# On Lagrangian Catenoids 

David E. Blair


#### Abstract

Recently I. Castro and F. Urbano introduced the Lagrangian catenoid. Topologically, it is $\mathbb{R} \times$ $S^{n-1}$ and its induced metric is conformally flat, but not cylindrical. Their result is that if a Lagrangian minimal submanifold in $\mathbb{C}^{n}$ is foliated by round $(n-1)$-spheres, it is congruent to a Lagrangian catenoid. Here we study the question of conformally flat, minimal, Lagrangian submanifolds in $\mathbb{C}^{n}$. The general problem is formidable, but we first show that such a submanifold resembles a Lagrangian catenoid in that its Schouten tensor has an eigenvalue of multiplicity one. Then, restricting to the case of at most two eigenvalues, we show that the submanifold is either flat and totally geodesic or is homothetic to (a piece of) the Lagrangian catenoid.


## 1 Introduction

In 1744 Euler showed that a catenoid is a minimal surface of revolution and in 1785 Meusnier proved the converse that the catenoid is the only surface of revolution that is minimal $[3,11]$.

For hypersurfaces of dimension $n \geq 4$ in Euclidean space, Cartan [2] showed that a conformally flat hypersurface is quasi-umbilical, i.e., the Weingarten map has an eigenvalue of multiplicity $\geq n-1$. Common examples, also due to Cartan, are the canal hypersurfaces, i.e., envelopes of one-parameter families of hyperspheres. Thus conformal flatness can be viewed as a natural generalization of a surface of revolution.

In [1] the author showed that a conformally flat minimal hypersurface $M^{n}, n \geq$ 4 of Euclidean space $\mathbb{R}^{n+1}$ is either totally geodesic or a hypersurface of revolution $S^{n-1} \times \gamma(s)$, where the profile curve $\gamma$ is determined by its curvature as a function of arc length by $\kappa=(1-n) / u^{n}$ and

$$
s=\int \frac{u^{n-1} d u}{\sqrt{C u^{2 n-2}-1}}
$$

where $C$ is a constant. For $n=3$, replacing conformal flatness by quasi-umbilicity gives the same result with the same proof. For $n=2$, the profile curve is a catenary and hence these hypersurfaces are called generalized catenoids. Jagy gave an independent study of this question by assuming that the minimal hypersurface is foliated by spheres from the outset [10].

Recently I. Castro and F. Urbano [4] (see also Castro [3]) introduced the Lagrangian catenoid. The manifold itself was introduced by Harvey and Lawson [9] as an example of a minimal Lagrangian submanifold and is defined by
$M_{0}=\left\{(x, y) \in \mathbb{R}^{n} \times \mathbb{R}^{n} \equiv \mathbb{C}^{n}:|x| y=|y| x, \Im(|x|+i|y|)^{n}=1,|y|<|x| \tan (\pi / n)\right\}$.

[^0]Topologically $M_{0}$ is $\mathbb{R} \times S^{n-1}$. To describe it precisely, let $S^{n-1}$ be the unit sphere in $\mathbb{R}^{n}$ and view a point $p \in S^{n-1}$ as an $n$-tuple in $\mathbb{R}^{n}$ giving its coordinates. Define a $\operatorname{map} \phi_{0}: \mathbb{R} \times S^{n-1} \rightarrow \mathbb{C}^{n} \equiv \mathbb{R}^{n} \times \mathbb{R}^{n}$ by

$$
\phi_{0}(u, p)=\cosh ^{1 / n}(n u) e^{i \beta(u)} p,
$$

where $\beta(u)=\frac{\pi}{2 n}-\frac{2}{n} \arctan \left(\tanh \frac{n u}{2}\right) \in\left(0, \frac{\pi}{n}\right)$ and the multiplication $e^{i \beta} p$ multiplies each coordinate of $p$ by $e^{i \beta}$ and lists the real and imaginary parts as a $2 n$-tuple in $\mathbb{C}^{n} \equiv \mathbb{R}^{n} \times \mathbb{R}^{n}$. Let $g_{0}$ be the standard metric of constant curvature 1 on $S^{n-1}$; then the metric induced on $\mathbb{R} \times S^{n-1}$ by $\phi_{0}$ is $d s^{2}=\cosh ^{2 / n}(n u)\left(d u^{2}+g_{0}\right)$, which is clearly conformally flat. This Lagrangian submanifold defined by the mapping $\phi_{0}: \mathbb{R} \times S^{n-1} \rightarrow \mathbb{C}^{n}$, together with its induced metric, is known as the Lagrangian catenoid. The second fundamental form is given in terms of the Weingarten maps as follows. Let $\left\{e_{1}, \ldots, e_{n}\right\}$ be a local orthonormal basis with $e_{1}$ in the $u$-direction and let $A_{i}$ denote the Weingarten map corresponding to the normal $J e_{i}$. Then the matrices of the Weingarten maps with respect to this basis are the following:

$$
A_{1}=\left(\begin{array}{cccc}
-(n-1) a & & & \bigcirc \\
& a & & \\
& & \ddots & \\
\bigcirc & & & a,
\end{array}\right) \quad A_{i}=\left(\begin{array}{ccccccc}
0 & \cdots & 0 & a & 0 & \cdots & 0 \\
\vdots & & & & & & \\
0 & & & & & & \\
a & & & \bigcirc & & & \\
0 & & & & & \\
\vdots & & & & & \\
0 & & & & &
\end{array}\right)
$$

where $a=\cosh ^{-\left(1+\frac{1}{n}\right)}(n u)$.
The main result of Castro and Urbano [4] is the following.
Theorem 1 Let $\phi: M^{n} \rightarrow \mathbb{C}^{n}$ be a minimal (non-flat) Lagrangian immersion. Then $M^{n}$ is foliated by pieces of round $(n-1)$-spheres in $\mathbb{C}^{n}$ if and only if $\phi$ is congruent (up to dilations) to an open subset of the Lagrangian catenoid.

