# CONSTRUCTION OF PRINGIPAL FUNGTIONS BY ORTHOGONAL PROJECTION 

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1. Normal operators. Given a point set $E$ on an open Riemann surface $V$ we denote by $H(E)$ the space of functions $u$ harmonic in open sets $O(u)$ containing $E$. Let $V_{0}$ be a regular region of $V$ with border $\alpha$, and consider restrictions $f$ to $\alpha$ of functions in $H(\alpha)$. For $V_{1}=V-\bar{V}_{0}$, an operator $L$ from $H(\alpha)$ to $H\left(\bar{V}_{1}\right)$ is, by definition, normal if

$$
\begin{align*}
& L f=f \text { on } \alpha,  \tag{1}\\
& L\left(c_{1} f_{1}+c_{2} f_{2}\right)=c_{1} L\left(f_{1}\right)+c_{2} L\left(f_{2}\right),  \tag{2}\\
& L 1=1,  \tag{3}\\
& L f \geqslant 0 \quad \text { for } f \geqslant 0,  \tag{4}\\
& \int_{\alpha} * d L f=0 . \tag{5}
\end{align*}
$$

For general properties of normal operators we refer to Ahlfors (1), Ahlfors and Sario(2), Oikawa (5, 6), Rodin (7), Sario (8, 9, 10), Sario, Schiffer, and Glasner (11), Sario and Weill (12), and Weill (13).

Let $\Omega$ be a regular region with border $\beta_{\Omega}$ such that $\bar{V}_{0} \subset \Omega$. For a given $f$ denote by $u_{\Omega}$ the harmonic function in $\bar{\Omega} \cap \bar{V}_{1}$ with $u_{\Omega}\left|\alpha=f, u_{\Omega}\right| \beta_{\Omega}=$ const, $\int_{\alpha}{ }^{*} d u_{\Omega}=0$. As $\Omega$ exhausts $V, u_{\Omega}$ tends to a harmonic limit $u=L_{1} f$ on $\bar{V}_{1}$, where $L_{1}$ is a normal operator. Using Royden's compactification we shall first show (Theorem 1) that a normal operator $L$ is $L_{1}$ if and only if, in a sense, $L f$ is constant on the ideal boundary of $V$.
2. Principal functions. The principal function problem consists in constructing, for a given $s \in H\left(\bar{V}_{1}\right)$ and given $L$, a function $p \in H(V)$ such that

$$
\begin{equation*}
p \mid \bar{V}_{1}=s+L(p-s \mid \alpha) \tag{6}
\end{equation*}
$$

It is known that the condition

$$
\begin{equation*}
\int_{\alpha}^{*} d s=0 \tag{7}
\end{equation*}
$$

is necessary and sufficient for the solvability of the problem (9). The solution, called the principal function, is unique up to an additive constant. The function $s$ is interpreted as having a singularity on the ideal boundary of $V$ and is called the singularity function.

[^0]We shall show that for any given $V$ and $s$ with (7) the existence of the principal function corresponding to $L_{1}$ can be proved by the method of orthogonal projection (Theorem 2). It should be noted here that an inequality of the Poincaré type (Lemma 3) takes the place of the Harnack inequality in the existence proof. It is hoped that our study, methodological in nature, will also pave the road to the solution of the main problem, the construction of principal forms in Riemannian spaces.

Reference is made here to the recent interesting Research Announcement by Browder (3). Although there were some technical difficulties in applying his approach to prove earlier results or to extend them in the original direction, his Announcement threw new light on the entire principal function problem and was indeed the immediate incentive to the present study.
3. Weyl's lemma. We denote by $\Gamma=\Gamma(V)$ the space of real measurable 1 -forms $\omega$ on $V$ with

$$
\int_{V} \omega \wedge{ }^{*} \omega<\infty
$$

With the inner product

$$
\left(\omega_{1}, \omega_{2}\right)=\int_{V} \omega_{1} \wedge^{*} \omega_{2}
$$

and the norm $\|\omega\|=\sqrt{(\omega, \omega)}, \Gamma$ becomes a Hilbert space. Let $\Gamma_{e}{ }^{1}$ be the subspace of $\Gamma$ of continuous exact differentials:

$$
\Gamma_{e}{ }^{1}=\left\{d f \mid f \in C^{1}(V), d f \in \Gamma\right\} .
$$

The closure of $\Gamma_{e}{ }^{1}$ in $\Gamma$ is denoted by $\Gamma_{e}$. We also consider the space $\Gamma_{e 0}{ }^{1}$ of continuous exact differentials with compact supports in $V$, i.e.

