

CERTAIN BILATERAL HYPERGEOMETRIC IDENTITIES OF CAYLEY AND ORR TYPE

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1. Recently I (1) gave some new basic hypergeometric identities of the Cayley and Orr type with the help of a certain basic differential operator. The present paper deals with some bilateral generalizations of those identities together with certain new identities of the same type. In § 4 is indicated how the generalizations of Orr's identities given recently by Shukla (8, Theorems I, II) may be connected with each other. Later in § 5 certain general expansions of hypergeometric functions are deduced. The following notation has been used throughout this paper:

$$\begin{aligned} (q^a; n) &= (a; n) = (1 - q^a)(1 - q^{a+1}) \dots (1 - q^{a+n-1}), \quad (a; 0) = 1 \\ (q^a; -n) &= (a; -n) = (-)^n q^{\frac{1}{2}n(n+1)}/q^{na} (q^{1-a}; n), \quad |q| < 1 \\ (a)_{-n} &= (-)^n / (1 - a)_n, \\ {}_r\Psi_r \left(\begin{matrix} a_1, a_2, \dots, a_r \\ b_1, b_2, \dots, b_r \end{matrix}; x \right) &= \sum_{n=-\infty}^{\infty} \frac{(a_1; n)(a_2; n) \dots (a_r; n) x^n}{(b_1; n)(b_2; n) \dots (b_r; n)}, \end{aligned}$$

where, for convergence, $|x| < 1$, $|b_1 b_2 \dots b_r| < |a_1 a_2 \dots a_r x| < 1$.

$$\begin{aligned} {}_rH_r \left(\begin{matrix} a_1, a_2, \dots, a_r \\ b_1, b_2, \dots, b_r \end{matrix}; x \right) &= \sum_{n=-\infty}^{\infty} \frac{(a_1)_n (a_2)_n \dots (a_r)_n x^n}{(b_1)_n (b_2)_n \dots (b_r)_n}, \quad |x| = 1 \\ {}_r\Phi_s \left(\begin{matrix} a_1, a_2, \dots, a_r \\ b_1, b_2, \dots, b_s \end{matrix}; x \right) &= \sum_{n=0}^{\infty} \frac{(a_1; n)(a_2; n) \dots (a_r; n)}{(1; n)(b_1; n)(b_2; n) \dots (b_s; n)} x^n, \quad |x| < 1 \\ {}_rF_s \left(\begin{matrix} a_1, a_2, \dots, a_r \\ b_1, b_2, \dots, b_s \end{matrix}; x \right) &= \sum_{n=0}^{\infty} \frac{(a_1)_n (a_2)_n \dots (a_r)_n x^n}{n! (b_1)_n (b_2)_n \dots (b_s)_n}, \quad |x| < 1 \\ \Gamma \left(\begin{matrix} a_1, a_2, \dots, a_m \\ b_1, b_2, \dots, b_m \end{matrix}; \right) &= \frac{\Gamma(a_1)\Gamma(a_2) \dots \Gamma(a_m)}{\Gamma(b_1)\Gamma(b_2) \dots \Gamma(b_m)}, \\ \Gamma_m \left(\begin{matrix} a, b \\ c, d \end{matrix}; \right) &= \frac{(a)_m (b)_m}{(c)_m (d)_m}, \\ (\delta; n) &= (1 - q^\delta)(1 - q^{\delta+1}) \dots (1 - q^{\delta+n-1}), \end{aligned}$$

where $\delta \equiv x \partial/\partial x$, and

$$\Delta_q(h) = \frac{\Gamma_q(q^{\delta+h})}{\Gamma_q(q^h)} = \frac{\Gamma_q(\delta + h)}{\Gamma_q(h)}$$

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where $\Gamma_q(x)$ is a basic gamma function defined by Jackson (7) and

$$\Delta(h) = \frac{\Gamma(\delta + h)}{\Gamma(h)}.$$

2. Three bilateral hypergeometric identities. We now proceed to prove the following three identities that:

$$(2.1) \quad (c; m)(d - b; m)(a + c - 2m; m) \Delta_q(c + m) \\ \times {}_1\Psi_1\left(\begin{matrix} d - b + m; x \\ 1 + m \end{matrix}\right) {}_2\Psi_2\left(\begin{matrix} a + c - m, b; xq^{-a-b+2m} \\ 1 + m, c \end{matrix}\right)$$

= the same expression with c and d interchanged.

$$(2.2) \quad (f - b; m)(d; m)(1 - a; m)(e - c; m) \Delta_q(d + m) \\ \times {}_1\Psi_1\left(\begin{matrix} f - b + m; x \\ 1 + m \end{matrix}\right) {}_3\Psi_3\left(\begin{matrix} a, b, e - c + m; xq^{f-b} \\ 1 + m, d, e \end{matrix}\right) \\ = (d - b; m)(f; m)(1 - c; m)(e - a; m) \Delta_q(f + m) \\ \times {}_1\Psi_1\left(\begin{matrix} d - b + m; x \\ 1 + m \end{matrix}\right) {}_3\Psi_3\left(\begin{matrix} c, b, e - a + m; xq^{d-b} \\ 1 + m, e, f \end{matrix}\right)$$

provided $a + f = c + d$.

$$(2.3) \quad \frac{(c; m)(c'; m)}{(c - a; m)(c' - b'; m)} \Delta_q(c + m) \Delta_q(c' + m) \\ \times {}_2\Psi_2\left(\begin{matrix} a, b; xq^{c-a-b+m} \\ 1 + m, c \end{matrix}\right) {}_2\Psi_2\left(\begin{matrix} a', b'; x \\ 1 + m, c' \end{matrix}\right) \\ = \frac{(1 - a - a'; m)(1 - b - b'; m)}{(1 - a'; m)(1 - b; m)} q^{(a+b')m} \Delta_q(a + a') \Delta_q(b + b') \\ \times {}_2\Psi_2\left(\begin{matrix} b', c - a + m; x \\ 1 + m, b + b' - m \end{matrix}\right) {}_2\Psi_2\left(\begin{matrix} c' - b' + m, a; xq^{a'+b'-c'-m} \\ 1 + m, a + a' - m \end{matrix}\right)$$

provided $a + a' + b + b' = c + c' + 2m$.