In view of the conformal flatness of the Lagrangian catenoid and the author's result on conformally flat minimal hypersurfaces in $\mathbb{R}^{n+1}$, it is natural to ask if, aside from the flat totally geodesic case, the Lagrangian catenoids are the only conformally flat minimal Lagrangian submanifolds in $\mathbb{C}^{n}$. We first show in Proposition 4 that the Schouten tensor of such a submanifold has an eigenvalue of multiplicity 1. Then, noting that the Schouten tensor of the Lagrangian catenoid has only two eigenvalues, we restrict ourselves to the case that Schouten tensor, equivalently the Ricci tensor, has at most two eigenvalues. In particular we prove the following theorem.

Theorem 2 Let $M^{n}, n \geq 4$, be a conformally flat, minimal, Lagrangian submanifold of $\mathbb{C}^{n}$. If the Schouten tensor has at most two eigenvalues, then either $M^{n}$ is flat and totally geodesic or is homothetic to (a piece of) the Lagrangian catenoid.

We comment that our proof is not a reduction to the result of Castro and Urbano. Even when one is at the stage of having essentially a warped product of a 1-dimensional manifold and a space of constant curvature, it does not follow immediately that the spaces of constant curvature which now have codimension $n+1$ are round spheres in $\mathbb{C}^{n}$. Instead we establish both the first and second fundamental forms of $M^{n}$.

We close the paper with a remark showing that the Lagrangian catenoid is not a quasi-umbilical submanifold.

## 2 Preliminaries

Let $\left(x_{1}+i y_{1}, \ldots, x_{n}+i y_{n}\right)$ be the coordinates on $\mathbb{C}^{n}$ and $J$ the almost complex structure. An $n$-dimensional submanifold $M^{n}$ of $\mathbb{C}^{n}$ is said to be Lagrangian if the restriction of the canonical symplectic form $\Omega=\sum_{i=1}^{n} d x_{i} \wedge d y_{i}$ to $M^{n}$ vanishes. Note also that if a vector $X$ is tangent to $M^{n}, J X$ is normal.

For an isometrically immersed submanifold $(M, g)$ of $\left(\mathbb{C}^{n},\langle\rangle,\right)$ the Levi-Civita connection $\nabla$ of $g$ and the second fundamental form $\sigma$ are related to the ambient Levi-Civita connection $\bar{\nabla}$ by $\bar{\nabla}_{X} Y=\nabla_{X} Y+\sigma(X, Y)$. For a normal vector field $\zeta$, let $A_{\zeta}$ denote the corresponding Weingarten map and let $D$ denote the connection in the normal bundle; in particular $A_{\zeta}$ and $D$ are defined by $\bar{\nabla}_{X} \zeta=-A_{\zeta} X+D_{X} \zeta$. The Gauss equation is

$$
R(X, Y, Z, W)=\langle\sigma(Y, Z), \sigma(X, W)\rangle-\langle\sigma(X, Z), \sigma(Y, W)\rangle
$$

Defining the covariant derivative of $\sigma$ by $\left(\nabla^{\prime} \sigma\right)(X, Y, Z)=D_{X} \sigma(Y, Z)-\sigma\left(\nabla_{X} Y, Z\right)-$ $\sigma\left(Y, \nabla_{X} Z\right)$, the Codazzi equation is

$$
\left(R_{X Y} Z\right)^{\perp}=\left(\nabla^{\prime} \sigma\right)(X, Y, Z)-\left(\nabla^{\prime} \sigma\right)(Y, X, Z)
$$

Now for a Lagrangian submanifold, the almost complex structure $J$ of the Kähler manifold $\mathbb{C}^{n}$ is an isometry between the tangent bundle and the normal bundle, and hence the equation of Ricci-Kühne, $R^{\perp}(X, Y, \xi, \zeta)=g\left(\left[A_{\xi}, A_{\zeta}\right] X, Y\right)$, gives no new information; in particular we can find the commutators of the Weingarten maps directly.

Also for a Lagrangian submanifold of a Kähler manifold with local orthonormal basis $\left\{e_{1}, \ldots, e_{n}\right\}$, let $A_{i}$ denote the Weingarten map corresponding to the normal $J e_{i}$. Then we have from

$$
-A_{i} e_{j}+D_{e_{j}} J e_{i}=\bar{\nabla}_{e_{j}} J e_{i}=J \bar{\nabla}_{e_{j}} e_{i}=J \nabla_{e_{j}} e_{i}+J \sigma\left(e_{i}, e_{j}\right)
$$

and the symmetry of the second fundamental form that

$$
A_{i} e_{j}=-J \sigma\left(e_{i}, e_{j}\right)=A_{j} e_{i}
$$

an important relation that we will use frequently.
In the course of our work we will use several properties of Codazzi tensors as developed by Derdziński [7]. A symmetric tensor field $L$ of type $(1,1)$ is a Codazzi tensor if it satisfies $\left(\nabla_{X} L\right) Y-\left(\nabla_{Y} L\right) X=0$. Derdziński proves the following Lemma.

Lemma 3 If a Codazzi tensor has more than one eigenvalue, then the eigenspaces for each eigenvalue form an integrable subbundle on open sets of constant multiplicity. If an eigenvalue has multiplicity greater than 1 , then the eigenvalue is constant on its integral submanifolds. Moreover the integral submanifolds are umbilical submanifolds, and if the eigenvalue is constant on the manifold then the integral submanifolds are totally geodesic.

As a matter of notation, for a Riemannian metric $g$, let $Q$ denote its Ricci operator and $\tau$ its scalar curvature.