$$
\Gamma_{e 0}{ }^{1}=\left\{d f \mid f \in C_{0}{ }^{1}(V)\right\} .
$$

Then $\Gamma_{e 0}{ }^{1} \subset \Gamma$, and we denote by $\Gamma_{e 0}$ the closure of $\Gamma_{e 0}{ }^{1}$ in $\Gamma$. We shall use Weyl's lemma in the following form (2):

Lemma 1. If an element $\alpha$ in $\Gamma_{e}$ is orthogonal to $\Gamma_{e 0}{ }^{1}$, then there exists a function $u$ in $H D(V)$ such that $\alpha=d u$, and vice versa.
4. Royden's boundary. A real-valued continuous function $f$ on $V$ is said to be a continuous Dirichlet function if there exists an $\omega$ in $\Gamma(V)$ such that

$$
\int_{V} \omega \wedge \omega_{0}=-\int_{V} f d \omega_{0}
$$

for any $C^{2}$-form $\omega_{0}$ on $V$ with compact support; we set $\omega=d f$. Denote by $R(V)$ the family of continuous Dirichlet functions on $V$ and by $R_{0}(V)$ the subfamily of functions with compact supports in $V$. For $f, g \in R(V)$ we set

$$
\rho(f, g)=\|d f-d g\|+\sum_{n=1}^{\infty} 2^{-n} \sup _{K_{n}}|f-g| \cdot(1+|f-g|)^{-1}
$$

where $\left\{K_{n}\right\}$ is an exhaustion of $V$. Endowed with $\rho, R(V)$ is a complete metric space. Let $R_{\delta}(V)$ be the closure of $R_{0}(V)$ in $R(V)$ in terms of this metric $\rho$.

The Royden compactification $V^{*}$ of $V$ is the compact Hausdorff space with the following two properties: (a) it contains $V$ as its open and dense subspace such that every function in $R\left(V^{\prime}\right)$ can be extended continuously to $V^{*}$ with infinite values admitted; (b) $R(V)$, considered as a family of functions on $V^{*}$, separates points in $V^{*}$. The compact set $\beta=V^{*}-V$ is the Royden boundary of $V$. The set

$$
\delta=\left\{p \in V^{*} \mid f(p)=0 \text { for every } f \in R_{\delta}(V)\right\}
$$

is a compact subset of $\beta$. We call $\delta$ the Royden harmonic boundary of $V$. For details and fundamental properties of these concepts we refer to (4).

Lemma 2. Let $u$ be a function in $H D\left(V_{1}\right)$ such that du can be continued to all of $V$ so as to be a 1 -form $\omega$ in $\Gamma_{e 0}(V)$. Then $u$ is finitely continuous on

$$
V_{1} \cup \beta=V^{*}-\bar{V}_{0}
$$

and a constant on $\delta$.
Proof. We take a smaller boundary neighbourhood $V^{\prime}{ }_{1}$ with $\bar{V}^{\prime}{ }_{1} \subset V_{1}$, if necessary, to assume that $u \in H D\left(\bar{V}_{1}\right)$. Then we can continue $u$ to $V$ as a function $u_{0}$ in $C^{1}(V)$. Clearly $u_{0} \in R(V)$; thus $u_{0}$ is continuous on $V^{*}$ and a fortiori $u$ is continuous on $V_{1} \cup \beta$, with infinite values admitted. If $V \in O_{G}$, then $\beta=\emptyset, H D\left(\bar{V}_{1}\right)=H B\left(\bar{V}_{1}\right)$, and the assertion is trivial. Therefore we may assume that $V \notin O_{G}$.

