ON THE REPRESENTATION OF FUNCTIONS AS FOURIER TRANSFORMS

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If $f \in L_p$ ($-\infty$, ∞), $1 < p \le 2$, then f has a Fourier-Plancherel transform $F \in L_q$ (- ∞ , ∞) where $p^{-1} + q^{-1} = 1$. Also if $|x|^{1-2/q} f(x) \in L_q$ (- ∞ , ∞), $q \ge 2$, then f has a Fourier-Plancherel transform $F \in L_q$ ($-\infty$, ∞). These results can be found in (2, Theorems 74 and 79). In neither case, however, does the collection of transforms cover L_q , except when $p = q = 2$, and in neither case, with the same exception, has the collection of transforms been characterized.

Further, if $f \in L_p(-\infty, \infty)$, $1 < p \le 2$, then its transform *F* has the property $|x|^{1-2/p} F(x) \in L_p(-\infty, \infty)$ (see 2, Theorem 80) but, except when $p = 2$, the collection of transforms does not cover the set of functions with this property, and again, except when $p = 2$, the collection of transforms has not been characterized.

Our object here is to find such characterizations, and this is done for the various cases in Theorems 1, 2, and 3 below. This characterization is given in terms of an operator

$$
\mathfrak{F}_{k,i}[F] = \frac{(-ik/t)^{k+1}}{(2\pi)^k} \int_{-\infty}^{\infty} \frac{1}{(x-ik/t)^{k+1}} F(x) dx, \quad k = 1, 2, \ldots.
$$

It transpires that this operator is an inversion operator for the Fourier transform, and its inversion theory will be the subject of another paper.

THEOREM 1. A necessary and sufficient condition that a function $F \in$ $L_q(-\infty, \infty), q \geqslant 2$, be the Fourier transform of a function in $L_p(-\infty, \infty)$, *with* $p^{-1} + q^{-1} = 1$, is that there exist a constant *M* such that

$$
\int_{-\infty}^{\infty} |\mathfrak{F}_{k, l}[F]|^p dt \leq M, \qquad k = 1, 2, \ldots.
$$

Proof of necessity. Suppose *F* is the Fourier transform of $f \in L_p$ ($-\infty, \infty$).

Now an easy calculation shows that for $k = 1, 2, \ldots$,

$$
\frac{1}{(2\pi)^{\frac{1}{2}}}\int_{-\infty}^{\infty}\frac{1}{(x-ik/t)^{k+1}}e^{-ixy}dx = \begin{cases} -(2\pi)^{\frac{1}{2}}(-i)^{k+1}y^{k}e^{ky/t}/k!, & y < 0, t > 0, \\ (2\pi)^{\frac{1}{2}}(-i)^{k+1}y^{k}e^{ky/t}/k!, & y > 0, t < 0, \\ 0, & yt > 0. \end{cases}
$$

Hence, since for each $t \neq 0$ and each $k = 1, 2, ..., (x - ik/t)^{-(k+1)} \in$ L_p ($-\infty$, ∞), we have from (2, Theorem 75) that

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$$
\mathfrak{F}_{k, t}[F] = \frac{(-ik/t)^{k+1}}{(2\pi)^k} \int_{-\infty}^{\infty} \frac{1}{(x-ik/t)^{k+1}} F(x) dx
$$

=
$$
\begin{cases} (k/t)^{k+1} (k!)^{-1} \int_{0}^{\infty} e^{-ky/t} y^k f(y) dy, & t > 0 \\ (k/|t|)^{k+1} (k!)^{-1} \int_{-\infty}^{0} e^{-ky/t} |y|^k f(y) dy, & t < 0. \end{cases}
$$

Thus, using Hölder's inequality, we have for $t > 0$

$$
|\mathfrak{F}_{k, t}[F]| \leq (k/t)^{k+1}(k!)^{-1} \Biggl\{ \int_0^{\infty} e^{-ky/t} y^k |f(y)|^p dy \Biggr\}^{1/p} \Biggl\{ \int_0^{\infty} e^{-ky/t} y^k dy \Biggr\}^{1/q} \\
= \Biggl\{ (k/t)^{k+1}(k!)^{-1} \int_0^{\infty} e^{-ky/t} |f(y)|^p dy \Biggr\}^{1/p},
$$

and consequently,

$$
\int_0^{\infty} |\mathfrak{F}_{k, t}[F]|^p dt \leq \frac{k^{k+1}}{k!} \int_0^{\infty} t^{-(k+1)} dt \int_0^{\infty} e^{-ky/t} y^k |f(y)|^p dy
$$