Recall that if the Weyl conformal curvature tensor of a Riemannian manifold $M^{n}$ vanishes, the curvature tensor is given in terms of the Schouten tensor

$$
L=-\frac{Q}{n-2}+\frac{\tau}{2(n-1)(n-2)} I
$$

by

$$
\begin{aligned}
g\left(R_{X Y} Z, V\right)=-g(L X, V) g(Y, Z)+g(L X, & Z) g(Y, V) \\
& -g(L Y, Z) g(X, V)+g(L Y, V) g(X, Z)
\end{aligned}
$$

It is well known that for $n \geq 4, M^{n}$ is conformally flat if and only if the Weyl conformal curvature tensor vanishes, and that this implies that $L$ is a Codazzi tensor. For $n=3$, the Weyl conformal curvature tensor vanishes identically, and the manifold is conformally flat if and only if the Schouten tensor is a Codazzi tensor.

## 3 Conformally Flat, Minimal, Lagrangian Submanifolds

In this section we develop to some extent the theory of conformally flat, minimal, Lagrangian submanifolds in general and specifically prove our theorem on the case of the Schouten tensor having two eigenvalues.

Let $\left\{e_{1}, \ldots, e_{n}\right\}$ be a local orthonormal eigenvector basis of $L$ and let $A_{i}$ denote the Weingarten map corresponding to the normal $J e_{i}$. Contracting the Gauss equation and using the minimality, we see that the Ricci operator and the scalar curvature are given by

$$
\begin{equation*}
Q=-\sum_{i} A_{i}^{2}, \quad \tau=-|\sigma|^{2} \tag{3.1}
\end{equation*}
$$

and hence the Schouten tensor becomes

$$
\begin{equation*}
L=\frac{1}{n-2} \sum_{i} A_{i}^{2}-\frac{|\sigma|^{2}}{2(n-1)(n-2)} I . \tag{3.2}
\end{equation*}
$$

We also note that

$$
\operatorname{tr} L=\frac{-\tau}{2(n-1)}=\frac{|\sigma|^{2}}{2(n-1)}
$$

We will approach our work by studying the eigenvalues of $L$. We will assume that not all the eigenvalues are equal on any neighborhood, for if they were, the submanifold would be of constant curvature. Ejiri proved that a minimal Lagrangian submanifold of constant curvature in a complex space form must be totally geodesic or flat [8]. Thus with the ambient space being $\mathbb{C}^{n}$ we have both these properties, giving the trivial case of the theorem.

Let $\nu_{j}$ denote the $j$-th eigenvalue of $L$ and recall that for a Lagrangian submanifold $A_{i} e_{j}=A_{j} e_{i}$. Then computing $g\left(L e_{j}, e_{m}\right)$, we see that

$$
\nu_{j} \delta_{j m}=\frac{1}{n-2} \operatorname{tr} A_{j} A_{m}-\frac{|\sigma|^{2}}{2(n-1)(n-2)} \delta_{j m}
$$

Consequently we have that $\operatorname{tr} A_{j} A_{m}=0$ for $j \neq m$ and that the eigenvalues of $L$ are given by

$$
\nu_{j}=\frac{1}{n-2} \operatorname{tr} A_{j}^{2}-\frac{|\sigma|^{2}}{2(n-1)(n-2)}
$$

Denote by $\nu_{\nu_{j}}$ the eigenspace of $\nu_{j}$. If $X, V \in \mathcal{V}_{\nu_{j}}$ and $Y, Z \in \mathcal{V}_{\nu_{k}}$, direct computation of the curvature gives $g\left(R_{X Y} Z, V\right)=-\left(\nu_{j}+\nu_{k}\right) g(X, V) g(Y, Z)$. On the other hand, by the Gauss equation,

$$
\begin{aligned}
g\left(R_{e_{j} e_{k}} e_{l}, e_{m}\right) & =\sum_{i} g\left(A_{k} e_{l}, e_{i}\right) g\left(A_{j} e_{m}, e_{i}\right)-\sum_{i} g\left(A_{j} e_{l}, e_{i}\right) g\left(A_{k} e_{m}, e_{i}\right) \\
& =g\left(A_{k} e_{l}, A_{j} e_{m}\right)-g\left(A_{j} e_{l}, A_{k} e_{m}\right) \\
& =g\left(\left[A_{j}, A_{k}\right] e_{l}, e_{m}\right)
\end{aligned}
$$

From the form of the curvature tensor of a conformally flat space, this is zero for $j, k, l, m$ distinct. Similarly, for $j, k, m$ distinct and $k=l$,

$$
\begin{aligned}
& g\left(\left[A_{j}, A_{k}\right] e_{k}, e_{m}\right)=-\frac{1}{n-2} \operatorname{tr} A_{j} A_{m}=0 \\
& g\left(\left[A_{j}, A_{k}\right] e_{k}, e_{j}\right)=g\left(R_{e_{j} e_{k}} e_{k}, e_{j}\right)=-\left(\nu_{j}+\nu_{k}\right)
\end{aligned}
$$

Note that $g\left(\left[A_{j}, A_{k}\right] e_{k}, e_{j}\right)$ is the $(j, k)$-component of the matrix of $\left[A_{j}, A_{k}\right]$. Thus, the only non-zero components of $\left[A_{j}, A_{k}\right]$ are the $(j, k)$ and $(k, j)$ components, and if $\left\{\omega^{1}, \ldots, \omega^{n}\right\}$ is the dual basis of $\left\{e_{1}, \ldots, e_{n}\right\}$,

$$
\begin{equation*}
\left[A_{j}, A_{k}\right]=\left(\nu_{j}+\nu_{k}\right)\left(\omega^{j} \otimes e_{k}-\omega^{k} \otimes e_{j}\right) \tag{3.3}
\end{equation*}
$$

We next make the simple observation that from the minimality $g\left(\sum_{k} A_{k} e_{k}, e_{l}\right)=$ $\sum_{k} g\left(e_{k}, A_{k} e_{l}\right)=\sum_{k} g\left(e_{k}, A_{l} e_{k}\right)=\operatorname{tr} A_{l}=0$, and hence

$$
\begin{equation*}
\sum_{k} A_{k} e_{k}=0 \tag{3.4}
\end{equation*}
$$

We now compute the commutator of $L$ and $A_{j}$. From equation (3.2) we have

$$
L A_{j}-A_{j} L=\frac{1}{n-2} \sum_{i}\left(A_{i}^{2} A_{j}-A_{j} A_{i}^{2}\right)=\frac{1}{n-2} \sum_{i}\left(A_{i}\left[A_{i}, A_{j}\right]+\left[A_{i}, A_{j}\right] A_{i}\right) .
$$