The identities (2.1, 2.2, and 2.3) are generalizations of certain earlier results (1, 2.1, 2.2, 2.3).

Proof of 2.1. Consider the known identity (1, 2.1), namely, that

$$(2.4) \quad \Delta_q(c) {}_1\Phi_0(d - b; x) {}_2\Phi_1\left(\begin{matrix} a + c, b; xq^{-a-b} \\ c \end{matrix}\right)$$

equals the same expression with c and d interchanged.

Comparing the coefficients of x^n on both the sides we get the transformation

$$(2.5) \quad (c; n)(d - b; n) {}_3\Phi_2\left(\begin{matrix} a + c, b, -n; q^{1-a-d} \\ c, 1 + b - d - n \end{matrix}\right) \\ = (d; n)(c - b; n) {}_3\Phi_2\left(\begin{matrix} a + d, b, -n; q^{1-a-c} \\ d, 1 + b - c - n \end{matrix}\right).$$

Replacing $n, a, b, c,$ and d respectively by $2m + n, a - m, b - m, c - m,$ $d - m$ in (2.5) we find that

$$(2.6) \quad (c; m + n)(d - b; m + n)(a + c - 2m; m) \\ \times {}_3\Psi_3\left(\begin{matrix} a + c - m, b, -m - n; q^{1-a-d+2m} \\ 1 + m, c, 1 + b - d - m - n \end{matrix}\right)$$

equals the same expression with c and d interchanged.* Hence on comparing the coefficients of x^n and using the relation (2.6) we find that

$$(c; m)(d - b; m)(a + c - 2m; m) \Delta_q(c + m) \\ \times {}_1\Psi_1\left(\begin{matrix} d - b + m; x \\ 1 + m \end{matrix}\right) {}_2\Psi_2\left(\begin{matrix} a + c - m, b; x q^{-a-b+2m} \\ 1 + m, c \end{matrix}\right)$$

equals the same expression with c and d interchanged, which proves (2.1). Putting $m = 0$ (2.1) gives (2.4).

Proceeding exactly in the above manner we can prove the identities (2.2) and (2.3). Putting $m = 0$ in (2.2) and (2.3) we get back to the known identities due to Agarwal (1, 2.2, and 2.3).

3. Certain new identities and their generalizations. In this section we prove three new identities and later deduce their bilateral generalizations. The identities are

$$(3.1) \quad \Delta(1 + a - c) {}_1F_0(e - c; x) {}_3F_2\left(\begin{matrix} a, b, c; x \\ 1 + a - b, 1 + a - c \end{matrix}\right) \\ = \Delta(e) {}_1F_0(1 + a - 2c; x) {}_4F_3\left(\begin{matrix} a - 2b, 1 + \frac{1}{2}a - b, -b, c; x \\ \frac{1}{2}a - b, 1 + a - b, e \end{matrix}\right)$$

and

$$(3.2) \quad \Delta(1 + a - c) {}_1F_0(e - c; x) {}_4F_3\left(\begin{matrix} a, \frac{1}{2}a + 1, b, c; x \\ \frac{1}{2}a, 1 + a - b, 1 + a - c \end{matrix}\right) \\ = \Delta(e) {}_1F_0(1 + a - 2c; x) {}_3F_2\left(\begin{matrix} a - 2b, -b, c; x \\ 1 + a - b, e \end{matrix}\right),$$

*The limiting case as $q \rightarrow 1$ of (2.6) can be obtained directly by putting $c = m + n + 1$ and replacing $2 - e, 2 - f, b, d,$ and a respectively by $c, d, 1 - b, 1 - m,$ and $1 + a - 2m$ in a known result due to M. Jackson (6, p. 34).

provided $1 + a - c - e = 2b$ in (3.1) and (3.2), and

$$(3.3) \quad \Delta(1 + a - c) {}_1F_0(e - c; x) {}_4F_3\left(a, \frac{1}{2}a + 1, b, c; x\right. \\ \left. \frac{1}{2}a, 1 + a - b, 1 + a - c\right) \\ = \Delta(e) {}_1F_0(1 + a - 2c; x) {}_4F_3\left(a - 2b - 1, \frac{1}{2}a + \frac{1}{2} - b, -b - 1, c; x\right. \\ \left. \frac{1}{2}a - \frac{1}{2} - b, 1 + a - b, e\right)$$

provided $a - c - e = 2b$.