= $\frac{k^{k+1}}{k!} \int_0^{\infty} y^k |f(y)|^p dy \int_0^{\infty} t^{-(k+1)} e^{-ky/t} dt = \int_0^{\infty} |f(y)|^p dy.$

A similar calculation for *t <* 0 shows that

$$
\int_{-\infty}^0 |\mathfrak{F}_{k,\,i}[F]|^p dt \leqslant \int_{-\infty}^0 |f(y)|^p dy,
$$

and hence

$$
\int_{-\infty}^{\infty} |\mathfrak{F}_{k. t}[F]|^p dt \leqslant \int_{-\infty}^{\infty} |f(y)|^p dy = M.
$$

Proof of sufficiency. For $s > 0$ let

$$
g_{+}(s) = -(2\pi)^{-1} i \int_{-\infty}^{\infty} \frac{1}{x - i s} F(x) dx,
$$

and

$$
g_{-}(s) = (2\pi)^{-\frac{1}{2}} i \int_{-\infty}^{\infty} \frac{1}{x + i s} F(x) dx,
$$

and denote by $L_{k,t}$ the Widder-Post inversion operator for the Laplace transformation; that is

$$
L_{k, i}[g] = (-1)^{k} (k/t)^{k+1} g^{(k)}(k/t)/k!, \qquad k = 1, 2, \ldots.
$$

Now if $s \ge \delta > 0$, and $k = 1, 2, \ldots$, then

$$
|(x\pm is)^{-(k+1)} F(x)| \leq (x^2+\delta^2)^{-(k+1)/2}|F(x)| \in L_1 \; (-\infty, \infty),
$$

since from Hölder's inequality

$$
\int_{-\infty}^{\infty} (x^2 + \delta^2)^{-(k+1)/2} |F(x)| dx
$$

\$\leqslant \left\{ \int_{-\infty}^{\infty} (x^2 + \delta^2)^{-p(k+1)/2} dx \right\}^{1/p} \cdot \left\{ \int_{-\infty}^{\infty} |F(x)|^q dx \right\}^{1/q} < \infty .

 \bullet

Hence by (1, Corollary 39.2), $g_{+}(s)$ has derivatives of all orders in $0 < s < \infty$, and these derivatives can be calculated by differentiating under the integral sign. Thus for $t > 0$,

$$
L_{k, t}[g_{+}] = \frac{(-ik/t)^{k+1}}{(2\pi)^{k}} \int_{-\infty}^{\infty} \frac{1}{(x - ik/t)^{k+1}} F(x) dx = \mathfrak{F}_{k, t}[F],
$$

and

$$
L_{k, t}[g_{-}] = \frac{(ik/t)^{k+1}}{(2\pi)^{\frac{k}{2}}} \int_{-\infty}^{\infty} \frac{1}{(x+ik/t)^{k+1}} F(x) dx = \mathfrak{F}_{k, -t}[F],
$$

so that

$$
\int_0^\infty |L_{k,\,t}[g_+]|^p dt = \int_0^\infty |\mathfrak{F}_{k,\,t}[F]|^p dt \leqslant M, \qquad k = 1, 2, \ldots,
$$

and

$$
\int_0^\infty |L_{k,\,t}[g_-]|^p dt = \int_0^\infty |\mathfrak{F}_{k,-t}[F]|^p dt \leqslant M, \qquad k = 1, 2, \ldots.
$$

Further $g_{\pm}(s) \rightarrow 0$ as $s \rightarrow \infty$. For from Hölder's inequality we have

$$
|g_{\pm}(s)| \leq (2\pi)^{-\frac{1}{2}} \left\{ \int_{-\infty}^{\infty} (x^2 + s^2)^{-p/2} dx \right\}^{1/p} \left\{ \int_{-\infty}^{\infty} |F(x)|^q dx \right\}^{1/q} = 0(s^{-1/q}).
$$

Hence by (3, Chapter 7, Theorem 15a) there are functions f_+ and f_- in $L_p(0, \infty)$ such that

$$
g_{+}(s) = \int_{0}^{\infty} e^{-st} f_{+}(t) dt, \qquad s > 0,
$$

and

$$
g_{-}(s) = \int_{0}^{\infty} e^{-st} f_{-}(t) dt, \qquad s > 0.
$$

Let

$$
f(t) = \begin{cases} f_+(t), & t > 0, \\ f_-(-t), & t < 0. \end{cases}
$$

Then clearly $f \in L_p$ ($-\infty$, ∞) and hence by (2, Theorem 74) f has a Fourier transform $F^* \in L_q$ ($-\infty$, ∞). We now show $F = F^*$ a.e.

