

Extreme laser pulses for non-thermal fusion ignition of hydrogen–boron for clean and low-cost energy

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Review Article

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Abstract

After achieving significant research results on laser-driven boron fusion, the essential facts are presented how the classical very low-energy gains of the initially known thermal ignition conditions for fusion of hydrogen (H) with the boron isotope 11 (HB11 fusion) were bridged by nine orders of magnitudes in agreement with experiments. This is possible under extreme non-thermal equilibrium conditions for ignition by >10 PW-ps laser pulses of extreme power and nonlinear conditions. This low-temperature clean and low-cost fusion energy generation is in crucial contrast to local thermal equilibrium conditions with the advantage to avoid the difficulties of the usual problems with extremely high temperatures.

Introduction

The conversion between different energies (mechanical, electric, chemical or nuclear reaction, thermal, radiation, gravitation, hydro-mechanical) of each one into another follows the principle of the energy conservation in physics. One exception can be by restrictions that thermal energy cannot completely flow on its own to higher temperatures (second law of thermodynamics). The changing of chemical energy per reaction of about electron volt (eV) into mechanic energy by the heat from burning is used since the steam engine or the combustion motor. The processes at local thermal equilibrium (LTE) reduces the efficiency from few to not more than about the range of a dozen of percents at temperatures in the order of 0.1 eV. Without needing heat and combustion, the direct change of chemical energy into electric and then into mechanical energy of motion is possible by batteries or accumulators.

Can this be done with nuclear fusion energy? This was a fundamental question at the Third International Symposium on High Power Laser Science and Technology (HPLST) in a Mini-Workshop by the organizer (Zhou, 2018) in order to demonstrate that pulsed fusion with lasers in plasmas can offer the low-temperature initiation of fusion even at the important high plasma density of about the solid state! This has the advantage to avoid the thermal equilibrium temperature fusion much above 10 million degrees and the unfavorable conditions at extremely low densities at magnetic confinement fusion at more than ten orders of magnitudes lower densities for continuously working reactors like ITER (Bigot, 2017). These developments were fundamental for the special new initiative by the Chinese Academy of Science (Zhang, 2018) for a research program on laser fusion with picoseconds laser pulses at very extreme, ultrahigh power. This is possible now by using the latest developments of picosecond laser pulses for initiation of the ignition of fusion for providing plasma conditions of extreme non-LTE. In addition, it opens for the very first time the fusion of hydrogen with the boron isotope 11 (HB11 fusion) (Hora *et al.*, 2015, 2017a). The difference by non-LTE conditions refers also to laser fusion with nanosecond laser pulses (Hurricane *et al.*, 2014) where the differing LTE condition arrived at other kinds of respectable results yet below break-even.

The energy of nuclear reactions is in the range of 10 MeV in contrast to the eV of chemical reactions. For thermal equilibrium reactions at LTE processes, higher temperatures of >100 million degrees or >10 keV are needed. These are the temperatures needed for fusion power reactors at conditions of thermal equilibrium in contrast to thermal burning of chemicals. A splendid example is the stellarator reactor. After more than 20 years of experiments for continuous fusion of deuterium D in a stellarator, it was possible to produce the very first DD reactions in the Wendelstein experiment (Grieger and Wendelstein Team 1981), see also Section 2.6 of Hora (1991), where equilibrium temperatures of 800 eV (close to 10 million degrees) were reached and continued to arrive at to considerably higher values with the ITER option (Bigot, 2017).

The following summary about the results for the HB11 fusion with picosecond laser pulses was based on the measured ultrahigh acceleration of plasma blocks using lasers by Sauerbrey (1996). This is possible only based on the most exceptional related measurement by Zang et al. (1998) that the picosecond laser pulses had to have a very difficult high contrast ratio such that relativistic self-focusing could be avoided. This was needed for the theory of plasma-block ignition (Hora 1981, 2016; Osman et al., 2004; Hora et al., 2015, 2017a)

Extreme thermal non-equilibrium conditions for boron fusion

The following evaluation to reach the very high-energy gain from laser-driven fusion initiated by ultrahigh power picosecond laser pulses has to be explained, why a change of the HB11 fusion energy gains by nine orders of magnitudes above the classical values was needed. It is easy to understand that bridging such a big gap may be considered as good as impossible. But it has been just measured by laser technology.

