

LETTER

35 W continuous-wave Ho:YAG single-crystal fiber laser

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Abstract

We report on a high-power Ho:YAG single-crystal fiber (SCF) laser inband pumped by a high-brightness Tm-fiber laser at 1908 nm. The Ho:YAG SCF grown by the micro-pulling-down technique exhibits a propagation loss of $0.05 \pm 0.005 \text{ cm}^{-1}$ at $2.09 \mu\text{m}$. A continuous-wave output power of 35.2 W is achieved with a slope efficiency of 42.7%, which is to the best of our knowledge the highest power ever reported from an SCF-based laser in the $2 \mu\text{m}$ spectral range.

Keywords: high-power laser; Ho:YAG; mid-IR laser; single-crystal fiber

1. Introduction

High-power solid-state lasers in the $2 \mu\text{m}$ spectral region are actively investigated due to their widespread applications in various fields such as eye-safe coherent laser radar for long-range hard-target measurement^[1], hard or soft tissue ablation owing to strong laser absorption by water molecules^[2, 3], high-sensitivity global wind and active atmospheric remote sensing^[4], and pumping optical parametric oscillators for the development of mid-IR sources^[5]. Current concepts for the power scaling of such $2 \mu\text{m}$ laser sources tend to use Tm- or Ho-doped glass fibers^[6, 7]. However, stimulated Brillouin scattering, stimulated Raman scattering, optical damage and photo-darkening present serious limiting factors^[8]. Single-crystal fibers (SCFs), typically with a diameter below 1 mm and a length of a few centimeters, are an alternative for high-power laser systems owing to the fiber-like geometry, higher

thermal conductivity, higher optical damage threshold and much lower Brillouin gain coefficient^[9]. The amplification of 800-fs pulses to an average power of 160 W^[10] and 380-fs pulses to a peak power of 2.2 GW^[11] using Yb:YAG SCFs indicates the inherent capability for simultaneous handling of high average and peak powers. With good thermal management provided by the large surface-to-volume ratio (i.e., small temperature difference between the center and the edge), output powers in the continuous-wave (CW) regime of 34 W at $0.95 \mu\text{m}$ and $>250 \text{ W}$ at $1 \mu\text{m}$ have been, respectively, demonstrated with Nd-doped^[12] and Yb-doped YAG SCFs^[13], in which the high overlap between the pump and the laser beam was brought about by pump guiding. Alternatively, high beam quality pump sources (e.g., high-brightness laser diodes or fiber lasers offering a sufficiently high depth of focus to cover the entire SCF) make this possible for freely propagating pumps. In comparison to a waveguide pumped Yb:YAG SCF, a conventional Yb:YAG rod with an unguided pump using a high-brightness laser diode has been demonstrated to exhibit even a higher gain in

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a two-pass chirped pulse amplifier^[14]. On the other side, Tm- or Ho-SCF based laser systems in the 2 μm spectral range have been seldom reported. In 2016, such a laser at 2.09 μm was demonstrated for the first time employing a Ho:YAG SCF fabricated by laser heated pedestal growth (LHPG)^[15]. A quasi-CW output power of 23.5 W was obtained under pulsed pumping. However, output power roll-off was observed at ~ 10 W in the pure CW regime. Most recently, we have reported on the first Q -switched Ho:YAG SCF laser delivering 7.5 ns pulses with 0.19 MW peak power^[16]. No optical damage or spectral broadening was observed in this experiment, indicating the high damage resistivity and suppressed nonlinear effects in the SCF. Nevertheless, the CW laser power of 2 μm SCF oscillators/amplifiers was clamped to ~ 10 W due to either available pump power or thermal issues under rigid focus conditions for realizing waveguide pumping.

The laser performance in the CW regime is investigated in the present work by utilizing a high-brightness Tm-fiber laser at 1908 nm for the resonant pumping of a Ho:YAG SCF grown by the micro-pulling-down (μ -PD) method. The laser power of 35 W produced at 2.09 μm is to the best of our knowledge the highest power ever reported from an SCF-based laser in the 2 μm spectral range.