Using (3.3), this gives

$$
\begin{aligned}
&\left(L A_{j}-A_{j} L\right) e_{k}=\frac{-1}{n-2} \sum_{i}\left(A_{i}\left(\nu_{i}+\nu_{j}\right)\left(\omega^{j} \otimes e_{i}-\omega^{i} \otimes e_{j}\right) e_{k}\right. \\
&\left.+\left(\nu_{i}+\nu_{j}\right)\left(\omega^{j} \otimes e_{i}-\omega^{i} \otimes e_{j}\right) A_{i} e_{k}\right)
\end{aligned}
$$

Noting that $\sum_{i} g\left(A_{i} e_{k}, e_{j}\right) e_{i}=\sum_{i} g\left(A_{k} e_{i}, e_{j}\right) e_{i}=\sum_{i} g\left(A_{k} e_{j}, e_{i}\right) e_{i}=A_{k} e_{j}$ and using (3.4),

$$
\begin{aligned}
&\left(L A_{j}-A_{j} L\right) e_{k}=\frac{-1}{n-2}\left(\delta_{j k} \sum_{i} \nu_{i} A_{i} e_{i}-\nu_{k} A_{k} e_{j}\right. \\
&\left.+\sum_{i} \nu_{i} g\left(A_{i} e_{k}, e_{j}\right) e_{i}-\sum_{i} \nu_{i} g\left(A_{i} e_{k}, e_{i}\right) e_{j}\right)
\end{aligned}
$$

Taking the inner product of both sides with $e_{m}$ and evaluating, we have

$$
\begin{aligned}
g\left(\left(L A_{j}-A_{j} L\right) e_{k}, e_{m}\right)= & \left(\nu_{m}-\nu_{k}\right) g\left(A_{j} e_{k}, e_{m}\right) \\
= & \frac{-1}{n-2}\left(\delta_{j k} \sum_{i} \nu_{i} g\left(A_{i} e_{i}, e_{m}\right)\right. \\
& \left.-\nu_{k} g\left(A_{k} e_{j}, e_{m}\right)+\nu_{m} g\left(A_{m} e_{k}, e_{j}\right)-\delta_{j m} \sum_{i} \nu_{i} g\left(A_{i} e_{k}, e_{i}\right)\right) \\
= & -\frac{\nu_{m}-\nu_{k}}{n-2} g\left(A_{k} e_{j}, e_{m}\right) \\
& \quad-\frac{1}{n-2}\left(\delta_{j k} \sum_{i} \nu_{i} g\left(A_{i} e_{i}, e_{m}\right)-\delta_{j m} \sum_{i} \nu_{i} g\left(A_{i} e_{i}, e_{k}\right)\right) .
\end{aligned}
$$

In particular,

$$
(n-1)\left(\nu_{m}-\nu_{k}\right) g\left(A_{j} e_{k}, e_{m}\right)=\delta_{j m} \sum_{i} \nu_{i} g\left(A_{i} e_{i}, e_{k}\right)-\delta_{j k} \sum_{i} \nu_{i} g\left(A_{i} e_{i}, e_{m}\right)
$$

Thus for $j, k, m$ distinct,

$$
\begin{equation*}
\left(\nu_{m}-\nu_{k}\right) g\left(A_{j} e_{k}, e_{m}\right)=0 \tag{3.5}
\end{equation*}
$$

and for $j=k \neq m$,
(3.6) $(n-1)\left(\nu_{m}-\nu_{j}\right) g\left(A_{j} e_{j}, e_{m}\right)=-\sum_{i} \nu_{i} g\left(A_{i} e_{i}, e_{m}\right)=(n-1)\left(\nu_{m}-\nu_{l}\right) g\left(A_{l} e_{l}, e_{m}\right)$
for $l \neq m$.
We now prove the following proposition, which does not use the hypothesis on the number of eigenvalues of the Schouten tensor.

Proposition 4 Let $M^{n}$ be a conformally flat, minimal, Lagrangian submanifold of $\mathbb{C}^{n}$. Then $L$ has an eigenvalue of multiplicity 1.

Proof From equation (3.6) if $\nu_{m}=\nu_{j}$, then $g\left(\sum_{i} \nu_{i} A_{i} e_{i}, e_{m}\right)=0$, and hence $\sum_{i} \nu_{i} A_{i} e_{i}$ is orthogonal to each eigenspace corresponding to an eigenvalue of multiplicity $>1$. Therefore either $\sum_{i} \nu_{i} A_{i} e_{i}=0$ or there exists an eigenvalue of multiplicity 1.

Suppose now that $L$ has no eigenvalue of multiplicity 1 . Order the $e_{1}, \ldots, e_{n}$ corresponding to the multiplicity of eigenvalues of $L$, i.e., $e_{1}, \ldots, e_{p}$ correspond to $\nu_{1}$ having multiplicity $p$, etc., and consider the corresponding block decomposition of a Weingarten map. Using $\sum_{i} \nu_{i} A_{i} e_{i}=0$ and (3.6) with $\nu_{m}$ and $\nu_{j}$ corresponding to different eigenspaces, we see that $g\left(A_{j} e_{j}, e_{m}\right)=0$. This together with equation (3.5) shows that the off-diagonal blocks of the block decomposition are zero. Also from $g\left(A_{j} e_{j}, e_{m}\right)=0$ we have $g\left(A_{m} e_{j}, e_{j}\right)=0$ and hence in the matrix of any $A_{m}$, the diagonal blocks not corresponding to $\nu_{m}$ have zeros along their diagonals. Now from the property (3.3) of $\left[A_{j}, A_{k}\right]$, we see that the sum of any two distinct eigenvalues is zero and hence there exist at most two distinct eigenvalues, one being the negative of the other, say $\nu$ and $-\nu$. From Lemma 3, we see that $\nu$ is constant on $M$ and hence the leaves of both foliations are totally geodesic submanifolds with curvatures of opposite sign. Now from $\sum_{i} \nu_{i} A_{i} e_{i}=0$, we see that $\sum_{i=1}^{p} A_{i} e_{i}=\sum_{i=p+1}^{n} A_{i} e_{i}$, but the sum on the left is tangent to one foliation and the sum on the right is tangent to the other, and hence both sides vanish. Therefore the leaves of positive curvature are minimal submanifolds in $\mathbb{C}^{n}$, a contradiction.

