

# An Automorphic Theta Module for Quaternionic Exceptional Groups

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*Abstract.* We construct an automorphic realization of the global minimal representation of quaternionic exceptional groups, using the theory of Eisenstein series, and use this for the study of theta correspondences.

## 1 Introduction

Let  $F = \mathbb{Q}$  be the rational number field, and let  $D$  be either a definite quaternion algebra  $\mathbb{H}$ , or a definite octonion algebra  $\mathbb{O}$  over  $F$ . To such a  $D$ , there corresponds an exceptional group  $H$  of relative rank 4, and type  $E_7$  or  $E_8$  accordingly. In  $H$ , there is a reductive dual pair  $G \times G'$ . Here  $G$  is a split group of type  $G_2$ , and  $G'$  is the automorphism group of the Jordan algebra  $J$  of 3-by-3 hermitian matrices with coefficients in  $D$ , which is anisotropic over  $F$ . In this paper, we study the global theta correspondence which arises from the dual pair  $G \times G'$ . The local analogue of this correspondence has been studied in [MS] and [SG], and the reason for writing  $F$  in place of  $\mathbb{Q}$  is the expectation that the results here hold for any totally real number field.

The first part of the paper is devoted to the construction of an automorphic realization of the global minimal representation  $\Pi$  of  $H$ , and follows the approach of [GRS1] for the split case. Let  $P$  be the Heisenberg parabolic subgroup of  $H$ , with modulus character  $\delta_P$ . For a standard section  $f_s \in \text{Ind}_{P(\Lambda)}^{H(\Lambda)} \delta_P^{\frac{1}{2}+s}$ , let  $E(g, f_s)$  be the usual Eisenstein series. Then we have:

**Theorem 1.1** *For any standard section  $f_s$ ,  $E(g, f_s)$  has at most a simple pole at  $s = s_0$  (for a certain specific  $s_0$ ). This pole is actually attained by some standard section. Moreover, the space  $\Theta$  of automorphic forms spanned by the residues of  $E(g, f_s)$  at  $s = s_0$  is an irreducible square-integrable automorphic representation isomorphic to  $\Pi$ .*

The proof of this theorem follows the method of [GRS1]; the main difference being that, unlike the split case, the local representation  $\Pi_v$  may be non-spherical here for some place  $v$ . Though the above result is to be expected, we find it useful and necessary to have the details properly worked out, because of the importance of the automorphic theta module  $\Theta$  in the theory of automorphic forms. For example, combined with the results of [Li1] and [Li2], the existence of  $\Theta$  implies the non-vanishing of  $L^2$ -cohomology groups in low degrees of certain locally symmetric spaces. We remark also that an automorphic theta module was constructed in [R] for the rank 2 form of  $E_6$  associated to a 9-dimensional division algebra.

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Received by the editors April 30, 1999; revised July 2, 1999.  
AMS subject classification: 11F27, 11F70.  
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In the second part of the paper, we study the theta lifts of automorphic forms from  $G'$  to  $G$  using  $\Theta$ . If  $\pi$  is an automorphic representation of  $G'$ , let  $\Theta(\pi)$  be its theta lift, which is a subspace of the space of automorphic forms on  $G$ . We characterize those  $\pi$  for which  $\Theta(\pi)$  is non-zero and cuspidal in terms of period integrals over certain (anisotropic) subgroups of  $G'$ , in the spirit of [GRS2] and [GrS]. More precisely, each totally real étale cubic algebra  $E$  determines a subgroup  $C_E$  of  $G'$ ; also let  $C$  be the stabilizer in  $G'$  of a primitive idempotent in  $J$ . Then we have:

**Theorem 1.2** *Let  $\pi$  be a non-trivial automorphic representation of  $G'$ . Then  $\Theta(\pi)$  is non-zero and cuspidal if and only if the following two conditions hold:*

- (i)  $\pi$  is not  $C$ -distinguished;
- (ii)  $\pi$  is  $C_E$ -distinguished for some totally real étale cubic algebra  $E$ .

Essentially, the groups  $C_E$  are related to certain generic Fourier coefficients of the theta lift, whereas the group  $C$  is related to degenerate Fourier coefficients.

When  $D = \mathbb{H}$ ,  $G'$  is an inner form of  $\mathrm{PGSp}_6$ , and the functoriality principle predicts that an automorphic representation  $\pi$  of  $G'$  is a functorial lift from  $G$  if and only if its Spin  $L$ -function  $L(s, \pi, \mathrm{Spin})$  has a pole at  $s = 1$ . On the other hand, in [MS] and [SG], the local theta correspondence is found to be a functorial lifting. In view of Theorem 1.2, one expects that  $L(s, \pi, \mathrm{Spin})$  has a pole at  $s = 1$  if and only if  $\pi$  is  $C_E$ -distinguished for some  $E$ . Can this be established independently?

