

Improving the STEM Spatial Resolution Limit

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Professor Hashimoto pioneered several new electron microscope capabilities and applied them to many key problems. His dedication to making major contributions in every field he entered was a great inspiration, and it is very appropriate to dedicate this short paper on STEM resolution in his memory.

Scanning transmission electron microscopy (STEM) has progressed very significantly in recent years. Improving the spatial resolution by aberration correction has led to 1 Å (100 pm) spatial resolution becoming routine at primary energies as low as 60 keV in cold field emission (CFEG) STEM [1-3], and probe currents of ~1 nA have become available in sub-2 Å probes at higher keVs. Another important development was the introduction of ultra-high energy resolution monochromators and spectrometers, which have made 5 meV energy resolution electron energy loss spectroscopy (EELS) possible [4].

The quest for improved spatial resolution over the last two decades is illustrated in Fig. 1. The resolution and the probe size are closely related, and we prefer to focus on the probe size, which can be defined without needing a test object with suitable spacings. The probe size depends on the electron wavelength λ and therefore the primary energy and also on the source size contribution, which is proportional to $(I_p/I_c)^{0.5}$, where I_p is the probe current and I_c the coherent current (=current in a probe that's 50% coherent), and the contribution is added in quadrature to the diffraction-limited probe size [3]. The orange curves in Fig. 1 show the probe size for uncorrected STEM computed using formulas described in [3], for $I_p/I_c = 0.1$ (solid curve) and $I_p/I_c = 0.5$ (dashed curve), i.e. for about 15 and 75 pA probe current in the CFEG case. They show that for $C_s = 0.5$ mm, probe size of 135 pm was attainable at 200 keV in uncorrected STEM, with a severely limited probe current. At 100 keV with $C_s = 0.7$ mm (VG STEM with ultra-high resolution OL polepiece), the minimum probe size was about 200 pm.

STEM aberration correctors initially improved this performance about 2x, with 100 pm probe sizes becoming possible at 100-120 keV [1], and <78 pm probe sizes at 300 keV [5]. The first successful STEM correctors were limited by 5th order 4-fold astigmatism $C_{5,4}$, illustrated by the teal curves in Fig. 1. Next came the correction of all 5th order aberrations and lower-order parasitic aberrations [6]. The probe size then became limited by chromatic aberration (red curves), with a stronger dependence on the primary energy than geometric aberrations, making the C_c limit especially severe at lower keVs.

Overcoming the chromatic aberration limit can be achieved by eliminating C_c with a chromatic aberration corrector, or by reducing the energy spread δE substantially with a monochromator. The second option was tried in 2012 [7], without convincing results. We have chosen this option, which reduces the available beam current, but has the major advantage that it leads to higher energy resolution in EELS, unlocking new experimental capabilities [4]. At 30 keV, the non-monochromated probe size in our microscope system is about 165 pm, and reducing δE from 350 to 100 meV lowers this limit 1.8x. Capturing the 107 pm spacing in graphene then becomes possible at 30 kV, provided that geometric

aberrations are controlled well enough for 50 mrad illumination half-angle. This performance is demonstrated by the graphene image and FFT in Fig. 2, and marked by the blue diamond in Fig. 1. The next fundamental limit for quadrupole-octupole correctors is then posed by 8-fold astigmatism $C_{7,8}$, indicated by the blue curves in Fig. 1.

Recording a 107 pm spacing at 30 keV ($\lambda = 6.98$ pm) demonstrates spatial resolution of 15λ . This is, as far as we know, the best resolution performance for any transmission electron microscope. It has also been reached by a C_c -corrected fixed-beam TEM [7] and by a low C_c STEM [8], but with lower contrast levels. Progressing further with the monochromation approach will become possible if brighter electron sources, which would also be very useful for monochromated EELS, are developed. If seventh-order aberrations are also overcome, STEM will become limited by 9th order aberrations such as $C_{9,10}$, whose influence is estimated by the green curves in Fig. 1. Spatial resolution may then progress to $\sim 7 \lambda$, which would make sub-100 pm resolution possible at 10 keV primary energy, and < 30 pm at 100 keV.

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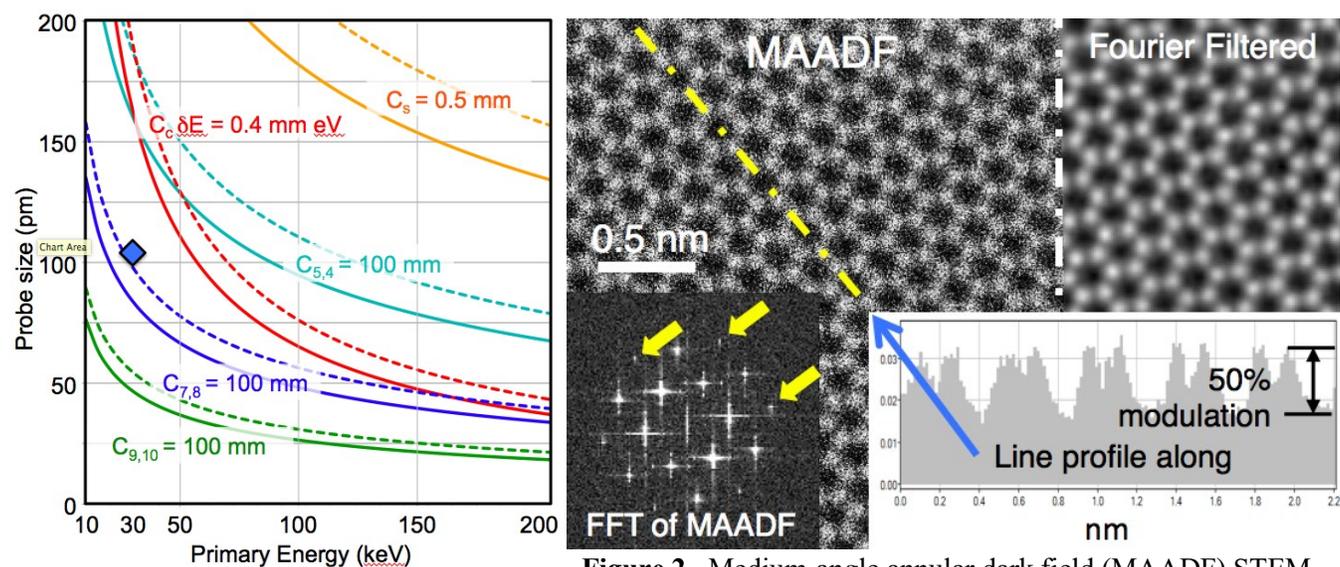


Figure 1. Theoretical STEM probe sizes for different types of limiting aberrations. Solid curves correspond to $I_p/I_c = 0.1$, dashed curves to $I_p/I_c = 0.5$. The blue diamond shows the performance demonstrated in Fig. 2.

Figure 2. Medium-angle annular dark field (MAADF) STEM image of graphene obtained with Nion U-HERMES™ 100 at 30 keV, with FFT and a line profile taken along the indicated direction. C_3 and C_5 -corrected, energy spread $\delta E \sim 100$ meV, $C_c = 0.96$ mm. Arrows in the FFT point to strongly transferred [20-20] 107 pm graphene spacings.