Since $R(V)=H D(V)+R_{\delta}(V)$, there exists a function $v \in H D(V)$ such that $u_{0}-v \in R_{\delta}(V)$. Then $d\left(u_{0}-v\right) \in \Gamma_{e 0}$ and consequently $d v \in \Gamma_{e 0}$. By Lemma $1, d v$ is orthogonal to $\Gamma_{e 0}{ }^{1}$ and hence to $\Gamma_{e 0}$. In particular,

$$
\|d v\|^{2}=(d v, d v)=0
$$

which means that $u_{0}$ - const $\in R_{\delta}(V)$, or $u_{0}=u=$ const on $\delta$. By the maximum principle, $u$ is finitely continuous at $\beta$.
5. A characterization of the $L_{1}$-operator. We can now give a characterization of the operator $L_{1}$ in terms of the Royden compactification: it is a normal operator such that $L_{1} f$ is finitely continuous at the Royden boundary $\beta$ and a constant $c_{f}$ on the Royden harmonic boundary $\delta$ for every $f$ in $H(\alpha)$. Explicitly, we have for a given $u \in H\left(\bar{V}_{1}\right)$ :

Theorem 1. Necessary and sufficient for $L_{1} u=u$ is that $u$ satisfies the following conditions:

$$
\begin{align*}
& u \text { is finitely continuous on } V_{1} \cup \beta,  \tag{8}\\
& u \text { is a constant on } \delta,  \tag{9}\\
& \int_{\alpha}{ }^{*} d u=0 . \tag{10}
\end{align*}
$$

Proof. If $V \in O_{G}$, the properties $L_{1} u=u$ and (8) are each equivalent to $u \in H B\left(\bar{V}_{1}\right)=H D\left(\bar{V}_{1}\right)$, which includes (10). Condition (9) is trivially satisfied by every $u$ in $H\left(\bar{V}_{1}\right)$ because $\delta=\emptyset$. Thus, we have only to consider the case $V \notin O_{G}$.

First assume that $L_{1} u=u$. Let $\left\{U_{n}\right\}_{n=0}^{\infty}$ be an exhaustion of $V$ by regular regions, with $U_{0}=V_{0}$. For $f$ in $H(\alpha)$, let $S_{n} f$ be the harmonic function in $\bar{U}_{n}-U_{0}$ with continuous boundary value 0 on $\partial U_{n}$ and $f$ on $\alpha=\partial V_{0}=\partial U_{0}$. Let $f_{0} \in C^{1}\left(\bar{U}_{0}\right)$ such that $f_{0}=f$ on $\alpha$. We set $S_{n}^{\prime} f=S_{n} f$ on $\bar{U}_{n}-U_{0}$, $S_{n}^{\prime} f=f_{0}$ on $\bar{U}_{0}$, and $S_{n}^{\prime} f=0$ on $V-U_{n}$. Then $S_{n}^{\prime} f \in R_{0}(V)$ and since

$$
\left\|d\left(S_{n}^{\prime} f\right)-d\left(S_{n+m}^{\prime} f\right)\right\|^{2}=\left\|d\left(S_{n}^{\prime} f\right)\right\|^{2}-\left\|d\left(S_{n+m}^{\prime} f\right)\right\|^{2}
$$

we see that $\left\{S_{n}^{\prime} f\right\}_{n=1}^{\infty}$ converges to a function, say $S^{\prime} f$, on $V$ in the $\rho$-metric. Hence $S^{\prime} f \in H D\left(\bar{V}_{1}\right), S f \mid \alpha=f, S f$ is continuous on $V_{1} \cup \beta$, and $S f=0$ on $\delta$. Moreover $\left\|d\left(S^{\prime}{ }_{n} f\right)-d(S f)\right\|_{v_{1}} \rightarrow 0$ as $n \rightarrow \infty$, and $S_{n}^{\prime} f$ converges to $S f$ uniformly in compact sets of $\bar{V}_{1}=V_{1} \cup \alpha$. It follows that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{\alpha}^{*} d\left(S_{n} f\right)=\int_{\alpha}^{*} d(S f) \tag{11}
\end{equation*}
$$