Proof of (3.1). It is easy to see that for suitably restricted parameters we have (4)

$$\frac{\Gamma(\delta + c)\Gamma(d - c)}{\Gamma(\delta + d)} f(x) = \int_0^1 u^{c-1}(1 - u)^{d-c-1} f(xu) du.$$

Let us replace c and d respectively by $1 + a - c$ and e and take

$$f(x) = {}_1F_0(e - c; x) {}_3F_2\left(a, b, c; x\right. \\ \left. 1 + a - b, 1 + a - c\right).$$

Then the right-hand side becomes

$$\int_0^1 u^{a-c}(1 - u)^{e+c-a-2}(1 - ux)^{c-e} {}_3F_2\left(a, b, c; xu\right. \\ \left. 1 + a - b, 1 + a - c\right) du.$$

Expanding the ${}_3F_2$ series and interchanging the order of integration and summation (which is easily justifiable), we have

$$\sum_{r=0}^{\infty} \frac{(a)_r(b)_r(c)_r}{r!(1 + a - b)_r(1 + a - c)_r} x^r \int_0^1 u^{a+r-c}(1 - u)^{e+c-a-2}(1 - ux)^{c-e} du \\ = \sum_{r=0}^{\infty} \frac{(a)_r(b)_r(c)_r x^r}{r!(1 + a - b)_r(1 + a - c)_r} \cdot \frac{\Gamma(1 + a - c + r)\Gamma(e + c - a - 1)}{\Gamma(e + r)} \\ {}_2F_1\left(\begin{matrix} e - c, 1 + a - c + r; \\ e + r \end{matrix} x\right).$$

Using Euler's identity (Tract 1, 1.2, 2.) we get

$$(3.4) \quad \frac{\Gamma(1 + a - c)\Gamma(e + c - a - 1)}{\Gamma(e)} {}_1F_0(1 + a - 2c; x) \\ \sum_{r=0}^{\infty} \frac{(a)_r(b)_r(c)_r}{r!(1 + a - b)_r(e)_r} x^r {}_2F_1\left(\begin{matrix} c + r, e + c - a - 1; \\ e + r \end{matrix} x\right).$$

Now, we have the transformation

$${}_4F_3\left(a - 2b, \frac{1}{2}a + 1 - b, -b, c; x\right. \\ \left. \frac{1}{2}a - b, 1 + a - b, e\right) \\ = \sum_{r=0}^{\infty} \frac{(a)_r(b)_r(c)_r}{r!(1 + a - b)_r(e)_r} x^r {}_2F_1\left(\begin{matrix} c + r, e + c - a - 1; \\ e + r \end{matrix} x\right)$$

provided $1 + a - c - e = 2b$, which can easily be obtained by collecting the coefficients of x^n and using the known summation theorem (2, §§ 4.5, 1.2). Hence using this transformation in (3.4) we have the required identity (3.1).

To prove (3.2) and (3.3) we proceed exactly as above with

$$f(x) = {}_1F_0(e - c; x) {}_4F_3\left(a, \frac{1}{2}a + 1, b, c; x; \frac{1}{2}a, 1 + a - b, 1 + a - c\right),$$

and use the transformations (2, §§ 4.5, 1.3, 1.4) giving the sum of a nearly-poised ${}_4F_3$.

We can also obtain the basic analogue of the identity (3.3) in the form

$$(3.5) \quad \Delta_q(1 + a - c) {}_1\Phi_0(q^{e-c}; x) {}_5\Phi_4\left(\begin{matrix} q^a, q^{\frac{1}{2}a+1}, -q^{\frac{1}{2}a+1}, q^b, q^c; xq^{a-2b-2c} \\ q^{\frac{1}{2}a}, -q^{\frac{1}{2}a}, q^{1+a-b}, q^{1+a-c} \end{matrix}\right) \\ = \Delta_q(e) {}_1\Phi_0(q^{1+a-2c}; x) {}_5\Phi_4\left(\begin{matrix} q^{a-2b-1}, q^{\frac{1}{2}a+\frac{1}{2}-b}, -q^{\frac{1}{2}a+\frac{1}{2}-b}, q^{-b-1}, q^c; xq^{1+a-2c} \\ q^{\frac{1}{2}a-\frac{1}{2}-b}, -q^{\frac{1}{2}a-\frac{1}{2}-b}, q^{1+a-b}, q^e \end{matrix}\right)$$

provided $a - c - e = 2b$.

To prove (3.5) we use the basic integral

$$(3.6) \quad \frac{\Gamma_q(\delta + c)\Gamma_q(d - c)}{\Gamma_q(\delta + d)} \Phi(x) = \int_0^1 u^{c-1}(1 - uq)^{d-c-1} \Phi(xu) d(qu)$$

used in an earlier paper as well (1), where $(1 - q^a x)^{-a}$ means the basic binomial expansion

$$1 + \frac{(1 - q^a)}{(1 - q)} x + \frac{(1 - q^a)(1 - q^{a+1})}{(1 - q)(1 - q^2)} x^2 + \dots,$$

or ${}_1\Phi_0(a; x)$.

Replace c and d respectively by $1 + a - c$ and e in (3.6) and take

$$\Phi(x) = {}_1\Phi_0(q^{e-c}; x) {}_5\Phi_4\left(\begin{matrix} q^a, q^{\frac{1}{2}a+1}, -q^{\frac{1}{2}a+1}, q^b, q^c; xq^{a-2b-2c} \\ q^{\frac{1}{2}a}, -q^{\frac{1}{2}a}, q^{1+a-b}, q^{1+a-c} \end{matrix}\right).$$

Proceeding as for (3.1) we obtain, on using a known summation theorem due to Bailey (3, § 3 (3)), the required identity (3.5).