The enormous bridging by nine orders of magnitudes is just possible by choosing the non-LTE option of a nearly non-thermal laser fusion based on the non-linear (ponderomotive) force by most extreme laser pulses of picosecond duration (Hora, 1969, 1981) as measured with HB11 fuel (Picciotto et al., 2014; Margarone et al., 2015; Hora, 2015, 2017a).

This can be seen from the equation of motion for the force density \mathbf{f} in a plasma with the thermokinetic pressure p given by the temperature T and the particle density n , and by Maxwell's stress tensor \mathbf{M} of the electric and magnetic fields \mathbf{E} and \mathbf{H} of the laser with inclusion of the complex refractive index varying dynamically on time and space within the plasma (Hora et al., 2015, 2017a)

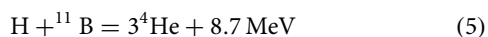
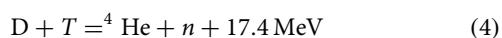
$$\mathbf{f} = -\nabla p + \mathbf{f}_{\text{NL}} = -\nabla p - \nabla \cdot \mathbf{M} \quad (1)$$

$$\begin{aligned} \mathbf{M} = & [\mathbf{E}\mathbf{E} + \mathbf{H}\mathbf{H} - 0.5(\mathbf{E}^2 + \mathbf{H}^2)\mathbf{1} \\ & + (1 + (\partial/\partial t)/\omega)(n^2 - 1)\mathbf{E}\mathbf{E}]/(4\pi) \\ & - (\partial/\partial t)\mathbf{E} \times \mathbf{H}/(4\pi c) \end{aligned} \quad (2)$$

with the unity tensor $\mathbf{1}$. For plane geometry interaction, the non-linear force \mathbf{f}_{NL} is for perpendicular irradiated plasma surface

$$\begin{aligned} f_{\text{NL}} = & -(\partial/\partial x)(\mathbf{E}^2 + \mathbf{H}^2)/(8\pi) \\ = & -(\omega_p/\omega)^2(\partial/\partial x)(E_v^2/\mathbf{n})/(16\pi) \end{aligned} \quad (3)$$

using the dielectric constant \mathbf{n} and the vacuum electric field amplitude E_v of the laser. At thermal equilibrium, the easiest fusion reactions of heavy hydrogen H with superheavy hydrogen D (DT fusion) need a temperature of about 100 million degrees and of HB11 fusion needs more than five times higher temperatures



where next to the harmless helium ${}^4\text{He}$, the generated neutrons n cause a problem of radioactivity in the reaction waste. This is primarily fully excluded in the HB11 reaction.

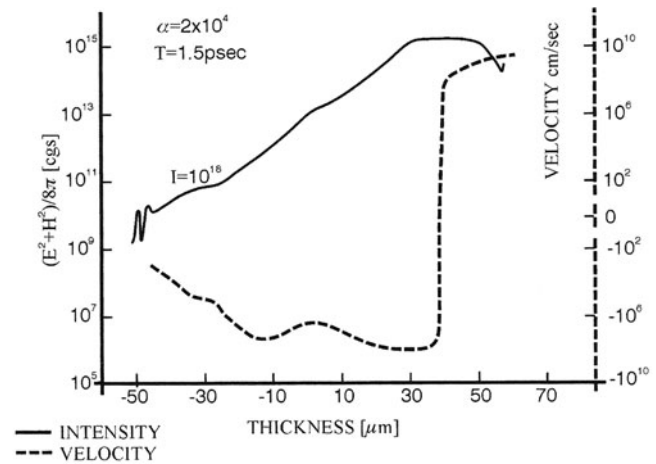


Fig. 1. The 10^{18} W/cm² neodymium glass laser intensity in one-dimensional geometry is incident from the right-hand side on an initially 100 eV hot deuterium plasma slab of initially 0.1 mm thickness whose initial density has a very low reflecting bi-Rayleigh profile, resulting in a laser energy density and a velocity distribution from plasma hydrodynamic computations at time $t = 1.5$ ps of interaction. The driving non-linear force is the negative of the energy density gradient of the laser field $(\mathbf{E}^2 + \mathbf{H}^2)/8\pi$. The dynamic development of temperature and density had accelerated the plasma block of about 15 vacuum wave length thickness of the dielectric enlarged skin layer moving against the laser (positive velocity) and another block into the plasma (negative velocity) showing ultrahigh $>10^{20}$ cm/s² acceleration of the deuterium plasma block to velocities above 10^9 cm/s within the 1.5 ps.