Let $w_{n}$ be harmonic in $\bar{U}_{n}-U_{0}$ with $w_{n}\left|\partial U_{0}=1, w_{n}\right| \alpha=0$. Similarly let $w$ be continuous on $V_{1} \cup \beta$ and harmonic in $V_{1}$ with boundary value 1 at $\delta$ and 0 at $\alpha$. By the same argument as above, $\left\|d w_{n}{ }_{n}-d w\right\|_{V_{1}} \rightarrow 0$ as $n \rightarrow \infty$, and $w^{\prime}{ }_{n}$ converges to $w$ uniformly in compact sets of $\bar{V}_{1}=V_{1} \cup \alpha$; here we set $w_{n}^{\prime}=w_{n}$ on $\bar{U}_{n}-U_{0}$ and $w_{n}^{\prime}=1$ on $V-U_{n}$. We have

$$
\int_{\alpha}^{*} *_{d w_{n}}=-\int_{\partial U_{n}}{ }^{*} d w_{n}=-\left\|d w_{n}^{\prime}\right\|_{V_{1}}^{2} .
$$

Hence

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{\alpha} *_{d} d w_{n}=\int_{\alpha}^{*} d w=-\lim _{n \rightarrow \infty}\left\|d w_{n}^{\prime}\right\|_{V_{1}}^{2}=-\|d w\|^{2}<0 . \tag{12}
\end{equation*}
$$

If we put

$$
L_{1}^{(n)} f=S_{n} f-\left(\int_{\alpha}^{*} d\left(S_{n} f\right) / \int_{\alpha}^{*} d w_{n}\right) w_{n}
$$

then $L_{1}{ }^{(n)} f$ is constant on $\partial U_{n}$ and

$$
\int_{\alpha}{ }^{*} d\left(L_{1}^{(n)} f\right)=0 .
$$

Thus by the definition of $L_{1} f, L_{1} f=\lim _{n} L_{1}{ }^{(n)} f$ on $V$. Using (11) and (12), we obtain the representation

$$
\begin{equation*}
L_{1} f=S f-\left(\int_{\alpha}^{*} d(S f) / \int_{\alpha} * d w\right) w \tag{13}
\end{equation*}
$$

By the properties of $S f$ and $w, L_{1} f$ is continuous on $V_{1} \cup \beta$ and a constant on $\delta$. Thus, in particular, $u=L_{1} u$ satisfies (8), (9), and (10).

Conversely assume that $u$ satisfies these three conditions. As above, the same is true of $L_{1} u$ and of $v=u-L_{1} u$. Let $v=c$ on $\delta$. Then $v=c w$ and, by (12), we must have $c=0$. Therefore $v \equiv 0$ on $V_{1}$.

These arguments are based on the fact that every function in $H D\left(\bar{V}_{1}\right)$ takes its maximum and minimum on $\alpha \cup \delta$.
6. Fundamental inequality. Before embarking on the existence proof by the method of orthogonal projection we need the following lemma; it plays a role in our proof of equal importance to the $q$-lemma in (9).

Lemma 3. Let $\omega$ be a fixed continuous 1 -form defined on $\alpha$ such that

$$
\int_{\alpha} \omega=0 .
$$

Then there exists a constant $c$ depending only on $V_{0}$ and $\omega$ such that

$$
\begin{equation*}
\left|\int_{\alpha} \phi \omega\right|^{2} \leqslant c \int_{V_{0}} d \phi \wedge{ }^{*} d \phi \tag{14}
\end{equation*}
$$

for every $\phi \in C^{1}\left(V_{0}\right) \cap C\left(\bar{V}_{0}\right)$.
Proof. First we note that we have only to prove (14) for $\phi$ in $H\left(V_{0}\right) \cap C\left(\bar{V}_{0}\right)$. In fact, for $\phi \in C^{1}\left(\bar{V}_{0}\right)$ we let $h_{\phi}$ be the harmonic function in $V_{0}$ with continuous boundary value $\phi$ at $\alpha=\partial V_{0}$. Since

$$
\int_{V_{0}} d h_{\phi} \wedge^{*} d h_{\phi} \leqslant \int_{V_{0}} d \phi \wedge{ }^{*} d \phi
$$

and

$$
\int_{\alpha} h_{\phi} \omega=\int_{\alpha} \phi \omega,
$$

the validity of (14) for $h_{\phi}$ gives that for $\phi$. We therefore may and will assume in the following that $\phi \in H\left(V_{0}\right) \cap C\left(\bar{V}_{0}\right)$.