2. Experimental details

YAG ($\text{Y}_3\text{Al}_5\text{O}_{12}$) is the most stable garnet phase in the Y_2O_3 – Al_2O_3 binary equilibrium diagram and melts congruently around $\sim 1980^\circ\text{C}$. Because of the small difference in the ionic radius between Ho^{3+} (104.1 pm) and Y^{3+} (104 pm), it is possible to substitute a high level of Y^{3+} by Ho^{3+} and the segregation issue is a minor problem in comparison to other dopants such as Nd^{3+} (112.3 pm) or Ce^{3+} (115 pm), thus maintaining the advantageous properties of the host crystal. The YAG SCF doped with 0.6% (atomic fraction) Ho^{3+} was grown by the μ -PD technique at a temperature slightly higher than the melting point. The fabricated SCF with a total length of 160 mm was cut into several segments and a 40-mm-long segment was used as a laser element in the present work. Its end faces were polished but left uncoated, whereas the lateral surface was not polished since it is inherently smooth. Figure 1(a) displays a photograph of the end facet of the SCF. The transverse cross-section presents a noncircular shape with a diameter of 960 and 920 μm in the two orthotropic directions, which was formed from the natural facets along the $\langle 111 \rangle$ direction during the growth process. This growth direction corresponds to the lowest surface energy. The growth of YAG fibers along it results in six facets with hexagonal shape, and this direction is considered to be the easiest growth orientation. To mitigate the thermal load, the SCF, tightly mounted on a specially designed copper holder, was water-cooled to 15.0°C .

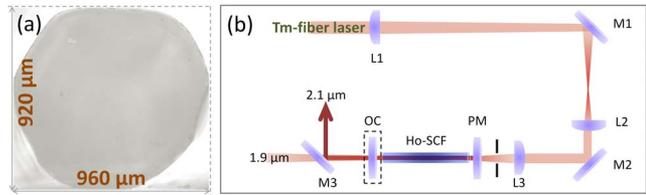


Figure 1. (a) Photograph of the end facet of the Ho:YAG SCF. (b) Schematic of the Ho:YAG SCF laser: L1–L3, lenses with $f = 750$, 100 and 150 mm, respectively; M1, M2, bending mirrors; M3, dichroic mirror; PM, pump mirror; OC, output coupler.

Figure 1(b) shows the schematic of the Ho:YAG SCF laser. The pump beam emitted from a Tm-fiber laser (IPG-TLM-100 W-1908 nm, $M^2 = 1.1$) passed through a 7.5 : 1 telescope system and thereafter was focused by a lens ($f = 150$ mm) into the SCF with a spot diameter of 450 μm . The beam diameter along the entire SCF was measured to be less than 600 μm , i.e., smaller than the aperture of the SCF. A plane mirror coated for high reflectivity at the laser wavelength was used as a pump mirror (PM). Plane mirrors with transmission of $T_{\text{oc}} = 30\%$, 50%, 60%, 72% and 92% (end-face of the SCF provided 8% Fresnel reflection at 2.1 μm) were used as output couplers (OCs). A dichroic mirror (M3) was used to separate the laser and pump beams.

3. Results and discussion

High-power CW operation of the Ho:YAG SCF laser was studied at a physical cavity length of 45 mm using different OCs. The low signal single-pass absorption efficiency of the SCF under nonlasing conditions amounted to 94.7% at 1.9 W of incident pump power. Further increasing the pump power to 71.3 W, the absorption efficiency dropped to 82.5% due to the bleaching effect. Under lasing conditions with $T_{\text{oc}} = 92\%$, the corresponding absorption efficiencies for the same pump levels were 95.4% and 87%. Figure 2(a) shows the output laser performance with respect to the absorbed pump power, and the laser beams in all the cases were measured to be unpolarized. A maximum output power of 35.2 W was achieved with $T_{\text{oc}} = 72\%$. No power roll-over was observed and further power scaling was limited just by the available pump power. The good laser performance is attributed to the high gain and the weak thermal load benefiting from the large surface area-to-volume ratio, i.e., the fiber-like geometry of the SCF. The highest slope efficiency of 53.3% was achieved by using the end facet of the SCF as OC, i.e., $T_{\text{oc}} = 92\%$. Unfortunately, further power scaling in this case was limited by a careless scratch point at one end facet of the SCF during the experiment. Nevertheless, the straightforward power scaling with almost linear pump power dependence indicates the huge potential of the SCF for high-power laser operation. Laser emission was not achieved without the input mirror, i.e., in a cavity formed by both end facets of the SCF in which case a much