With the aim to minimize heat by using non-LTE or non-equilibrium plasmas, the equation of motion (1) needs such high laser intensities with fields that the non-linear force \mathbf{f}_{NL} produces much higher pressures than the thermal pressure p . The result of an example calculated in 1977 is shown in Figure 1 (Figs 10.18a&b of Hora 1981 or drawn together in Fig. 8.4 of Hora 2016).

The motion of the plasma blocks after 1.5 ps interaction in Figure 1 can be seen in Figure 2. The ultrahigh acceleration of the plasma block when moving against the laser light were first measured by the blue Doppler shift of the spectral lines by Sauerbrey (1996) using the sub-picosecond laser pulses (Strickland and Mourou, 1995) of the comparable intensities as in Figure 1 in agreement with the numerical calculation of 1977 (Hora, 1981) with reconstructing a dielectric swelling of the laser intensity in the irradiated plasma of a usual value near 3 (Hora et al., 2017a, 2007). This measurement was repeated (Földes et al., 2000) after similar accelerations were measured (Badziak et al., 1999) which were in drastic difference to the measurements with a red Doppler shift. The blue shift was proved to be possible only by using most extreme contrast ratios (Zhang et al., 1998; Danson et al., 2018) of the laser pulse in order to avoid relativistic self-focusing. This was measured in a most sophisticated way by Zhang et al. (1998), see Chapter 8.3 of Hora (2016). The red shift happens always at plasma densities lower than critical when instead of the dielectric plasma-block explosion, only ordinary radiation pressure acceleration happens. The red shift was the result of numerous PIC computations where dielectric effects were not included. The first difficult inclusion of densities close to critical densities were successfully demonstrating the dielectric explosion of plasma blocks (Hora et al., 2018) (see Fig. 2).

The discovery of Sauerbrey (1996) of the blue shift on the basis of the measurements by Zhang et al. (1998), Badziak et al. (1999) and Földes et al. (2000) led to the exploring of the picosecond

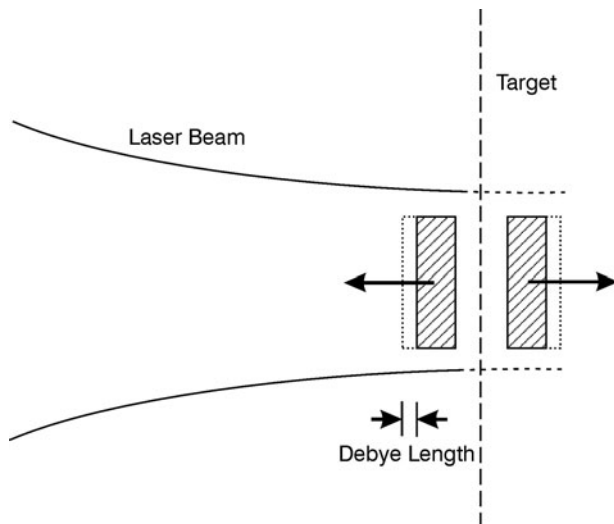


Fig. 2. Schematic representation of skin depth laser interaction where the non-linear force accelerates a plasma block against the laser light and another block toward the target interior. In front of the blocks are electron clouds of the thickness of the effective Debye lengths.

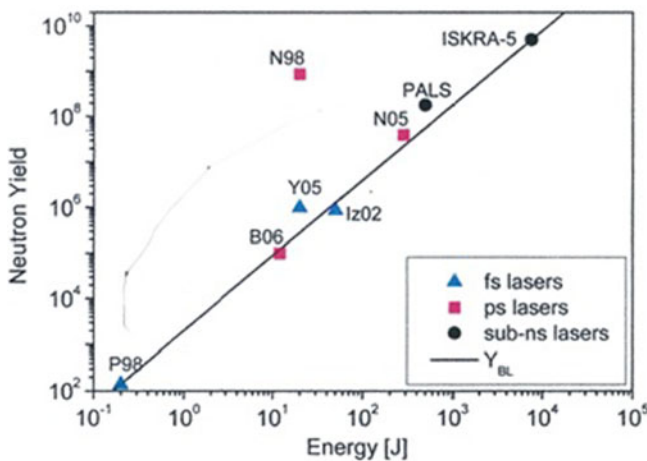


Fig. 3. Measured fusion neutrons emitted from solid targets containing deuterium irradiated by femto to 300 ns laser pulses depending on the energy of the pulses [compiled by Krasa *et al.* (2013)].