Let

$$
g_{+}^{*}(s) = -(2\pi)^{-\frac{1}{2}} i \int_{-\infty}^{\infty} \frac{1}{x - is} F^{*}(x) dx, \qquad s > 0,
$$

and

$$
g_{-}^{*}(s) = (2\pi)^{-\frac{1}{2}} i \int_{-\infty}^{\infty} \frac{1}{x + is} F^{*}(x) dx, \qquad s > 0.
$$

Then since for each $s > 0$, $(x - is)^{-1} \in L_p(-\infty, \infty)$, and

$$
(2\pi)^{-\frac{1}{2}}(P)\int_{-\infty}^{\infty}\frac{1}{(x-is)}e^{-ixy}dx = \begin{cases} (2\pi)^{\frac{1}{2}}i\ e^{iy}, & y < 0, s > 0, \\ -(2\pi)^{\frac{1}{2}}i\ e^{iy}, & y > 0, s < 0, \\ 0, & sy > 0, \end{cases}
$$

we have from $(2,$ Theorem 75) for $s > 0$,

$$
g_{+}^{*}(s) = -(2\pi)^{-\frac{1}{2}} i \int_{-\infty}^{\infty} \frac{1}{x - is} F^{*}(x) dx
$$

=
$$
\int_{0}^{\infty} e^{-sy} f(y) dy = \int_{0}^{\infty} e^{-sy} f_{+}(y) = g_{+}(s),
$$

and

$$
g_{-}^{*}(s) = (2\pi)^{-\frac{1}{2}} i \int_{-\infty}^{\infty} \frac{1}{x + is} F^{*}(x) dx
$$

=
$$
\int_{-\infty}^{0} e^{sy} f(y) dy = \int_{0}^{\infty} e^{-sy} f_{-}(y) dy = g_{-}(s).
$$

Consequently, for $s > 0$

$$
\int_{-\infty}^{\infty} \frac{1}{x - is} \left(F(x) - F^*(x) \right) dx = 0
$$

and

$$
\int_{-\infty}^{\infty} \frac{1}{x + is} \left(F(x) - F^*(x) \right) dx = 0.
$$

Letting $\phi(x) = F(x) - F^*(x)$, the last two equations yield

$$
\int_{-\infty}^{\infty} \frac{1}{x + is} \phi(x) dx = 0, \qquad s \neq 0.
$$

Then denoting the even and odd parts of ϕ by ϕ_e and ϕ_0 respectively, we have for $s \neq 0$

$$
\int_{-\infty}^{\infty} \frac{1}{x + is} \phi_e(x) dx = - \int_{-\infty}^{\infty} \frac{1}{x + is} \phi_0(x) dx.
$$

But the function on the left of this equation is an odd function of *s* while the function on the right is even. Hence each is zero, so that for $s^{\dagger} \neq 0$

$$
\int_0^\infty \frac{1}{x^2+s^2} \phi_\epsilon(x) dx = -\frac{1}{2is} \int_{-\infty}^\infty \frac{1}{x+is} \phi_\epsilon(x) dx = 0,
$$

and

$$
\int_0^{\infty} \frac{x}{x^2 + s^2} \phi_0(x) dx = -\frac{1}{2} \int_{-\infty}^{\infty} \frac{1}{x + is} \phi_0(x) dx = 0.
$$

Thus for each $3 \gt 0$,

$$
\int_0^\infty \frac{1}{x+s} x^{-\frac{1}{2}} \phi_\epsilon(x^{\frac{1}{2}}) dx = 2 \int_0^\infty \frac{1}{x^2+s} \phi_\epsilon(x) dx = 0,
$$

and

$$
\int_0^{\infty} \frac{1}{x+s} \phi_0(x^{\frac{1}{2}}) dx = 2 \int_0^{\infty} \frac{x}{x^2+s} \phi_0(x) dx = 0,
$$

10 $\frac{1}{2}$ **1** $\frac{1}{2$ chanter 8. Theorem 5a $\overline{}$ and $\overline{}$ are zero almost everywhere. Thus $\overline{}$ is zero chapter 8, Theorem 5a) *<t>e* and <£0 are zero almost everywhere. Thus <£ is zero almost everywhere so that $F = F^*$ almost everywhere, and *F* has the prescribed representation.