Let $z_{0}$ be a fixed point in $V_{0}$ and let $g_{0}\left(z, z_{0}\right)$ be Green's function on $V_{0}$. We put

$$
r(z)=\exp \left(-g_{0}\left(z, z_{0}\right)\right), \quad d \theta(z)=-{ }^{*} d g_{0}\left(z, z_{0}\right)
$$

and form Green's star domain $V^{\prime}{ }_{0}$ of $g_{0}\left(z, z_{0}\right)$ on $V_{0}$. Explicitly $V^{\prime}{ }_{0}$ is obtained from $V_{0}$ by removing all closures of Green's lines issuing from the branch point of $g_{0}\left(z, z_{0}\right)$ in $V_{0}$. Then $r(z) e^{i \theta(z)}$ maps $V^{\prime}{ }_{0}$ onto a unit disk with a finite number of radial slits issuing from some point in the disk different from the origin and terminating at the unit circumference in a one-to-one and conformal fashion, $\alpha$ corresponding to the unit circumference. Fix a positive number $a$ such that $0<2 a<1$ and $\left\{z \mid z \in V_{0}, r(z) \leqslant 2 a\right\}$ is the disk in $V^{\prime}{ }_{0}$. We write $\phi_{r}(\theta)=\phi\left(r e^{i \theta}\right)$, which can be considered as an element in $L^{2}(0,2 \pi)$ with norm

$$
\left|\phi_{r}\right|_{2}=\sqrt{\int_{0}^{2 \pi}\left|\phi_{r}(\theta)\right|^{2} d \theta}
$$

First assume that $\phi\left(z_{0}\right)=0$. Except for a finite number of values $\theta$,

$$
\phi\left(e^{i \theta}\right)-\phi\left(a e^{i \theta}\right)=\int_{a}^{1} \frac{\partial}{\partial r} \phi\left(r e^{i \theta}\right) d r
$$

By Schwarz's inequality,

$$
\left|\phi_{1}(\theta)-\phi_{a}(\theta)\right|^{2} \leqslant \int_{a}^{1}\left|\frac{\partial}{\partial r} \phi\left(r e^{i \theta}\right)\right|^{2} r d r \cdot \int_{a}^{1} r^{-1} d r .
$$

Hence on integrating both sides over $(0,2 \pi)$ and on observing that

$$
\left|\frac{\partial}{\partial r} \phi\left(r e^{i \theta}\right)\right|^{2} \leqslant\left|\frac{\partial}{\partial r} \phi\left(r e^{i \theta}\right)\right|^{2}+r^{-2}\left|\frac{\partial}{\partial \theta} \phi\left(r e^{i \theta}\right)\right|^{2}=|\operatorname{grad} \phi|^{2},
$$

we obtain

$$
\begin{equation*}
\left|\phi_{1}-\phi_{a}\right|_{2} \leqslant \sqrt{\log a^{-1}} \cdot \sqrt{\int_{V_{0}} d \phi \wedge^{*} d \phi} \tag{15}
\end{equation*}
$$

Let ${ }^{*} \phi(z)$ be the conjugate harmonic function of $\phi(z)$ in $\left|z-z_{0}\right|<2 a$ such that ${ }^{*} \phi\left(z_{0}\right)=0$. Consider the analytic function $f(z)=\phi(z)+i^{*} \phi(z)$ in $\left|z-z_{0}\right|<2 a$. Since $f\left(z_{0}\right)=0$, we have

$$
f(z)=\int_{z_{0}}^{z} f^{\prime}(z) d z
$$

and thus

$$
|\phi(z)| \leqslant|f(z)| \leqslant\left|z-z_{0}\right| \max _{\left|z-z_{0}\right| \leqslant a}\left|f^{\prime}(z)\right|
$$

for $\left|z-z_{0}\right| \leqslant a$. In particular

$$
\begin{equation*}
\left|\phi_{a}(\theta)\right|^{2} \leqslant a^{2} \max _{\left|z-z_{0}\right| \leqslant a}\left|f^{\prime}(z)\right|^{2} \tag{16}
\end{equation*}
$$