initiation of solid density fusion fuel (Hora *et al.*, 2002a, 2002b; Picciotto *et al.*, 2014; Margarone *et al.*, 2015) by updating (Hora, 2009) of the computations by Chu (1972) and Bobin (1974). This was leading to the five orders of magnitudes increased fusion gains of HB11 above the classical values (Hora *et al.*, 2010) and the further four orders increase by the avalanche of the α reaction [Eq. (5)] (Eliezer *et al.*, 2016) for arriving at the results (Hora *et al.*, 2015, 2017a). Only this rather unusual research path led to the design of the basically new type of a clean and low-cost LASER BORON FUSION reactor.

First measurement of low-temperature fusion reactions by laser pulses

Up to 1998, laser fusion was studied mostly by heating of the fuel under conditions of thermal equilibrium. Many measurements were known how the laser interaction with nuclear fusion fuel – mostly with heavy hydrogen (deuterium D) or mixed with

superheavy hydrogen (tritium T) – produced nuclear fusion (DT reactions) detected by the measured number of the generated neutrons. The experiments up to the year 1998 were mostly in the way that laser pulses of about nanosecond duration heated and compressed the fuel under thermal equilibrium. These reactions arrived now at the NIF-experiment in Livermore/California nearly at break-even but where the generated energy is still lower than the input laser energy but reached respectable values. These experiments are based on the conditions of LTE (Hurricane *et al.*, 2014).

Apart from these studies with plasmas at LTE conditions, a most exceptional experiment was in 1998 by Norreys *et al.* (1998) irradiating deuterated targets, and to apply similar shorter than picosecond laser pulses as by Sauerbrey (1996). Compared with very many usually laser-driven experiments, suddenly, the number of fusion neutrons (Krasa *et al.*, 2013), there was a neutron gain nearly four orders of magnitudes higher than under the known heating conditions (Fig. 3, see N98). It was clearly confirmed that the interaction area had a very low temperature exactly showing the conditions of non-LTE. In retrospect with what we know today, this was a typical non-linear force-driven plasma-block fusion (Hora, 1969; Hora *et al.*, 2015, 2017a). At different conditions when irradiating sandwich targets for thermally dominated reactions (N05 of Fig. 3), the measured fusion neutrons were going down to the values in full agreement with all the large number of measured usual fusion gains at thermal equilibrium.

The understanding of the four orders of magnitudes increased fusion gains without much heating is given by the force density in the plasma in the presence of the very high laser intensity. The thermokinetic pressure p can be only a small perturbation against the pressure due to non-linear force by the laser electric and magnetic fields E and H with a much higher energy density $(E^2 + H^2)/(8\pi)$ as in the case of Figure 1, where a first interpretation of a non-LTE, non-linear, force-dominated, low-temperature fusion was elaborated before (Hora *et al.*, 2002a). This is now endorsed in convincing details by the results of laser boron fusion (Hora *et al.*, 2015, 2017a, 2017b, 2018; Hora and Miley, 2018; Eliezer *et al.*, 2016).