For Theorems 2 and 3 let us denote by \mathscr{L}_r ($-\infty$, ∞) the collection of functions f such that $|x|^{1-2/r} f(x) \in L_r$ ($-\infty$, ∞).

THEOREM 2. A necessary and sufficient condition that a function $F \in$ L_q ($-\infty$, ∞), $q \geqslant 2$, be the Fourier transform of a function in \mathscr{L}_q ($-\infty$, ∞), $q \geqslant 2$, is that there exist a constant M such that

$$
\int_{-\infty}^{\infty} |t|^{q-2} |\mathfrak{F}_{k,\,t}[F]|^q dt \leqslant M, \qquad k > q-2.
$$

Proof of necessity. Suppose F is the Fourier transform of $f \in \mathscr{L}_{q}$ ($-\infty$, ∞). Then as in the proof of Theorem 1, for $t > 0$ and $k > q - 2$

$$
|\mathfrak{F}_{k, i}[F]| \leq \left\{ (k/t)^{k+1}(k!)^{-1} \int_0^\infty e^{-ky/t} y^k |f(y)|^q dy \right\}^{1/q}
$$

and consequently if $k > q - 2$

$$
\int_0^{\infty} t^{\sigma-2} |\mathfrak{F}_{k, t}[F]|^{\sigma} dt \leq \frac{k^{k+1}}{k!} \int_0^{\infty} t^{\sigma-k-3} dt \int_0^{\infty} e^{-ky/t} y^k |f(y)|^{\sigma} dy
$$

= $\frac{k^{k+1}}{k!} \int_0^{\infty} y^k |f(y)|^{\sigma} dy \int_0^{\infty} e^{-ky/t} t^{\sigma-k-3} dt$
= $K(k) \int_0^{\infty} y^{\sigma-2} |f(y)|^{\sigma} dy$,

where $K(k) = k^{q-1} \Gamma(k - q + 2)/k!$ Similarly

$$
\int_{-\infty}^0 |t|^{q-2} |\mathfrak{F}_{k, t}[F]|^q dt \leqslant K(k) \int_{-\infty}^0 |y|^{q-2} |f(y)|^q dy,
$$

so that

$$
\int_{-\infty}^{\infty} |t|^{q-2} |\mathfrak{F}_{k,\,l}[F]|^q dt \leqslant K(k) \int_{-\infty}^{\infty} |y|^{q-2} |f(y)|^q dy.
$$

But from Stirling's formula,

$$
\lim_{k\to\infty}K(k)=1,
$$

so that $K(k)$ is bounded for $k > q - 2$. Hence there is an *M* such that

$$
\int_{-\infty}^{\infty} |t|^{q-2} |\mathfrak{F}_{k,\,t}[F]|^q dt \leqslant M, \qquad k > q-2.
$$

Proof of sufficiency. Let g_+ and g_- be defined as in the proof of Theorem 1. Then as in that proof, for $t > 0$

$$
L_{k, t}[g_{+}] = \mathfrak{F}_{k, t}[F],
$$

and

$$
L_{k,i}[g_-] = \mathfrak{F}_{k,-i}[F],
$$

and hence

$$
\int_0^{\infty} t^{q-2} |L_{k, t}[g_+]|^q dt = \int_0^{\infty} t^{q-2} |\mathfrak{F}_{k, t}[F]|^q dt \leq M, \qquad k > q-2
$$

and

$$
\int_0^{\infty} t^{q-2} |L_{k, t}[g_-]|^q dt = \int_0^{\infty} t^{q-2} |\mathfrak{F}_{k, t}[F]|^q dt \leq M, \qquad k > q-2.
$$

Consider first g_+ . By (3, chapter 1, Theorem 17a), with $\alpha_k(t) = t^{1-2/q} L_{k,t}[g_+]$, there is a function f_+ with $t^{1-2/q} f_+(t) \in L_q(0, \infty)$, and an increasing unbounded sequence of integers $\{k_i\}$ such that for any function $\beta(t) \in L_p(0, \infty)$,

$$
\lim_{t\to\infty}\int_0^\infty\beta(t)\ t^{1-2/q}\ L_{k\epsilon,\ t}[g_+]dt=\int_0^\infty\beta(t)\ t^{1-2/q}f_+(t)dt.
$$

But for each $s > 0$, $t^{-(1-2/q)} e^{-st} \in L_p(0, \infty)$, and hence choosing this as our $\beta(t)$ we have for $s > 0$

$$
\lim_{t\to\infty}\int_0^\infty e^{-st}L_{k,i},[g_+]dt=\int_0^\infty e^{-st}f_+(t)dt.
$$