It may be noted that as a consequence of these identities we get certain interesting relations between two terminating nearly-poised series. From (3.1), (3.2), and (3.3) respectively we get

$$(3.7) \quad {}_4F_3\left(a, b, c, -n; 1 + a - b, 1 + a - c, 1 + c - e - n\right) \\ = \frac{(1 + a - 2c)_n(e)_n}{(1 + a - c)_n(e - c)_n} {}_5F_4\left(a - 2b, 1 + \frac{1}{2}a - b, -b, c, -n; \frac{1}{2}a - b, 1 + a - b, e, 2c - a - n\right)$$

provided $1 + a - c - e = 2b$,

$$(3.8) \quad {}_5F_4\left(a, \frac{1}{2}a + 1, b, c, -n; \frac{1}{2}a, 1 + a - b, 1 + a - c, 1 + c - e - n\right) \\ = \frac{(1 + a - 2c)_n(e)_n}{(1 + a - c)_n(e - c)_n} {}_4F_3\left(a - 2b, -b, c, -n; 1 + a - b, e, 2c - a - n\right)$$

provided $1 + a - c - e = 2b$, and

$$(3.9) \quad {}_5F_4\left(a, \frac{1}{2}a + 1, b, c, -n; \frac{1}{2}a, 1 + a - b, 1 + a - c, 1 + c - e - n\right) \\ = \frac{(1 + a - 2c)_n (e)_n}{(1 + a - c)_n (e - c)_n} {}_5F_4\left(a - 2b - 1, \frac{1}{2}a + \frac{1}{2} - b, -b - 1, c, -n; \frac{1}{2}a - \frac{1}{2} - b, 1 + a - b, e, 2c - a - n\right)$$

provided $a - c - e = 2b$.

The basic analogue of (3.9) may be written as (from 3.5)

$$(3.10) \quad {}_6\Phi_5\left(\begin{matrix} q^a, q^{\frac{1}{2}a+1}, -q^{\frac{1}{2}a+1}, q^b, q^c, q^{-n}; q \\ q^{\frac{1}{2}a}, -q^{\frac{1}{2}a}, q^{1+a-b}, q^{1+a-c}, q^{1+c-e-n} \end{matrix}\right) \\ = \frac{(q^{1+a-2c}; n)(q^e; n)}{(q^{1+a-c}; n)(q^{e-c}; n)} {}_6\Phi_5\left(\begin{matrix} q^{a-2b-1}, q^{\frac{1}{2}a+\frac{1}{2}-b}, -q^{\frac{1}{2}a+\frac{1}{2}-b}, q^{-b-1}, q^c, q^{-n}; q \\ q^{\frac{1}{2}a-\frac{1}{2}-b}, -q^{\frac{1}{2}a-\frac{1}{2}-b}, q^{1+a-b}, q^e, q^{2c-a-n} \end{matrix}\right)$$

provided $a - c - e = 2b$.

Next we deduce the bilateral generalizations of the identities (3.1), (3.2), and (3.3). In the known transformation due to Shukla (8, 2.2) let us take $M = 3, N = 1, a_1 = b_1 = 1$ and $c_4 = 0$. This gives us a relation between three non-terminating nearly-poised ${}_4H_4$ series and a terminating ${}_4F_3$ series viz.:

$$(3.11) \quad (1 - E)(1 - F) \\ \Gamma\left(\begin{matrix} 1 + E - D, 1 + F - D, D - E, D - F; \\ 2 + a - b - D, 2 + a - c - D, 2 + c - e - D, D - a, D - b, D - c \end{matrix}\right) \\ \times \frac{(D + e - c - 1)_n}{(D)_n} \\ {}_4H_4\left(\begin{matrix} 1 + a - D, 1 + b - D, 1 + c - D, 1 - D - n; \\ 2 - D, 2 + a - b - D, 2 + a - c - D, 2 + c - e - D - n \end{matrix}\right) \\ + \text{idem } (D; E, F) \\ = \Gamma\left(\begin{matrix} D, E, F, 2 - D, 2 - E, 2 - F; \\ 1 + a - b, 1 + a - c, 1 + c - e, 1 - a, 1 - b, 1 - c \end{matrix}\right) \\ \times \frac{(e - c)_n}{(1)_n} {}_4F_3\left(a, b, c, -n; 1 + a - b, 1 + a - c, 1 + c - e - n\right).$$

Transform the ${}_4F_3$ series on the right by (3.7) and then replace n, a, b, c , and e respectively by $2m + n, a - 2m, b - m, c - m$, and $e - m$. After some simplification we find that (3.11) may be written as

$$(3.12) \quad (1 - E)(1 - F) \\ \Gamma\left(\begin{matrix} 1 + E - D, 1 + F - D, D - E, D - F; \\ 2 + a - b - D - m, 2 + a - c - D - m, 2 + c - e - D, \\ D - a + 2m, D - b + m, D - c + m \end{matrix}\right) \\ \times \frac{(D + e - c - 1)_{2m}}{(D)_{2m}} {}_1H_1\left(\begin{matrix} D + e - c - 1 + 2m; x \\ D + 2m \end{matrix}\right) \times$$

$$\begin{aligned}
 & {}_3H_3\left(1+a-D-2m, 1+b-D-m, 1+c-D-m; x\right) \\
 & + \text{idem } (D; E, F) \\
 = & \Gamma\left(\begin{matrix} D, E, F, 2-D, 2-E, 2-F; \\ 1+a-b-m, 1+a-c-m, 1+c-e, 1-a+2m, \\ 1-b+m, 1-c+m \end{matrix}\right) \\
 & \times \Gamma_m\left(\begin{matrix} e, a-2b, \frac{1}{2}a+1-b, 1-c, -b+m, 1+a-2c; \\ 1, \frac{1}{2}a-b, b-a, c-a, 1+a-c \end{matrix}\right) \\
 & \times \frac{\Delta(e+m)}{\Delta(1+a-c+m)} \\
 & \times {}_4H_4\left(\begin{matrix} a-2b+m, \frac{1}{2}a+1-b+m, -b+2m, c; x \\ 1+m, \frac{1}{2}a-b+m, 1+a-b, e \end{matrix}\right) \\
 & {}_1H_1\left(\begin{matrix} 1+a-2c+m; x \\ 1+m \end{matrix}\right),
 \end{aligned}$$

provided $1+a-c-e = 2b-2m$. (3.12) reduces to (3.1) when $D = 1$ and $m = 0$.