## General Notations

In this paper,  $F$  will denote the rational number field  $\mathbb{Q}$ . A place of  $F$  will be denoted by  $v$ , with  $F_v$  the corresponding completion of  $F$ . Let  $\zeta_v$  be the local zeta factor of  $F$  at  $v$ . Hence,

$$\zeta_v(s) = \begin{cases} (1 - p^{-s})^{-1}, & \text{if } v = p \text{ is finite;} \\ \pi^{-\frac{s}{2}} \cdot \Gamma(\frac{s}{2}), & \text{if } v \text{ is real.} \end{cases}$$

The ring of adèles of  $F$  will be denoted by  $\mathbb{A}$ . As in the introduction,  $D$  will denote either a definite quaternion algebra  $\mathbb{H}$ , or a definite octonion algebra  $\mathbb{O}$  over  $F$ . Let  $\mathrm{SL}(D)$  be the group of norm 1 elements of  $D$ . Set  $D_v := D \otimes F_v$ , and let  $S$  be the finite set of places  $v$  where  $D_v$  is ramified. Hence, if  $D = \mathbb{H}$ ,  $S$  contains the real place and has even cardinality, whereas for  $D = \mathbb{O}$ ,  $S$  contains just the real place. For  $v \notin S$ , set:

$$\zeta_{D_v}(s) = \begin{cases} \zeta_v(s) \cdot \zeta_v(s - 1), & \text{if } D = \mathbb{H}; \\ \zeta_v(s) \cdot \zeta_v(s - 3), & \text{if } D = \mathbb{O}. \end{cases}$$

For any algebraic group  $H$  over  $F$ , we shall write  $H$  for  $H(F)$  and  $H_v$  for  $H(F_v)$ . If  $v$  is finite, an element of  $X^\bullet(H_v) \otimes_{\mathbb{Z}} \mathbb{C}$ , where  $X^\bullet(H_v)$  is the group of rational characters of  $H_v$ , gives rise to an unramified character of  $H_v$ , taking values in  $\mathbb{C}^\times$ . Hence we shall often identify an unramified character of  $H_v$  with an element of  $X^\bullet(H_v) \otimes_{\mathbb{Z}} \mathbb{C}$ . Similarly, an element of  $X^\bullet(H) \otimes_{\mathbb{Z}} \mathbb{C}$  gives rise to an unramified character of  $H(\mathbb{A})$ , which is trivial on  $H$ .

Assume now that  $H$  is reductive. Let  $P_0 = M_0 \cdot N_0$  be a fixed minimal parabolic subgroup, with modulus character  $\delta_0: P_0 \rightarrow \mathbb{R}_+^\times$ . Let  $A \subset M_0$  be a maximal split torus of  $H$ ,  $\Phi$  the set of roots of  $H$  relative to  $A$ , and  $\Phi^+$  the set of positive roots determined by  $P_0$ .

Let  $\Delta \subset \Phi^+$  be the set of simple roots, and let  $W := N_H(A)/M_0$  be the (relative) Weyl group. For any  $\alpha \in \Phi$ ,  $\alpha^\vee$  will denote the corresponding coroot, and  $U_\alpha$  the corresponding root subgroup. Moreover, let  $\langle \cdot, \cdot \rangle$  be the canonical pairing between the roots and the coroots. Sometimes  $\langle \cdot, \cdot \rangle$  will also denote the Killing form of various Lie algebras. We hope that this will not cause any confusion.

For a standard parabolic subgroup  $P = M \cdot N$  of  $H$ , let  $\Delta_M \subset \Delta$  be the set of simple roots of its Levi factor  $M$ ,  $\Phi_M$  the corresponding root system and  $W_M$  the Weyl group of  $M$ . The opposite parabolic is denoted by  $\bar{P} = M \cdot \bar{N}$ , and the modulus character of  $P$  by  $\delta_P$ . For each  $v$ , let  $K_v$  be a maximal compact subgroup of  $H_v$ , which is special if  $v$  is finite, so that the Iwasawa decomposition holds:  $H_v = P_{0,v} \cdot K_v$ . Then for almost all  $v$ ,  $K_v$  is hyperspecial, and  $K = \prod_v K_v$  is a maximal compact subgroup of  $H(\mathbb{A})$ .