However, for $x > 0$,

$$
\int_0^x |L_{k, t}[g_+]| dt \leq \left\{ \int_0^x t^{p-2} dt \right\}^{1/p} \left\{ \int_0^x t^{q-2} |L_{k, t}[g_+]|^q dt \right\}^{1/q}
$$

 $\leq (p-1)^{-1/p} M x^{1/q} = O(x)$ as $x \to \infty$,

and as in the proof of Theorem 1, $g_{+}(s) \rightarrow 0$ as $s \rightarrow \infty$. Hence by (3, chapter 7, Theorem lib),

$$
\lim_{t\to\infty}\int_0^\infty e^{-st} L_{kt, t}[g_+]dt = g_+(s), \qquad s > 0,
$$

and thus

$$
g_{+}(s) = \int_{0}^{\infty} e^{-st} f_{+}(t) dt, \qquad s > 0.
$$

Similarly f_{-} exists with $t^{1-2/q} f_{-}(t) \in L_q(0, \infty)$ such that

$$
g_{-}(s) = \int_{0}^{\infty} e^{-st} f_{-}(t) dt, \qquad s > 0.
$$

Let

$$
f(t) = \begin{cases} f_{+}(t), & t > 0, \\ f_{-}(-t), & t < 0. \end{cases}
$$

Then clearly $f \in \mathscr{L}_q$ ($-\infty$, ∞), and hence by (2, Theorem 79) f has a Fourier transform $F^* \in L_q$ ($-\infty$, ∞). It remains to show $F = F^*$ a.e., which now follows as in Theorem 1.

THEOREM 3. A necessary and sufficient condition that a function $F \in$ \mathscr{L}_p ($-\infty$, ∞), $1 < p \leqslant 2$, be the Fourier transform of a function in L_p ($-\infty$, ∞) *is that there exist a constant M such that*

$$
\int_{-\infty}^{\infty} |\mathfrak{F}_{k,i}[F]|^p dt \leqslant M, \qquad k = 1, 2, \ldots.
$$

Proof of necessity. If $F \in \mathscr{L}_p$ ($-\infty$, ∞) is the Fourier transform of $f \in L_p$ $(-\infty, \infty)$ then by (2, Theorem 74), $F \in L_q$ $(-\infty, \infty)$, and hence by Theorem 1, there is a constant *M* so that

$$
\int_{-\infty}^{\infty} |\mathfrak{F}_{k,\,l}[F]|^p dt \leqslant M, \qquad k = 1, 2, \ldots.
$$

Proof of sufficiency. Let $g_{+}(s)$ and $g_{-}(s)$ be defined as in Theorem 1. Then as in that theorem,

$$
\int_0^\infty |L_{k,\,l}[g_+(s)]|^p dt \leqslant M, \qquad k=1,2,\ldots,
$$

and

$$
\int_0^\infty |L_{k,\,t}[g_-(s)]|^p dt \leqslant M, \qquad k=1,2,\ldots.
$$

Further $g_{\pm}(s) \rightarrow 0$ as $s \rightarrow \infty$. For from Hölder's inequality we have for $s > 0$

$$
|g_{\pm}(s)| \leqslant \left\{ \int_{-\infty}^{\infty} \frac{|x|^{q-2}}{(x^2+s^2)^{q/2}} \, dx \right\}^{1/q} \left\{ \int_{-\infty}^{\infty} |x|^{p-2} |F(x)|^p dx \right\}^{1/p} = 0(s^{-1/q}).
$$

Hence by (3, chapter 7, Theorem 15a), there are functions f_+ and f_- in L_p (0, ∞) such that

$$
g_{+}(s) = \int_{0}^{\infty} e^{-st} f_{+}(t) dt, \qquad s > 0
$$

and

$$
g_{-}(s) = \int_{0}^{\infty} e^{-st} f_{-}(t) dt, \qquad s > 0.
$$

Let

$$
f(t) = \begin{cases} f_{+}(t), & t > 0, \\ f_{-}(-t), & t < 0. \end{cases}
$$

Then clearly $f \in L_p$ ($-\infty$, ∞) and hence by (2, Theorems 75 and 80) f has a Fourier transform $F^* \in \mathscr{L}_{\rho}(-\infty, \infty)$. It remains to show that $F = F^*$ a.e., and this follows as in Theorem 1.

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