In exactly the same manner bilateral generalizations of the identities (3.2) and (3.3) may be written in the form

$$\begin{aligned}
 (3.13) \quad & (1-E)(1-F)(1-G) \\
 & \Gamma\left[\begin{matrix} 1+E-D, 1+F-D, 1+G-D, D-E, D-F, \\ 1+\frac{1}{2}a-D-m, 2+a-b-D-m, 2+a-c-D-m, 2+c-e-D, \\ D-G; \\ D-a+2m, D-\frac{1}{2}a-1+m, D-b+m, D-c+m \end{matrix}\right] \\
 & \times \frac{(D+e-c-1)_{2m}}{(D)_{2m}} {}_1H_1\left(\begin{matrix} D+e-c-1+2m; x \\ D+2m \end{matrix}\right) \\
 & {}_4H_4\left(\begin{matrix} 1+a-D-2m, 2+\frac{1}{2}a-D-m, 1+b-D-m, 1+c-D-m; x \\ 2-D, 1+\frac{1}{2}a-D-m, 2+a-b-D-m, 2+a-c-D-m \end{matrix}\right) \\
 & + \text{idem } (D; E, F, G) \\
 = & \Gamma\left(\begin{matrix} D, E, F, G, 2-D, 2-E, 2-F, 2-G; \\ 1+a-b-m, 1+a-c-m, 1+c-e, \frac{1}{2}a-m, 1-a+2m, \\ -\frac{1}{2}a+m, 1-b+m, 1-c+m \end{matrix}\right) \\
 & \times \Gamma_m\left(\begin{matrix} e, a-2b, -b+m, 1-c, 1+a-2c \\ 1, b-a, c-a, 1+a-c \end{matrix}\right) \frac{\Delta(e+m)}{\Delta(1+a-c+m)} \\
 & \times {}_1H_1\left(\begin{matrix} 1+a-2c+m; x \\ 1+m \end{matrix}\right) {}_3H_3\left(\begin{matrix} a-2b+m, -b+2m, c; x \\ 1+m, 1+a-b, e \end{matrix}\right)
 \end{aligned}$$

provided $1 + a - c - e = 2b - 2m$,

(3.14)

$$\begin{aligned} &= \Gamma \left(\begin{matrix} D, E, F, G, 2 - D, 2 - E, 2 - F, 2 - G; \\ 1 + a - b - m, 1 + a - c - m, 1 + c - e, \frac{1}{2}a - m, 1 - a + 2m, \\ \qquad \qquad \qquad -\frac{1}{2}a + m, 1 - b + m, 1 - c + m \end{matrix} \right) \\ &\times \Gamma_m \left(\begin{matrix} e, 1 + a - 2c, a - 2b - 1, -b - 1 + m, 1 - c, \frac{1}{2}a + \frac{1}{2} - b \\ 1, b - a, c - a, 1 + a - c, \frac{1}{2}a - \frac{1}{2} - b \end{matrix} \right) \\ &\qquad \qquad \qquad \frac{\Delta(e + m)}{\Delta(1 + a - c + m)} \\ &\times {}_1H_1 \left(\begin{matrix} 1 + a - 2c + m; x \\ 1 + m \end{matrix} \right) \\ &\times {}_4H_4 \left(\begin{matrix} a - 2b + m - 1, \frac{1}{2}a + \frac{1}{2} - b + m, -b - 1 + 2m, c \\ 1 + m, \frac{1}{2}a - \frac{1}{2} - b + m, 1 + a - b, e \end{matrix} ; x \right) \end{aligned}$$

provided $a - c - e = 2b - 2m$. Putting $D = 1$ and $m = 0$ in (3.13) and (3.14) we get back to the identities (3.1) and (3.2).

4. Next we show how the bilateral generalizations due to Shukla (8, Theorems I, II) of Orr’s identities (Tract, § 10.1 (2 and 3)) may be deduced from each other by the use of the following identity:

(4.1)

$$\begin{aligned} &(1 - a)_m(1 - b)_m(1 - a')_m(1 - b')_m {}_2H_2 \left(\begin{matrix} a, b \\ 1 + m, c \end{matrix} ; x \right) {}_2H_2 \left(\begin{matrix} a', b' \\ 1 + m, c' \end{matrix} ; x \right) \\ &= (c - a)_m(c - b)_m(c' - a')_m(c' - b')_m \\ &{}_2H_2 \left(\begin{matrix} c - a + m, c - b + m \\ 1 + m, c \end{matrix} ; x \right) {}_2H_2 \left(\begin{matrix} c' - a' + m, c' - b' + m \\ 1 + m, c' \end{matrix} ; x \right) \end{aligned}$$

with $a + a' + b + b' = c + c' + 2m$.

The identity (4.1) may be obtained from the known identity due to Chaundy (4, 25).