For the theorem, since we now know that $L$ has an eigenvalue of multiplicity 1 , say $\nu_{1}$, the other eigenvalue, $\nu_{2}$, has multiplicity $n-1$. We show next that relative to the basis $\left\{e_{i}\right\}$, the Weingarten map $A_{1}$ is given by a diagonal matrix.

Lemma 5 The matrix of $A_{1}$ is diagonal. Moreover, setting $\Phi=\frac{-1}{n-1} g\left(\sum_{i} \nu_{i} A_{i} e_{i}, e_{1}\right)$, the $j$-th diagonal entry for $j>1$ is $\frac{\Phi}{\nu_{1}-\nu_{2}}$ and the first entry is $-\frac{(n-1) \Phi}{\nu_{1}-\nu_{2}}$.

Proof We first show that the lower right-hand $(n-1) \times(n-1)$ block of $A_{1}$ is diagonal. From equation (3.5) and the general property $A_{i} e_{j}=A_{j} e_{i}$, we have

$$
\left(\nu_{1}-\nu_{k}\right) g\left(A_{1} e_{j}, e_{k}\right)=0
$$

for $j \neq k$ and $j, k \neq 1$, giving the diagonal form of the lower right-hand block. From equation (3.6)

$$
(n-1)\left(\nu_{1}-\nu_{j}\right) g\left(A_{1} e_{j}, e_{j}\right)=-g\left(\sum_{i} \nu_{i} A_{i} e_{i}, e_{1}\right),
$$

giving the diagonal entry. In the proof of Proposition 4 we saw that $\sum_{i} \nu_{i} A_{i} e_{i}$ is orthogonal to each eigenspace corresponding to an eigenvalue of multiplicity $>1$. Therefore from equation (3.6) we see that $\left(\nu_{m}-\nu_{1}\right) g\left(A_{1} e_{1}, e_{m}\right)=0$. Therefore the matrix of $A_{1}$ is diagonal. Finally by the minimality,

$$
g\left(A_{1} e_{1}, e_{1}\right)=-\frac{(n-1) \Phi}{\nu_{1}-\nu_{2}}
$$

Proof of Theorem 2 With $A_{1}$ diagonal and $A_{1} e_{2}=A_{2} e_{1}$ we can easily compute the $(1,2)$ component of $\left[A_{1}, A_{2}\right]$; comparing with (3.3) we have two cases.

Case I

$$
\begin{gathered}
\nu_{1}+\nu_{2}=\frac{n \Phi^{2}}{\left(\nu_{1}-\nu_{2}\right)^{2}}>0 \\
\Phi=0\left(\nu_{1}+\nu_{2}=0\right)
\end{gathered}
$$

Case II
If for a conformally flat manifold, the Schouten tensor has an eigenvalue $\nu_{1}$ of multiplicity 1 and a second eigenvalue $\nu_{2}$ of multiplicity $n-1$, then by Lemma 3, the eigenspaces of $\nu_{2}$ form an integrable subbundle and $\nu_{2}$ is constant along the integral submanifolds. Moreover the integral submanifolds are umbilical in $M^{n}$. This and the fact that $n-1 \geq 3$ imply that the integral submanifolds are of constant curvature, and hence we can write the metric in the form

$$
\begin{equation*}
d s^{2}=e^{2 f\left(u_{1}\right)}\left(d u_{1}^{2}+\frac{d u_{2}^{2}+\cdots+d u_{n}^{2}}{\left(1+\frac{\varepsilon}{4} \sum_{i=2}^{n} u_{i}^{2}\right)^{2}}\right) \tag{3.7}
\end{equation*}
$$

where $\varepsilon=1,-1,0$ according as the submanifolds $u_{1}=$ const have positive, negative or zero constant curvature. With respect to the orthonormal basis

$$
e_{1}=e^{-f} \frac{\partial}{\partial u_{1}} \quad \text { and } \quad e_{j}=e^{-f}\left(1+\frac{\varepsilon}{4} \sum_{i=2}^{n} u_{i}^{2}\right) \frac{\partial}{\partial u_{j}}, j>1
$$

the Levi-Civita connection is given as follows, where $i, j>1, i \neq j$ :

$$
\begin{gathered}
\nabla_{e_{1}} e_{1}=0, \quad \nabla_{e_{1}} e_{j}=0, \quad \nabla_{e_{i}} e_{1}=\left(e_{1} f\right) e_{i}, \\
\nabla_{e_{i}} e_{i}=-\left(e_{1} f\right) e_{1}+e^{-f} \frac{\varepsilon}{2} \sum_{l \neq 1, i} u_{l} e_{l}, \quad \nabla_{e_{i}} e_{j}=-e^{-f} \frac{\varepsilon}{2} u_{j} e_{i} .
\end{gathered}
$$

The computation of the curvature is now straightforward, and $\left\{e_{1}, \ldots, e_{n}\right\}$ is an eigenvector basis of $L$.