Computations based on cases as in Figure 1 but with picosecond high-intensity block ignition were studied since 2000 for explanations of the results of Sauerbrey (1996) by related experiments and computations to analyse this ignition mechanism (Hora *et al.*, 2002a, 2002b, 2002c; Laska *et al.*, 2003; Badziak *et al.*, 2003; Wolowski *et al.*, 2003; Osman *et al.*, 2004; Hora, 2004) the results of which were derived by further international support including the IAEA (UN-International Atomic Energy Agency in Vienna) (Cang *et al.*, 2005; Hora *et al.*, 2007, 2008). The high-energy density non-thermal equilibrium (low-temperature) fusion ignition by plasma blocks from picosecond ultrahigh acceleration were studied also based on the measurements by Sauerbrey (1996) but also at other centers in Poland, Hungary, Shanghai, and few other “fast ignition” projects. When from 2008 (Hora *et al.*, 2008) instead of the most studied DT fusion, the usually five orders of magnitudes lower gaining HB11 (usually at classical thermal equilibrium) were used with non-thermal equilibrium for block ignition, it was most surprising that then gains like those of DT were the result (Hora *et al.*, 2010). This increase by five orders of magnitudes for HB11 was – similar to the measurements of Norreys *et al.* (1998) – due to the non-equilibrium conditions of the block ignition. This was then in advance resulting in the conclusions of Hora *et al.* (2015, 2017a) that could be concluded from the pioneering measurements of laser-produced HB11 fusion (Belyaev *et al.*, 2005; Labaune *et al.*, 2013). The needed non-thermal and non-linear force-dominating conditions

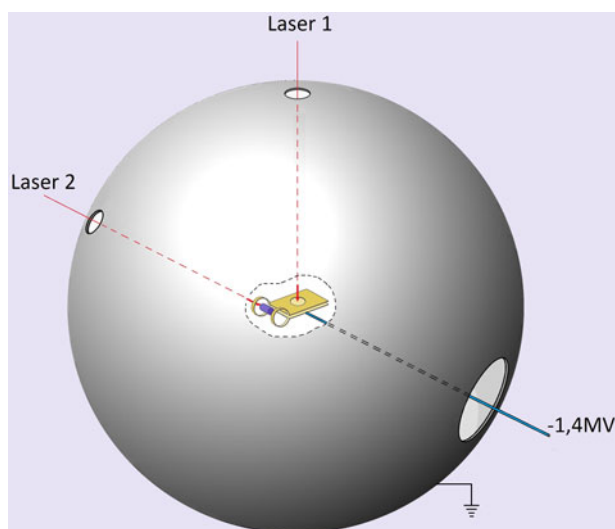


Fig. 4. Clean generator for electric power by laser boron fusion with nanosecond laser 1 to produce the kilotesla magnetic field in the capacitor-coil (Fujioka *et al.* 2013) reaction unit in the center of the spherical generator (Hora *et al.*, 2015, 2017a) and the >10 PW-ps laser pulse 2 to initiate end-on the non-thermal non-linear force-driven reaction in the HB11 fuel cylinder.

for HB11 fusion had been postulated earlier (Hora, 1988) at a conference in Princeton in 1987 for justifying alternative research on aneutronic fusion in contrast to LTE conditions.

Design of laser boron fusion reactor

Though details of the clean and low-cost laser boron fusion reactor have been discussed before (Hora *et al.*, 2017a, 2017b, 2017c), the following points should be underlined. The reaction unit in the center of the reaction sphere (Fig. 4) with the cylindrical solid density fuel of hydrogen and boron ^{11}B is ignited by the direct-drive pulses of the picosecond laser 2 of more than 10 PW power. There is a flexibility about the details of the pulse profile that on the one hand is a problem how to generate this (Danson *et al.*, 2018), but on the other hand, it permits modifications for producing optimized conditions. The same optimization is possible by the time dependence of the kilotesla magnetic field in the capacitor coil depending on the quality of the nanosecond laser pulse 1 of at least kJ energy. The laser technology permits a wide range of variations.

Another flexibility in modifications is the preparation of the fuel density within the first 10 μm at the target for the direct-drive interaction area of laser 2 at the end area of the fuel. This influences the initiation of the picosecond ignition process for the propagating reaction plane through the fuel as it was the result of a very large number of cases for calculation with the genuine multi-fluid computations (Lalousis and Hora, 1983; Hora *et al.*, 1984, 2015). These details are also of importance when the non-linear force acceleration process in target layers is used for the generation of space charge neutral ion beams, with more than million times higher ion densities than the best classical accelerator can produce (Hora *et al.*, 2017a, 2007) for ion energies of few hundred MeV energy cancer treatment or for space craft propulsion (Hora *et al.*, 2011; Hora and Miley 2018; Lalousis *et al.*, 2013; Hoffmann *et al.*, 2018). Even nearly monoenergetic proton beams above GeV (Xu *et al.*, 2016; 2018) are related to the results of Hora *et al.* (2018).

