

# Scattering Length and the Spectrum of $-\Delta + V$

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*Abstract.* Given a non-negative, locally integrable function  $V$  on  $\mathbb{R}^n$ , we give a necessary and sufficient condition that  $-\Delta + V$  have purely discrete spectrum, in terms of the scattering length of  $V$  restricted to boxes.

## 1 Introduction

It is a classical result of K. Friedrichs [F] that if  $V \in L^1_{\text{loc}}(\mathbb{R}^n)$  and  $V \geq 0$ , then  $-\Delta + V$  yields a positive self-adjoint operator on  $L^2(\mathbb{R}^n)$ , and its spectrum is discrete if  $V(x) \rightarrow +\infty$  as  $|x| \rightarrow \infty$ . A. Molchanov [Mol] produced a necessary and sufficient condition for such an operator to have discrete spectrum. His condition takes the form

$$(1.1) \quad \inf_F \int_{Q_{b,\xi} \setminus F} V(x) dx \rightarrow \infty, \quad \text{as } |\xi| \rightarrow \infty,$$

for each  $b \in (0, 1]$ , where  $Q_{b,\xi}$  is the  $n$ -dimensional cube of the form

$$(1.2) \quad Q_{b,\xi} = \left\{ x \in \mathbb{R}^n : \xi_j - \frac{b}{2} \leq x_j \leq \xi_j + \frac{b}{2} \right\}.$$

(We henceforth say  $Q_{b,\xi}$  is the cube with sidelength  $b$  and center  $\xi$ .) In (1.1),  $F$  runs over the “negligible” subsets of  $Q_{b,\xi}$ , defined by the condition  $\text{cap } F \leq \gamma \text{cap } Q_{b,\xi}$ . In [Mol],  $\gamma$  was taken to be a particular (small) constant  $\gamma_n$ .

Recent important work of V. Maz’ya and M. Shubin [MS] provides a cleaner form for the necessary and sufficient condition. In particular,  $\gamma$  can be given any value in  $(0, 1)$ . Furthermore, they allow  $\gamma = \gamma(b)$ , possibly decaying to 0 as  $b \rightarrow 0$ , as long as  $b^{-2}\gamma(b) \rightarrow \infty$ .

Our purpose here is to produce an alternative formulation of a necessary and sufficient condition that  $-\Delta + V$  have discrete spectrum (given  $V \geq 0$ ,  $V \in L^1_{\text{loc}}(\mathbb{R}^n)$ ). Our result is phrased in terms of “scattering length,” a quantity  $\Gamma(\nu)$  associated to integrable  $\nu \geq 0$  that is somewhat parallel to the notion of capacity of a set. In fact, if  $K$  is a compact set satisfying a mild regularity condition,

$$(1.3) \quad \text{cap } K = \lim_{r \rightarrow +\infty} \Gamma(r\chi_K),$$

where  $\chi_K$  denotes the characteristic function of  $K$ . We will recall the definition of  $\Gamma(\nu)$  in §2. Our main result is the following.

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**Theorem 1.1** Given  $V \geq 0$ ,  $V \in L^1_{\text{loc}}(\mathbb{R}^n)$ , the following three conditions are equivalent.

- (1)  $-\Delta + V$  has purely discrete spectrum on  $L^2(\mathbb{R}^n)$ .
- (2) Given  $A \in (0, \infty)$ , there exists  $b = b(A) \in (0, 1]$  and  $R \in (0, \infty)$  such that

$$\Gamma(b^2V_{b,\xi}) \geq Ab^2, \quad \text{for } |\xi| \geq R.$$

- (3) Given  $A \in (0, \infty)$ , there exists  $b_0 = b_0(A) \in (0, 1]$  and  $R: (0, b_0] \rightarrow (0, \infty)$  such that

$$\Gamma(b^2V_{b,\xi}) \geq Ab^2, \quad \text{for } b \in (0, b_0], |\xi| \geq R(b).$$

Here  $V_{b,\xi}$  is a positive function supported on the unit cube  $Q = Q_{1,0}$ , given by

$$(1.4) \quad V_{b,\xi}(x) = V(bx + \xi), \quad x \in Q.$$

The rest of this paper is structured as follows. In §2 we define  $\Gamma(\nu)$  for positive, integrable  $\nu$  and review some of its crucial properties. In §3 we prove that (2)  $\Rightarrow$  (1) in Theorem 1.1, and in §4 we prove that (1)  $\Rightarrow$  (3). Clearly (3)  $\Rightarrow$  (2), so this will prove Theorem 1.1. There is one result in §4, Lemma 4.2, whose proof is presented in §5.