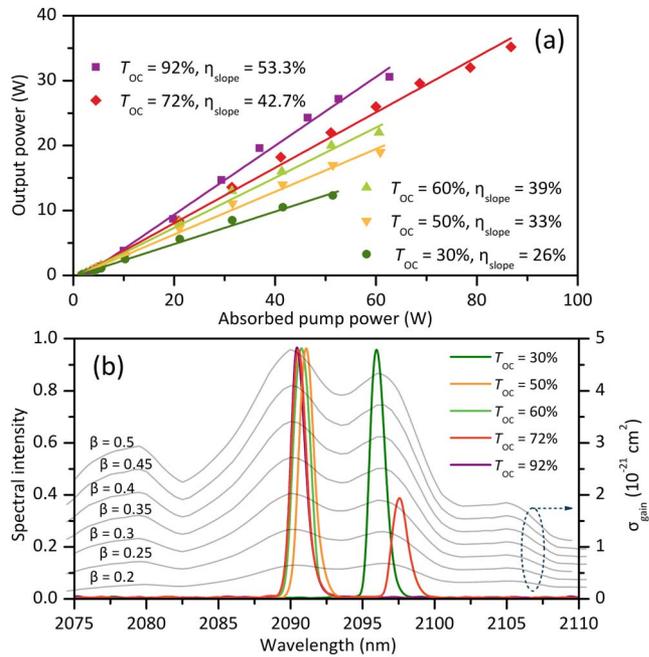


Figure 2. (a) Laser performance of the Ho:YAG SCF laser with different OCs and (b) the corresponding optical spectra. The gray curves are the calculated gain spectra of Ho:YAG^[17] with different population inversion parameters β .

higher population inversion is required. This is an indication of the nonparallel faces of the used SCF. Figure 2(b) shows the optical spectra of the Ho:YAG SCF laser with different OCs (for comparison the gain spectra of Ho:YAG^[17] are also depicted). The peak wavelength of the laser emission changed from 2090.3 nm to 2095.7 nm with increasing reflectivity of the OC. This wavelength red-shift is caused by the enhanced reabsorption effect with low OC transmission, where a lower population inversion is needed, resulting in a drift of gain spectrum. Note that the two emission peaks for $T_{oc} = 72\%$ could be attributed to the same gain cross-sections in this case.

To evaluate the propagation loss of the SCF, Caird analysis that was modified for the case of high T_{oc} was performed^[18]. The dependence of the slope efficiency η on T_{oc} and the internal loss L per pass could be expressed as^[19]

$$\frac{1}{\eta} = \frac{1}{\eta_0 \left(1 + \frac{2\gamma}{\gamma_{oc}}\right)}, \quad (1)$$

where η_0 is an intrinsic slope efficiency, $\gamma = -\ln(1 - L)$ and $\gamma_{oc} = -\ln(1 - T_{oc})$. Note that the emission wavelength and the corresponding stimulated emission cross-section were regarded as constant since the laser wavelength depended only weakly on T_{oc} (only 5 nm drift as shown in Figure 2(b)). Figure 3 shows the inverse slope efficiency versus the inverse output-coupling loss. Best fit of the experimental data yielded $\eta_0 \approx 60\%$ and a propagation loss