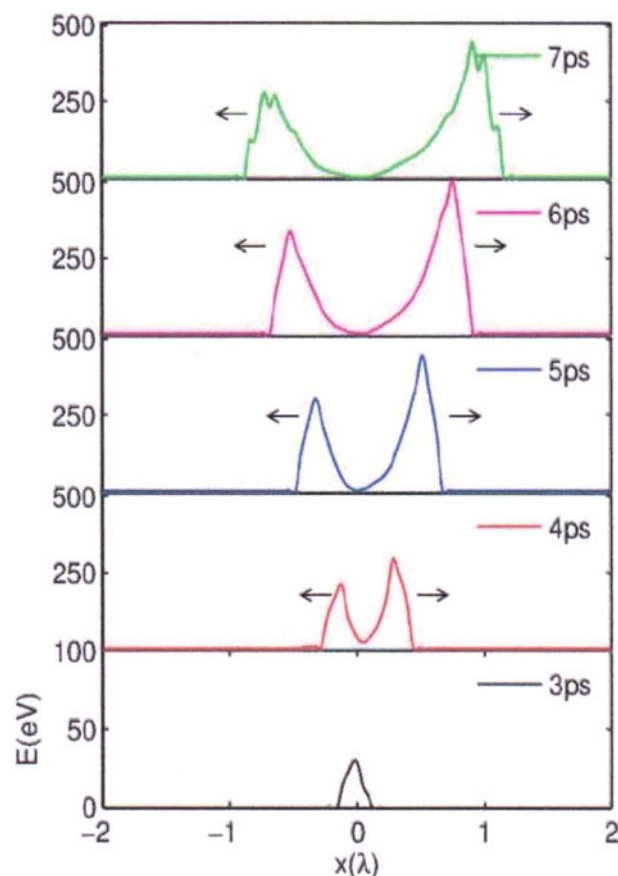


Fig. 5. PIC computation of the dielectric explosion of the plasma blocks at plasma densities close to the critical density (Hora *et al.*, 2018) confirming of the result of hydrodynamic computations (Fig. 1) as necessary process for the very rare measurement of the blue Doppler shift in the reflected light (Sauerbrey, 1996; Zhang *et al.*, 1998; Hora, 1969, 1981).

It is shown in the 2nd and 3rd Sections that the ultrahigh plasma-block acceleration is the essential mechanism for the non-LTE process of the low-temperature ignition (Hora *et al.*, 2015, 2017a, 2017b, 2017c, 2018) of the laser boron fusion for the new reactor design. This was possible only by the extreme contrast ratio (Danson *et al.*, 2018) of the laser pulses as it was realized in retrospect only (Hora, 2016) by the blue Doppler shift of spectral lines (Sauerbrey, 1996; Zhang *et al.*, 1998; Földes *et al.*, 2000). The PIC computation with inclusion of dielectric response was difficult and succeeded only recently (Xu *et al.*, 2016, 2018; Hora *et al.*, 2018) to confirm the earlier hydrodynamic derivation of basic requirements for the operation of the reactor design of Figure 5.

A further problem had to be solved with respect to the secondary neutron production of HB11 fusion. The primary HB11 reaction [Eq. (5)] is indeed free from neutron generation, but it is well known that secondary reactions produce about a neutron per thousand generated α particles of helium. The elimination of these neutrons is possible by screening the reactor sphere of Figure 4 by an equipment of about 10 cm thickness to the level that the reactor works clearly below the level of any radioactive pollution problem (Eliezer *et al.*, 2017). This neutron capturing, for example, for a reactor for 100 PW generation of electricity does not need a replacement of the screening material for a few years. The advantage compared with the radioactive waste

problem of fission reactors is that the neutrons decay with a half-life of 14.69 min into harmless electrons and protons.

Patenting procedures (Hora, 2014) are based on HB11 reactions with some similarity to (Margarone *et al.*, 2013) based on non-linear force ultrahigh acceleration processes. The essential difference consists in the fact that the patent (Hora, 2014) is a combination with the otherwise earlier known kilotesla magnetic field generated by capacitor coils (Fujioka *et al.*, 2013) for the cylindrical trapping of the reaction that is necessary for the reactor of Figure 4 to provide non-LTE low-temperature laser fusion.

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