**Remark** In the formal limit  $V = +\infty$  on  $K = \mathbb{R}^n \setminus \Omega$ , where one considers  $-\Delta$  on  $L^2(\Omega)$ , with the Dirichlet boundary condition on  $\partial\Omega$ , the condition (3) of Theorem 1.1 becomes that for each  $A \in (0, \infty)$ , there exists  $b_0 = b_0(A) \in (0, 1]$  and  $R: (0, b_0] \rightarrow (0, \infty)$  such that

$$(1.5) \quad \text{cap } K_{b,\xi} \geq Ab^2(\text{cap } Q_{b,\xi}), \quad \forall b \in (0, b_0], |\xi| \geq R(b),$$

where  $K_{b,\xi} = K \cap Q_{b,\xi}$ . This coincides with one of the criteria (necessary and sufficient) for discreteness presented in [MS, Remark 2.7].

## 2 Scattering Length

Here we define the scattering length  $\Gamma(\nu)$  of a positive integrable potential  $\nu$  and review some of its properties. Our material is taken from [T], which in turn was influenced by results on scattering length presented in [K, KL]. For simplicity we take  $n \geq 3$ .

To such  $\nu$  we associate the capacity potential  $U_\nu$  and the scattering length  $\Gamma(\nu)$  as follows. First assume that  $\nu \in L^2(\mathbb{R}^n)$  and has support in a compact set  $K$ , as well as  $\nu \geq 0$ . We define  $U_\nu$  by

$$(2.1) \quad U_\nu(x) = \lim_{\varepsilon \searrow 0} (\varepsilon + \nu - \Delta)^{-1} \nu(x).$$

It is shown that this limit exists in  $L^2_{\text{loc}}(\mathbb{R}^n)$  and satisfies

$$(2.2) \quad 0 \leq U_\nu \leq 1, \quad \nu \leq w \Rightarrow U_\nu \leq U_w.$$

The existence proof in [K, KL] involves producing the formula

$$(2.3) \quad U_\nu(x) = E_x \left\{ 1 - \exp \left( - \int_0^\infty \nu(b(\tau)) d\tau \right) \right\},$$

where  $E_x$  is expectation with respect to Wiener measure on Brownian paths  $b$  starting at  $x$ ; see also [T, p. 292] for a derivation of this formula.

The function  $U_\nu$  solves the PDE

$$(2.4) \quad \Delta U_\nu = -\nu(1 - U_\nu).$$

It follows that  $-\Delta U_\nu = \mu_\nu$  is a positive measure on  $\mathbb{R}^n$ . We set

$$(2.5) \quad \Gamma(\nu) = \int d\mu_\nu(x).$$

Some basic results on  $\Gamma(\nu)$  include:

$$(2.6) \quad \begin{aligned} \nu \leq w &\implies \Gamma(\nu) \leq \Gamma(w), \\ \Gamma(\nu + w) &\leq \Gamma(\nu) + \Gamma(w), \\ \nu_n \nearrow \nu &\implies \Gamma(\nu_n) \nearrow \Gamma(\nu), \\ \Gamma(\nu) &\leq \|\nu\|_{L^1}. \end{aligned}$$

We also have

$$(2.7) \quad \|\nabla U_\nu\|_{L^2(\mathbb{R}^n)}^2 = \int_{\mathbb{R}^n} U_\nu(x) d\mu_\nu(x) \leq \Gamma(\nu),$$

and, for any ball  $B \subset \mathbb{R}^n$ ,

$$(2.8) \quad \|U_\nu\|_{L^1(B)} \leq \alpha(B)\Gamma(\nu).$$

These results are established in [T, Propositions 1.2–1.6]. They allow us to define  $U_\nu$  and  $\Gamma(\nu)$  for positive  $\nu \in L^1(\mathbb{R}^n)$ , having

$$(2.9) \quad \nu_n \nearrow \nu, \nu_n \in L^2_{\text{comp}}(\mathbb{R}^n) \implies U_{\nu_n} \nearrow U_\nu, \Gamma(\nu_n) \nearrow \Gamma(\nu).$$

We now give two key estimates, established in [T], which connect scattering length to eigenvalue estimates. Suppose  $\nu \geq 0$  is an integrable function supported on  $Q$ , the cube of sidelength 1 centered at 0. Let  $\lambda_1(\nu) \in [0, \infty)$  denote the smallest eigenvalue of  $-\Delta + \nu$ , with the Neumann boundary condition, on  $L^2(Q)$ . The following result summarizes [T, Propositions 2.2–2.3].