The identities due to Shukla can also be written in the form

$$\begin{aligned} (4.2) \quad &(1 - E) \Gamma \left(\begin{matrix} 1 + E - D, D - E; \\ 1 + 2c - D - 2m, 3/2 + a + b - c - D - m, \\ \qquad \qquad \qquad D - 2b + 2m, D - 2a + 2m \end{matrix} \right) \\ &\times \frac{(c + D - a - b + \frac{1}{2} + m)_{2m}}{(D)_{2m}} {}_1H_1 \left(\begin{matrix} c + D - a - b - \frac{1}{2} + 3m; x \\ D + 2m \end{matrix} \right) \\ &{}_2H_2 \left(\begin{matrix} 1 + 2a - D - 2m, 1 + 2b - D - 2m; x \\ 2 - D, 1 + 2c - D - 2m \end{matrix} \right) + \text{idem } (D; E) \end{aligned}$$

$$\begin{aligned}
 &= \Gamma \left[\begin{matrix} D, E, 2 - D, 2 - E; \\ 2c - 2m, \frac{1}{2} + a + b - c - m, 1 - 2a + 2m, 1 - 2b + 2m \end{matrix} \right] \\
 &\Gamma_m \left[\begin{matrix} \frac{1}{2} + c - a, \frac{1}{2} + c - b, 1 + c, 1 - a, 1 - b; \\ 1, 1, \frac{1}{2} + c, \frac{1}{2} + c, 1 - c, \frac{1}{2} - c \end{matrix} \right] \\
 &\frac{\Delta(c + 1 + m)}{\Delta(c + \frac{1}{2} + m)} {}_2H_2 \left(\begin{matrix} a, b \\ 1 + m, c \end{matrix}; x \right) {}_2H_2 \left(\begin{matrix} c - a + \frac{1}{2} + m, c - b + \frac{1}{2} + m; \\ 1 + m, c + 1 \end{matrix}; x \right)
 \end{aligned}$$

and

$$\begin{aligned}
 (4.3) \quad &(1 - E) \Gamma \left(\begin{matrix} 1 + E - D, D - E; \\ 2c - D - 2m, 3/2 + a + b - c - D - m, \\ 1 + D - 2a + 2m, D - 2b + 2m \end{matrix} \right) \\
 &\times \frac{(c + D - a - b - \frac{1}{2} + m)_{2m}}{(D)_{2m}} {}_1H_1 \left(\begin{matrix} c + D - a - b - \frac{1}{2} + 3m; \\ D + 2m \end{matrix}; x \right) \\
 &{}_2H_2 \left(\begin{matrix} 2a - 2m - D, 2b + 1 - D - 2m; \\ 2 - D, 2c - 2m - D \end{matrix}; x \right) + \text{idem } (D; E) \\
 &= \Gamma \left(\begin{matrix} D, E, 2 - D, 2 - E; \\ 2c - 2m - 1, \frac{1}{2} + a + b - c - m, 2 - 2a + 2m, 1 - 2b + 2m \end{matrix} \right) \\
 &\Gamma_m \left(\begin{matrix} \frac{1}{2} + c - a, c - \frac{1}{2} - b, 1 - a, 1 - b; \\ 1, 1, 3/2 - c, c - \frac{1}{2}, 1 - c \end{matrix} \right) \frac{\Delta(c + m)}{\Delta(c - \frac{1}{2} + m)} {}_2H_2 \left(\begin{matrix} a, b \\ 1 + m, c \end{matrix}; x \right) \\
 &{}_2H_2 \left(\begin{matrix} c - a + \frac{1}{2} + m, c - b - \frac{1}{2} + m; \\ 1 + m, c \end{matrix}; x \right).
 \end{aligned}$$

Now let $q \rightarrow 1$ in (2.3) and then use the identity (4.1); to transform the right-hand side put $a + a' = b + b' = c + \frac{1}{2} + m$, and $c' = c + 1$. This gives us the transformation

$$\begin{aligned}
 &\Gamma_m(c, c + 1, \frac{1}{2} + c - b, 1 - a) \frac{\Delta(c + m + 1)}{\Delta(c + m + \frac{1}{2})} \\
 &\times {}_2H_2 \left(\begin{matrix} a, b \\ 1 + m, c \end{matrix}; x \right) {}_2H_2 \left(\begin{matrix} c - a + \frac{1}{2} + m, c - b + \frac{1}{2} + m; \\ 1 + m, c + 1 \end{matrix}; x \right) \\
 &= \Gamma_m(\frac{1}{2} + c - a, \frac{1}{2} + c, \frac{1}{2} + c, \frac{1}{2} - a, c - b) \frac{\Delta(c + \frac{1}{2} + m)}{\Delta(c + m)} \\
 &\times {}_2H_2 \left(\begin{matrix} b, a + \frac{1}{2} \\ 1 + m, c + \frac{1}{2} \end{matrix}; x \right) {}_2H_2 \left(\begin{matrix} c - a + \frac{1}{2} + m, c - b + m; \\ 1 + m, c + \frac{1}{2} \end{matrix}; x \right).
 \end{aligned}$$

Using (4.2) on the left of the above equation we find that it reduces to (4.3) with $a + \frac{1}{2}, c + \frac{1}{2}$ for a and c . Thus (4.3) is connected with (4.2).

5. Certain general hypergeometric expansions. In this section we obtain certain general expansions both for ordinary and bilateral hyper-

geometric series with the help of transformations deduced in previous sections. Before proceeding to the actual deduction of the general expansions we prove a lemma which is a bilateral generalization of a known transformation due to \ast Chaundy (4, 42).