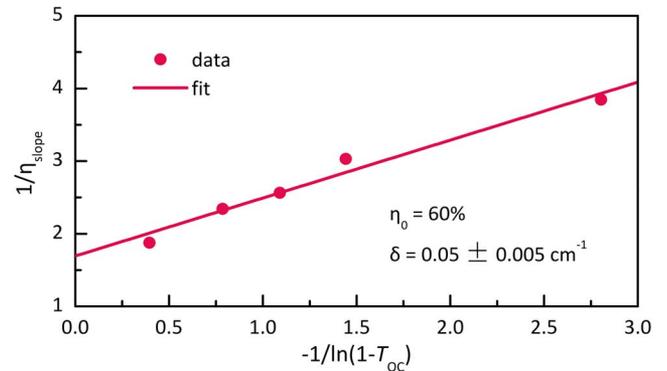


Figure 3. Caird plot for the high-power Ho:YAG SCF laser: inverse slope efficiency with respect to the inverse output-coupling loss, i.e., $-\ln(1 - T_{oc})$.

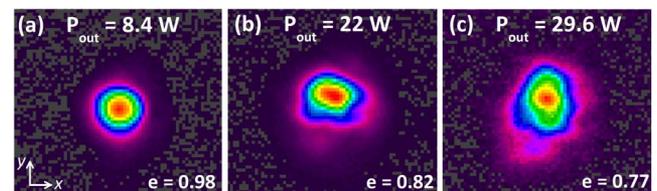


Figure 4. Beam intensity profiles recorded at different output powers with $T_{oc} = 72\%$; e is the calculated beam ellipticity from each image.

at the laser wavelength of $\delta = L/l = 0.05 \pm 0.005 \text{ cm}^{-1}$. The calculated limiting slope efficiency is considerably less than the quantum efficiency of $\sim 91\%$, which is caused by the following two reasons: first, the noncollinear pumping configuration, i.e., about 2° between the pump and laser beams at the front facet of the SCF used to avoid the unwanted feedback to the Tm-fiber pump laser, and second, the energy-transfer upconversion of Ho³⁺ ions in the ⁵I₇ manifold under pump bleaching^[20]. The propagation loss is larger than the previously reported value of 0.02 cm^{-1} , ‘directly measured’ in a 400 μm -diameter Ho:YAG SCF fabricated by LHPG method^[21]. Anti-reflection coating on the SCF and optimizing the crystal quality would decrease the overall losses, thus contributing to further scaling of the laser power.

Far-field beam intensity profiles of the Ho:YAG SCF laser with $T_{oc} = 72\%$ were recorded at different output powers by using a CCD camera (PY-III-HR-C-A, Ophir). As shown in Figure 4(a), the beam profile at 8.4 W output power was a perfect circle with a beam ellipticity (defined as a ratio of the beam diameters in the x - and y -directions) of $e = 0.98$. This implies uniform thermal distribution over the cross-section of the SCF owing to the large surface area-to-volume ratio. The corresponding near-field beam profile presented a similar intensity distribution which was obviously not affected by the noncircular shape of the SCF, indicating a free-space propagation of the laser beam in the SCF. This could be further confirmed by calculating the beam size on

the end facet of the SCF with $T_{oc} = 92\%$ for the case when the end facet itself was employed as the output coupler. Using a $2f$ imaging system and knife-edge technique, the beam diameter was roughly estimated to be $300\ \mu\text{m}$ on the SCF end facet, which equals the simulated value using the ABCD-matrix analysis. When the output power reached 22 W, as shown in Figure 4(b), the beam shape was no more circular because of the generation of high-order modes, giving an e value of 0.82. The excitation of high-order modes is easily understood since the noncollinear pumping configuration and no additional transverse-mode-limiting optical elements were applied. Further power scaling to nearly 30 W produced a deteriorated beam profile (see Figure 4(c)) with an elliptic shape ($e = 0.77$). The smaller central intensity distribution (the area of the red part) is an indication of thermal lens effect and the asymmetric thermal distribution over the transverse cross-section of the SCF. To mitigate the thermal issues, longer SCF with lower active ion doping level will be preferable for further improving the laser performance. Moreover, direct water cooling around the SCF will mitigate thermal effect and favor uniform thermal diffusion, since the nonoptimized thermal contact is responsible for the main thermal difference between the mount and SCF center^[22]. Nevertheless, the thermal issue in present work is more mitigated than in the previous report in which the CW laser power began to roll off at 10 W output^[15].