**Proposition 2.1** *There exists  $C_n \in (0, \infty)$  such that*

$$(2.10) \quad \lambda_1(\nu) \geq C_n \Gamma(\nu).$$

*Furthermore, there exist  $E_n, \tilde{C}_n \in (0, \infty)$  such that*

$$(2.11) \quad \Gamma(\nu) \leq E_n \implies \lambda_1(\nu) \leq \tilde{C}_n \Gamma(\nu).$$

We refer to [T, pp. 295–297] for proofs of these results. We mention that (2.11) is proven by an apt choice of test function in the variational characterization of  $\lambda_1(\nu)$ , while (2.10) is proven by examining the decay rate for  $e^{-tL_N}$ , where  $L_N$  denotes  $-\Delta + \nu$ , with the Neumann boundary condition, on  $L^2(Q)$ .

### 3 Sufficient Condition for Discrete Spectrum

The following result yields the implication (2)  $\Rightarrow$  (1) in Theorem 1.1.

**Proposition 3.1** *Take  $A \in (0, \infty)$  and let  $C_n$  be as in (2.10). Assume that there exists  $b = b(A) \in (0, 1]$  and  $R = R(A) \in (0, \infty)$  such that*

$$(3.1) \quad C_n \Gamma(b^2 V_{b,\xi}) \geq Ab^2, \quad \text{for } |\xi| \geq R.$$

Then

$$(3.2) \quad \text{ess spec}(-\Delta + V) \subset [A, \infty).$$

**Proof** Let  $Q_{b,\xi}$  denote the cube of edge  $b$ , center  $\xi$ , as in (1.2), and let  $L_{b,\xi}$  denote the operator  $-\Delta + V$  on  $L^2(Q_{b,\xi})$ , with the Neumann boundary condition. A standard argument involving Rellich's theorem shows that, if there exists  $R = R(A)$  such that

$$(3.3) \quad \text{spec } L_{b,\xi} \subset [A, \infty), \quad \text{for } |\xi| \geq R,$$

then (3.2) holds. Now  $L_{b,\xi}$  is unitarily equivalent to the operator

$$(3.4) \quad -b^{-2}\Delta + V_{b,\xi} = b^{-2}(-\Delta + b^2 V_{b,\xi}),$$

on  $L^2(Q)$  ( $Q$  denoting the cube of edge 1, center 0), where

$$(3.5) \quad V_{b,\xi}(x) = V(bx + \xi), \quad x \in Q,$$

and one places the Neumann boundary condition on the operator (3.4). Now, by Proposition 2.1, the spectrum of this operator is bounded below by

$$(3.6) \quad C_n b^{-2} \Gamma(b^2 V_{b,\xi}),$$

so Proposition 3.1 is proven. ■

### 4 Necessary Condition for Discrete Spectrum

It is convenient to set up some notation. Given a cube  $Q_\nu \subset \mathbb{R}^n$ , we denote by

$$(4.1) \quad \lambda_D^{Q_\nu}(-\Delta + V), \quad \text{resp.,} \quad \lambda_N^{Q_\nu}(-\Delta + V),$$

the smallest eigenvalue of  $-\Delta + V$  on  $L^2(Q_\nu)$ , where we impose, respectively, the Dirichlet or Neumann boundary condition on  $\partial Q_\nu$ . As before, let  $Q_{b,\xi}$  denote the cube of edge  $b$ , center  $\xi$ , as in (1.2). We continue to assume  $V \geq 0$  and  $V \in L^1_{\text{loc}}(\mathbb{R}^n)$ .

**Lemma 4.1** *If  $-\Delta + V$  has discrete spectrum on  $L^2(\mathbb{R}^n)$ , then for each  $b \in (0, 1]$ ,*

$$(4.2) \quad \lambda_D^{Q_{b,\xi}}(-\Delta + V) \longrightarrow +\infty, \quad \text{as } |\xi| \rightarrow \infty.$$

**Proof** As is well known,  $-\Delta + V$  has discrete spectrum on  $L^2(\mathbb{R}^n)$  if and only if the set

$$(4.3) \quad X = \{u \in H^1(\mathbb{R}^n) : \|\nabla u\|_{L^2}^2 + \|V^{1/2}u\|_{L^2}^2 \leq 1\}$$

is compact in  $L^2(\mathbb{R}^n)$ . In turn, such compactness implies

$$(4.4) \quad \int_{|x| \geq R} |u(x)|^2 dx \leq \varepsilon(R), \quad \forall u \in X,$$

where  $\varepsilon(R) \rightarrow 0$  as  $R \rightarrow \infty$ . If we restrict attention to  $u \in H_0^1(Q_{b,\xi})$ , this gives (4.2). ■

In the following lemma,  $Q = Q_{1,0}$ , the unit cube centered at 0.