As $\left|f^{\prime}(z)\right|^{2}$ is subharmonic in $\left|z-z_{0}\right|<2 a$,

$$
\left|f^{\prime}(z)\right|^{2} \leqslant\left(\pi a^{2}\right)^{-1} \int_{|\zeta-z| \leqslant a}\left|f^{\prime}(\zeta)\right|^{2} d \xi d \eta
$$

for $z$ in $\left|z-z_{0}\right| \leqslant a$, with $\zeta=\xi+i \eta$. Since $\left|f^{\prime}(\zeta)\right|^{2} d \xi d \eta=d \phi \wedge{ }^{*} d \phi$ in $|z-\zeta| \leqslant a$, we conclude that

$$
\left|f^{\prime}(z)\right|^{2} \leqslant\left(\pi a^{2}\right)^{-1} \int_{V_{0}} d \phi \wedge^{*} d \phi
$$

This with (16) gives

$$
\left|\phi_{a}(\theta)\right|^{2} \leqslant \pi^{-1} \int_{V_{0}} d \phi \wedge^{*} d \phi
$$

and therefore

$$
\begin{equation*}
\left|\phi_{a}\right|_{2} \leqslant \sqrt{2 \int_{V_{0}} d \phi \wedge^{*} d \phi} \tag{17}
\end{equation*}
$$

Since $\left|\phi_{1}\right|_{2} \leqslant\left|\phi_{1}-\phi_{a}\right|_{2}+\left|\phi_{a}\right|_{2}$, we infer from (15) and (17) that

$$
\begin{equation*}
\int_{0}^{2 \pi}\left|\phi\left(e^{i \theta}\right)\right|^{2} d \theta \leqslant c_{1} \int_{V_{0}} d \phi \wedge^{*} d \phi \tag{18}
\end{equation*}
$$

where $c_{1}=2+\log a^{-1}$ depends only on $V_{0}$. By a piecewise analytic representation of $\alpha$ with parameter $\theta, \omega$ can be expressed as $\omega=\Omega(\theta) d \theta$ on $\alpha$. Here $\Omega(\theta)$ is bounded, say $|\Omega(\theta)| \leqslant c_{2}$, and piecewise continuous on $\alpha$, with $c_{2}$ depending only on $\omega$ on $V_{0}$. By Schwarz's inequality and (18) we obtain

$$
\begin{aligned}
\left|\int_{\alpha} \phi \omega\right|^{2} & =\left|\int_{0}^{2 \pi} \phi\left(e^{i \theta}\right) \cdot \Omega(\theta) d \theta\right|^{2} \\
& \leqslant \int_{0}^{2 \pi}\left|\phi\left(e^{i \theta}\right)\right|^{2} d \theta \cdot \int_{0}^{2 \pi}|\Omega(\theta)|^{2} d \theta \\
& \leqslant c \int_{0}^{2 \pi} d \phi \wedge^{*} d \phi
\end{aligned}
$$

where $c=c_{1} \cdot c_{2}{ }^{2}$ depends on $\omega$ and $V_{0}$.
Next consider $\phi \in H\left(V_{0}\right) \cap C\left(\bar{V}_{0}\right)$, with not necessarily $\phi\left(z_{0}\right)=0$; then $\tilde{\phi}=\phi-\phi\left(z_{0}\right)$ satisfies $\tilde{\phi}\left(z_{0}\right)=0$. By the above reasoning,

$$
\begin{equation*}
\left|\int_{\alpha} \tilde{\phi} \omega\right|^{2} \leqslant c \int_{V_{0}} d \tilde{\phi} \wedge{ }^{*} d \tilde{\phi} \tag{19}
\end{equation*}
$$

In view of the assumption

$$
\int_{\alpha} \omega=0
$$

we have

$$
\int_{\alpha} \tilde{\phi} \omega=\int_{\alpha} \phi \omega-\phi\left(z_{0}\right) \int_{\alpha} \omega=\int_{\alpha} \phi
$$

Obviously $d \tilde{\phi}=d \phi$ and substitution in (19) gives (14) for $\phi$.
7. Existence proof by orthogonal projection. We proceed to the proof by orthogonal projection of the existence of principal functions:

Theorem 2. Let s be a harmonic function on $V_{1} \cup \alpha$ with property (7). Then there exists a harmonic function $p$ on $V$ which satisfies equation (6).