Now using the fact that $L$ is a Codazzi tensor, we have $\left(\nabla_{e_{1}} L\right) e_{i}-\left(\nabla_{e_{i}} L\right) e_{1}=0$; expanding and taking inner products with $e_{1}, e_{i}, i \neq 1$, we have

$$
\left(\nu_{2}-\nu_{1}\right) g\left(\nabla_{e_{1}} e_{i}, e_{1}\right)-e_{i} \nu_{1}=0, \quad e_{1} \nu_{2}+\left(\nu_{2}-\nu_{1}\right) g\left(\nabla_{e_{i}} e_{1}, e_{i}\right)=0
$$

Continuing by direct use of the Levi-Civita connection of the metric (3.7), we see first that since the integral curves of $e_{1}$ are geodesics, $e_{i} \nu_{1}=0$, and hence that $\nu_{1}$
is a function of $u_{1}$ alone. Secondly, since $\nabla_{e_{i}} e_{1}=\left(e_{1} f\right) e_{i}, i>1$, we have $e_{1} \nu_{2}=$ $\left(e_{1} f\right)\left(\nu_{1}-\nu_{2}\right)$, or

$$
\begin{equation*}
\nu_{2}^{\prime}=\left(\nu_{1}-\nu_{2}\right) f^{\prime} \tag{3.8}
\end{equation*}
$$

where the prime denotes differentiation with respect to $u_{1}$. Computing the curvature, we see that $\nu_{1}+\nu_{2}=-g\left(R_{e_{1} e_{2}} e_{2}, e_{1}\right)=e_{1} e_{1} f+\left(e_{1} f\right)^{2}$, giving

$$
\begin{equation*}
\nu_{1}+\nu_{2}=e^{-2 f} f^{\prime \prime} \tag{3.9}
\end{equation*}
$$

We now use the Codazzi equation

$$
\begin{equation*}
g\left(\left(\nabla_{e_{j}} A_{i}\right) e_{k}, e_{m}\right)-g\left(A_{k} e_{m}, \nabla_{e_{j}} e_{i}\right)=g\left(\left(\nabla_{e_{k}} A_{i}\right) e_{j}, e_{m}\right)-g\left(A_{j} e_{m}, \nabla_{e_{k}} e_{i}\right) \tag{3.10}
\end{equation*}
$$

Taking $i=j=1, k=m=2$ and expanding, we have

$$
g\left(\nabla_{e_{1}} \frac{\Phi}{\nu_{1}-\nu_{2}} e_{2}, e_{2}\right)=-(n-1) \frac{\Phi}{\nu_{1}-\nu_{2}} g\left(\left(e_{1} f\right) e_{2}, e_{2}\right)-2 g\left(A_{1}\left(e_{1} f\right) e_{2}, e_{2}\right)
$$

and simplifying,

$$
\begin{equation*}
e_{1} \frac{\Phi}{\nu_{1}-\nu_{2}}=-(n+1) \frac{\Phi}{\nu_{1}-\nu_{2}} e_{1} f \tag{3.11}
\end{equation*}
$$

From the statement of Case I,

$$
\frac{\Phi}{\nu_{1}-\nu_{2}}=\sqrt{\frac{\nu_{1}+\nu_{2}}{n}}
$$

from which follows

$$
e_{1} \frac{\Phi}{\nu_{1}-\nu_{2}}=\frac{1}{2 n}\left(\frac{\Phi}{\nu_{1}-\nu_{2}}\right)^{-1}\left(e_{1} \nu_{1}+\left(\nu_{1}-\nu_{2}\right) e_{1} f\right)
$$

Comparing with (3.11), we have $-(n+1)\left(\nu_{1}+\nu_{2}\right) e_{1} f=\frac{1}{2}\left(e_{1} \nu_{1}+\left(\nu_{1}-\nu_{2}\right) e_{1} f\right)$, and simplifying,

$$
\begin{equation*}
\nu_{1}^{\prime}=-\left((2 n+3) \nu_{1}+(2 n+1) \nu_{2}\right) f^{\prime} \tag{3.12}
\end{equation*}
$$

Our task will be to solve equations (3.8), (3.9), and (3.12). Differentiating (3.9) and comparing with the sum of (3.8) and (3.12), we obtain $f^{\prime \prime \prime}=-2 n f^{\prime \prime} f^{\prime}$. Let $v=f^{\prime}$ and $w=v^{\prime}$. Then $w \frac{d w}{d v}=-2 n v w$, from which follows either

Case Ia

$$
w=-n v^{2}+C_{1}
$$

where $C_{1}$ is the constant of integration, or
Case Ib

$$
w=0
$$

Integrating the first case:

$$
\begin{equation*}
u_{1}=\int \frac{d v}{C_{1}-n v^{2}} \tag{3.13}
\end{equation*}
$$

A second constant of integration is unnecessary as it represents a change of position of the origin in the $u_{1}$ variable.

We remark at this stage that if $C_{1}=n$, then $v=\tanh \left(n u_{1}\right)$ and therefore

$$
f=\frac{1}{n} \ln \cosh \left(n u_{1}\right)
$$

giving the Lagrangian catenoid; note again that another constant of integration is unnecessary as it would give a homothetic change of metric and the problem is homothetic invariant.

Of course we must deal with a general $C_{1}$ which we now do in cases by sign: $C_{1}>0, C_{1}<0, C_{1}=0$.