**Lemma 4.2** *There exists  $A_n \in (0, \infty)$  and  $B_n: [A_n, \infty) \rightarrow (0, \infty)$ , such that  $B_n(A) \rightarrow \infty$  as  $A \rightarrow \infty$ , and such that whenever  $v \in L^1(Q)$  is non-negative and whenever  $A \geq A_n$ ,*

$$(4.5) \quad \lambda_D^Q(-\Delta + v) \geq A \implies \lambda_N^Q(-\Delta + v) \geq B_n(A).$$

Such a result is established in [Mol]; a proof is also given in [KS, Lemma 2.9]. For the convenience of the reader, we present yet another proof of Lemma 4.2 in §5. Granted the result, we deduce from Lemma 4.1 the following.

**Corollary 4.3** *If  $-\Delta + V$  has discrete spectrum on  $L^2(\mathbb{R}^n)$ , then, for each  $b \in (0, 1]$ ,*

$$(4.6) \quad \lambda_N^{Q_{b,\xi}}(-\Delta + V) \longrightarrow +\infty, \quad \text{as } |\xi| \rightarrow \infty.$$

Note that the left side of (4.6) is equal to

$$(4.7) \quad b^{-2} \lambda_N^Q(-\Delta + b^2 V_{b,\xi}).$$

We are now ready to prove the implication (1)  $\implies$  (3) in Theorem 1.1. Given  $A \in (0, \infty)$ , pick  $b_0 = b_0(A)$  so small that (2.11) applies, so that

$$(4.8) \quad \Gamma(v) \leq b_0^2 A \implies \lambda_N^Q(-\Delta + v) \leq \tilde{C}_n \Gamma(v).$$

Consequently, for  $b \in (0, b_0]$ ,

$$(4.9) \quad \Gamma(b^2 V_{b,\xi}) \leq Ab^2 \implies \lambda_N^Q(-\Delta + b^2 V_{b,\xi}) \leq \tilde{C}_n \Gamma(b^2 V_{b,\xi}) \leq \tilde{C}_n Ab^2.$$

Now, by (4.6)–(4.7), we cannot have the bound  $\lambda_N^Q(-\Delta + b^2 V_{b,\xi}) \leq \tilde{C}_n Ab^2$  for large  $|\xi|$ , so consequently we cannot have the bound  $\Gamma(b^2 V_{b,\xi}) \leq Ab^2$  for large  $|\xi|$ . The proof of Theorem 1.1 is complete, modulo the proof of Lemma 4.2, which will be given in the next section.

### 5 Proof of Lemma 4.2

Given non-negative  $v \in L^1(Q)$ , let us denote by  $L_D$  the operator  $L = -\Delta + v$  on  $L^2(Q)$  with the Dirichlet boundary condition and by  $L_N$  the operator with the Neumann boundary condition. We assume

$$(5.1) \quad \lambda = \lambda_D^Q(-\Delta + v),$$

the smallest eigenvalue of  $L_D$ , is large, and we want to estimate the smallest eigenvalue of  $L_N$ . We will estimate various “heat semigroups.” For  $x, y \in Q, t > 0$ , set

$$(5.2) \quad \begin{aligned} p_D(t, x, y) &= e^{-tL_D} \delta_y(x), & p_N(t, x, y) &= e^{-tL_N} \delta_y(x), \\ p_Q(t, x, y) &= e^{t\Delta_N} \delta_y(x), & p_0(t, x, y) &= (4\pi t)^{-n/2} e^{-|x-y|^2/4t}. \end{aligned}$$

Here  $\Delta_N$  denotes the Laplace operator on  $L^2(Q)$ , with the Neumann boundary condition. It will be convenient to note the following inequalities:

$$(5.3) \quad p_D(t, x, y) \leq p_0(t, x, y), \quad p_N(t, x, y) \leq p_Q(t, x, y).$$