Proof. First we extend $s$ to all of the surface $V$ as a function $s_{0} \in C^{2}(V)$. For $\omega \in \Gamma_{e 0}{ }^{1}$ we put

$$
\begin{equation*}
T(\omega)=-\int_{V} \omega \wedge^{*} d s_{0}=-\int_{V} d s_{0} \wedge^{*} \omega \tag{20}
\end{equation*}
$$

This is well-defined because $\omega$ has compact support and $\omega \wedge^{*} d s_{0}=d s_{0} \wedge^{*} \omega$ is a continuous 2 -form with compact support in $V$. Clearly $T$ gives rise to a linear operator on $\Gamma_{e 0}{ }^{1}$.

We shall show next that $T$ is a bounded linear operator on $\Gamma_{e 0}{ }^{1}$. Take a regular region $W \subset V$ which contains $\bar{V}_{0}$ and the support of $\omega \in \Gamma_{e 0^{1}}$. Then

$$
\begin{equation*}
-T(\omega)=\int_{\bar{v}_{0}} \omega \wedge * d s_{0}+\int_{W-\bar{v}_{0}} \omega \wedge * d s_{0} . \tag{21}
\end{equation*}
$$

By Schwarz's inequality

$$
\begin{equation*}
\left|\int_{\bar{V}_{0}} \omega \wedge * d s_{0}\right| \leqslant \sqrt{\int_{\bar{V}_{0}} d s_{0} \wedge^{*} d s_{0}} \cdot\|\omega\| \tag{22}
\end{equation*}
$$

Since $\omega \in \Gamma_{e 0}{ }^{1}$, there exists a function $\phi$ in $C_{0}{ }^{1}(V)$ with its support in $W$ and such that $\omega=d \phi$ on $W$. As

$$
\omega \wedge * d s_{0}=d \phi \wedge * d s_{0}=d\left(\phi^{*} d s_{0}\right)-\phi d^{*} d s_{0}=d\left(\phi^{*} d s_{0}\right)
$$

in $V-\bar{V}_{0}$ and thus in $W-\bar{V}_{0}$, we have by Green's formula

$$
\begin{equation*}
\int_{W-\bar{v}_{0}} \omega \wedge{ }^{*} d s_{0}=\int_{W-\bar{v}_{0}} d\left(\phi^{*} d s_{0}\right)=\int_{\partial W+\alpha} \phi^{*} d s_{0}=\int_{\alpha} \phi^{*} d s_{0} \tag{23}
\end{equation*}
$$

By virtue of

$$
\int_{\alpha}^{*} d s_{0}=\int_{\alpha}^{*} d s=0
$$

and $\phi \in C^{1}\left(V_{0}\right) \cap C\left(\bar{V}_{0}\right)$, we can apply Lemma 3 to

$$
\int_{\alpha} \phi^{*} d s_{0}
$$

so as to obtain a constant $c$ depending only on ${ }^{*} d s_{0}$ and $V_{0}$ and such that

$$
\left|\int_{\alpha} \phi^{*} d s_{0}\right|^{2} \leqslant c \int_{V_{0}} d \phi \wedge^{*} d \phi
$$

Because of $d \phi=\omega$ and (23) we infer that

$$
\begin{equation*}
\left|\int_{W-\bar{v}_{0}} \omega \wedge{ }^{*} d s_{0}\right| \leqslant \sqrt{c}\|\omega\| . \tag{24}
\end{equation*}
$$