The beam propagation factors (M^2) at different absorbed pump powers were measured using the same CCD camera. As depicted in Figure 5(a), the measured M^2 -factors in both x - and y -directions were around 1.1 at low pump levels, i.e., <20 W absorbed pump power. However, excitation of high-order transverse modes when further increasing the pump power would gradually degrade the beam quality. On the other side, beam quality along the x -axis exhibits a stronger deterioration than that along the y -axis, which is attributed to the unidirectional heat dissipation in the SCF at high power levels. At 25 W of absorbed pump power, the M^2 -factors along the x - and y -directions were measured to be 1.32 and 1.24, respectively. The M^2 -factors were not measured at higher output powers because of the obvious multi-mode emission. Figure 5(b) shows a typical M^2 -factor measurement at an output power of 7 W corresponding to an absorbed power of 15 W. In this case, the measured M^2 -factor is comparable with the value of 1.1 that was obtained by employing a Ho:YAG thick rod^[23].

4. Conclusion

In conclusion, we have experimentally demonstrated a high-power Ho:YAG SCF laser inband pumped by a Tm-fiber laser at 1908 nm. With a modified Caird analysis, the propagation loss coefficient of this Ho:YAG SCF was estimated to

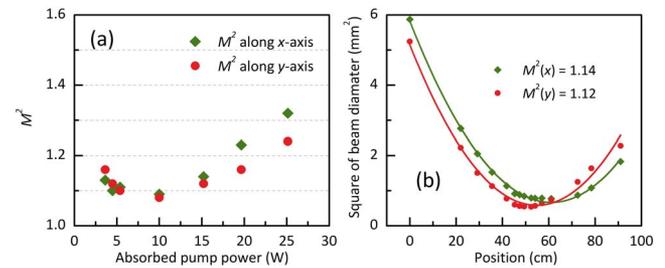


Figure 5. (a) The measured M^2 -factor of the Ho:YAG SCF laser ($T_{oc} = 92\%$) at different absorbed pump powers and (b) a typical M^2 measurement at 7 W output laser power.

be $0.05 \pm 0.005\ \text{cm}^{-1}$. A maximum output power of 35.2 W was achieved with $T_{oc} = 72\%$, corresponding to a slope efficiency of 42.7%. Without the traditional spatial mode matching, excitation of high-order transverse modes seems unavoidable under high powers, which will result in a serious deterioration of the beam quality. Nevertheless, the measured M^2 -factors are below 1.2 for output powers up to 10 W. Thermal lens effects were obviously enhanced at a power level approaching ~ 30 W. Therefore, a longer SCF with a lower doping level and direct water cooling will be essential for further power scaling and improving the beam quality, whereas anti-reflection coating and limitation of the high-order transverse modes are expected to improve the overall laser performance.

The present work is a demonstration of the high quality of the fabricated Ho:YAG SCF. Amplification of ultrashort pulses (including picosecond and femtosecond) near $2\ \mu\text{m}$ to high average/peak power will be the next step, relying on advantageous features such as the high laser damage threshold, straightforward thermal management (large surface-to-volume ratio), excellent thermo-mechanical properties and weak nonlinear effects. Moreover, the encountered problems such as the excitation of high-order transverse modes and thermal lens effect at high power levels would be absent or effectively compensated by external lens when the SCF is employed as an amplifier.