LEMMA. *If*

$$\Psi\left(\begin{matrix} a, \dots \\ 1 + m, c, \dots \end{matrix}; xq^\lambda\right) \quad \text{and} \quad \Psi\left(\begin{matrix} a', \dots \\ 1 + m, c', \dots \end{matrix}; x\right)$$

are two basic bilateral hypergeometric functions (of any order) and h, k , any two suitable constants then

(5.1)

$$\begin{aligned} & \frac{(h + m; m)(k; m) \Delta_q(h + 2m)}{(k + m; m)(h; m) \Delta_q(k + 2m)} \Psi\left[\begin{matrix} a, \dots \\ 1 + m, c, \dots \end{matrix}; xq^\lambda\right] \Psi\left[\begin{matrix} a', \dots \\ 1 + m, c', \dots \end{matrix}; x\right] \\ &= \sum_{r=0}^{\infty} \frac{(k - 1; r)(k; 2r)(k - h; r)[(h + m; r)]^2(a; r) \dots (a'; r) \dots}{(1; r)(h; r)(k - 1; 2r)[(k + m; 2r)]^2(c; r) \dots (c'; r) \dots} \\ & \times x^{2r} q^{\tau(r-1) + (\lambda+h)r} \Psi\left(\begin{matrix} h + m + r, a + r, \dots \\ k + m + 2r, 1 + m, c + r, \dots \end{matrix}; xq^\lambda\right) \\ & \qquad \qquad \qquad \Psi\left(\begin{matrix} h + m + r, a' + r, \dots \\ k + m + 2r, 1 + m, c' + r, \dots \end{matrix}; x\right). \end{aligned}$$

Proof: It is easily seen that

$$\frac{\Delta_q(h)}{\Delta_q(k)} \Psi\left(\begin{matrix} a, \dots \\ c, \dots \end{matrix}; x\right) = \Psi\left(\begin{matrix} h, a, \dots \\ k, c, \dots \end{matrix}; x\right)$$

and

$$\begin{aligned} & q^{(\delta+m)r} \frac{(-\delta - m; r) \Delta_q(h + m)}{(\delta + m + k; r) \Delta_q(k + m)} \Psi\left(\begin{matrix} a, \dots \\ 1 + m, c, \dots \end{matrix}; xq^\lambda\right) \\ &= \frac{(a; r) \dots (h + m; r)}{(c; r) \dots (k + m; 2r)} (-x)^r q^{\frac{1}{2}r(\tau-1) + \lambda r} \\ & \qquad \qquad \qquad \Psi\left(\begin{matrix} h + m + r, a + r, \dots \\ k + m + 2r, 1 + m, c + r, \dots \end{matrix}; xq^\lambda\right). \end{aligned}$$

Using the above transformation on the right-hand side of (5.1), it becomes

$$\begin{aligned} (5.2) \quad & \frac{\Gamma_q(\theta + h + m) \Gamma_q(\phi + h + m)}{\Gamma_q(\theta + k + m) \Gamma_q(\phi + k + m)} \left[\frac{\Gamma_q(k + m)}{\Gamma_q(h + m)} \right]^2 \\ & \times \sum_{r=0}^{\infty} \frac{(k - 1; r)(k; 2r)(k - h; r)(-\theta - m; r)(-\phi - m; r)}{(1; r)(k - 1; 2r)(h; r)(\theta + k + m; r)(\phi + k + m; r)} q^{(\theta + \phi + 2m + h)r} \\ & \times \Psi\left(\begin{matrix} a, \dots \\ 1 + m, c, \dots \end{matrix}; xq^\lambda\right) \Psi\left(\begin{matrix} a', \dots \\ 1 + m, c', \dots \end{matrix}; x\right), \end{aligned}$$

where $\theta \equiv x \partial/\partial x$ and $\phi \equiv x \partial/\partial x$ operate on the first and second series alone respectively and hence $q^{\theta+\phi} = q^{\theta}$. Summing up the above well-poised ${}_6\Phi_5$ series we get the required result (5.1). If we take the limit $q \rightarrow 1$ we get the corresponding transformation for ordinary bilateral hypergeometric series, viz.:

$$(5.3) \quad \frac{(h+m)_m (k)_m \Delta(h+2m)}{(k+m)_m (h)_m \Delta(k+2m)} H \left[\begin{matrix} a, \dots \\ 1+m, c, \dots \end{matrix}; x \right] H \left[\begin{matrix} a', \dots \\ 1+m, c', \dots \end{matrix}; x \right]$$

$$= \sum_{r=0}^{\infty} \frac{(k-1)_r (k)_{2r} (k-h)_r ((h+m)_r)^2 (a)_r \dots (a')_r \dots x^{2r}}{r! (k-1)_{2r} ((k+m)_{2r})^2 (h)_r (c)_r \dots (c')_r \dots} x^{2r}$$

$$\times H \left(\begin{matrix} h+m+r, a+r, \dots \\ 1+m, k+m+2r, c+r, \dots \end{matrix}; x \right) H \left(\begin{matrix} h+m+r, a'+r, \dots \\ 1+m, k+m+2r, c'+r, \dots \end{matrix}; x \right).$$