From (21), (22), and (24), we obtain

$$
\begin{equation*}
|T(\omega)| \leqslant K\|\omega\| \tag{25}
\end{equation*}
$$

for $\omega \in \Gamma_{e 0^{1}}$, where

$$
K=\sqrt{c}+\int_{\bar{v}_{0}} d s_{0} \wedge^{*} d s_{0}
$$

depends only on $s_{0}$ and $V_{0}$.
By the general Hilbert space theory, $T$ can be extended to

$$
\Gamma_{e 0}=\overline{\Gamma_{e 0}}
$$

so as to satisfy (25) again on $\Gamma_{e 0}$. Since $\Gamma_{e 0}$ is self-adjoint, there exists a unique element $\lambda$ in $\Gamma_{e 0}$ such that

$$
\begin{equation*}
T(\omega)=(\omega, \lambda) \tag{26}
\end{equation*}
$$

for $\omega \in \Gamma_{e 0}$. In particular, by (20) and (26),

$$
\begin{equation*}
\int_{V}\left(\lambda+d s_{0}\right) \wedge{ }^{*} \omega=0 \tag{27}
\end{equation*}
$$

for $\omega \in \Gamma_{e 0}{ }^{1}$.
Again let $\left\{U_{n}\right\}_{n=0}^{\infty}$ be an exhaustion of $V$ with $U_{0}=V_{0}$. Although $\lambda+d s_{0}$ is not an element of $\Gamma_{e}(V)$ in general, we can conclude that

$$
\lambda+d s_{0} \in \Gamma_{e}\left(U_{n}\right)
$$

because

$$
\lambda \in \Gamma_{e 0}(V) \subset \Gamma_{e}(V) \subset \Gamma_{e}\left(U_{n}\right)
$$

and $s_{0} \in C^{2}\left(\bar{U}_{n}\right)$ and hence $d s_{0} \in \Gamma_{e}\left(U_{n}\right)$. Since (27) holds for

$$
\omega \in \Gamma_{e 0}{ }^{1}\left(U_{n}\right) \subset \Gamma_{e 0}{ }^{1}(V)
$$

there exists by Lemma 1 a $q_{n} \in H D\left(U_{n}\right)$ such that $d q_{n}=\lambda+d s_{0}$ on $U_{n}$. Clearly $d q_{n+m}=d q_{n}$ on $U_{n}$ and therefore $q_{n+m}=q_{n}+$ const on $U_{n}$. Let $c_{n}$ be a constant such that $q_{n+1}=q_{n}+c_{n}$ on $U_{n}$ and set $p_{1}=q_{1}$ on $U_{1}, p_{n}=q_{n}-c_{n-1}$ on $U_{n}$ with $n>1$. Then $p_{n} \in H D\left(U_{n}\right)$ and $d p_{n}=\lambda+d s_{0}$ on $U_{n}$ and $p_{n+m}=p_{n}$ on $U_{n}$. Thus if we put

$$
\begin{equation*}
p(z)=p_{n}(z) \tag{28}
\end{equation*}
$$

for $z$ in $U_{n}$, then $p(z)$ does not depend on the choice of $U_{n}$ to which $z$ belongs. Therefore $p \in H(V)$ and

$$
\begin{equation*}
d p=\lambda+d s_{0} \tag{29}
\end{equation*}
$$

The function $u=p-s_{0}$ belongs to $C^{2}(V)$ and clearly

$$
\begin{equation*}
u \in H\left(\bar{V}_{1}\right) \tag{30}
\end{equation*}
$$

together with $p$ and $s_{0}$. Since

$$
\int_{\alpha}^{*} d p=\int_{V_{0}} d^{*} d p=\int_{V_{0}} \Delta p=0
$$

we have by (7)

$$
\begin{equation*}
\int_{\alpha}^{*} d u=0 \tag{31}
\end{equation*}
$$

From $d u=\lambda \in \Gamma_{e 0}, u \in H D\left(V_{1}\right)$, and Lemma 2, it follows that

$$
\begin{equation*}
u \in C\left(V_{1} \cup \beta\right), \quad u=\text { const on } \delta . \tag{32}
\end{equation*}
$$

On applying Theorem 1 to this $u$, we conclude by (30), (32), and (31) that $L_{1} u=u$ on $\bar{V}_{1}=V_{1} \cup \alpha$. In view of $u=p-s_{0}=p-s$ on $\alpha \cup V_{1}$, we have $L_{1}(p-s)=p-s$ on $V_{1}$, i.e. $p$ satisfies (6).

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