If $C_{1}>0$,

$$
u_{1}=\frac{1}{C_{1}} \int \frac{d v}{1-\frac{n}{C_{1}} v^{2}}= \begin{cases}\frac{1}{\sqrt{C_{1} n}} \tanh ^{-1} \sqrt{\frac{n}{C_{1}}} v, & |v|<\sqrt{\frac{n}{C_{1}}} \\ \frac{1}{\sqrt{C_{1} n}} \operatorname{coth}^{-1} \sqrt{\frac{n}{C_{1}}} v, & |v|>\sqrt{\frac{n}{C_{1}}}\end{cases}
$$

where again we may take the constant of integration equal to zero. From these we see that

$$
f=\ln \cosh ^{\frac{1}{n}}\left(\sqrt{C_{1} n} u_{1}\right) \quad \text { or } \quad f=\ln \sinh ^{\frac{1}{n}}\left(\sqrt{C_{1} n} u_{1}\right)
$$

For the first of these we have from $-2 \nu_{2}=g\left(R_{e_{i} e_{j}} e_{j}, e_{i}\right)=e^{-2 f}-\left(e_{1} f\right)^{2}, i, j>1$, that

$$
\nu_{2}=-\frac{C_{1}}{2 n} \cosh ^{-2-\frac{2}{n}}\left(\sqrt{C_{1} n} u_{1}\right)+\frac{C_{1}-n}{2 n} \cosh ^{\frac{-2}{n}}\left(\sqrt{C_{1} n} u_{1}\right)
$$

and from (3.8),

$$
\nu_{1}=\nu_{2}+\frac{\nu_{2}}{f^{\prime}}=\frac{(2 n+1) C_{1}}{2 n} \cosh ^{-2-\frac{2}{n}}\left(\sqrt{C_{1} n} u_{1}\right)-\frac{C_{1}-n}{2 n} \cosh ^{\frac{-2}{n}}\left(\sqrt{C_{1} n} u_{1}\right)
$$

We noted the Ricci operator (3.1) at the outset and therefore

$$
\begin{aligned}
g\left(Q e_{2}, e_{2}\right) & =-g\left(\sum_{i=1}^{n} A_{i}^{2} e_{2}, e_{2}\right)=-\sum_{i=1}^{n} g\left(A_{2}^{2} e_{i}, e_{i}\right)=-\operatorname{tr} A_{2}^{2}=-2 \frac{\Phi^{2}}{\left(\nu_{1}-\nu_{2}\right)^{2}} \\
& =-2 \frac{\nu_{1}+\nu_{2}}{n}=-\frac{2 C_{1}}{n} \cosh ^{-2-\frac{2}{n}}\left(\sqrt{C_{1} n} u_{1}\right)
\end{aligned}
$$

On the other hand, we can compute $g\left(Q e_{2}, e_{2}\right)$ intrinsically:

$$
\begin{aligned}
g\left(Q e_{2}, e_{2}\right) & =(n-2) e^{e f}-e_{1} e_{1} f-(n-1)\left(e_{1} f\right)^{2} \\
& =-\frac{2 C_{1}}{n} \cosh ^{-2-\frac{2}{n}}\left(\sqrt{C_{1} n} u_{1}\right)+\frac{(n-2)\left(n-C_{1}\right)}{n} \cosh ^{-\frac{2}{n}}\left(\sqrt{C_{1} n} u_{1}\right) .
\end{aligned}
$$

Comparing, we have that $C_{1}=n$, the case of the Lagrangian catenoid.
To complete the analysis of the case $C_{1}>0$, we treat $f=\ln \sinh ^{\frac{1}{n}}\left(\sqrt{C_{1} n} u_{1}\right)$.
Here we have from equation (3.9),

$$
\nu_{1}+\nu_{2}=e^{-2 f} f^{\prime \prime}=-C_{1} \sinh ^{-2-\frac{2}{n}}\left(\sqrt{C_{1} n} u_{1}\right)<0
$$

contradicting the fact that in Case I $\nu_{1}+\nu_{2}>0$.
We now turn to the case $C_{1}<0$; let $C=-C_{1}$. Equation (3.13) now takes the form

$$
u_{1}=-\int \frac{d v}{C+n v^{2}}=-\frac{1}{\sqrt{C n}} \tan ^{-1} \sqrt{\frac{n}{c}} v
$$

From this we obtain $f=\ln \cos ^{1 / n}\left(\sqrt{C n} u_{1}\right)$, and $\nu_{1}+\nu_{2}$ becomes

$$
\nu_{1}+\nu_{2}=e^{-2 f} f^{\prime \prime}=-C \cos ^{-2-\frac{2}{n}}\left(\sqrt{C n} u_{1}\right)<0
$$

Again this contradicts $\nu_{1}+\nu_{2}>0$.
Finally for $C_{1}=0, v^{\prime}=-n v^{2}$ from which $f\left(u_{1}\right)=\frac{1}{n} \ln \left(n u_{1}\right)$. From this

$$
\nu_{1}+\nu_{2}=-\frac{1}{n}\left(u_{1}\right)^{-2-\frac{2}{n}}<0
$$

a contradiction.
Thus we have shown that in Case Ia, $M^{n}$ is (to within homothety) locally isometric to the Lagrangian catenoid. To show that the second fundamental form agrees with that of the Lagrangian catenoid, it remains to show that for the Weingarten maps, $A_{k}, k \geq 2$, the lower right-hand $(n-1) \times(n-1)$ blocks vanish. This now is immediate from $\tau=-|\sigma|^{2}$, equation (3.1). The left side can be computed intrinsically from the curvature of (3.7) with the now known function $f$. The terms in the right-hand side coming from $A_{1}$ and the first columns and rows of the other $A_{k}$ contribute the same as $\tau$ on the left. Thus the lower right-hand $(n-1) \times(n-1)$ blocks vanish.

Turning to Case Ib and Case II, Case Ib leads immediately to $f^{\prime \prime}=0$, and therefore by equation (3.9) to $\nu_{1}+\nu_{2}=0$, and hence these two cases are the same. Thus we have $\Phi=0$, and hence $A_{1}=0$. Since $f^{\prime \prime}=0, f$ is linear, say $f=a u_{1}$, a second constant of integration being unnecessary as before. An intrinsic computation yields $\nu_{2}=\frac{a^{2}+1}{2} e^{-2 a u_{1}}$. Using the Codazzi equation (3.10) with index $k=1$, we get $-a g\left(A_{i} e_{j}, e_{m}\right)=\frac{\partial}{\partial u_{1}} g\left(A_{i} e_{j}, e_{m}\right)$, and therefore each $g\left(A_{i} e_{j}, e_{m}\right)$ is $e^{-a u_{1}}$ times a function of $u_{2}, \ldots, u_{n}$.