We want to estimate  $p_N(t, x, y)$ , but first we will estimate  $p_D(t, x, y)$ . Let us fix  $a \in (0, 1)$  and set

$$(5.4) \quad \tau = \lambda^{-a}.$$

Using (5.3) we have

$$(5.5) \quad \|e^{-\tau L_D}\|_{\mathcal{L}(L^1, L^2)} \leq (4\pi\tau)^{-n/4}, \quad \|e^{-\tau L_D}\|_{\mathcal{L}(L^2, L^\infty)} \leq (4\pi\tau)^{-n/4},$$

while (5.1) gives

$$(5.6) \quad \|e^{-\tau L_D}\|_{\mathcal{L}(L^2, L^2)} \leq e^{-\tau\lambda}.$$

Hence

$$(5.7) \quad \|e^{-3\tau L_D} \delta_y\|_{L^\infty} \leq C\tau^{-n/2} e^{-\tau\lambda} = C\lambda^{an/2} e^{-\lambda^{1-a}}.$$

In other words,

$$(5.8) \quad 0 \leq p_D(3\tau, x, y) \leq C\lambda^{an/2} e^{-\lambda^{1-a}}, \quad \forall x, y \in Q.$$

Next, we estimate  $V_y(t, x) = p_N(t, x, y) - p_D(t, x, y)$ , for  $t \in [0, 3\tau]$ . We have

$$(5.9) \quad (\partial_t - L)V_y = 0 \text{ on } \mathbb{R}^+ \times Q, \quad V_y(0, x) = 0,$$

and  $x \in \partial Q \implies V_y(t, x) = p_N(t, x, y)$ . Hence

$$(5.10) \quad x \in \partial Q \implies 0 \leq V_y(t, x) \leq p_Q(t, x, y).$$

Let us define the set  $\Omega_\tau \subset Q$  by

$$(5.11) \quad \Omega_\tau = \{y \in Q : \text{dist}(y, \partial Q) \geq \tau^{1/3}\}.$$

It is clear that, if  $\lambda$  is sufficiently large, so  $\tau$  is sufficiently small,

$$(5.12) \quad x \in \partial Q, y \in \Omega_\tau, t \in [0, 3\tau] \Rightarrow p_Q(t, x, y) \leq Ce^{-\lambda^{a/4}},$$

so applying the maximum principle to (5.9)–(5.10) gives

$$(5.13) \quad V_y(t, x) \leq Ce^{-\lambda^{a/4}}, \quad \text{for } x \in Q, y \in \Omega_\tau, t \in [0, 3\tau],$$

and hence, by (5.8), if we take  $a = 4/5$  and assume  $\lambda$  is sufficiently large,

$$(5.14) \quad 0 \leq p_N(3\tau, x, y) \leq C\lambda^{2n/5}e^{-\lambda^{1/5}}, \quad \forall x \in Q, y \in \Omega_\tau.$$

Now, using the semigroup property of  $e^{tL_N}$  and the fact that  $\|e^{-tL_N}\|_{\mathcal{L}(L^\infty, L^\infty)} \leq 1$ , we deduce that

$$(5.15) \quad 0 \leq p_N(t, x, y) \leq C\lambda^{2n/5}e^{-\lambda^{1/5}}, \quad \forall x \in Q, y \in \Omega_\tau, t \geq 3\tau.$$

In particular, if  $\lambda$  is large enough that  $3\tau = 3\lambda^{-4/5} < 1$ , the estimate (5.15) applies with  $t = 1$ . On the other hand, we can use (5.3) to obtain

$$(5.16) \quad p_N(1, x, y) \leq p_Q(1, x, y) \leq C, \quad \forall x \in Q, y \in Q \setminus \Omega_\tau.$$

It follows that

$$(5.17) \quad \begin{aligned} \int_Q p_N(1, x, y) dy &\leq C\lambda^{2n/5}e^{-\lambda^{1/5}} + C \text{Vol}(Q \setminus \Omega_\tau) \\ &\leq C\lambda^{2n/5}e^{-\lambda^{1/5}} + C\lambda^{-4/15}. \end{aligned}$$

Of course  $p_N(1, x, y) = p_N(1, y, x)$ , so there is a similar bound on  $\int_Q p_N(1, x, y) dx$ . Hence we deduce that

$$(5.18) \quad \|e^{-L_N}\|_{\mathcal{L}(L^2, L^2)} \leq C\lambda^{2n/5}e^{-\lambda^{1/5}} + C\lambda^{-4/15} = \Phi(\lambda).$$

It follows that

$$(5.19) \quad \lambda_N^Q(-\Delta + \nu) \geq \log \frac{1}{\Phi(\lambda)},$$

and Lemma 4.2 is proven.

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