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References

1. S. W. Henderson, C. P. Hale, J. R. Magee, M. J. Kavaya, and A. V. Huffaker, *Opt. Lett.* **16**, 773 (1991).
2. C. Gottlob, G. E. Kopchok, S.-K. Peng, M. Tabbara, D. Cavaye, and R. A. White, *Lasers Surg. Med.* **12**, 86 (1992).
3. T. Sumiyoshi, H. Sekita, T. Arai, S. Sato, M. Ishihara, and M. Kikuchi, *IEEE J. Sel. Top. Quantum Electron.* **5**, 936 (1999).
4. U. N. Singh, B. M. Walsh, J. Yu, M. Petros, M. J. Kavaya, T. F. Refaat, and N. P. Barnes, *Opt. Mater. Express* **5**, 827 (2015).
5. A. Dergachev, D. Armstrong, A. Smith, T. Drake, and M. Dubois, *Opt. Express* **15**, 14404 (2007).
6. T. Walbaum, M. Heinzig, T. Schreiber, R. Eberhardt, and A. Tünnermann, *Opt. Lett.* **41**, 2632 (2016).
7. S. D. Jackson, A. Sabella, A. Hemming, S. Bennetts, and D. Lancaster, *Opt. Lett.* **32**, 241 (2007).
8. M. N. Zervas and C. A. Codemard, *IEEE J. Sel. Top. Quantum Electron.* **20**, 0904123 (1999).
9. X. Délen, A. Aubourg, L. Deyra, F. Lesparre, I. Martial, J. Didierjean, F. Balembois, and P. Georges, *Proc. SPIE* **9342**, 934202 (2015).
10. V. Markovic, A. Rohrbacher, P. Hofmann, W. Pallmann, S. Pierrot, and B. Resan, *Opt. Express* **23**, 25883 (2015).
11. X. Délen, Y. Zaouter, I. Martial, N. Aubry, J. Didierjean, C. Hönninger, E. Mottay, F. Balembois, and P. Georges, *Opt. Lett.* **38**, 109 (2013).
12. X. Délen, I. Martial, J. Didierjean, N. Aubry, D. Sangla, F. Balembois, and P. Georges, *Appl. Phys. B* **104**, 1 (2011).
13. X. Délen, S. Piehler, J. Didierjean, N. Aubry, A. Voss, M. A. Ahmed, T. Graf, F. Balembois, and P. Georges, *Opt. Lett.* **37**, 2898 (2012).
14. A. M. Rodin and M. Zopelis, in *Conference on Lasers and Electro-Optics Pacific Rim (CLEO-PR)* (2017), p. 1.
15. Y. Li, K. Miller, E. G. Johnson, C. D. Nie, S. Bera, J. A. Harrington, and R. Shori, *Opt. Express* **24**, 9751 (2016).
16. J. Wang, Q. Song, Y. Sun, Y. Zhao, W. Zhou, D. Li, X. Xu, C. Shen, W. Yao, L. Wang, J. Xu, and D. Shen, *Opt. Lett.* **44**, 455 (2019).
17. Y. Wang, R. Lan, X. Mateos, J. Li, C. Hu, C. Li, S. Suomalainen, A. Härkönen, M. Guina, V. Petrov, and U. Griebner, *Opt. Express* **24**, 18003 (2016).
18. J. Morris, N. K. Stevenson, H. T. Bookey, A. K. Kar, C. T. A. Brown, J.-M. Hopkins, M. D. Dawson, and A. A. Lagatsky, *Opt. Express* **25**, 14910 (2017).
19. P. Loiko, R. Soulard, G. Brasse, J. Doualan, B. Guichardaz, A. Braud, A. Tyazhev, A. Hideur, and P. Camy, *Opt. Express* **26**, 24653 (2018).
20. N. P. Barnes, B. M. Walsh, and E. D. Filer, *J. Opt. Soc. Am. B* **20**, 1212 (2003).
21. Y. Li, E. G. Johnson, C. D. Nie, J. A. Harrington, and R. Shori, *Opt. Express* **22**, 14896 (2014).
22. D. Sangla, I. Martial, N. Aubry, J. Didierjean, D. Perrodin, F. Balembois, K. Lebbou, A. Brenier, P. Georges, O. Tillement, and J. M. Fourmigue, *Appl. Phys. B* **97**, 263 (2009).
23. D. Y. Shen, A. Abdolvand, L. J. Cooper, and W. A. Clarkson, *Appl. Phys. B* **79**, 559 (2004).