Now since $A_{1}=0$, the first row and column of $A_{2}$ vanish and hence $L$ and $A_{2}$ commute. We may therefore choose $e_{2}, \ldots, e_{n}$ such that $A_{2}$ is diagonal, say $e^{-a u_{1}}$ times a diagonal matrix with entries $b_{k k}, b_{11}$ being zero. The matrix of any other $A_{k}$ is $e^{-a u_{1}}$ times a matrix whose first row and column vanish and whose second row and column have $b_{k k}$ as the only possible non-zero entry. Now compute the $(2, k)$ entry of [ $\left.A_{2}, A_{k}\right]$ and use (3.3) to obtain $b_{22} b_{k k}-b_{k k}^{2}=a^{2}+1$. This is a quadratic in $b_{k k}$ and we obtain $b_{k k}=\left(b_{22} \pm D\right) / 2$, where $D=\sqrt{b_{22}^{2}-4\left(a^{2}+1\right)}$. However the submanifold is minimal and hence the trace of $A_{2}$ vanishes. Therefore

$$
b_{22}+\frac{(n-2)}{2} b_{22}+\frac{1}{2}(\text { a sum, difference of }(n-2) \text { copies of } D)=0
$$

Rearranging we have

$$
n b_{22}=\text { a sum, difference of }(n-2) \text { copies of } D
$$

which, since $D<\left|b_{22}\right|$, is impossible.

## 4 The Lagrangian Catenoid Is Not Quasi-Umbilical

We close with a remark showing that the Lagrangian catenoid is not a quasi-umbilical submanifold. It is known [6] that a quasi-umbilical submanifold of dimension $\geq 4$ of a conformally flat space is conformally flat, but in general not conversely. Generalizing the result of Cartan mentioned in the introduction, Moore and Morvan [12] showed that a conformally flat submanifold $M^{n}$ of Euclidean space $E^{n+p}$ is quasiumbilical if $p \leq \min \{4, n-3\}$. Chen and Verstraelen [5] showed that an $n$-dimensional submanifold of a conformally flat space of dimension $(n+p)$ with flat normal connection is quasi-umbilical for $p \leq n-3$. However here in the case of Lagrangian submanifolds, the codimension is $n$. We begin by recalling the definition of quasiumbilicity.

Consider an $n$-dimensional submanifold of an $(n+p)$-dimensional Riemannian manifold. A (local) normal vector field is a quasi-umbilical section of the normal bundle if the corresponding Weingarten map has at least $n-1$ eigenvalues equal. The submanifold is said to be quasi-umbilical if there exist $p$ mutually orthogonal quasi-umbilical normal sections.

We have already noted the Weingarten maps of the Lagrangian catenoid with respect to the normal fields $J e_{i}$, namely

$$
A_{1}=\left(\begin{array}{cccc}
-(n-1) a & & & \bigcirc \\
& a & & \\
& & \ddots & \\
\bigcirc & & & a
\end{array}\right), \quad A_{i}=\left(\begin{array}{ccccccc}
0 & \cdots & 0 & a & 0 & \cdots & 0 \\
\vdots & & & & & & \\
0 & & & & & \\
a & & & \bigcirc & & \\
0 & & & & & \\
\vdots & & & & & \\
0 & & & & &
\end{array}\right)
$$

where $a=\cosh ^{-\left(1+\frac{1}{n}\right)}(n u)$.
Suppose now that $\zeta_{i}=\sum_{j} \alpha_{i j} J e_{j}$ is an orthonormal basis of normal vectors with respect to which the Weingarten maps $B_{i}$ have $n-1$ eigenvalues equal. Then from the Weingarten equation $\bar{\nabla}_{X} \zeta_{i}=-B_{i} X+D_{X} \zeta_{i}$ we have

$$
-B_{i} X+D_{X} \zeta_{i}=\sum_{j}\left(X \alpha_{i j}\right) J e_{j}+\sum_{j} \alpha_{i j}\left(-A_{j} X+D_{X} J e_{j}\right)
$$

Thus the matrix of $B_{i}$ is

$$
\left(\begin{array}{cccc}
-(n-1) a \alpha_{i 1} & a \alpha_{i 2} & \cdots & a \alpha_{i n} \\
a \alpha_{i 2} & a \alpha_{i 1} & & \bigcirc \\
\vdots & & \ddots & \\
a \alpha_{i n} & \bigcirc & & a \alpha_{i 1}
\end{array}\right)
$$

We now compute the characteristic polynomial $P(\lambda)$ of this matrix. Since $\left(\alpha_{i j}\right)$ is an orthognal matrix, $\left(a \alpha_{i 1}\right)^{2}+\cdots+\left(a \alpha_{i n}\right)^{2}=a^{2}$, and we find that

$$
P(\lambda)=\left(a \alpha_{i 1}-\lambda\right)^{n-2}\left[\lambda^{2}+(n-2) a \alpha_{i 1} \lambda-(n-2) a^{2} \alpha_{i 1}^{2}-a^{2}\right] .
$$

The zeros of the quadratic factor are

$$
\frac{-(n-2) a \alpha_{i 1} \pm \sqrt{\left(n^{2}-4\right) a^{2} \alpha_{i 1}^{2}+4 a^{2}}}{2}
$$

If now the eigenvalue $a \alpha_{i 1}$ has multiplicity $n-1$, we have

$$
n a \alpha_{i 1}= \pm \sqrt{\left(n^{2}-4\right) a^{2} \alpha_{i 1}^{2}+4 a^{2}}
$$

Squaring gives $4 a^{2} \alpha_{i 1}^{2}=4 a^{2}$ or $\alpha_{i 1}^{2}=1$, thus each entry in the first column of $\left(\alpha_{i j}\right)$ is $\pm 1$, contradicting its orthogonality.

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Department of Mathematics
Michigan State University
East Lansing, MI 48824
U. S. A.
e-mail: blair@math.msu.edu


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