Applications. Applying the transformation (5.1) to the identity (2.1) we get

$$(5.4) \quad {}_1\Psi_1 \left(\begin{matrix} d-b+m \\ 1+m \end{matrix}; x \right) {}_2\Psi_2 \left(\begin{matrix} a+c-m, b \\ 1+m, c \end{matrix}; x q^{-a-b+2m} \right)$$

$$= \frac{(c-b; m)(a+d-2m; m)(d-m; m)}{(d-b; m)(a+c-2m; m)(c-m; m)}$$

$$\times \sum_{r=0}^{\infty} \frac{(c-m-1; r)(c-m; 2r)(c-d; r)(d; r)(c-b+m; r)(d+a-m; r)}{(1; r)(c-m-1; 2r)(c; 2r)(c; 2r)(d-m; r)}$$

$$\times (b; r) x^{2r} q^{r(r-1)+(d-a-b+m)r} {}_2\Psi_2 \left(\begin{matrix} d+r, c-b+m+r \\ 1+m, c+2r \end{matrix}; x \right)$$

$$\times {}_3\Psi_3 \left(\begin{matrix} d+r, a+d-m+r, b+r \\ 1+m, d+r, c+2r \end{matrix}; x q^{-a-b+2m} \right).$$

Similarly, applying the transformation (5.1) to (2.2) and (5.3) to (4.2) and (4.3) one can obtain three other expansions of similar type.

Next, applying the transformation (5.3) for $m = 0$ (4, 42) to the right-hand side of the identities (3.1), (3.2), and (3.3) respectively, we get the following three expansions

$$(5.5) \quad {}_3F_2 \left(\begin{matrix} a, b, c \\ 1+a-b, 1+a-c \end{matrix}; x \right)$$

$$= \sum_{r=0}^{\infty} \frac{(2b)_r (1+a-2c)_r (a-2b)_r (\frac{1}{2}a-b+1)_r (-b)_r (c)_r}{r! (a-c+r)_r (1+a-c)_{2r} (\frac{1}{2}a-b)_r (1+a-b)_r} x^{2r}$$

$$\times {}_2F_1 \left(\begin{matrix} c+r, 2b+r \\ 1+a-c+2r \end{matrix}; x \right)$$

$${}_4F_3 \left(\begin{matrix} a-2b+r, \frac{1}{2}a-b+1+r, -b+r, c+r \\ \frac{1}{2}a-b+r, 1+a-b+r, 1+a-c+2r \end{matrix}; x \right)$$

with the condition $1 + a - c - e = 2b$,

$$(5.6) \quad {}_4F_3\left(a, \frac{1}{2}a + 1, b, c; \frac{1}{2}a, 1 + a - b, 1 + a - c; x\right) \\ = \sum_{r=0}^{\infty} \frac{(2b)_r(1+a-2c)_r(a-2b)_r(-b)_r(c)_r}{r!(a-c+r)_r(1+a-c)_{2r}(1+a-b)_r} x^{2r} \\ \times {}_2F_1\left(c+r, 2b+r; 1+a-c+2r; x\right) {}_3F_2\left(a-2b+r, -b+r, c+r; 1+a-b+r, 1+a-c+2r; x\right)$$

with $1 + a - c - e = 2b$,

$$(5.7) \quad = \sum_{r=0}^{\infty} \frac{(2b-1)_r(1+a-2c)_r(a-2b-1)_r(\frac{1}{2}a+\frac{1}{2}-b)_r(-b-1)_r(c)_r}{r!(a-c+r)_r(1+a-c)_{2r}(\frac{1}{2}a-\frac{1}{2}-b)_r(1+a-b)_r} x^{2r} \\ \times {}_2F_1\left(c+r, 1+2b+r; 1+a-c+2r; x\right) \\ {}_4F_3\left(a-2b-1+r, \frac{1}{2}a+\frac{1}{2}-b+r, -b-1+r, c+r; \frac{1}{2}a-\frac{1}{2}-b+r, 1+a-c+2r, 1+a-b+r; x\right)$$

provided $a - c - e = 2b$.

The basic analogue of (5.7) may be written as below (by using (5.1) for $m = 0$ in (3.5))

$$(5.8) \quad {}_5\Phi_4\left(q^a, q^{\frac{1}{2}a+1}, -q^{\frac{1}{2}a+1}, q^b, q^c; q^{\frac{1}{2}a}, -q^{\frac{1}{2}a}, q^{1+a-b}, q^{1+a-c}; x q^{a-2b-2c}\right) \\ = \sum_{r=0}^{\infty} \frac{(q^{2b-1}; r)(q^{1+a-2c}; r)(q^{a-2b-1}; r)(q^{\frac{1}{2}a+\frac{1}{2}-b}; r)}{(q; r)(q^{a-c+r}; r)(q^{1+a-c}; 2r)(q^{\frac{1}{2}a-\frac{1}{2}-b}; r)} \\ \times \frac{(-q^{\frac{1}{2}a+\frac{1}{2}-b}; r)(q^{-b-1}; r)(q^c; r)}{(-q^{\frac{1}{2}a-\frac{1}{2}-b}; r)(q^{1+a-b}; r)} x^{2r} q^{r^2+(e+a-2c)r} \\ \times {}_2\Phi_1\left(q^{c+r}, q^{2b+1+r}; q^{1+a-c+2r}; x q^{e-c}\right) \\ \times {}_5\Phi_4\left(q^{a-2b-1+r}, q^{\frac{1}{2}a+\frac{1}{2}-b+r}, -q^{\frac{1}{2}a-b+r+\frac{1}{2}}, q^{-b-1+r}, q^{c+r}; q^{\frac{1}{2}a-\frac{1}{2}-b+r}, -q^{\frac{1}{2}a-\frac{1}{2}-b+r}, q^{1+a-b+r}, q^{1+a-c+2r}; x q^{1+a-2c}\right)$$

with $a - c - e = 2b$.

Similar expansions could also be obtained from results due to Shukla (8, vii) and Henrici (5, $a \sim b, a \sim c$).

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