Suppose that  $H$  is not split over  $F$ , but  $H_v$  is split. Then, in such cases,  $P_{0,v}$  is no longer the minimal parabolic subgroup. Let  $B_v$  be a fixed Borel subgroup of  $H_v$  contained in  $P_{0,v}$ , with modulus character  $\delta_{B_v}$ , and let  $B_v \supset T_v \supset A_v$  be a maximal torus. If  $\Phi^0$  is the (absolute) root system of  $H_v$  relative to  $T_v$ , then we have a canonical map,  $\Phi^0 \rightarrow \Phi$ , given by restriction of characters. For  $\beta \in \Phi^0$  and  $\alpha \in \Phi$ , we write  $\beta \mapsto \alpha$  if  $\beta$  restricts to  $\alpha$  under the above map.

## 2 Quaternionic Exceptional Groups

In this section, we describe the groups which we will study in this paper. For more details, see [SG, Section 3].

Let  $J$  be the Jordan algebra of 3-by-3 hermitian matrices with coefficients in  $D$ . Then the dimension of  $J$  is:

$$(2.1) \quad d = \begin{cases} 15, & \text{if } D = \mathbb{H}; \\ 27, & \text{if } D = \mathbb{O}. \end{cases}$$

There is a cubic form  $\det$  on  $J$ , giving rise to a symmetric trilinear form  $(\cdot, \cdot, \cdot)$  normalized by  $(X, X, X) = \det(X)$ , and a symmetric bilinear trace form  $(\cdot, \cdot)$ , given by:  $(X, Y) = \text{Tr}(XY)$ . Then, for  $X, Y \in J$ , we have the element  $X \times Y \in J$ , uniquely determined by:

$$(2.2) \quad (X \times Y, Z) = 3(X, Y, Z)$$

for all  $Z \in J$ . Recall that  $X$  has rank one if  $X \neq 0$  but  $X \times X = 0$ . Equivalently,  $X^2 = \text{Tr}(X)X$ . Note that if  $F$  is a totally real number field, as is our case, there are no rank one elements with trace zero.

Let  $L$  be the algebraic group of linear transformations on  $J$  which preserve the determinant form  $\det$ . Then

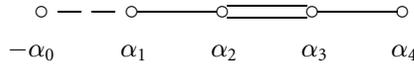
$$(2.3) \quad L \cong \begin{cases} \text{SL}_3(D)/\mu_2, & \text{if } D = \mathbb{H}; \\ E_{6,2}^{\text{sc}}, & \text{if } D = \mathbb{O}. \end{cases}$$

Here,  $E_{6,2}^{sc}$  is a simply-connected group of type  $E_6$  and relative rank 2.

Now associated to  $D$  is a simple adjoint algebraic group  $H$  of relative rank 4, and type  $E_7$  (respectively  $E_8$ ) if  $D = \mathbb{H}$  (respectively  $\mathbb{O}$ ). The Satake diagram of  $H$  is:



Moreover, the relative Dynkin diagram is of type  $F_4$ :



Here  $\alpha_0$  denotes the highest root. In particular,

$$(2.4) \quad \alpha_0 = 2\alpha_1 + 3\alpha_2 + 4\alpha_3 + 2\alpha_4.$$

For simplicity, we shall represent a root  $\sum_{i=1}^4 a_i \alpha_i$  as a 4-tuple  $(a_1, a_2, a_3, a_4)$ . Let  $w_i$  be the simple reflection in  $W$  corresponding to the simple root  $\alpha_i$ . If  $w \in W$  has a minimal length expression  $w = w_{i_1} \cdot w_{i_2} \cdot \dots \cdot w_{i_k}$ , then we shall write  $w = (i_1, i_2, \dots, i_k)$ .

Let  $P = M \cdot N$  be the Heisenberg maximal parabolic subgroup of  $H$ , which corresponds to the vertex  $\alpha_1$  in the relative Dynkin diagram. In particular, its unipotent radical  $N$  is a Heisenberg group with center  $Z$ , and the abelian group  $V = N/Z$  has a natural structure of a symplectic vector space. Let  $\bar{V} = \bar{N}/\bar{Z}$ , where  $\bar{Z}$  is the center of  $\bar{N}$ . Then there is a natural identification [MS, Section 6]

$$(2.5) \quad \bar{V} \cong F \oplus J \oplus J \oplus F.$$

The Levi factor  $M$  has derived group  $M^1$  of type  $D_6$  (respectively  $E_7$ ) if  $D = \mathbb{H}$  (respectively  $\mathbb{O}$ ). The action of  $M^1$  on  $\bar{V}$  is the half-spin representation of dimension 32 if  $D = \mathbb{H}$ , and the 56-dimensional miniscule representation if  $D = \mathbb{O}$ . Moreover, the minimal non-trivial  $M$ -orbit  $\Omega$  is the orbit of a highest weight vector, which can be chosen to be:

$$(2.6) \quad v_0 = (0, 0, 0, 1) \in F \oplus J \oplus J \oplus F.$$

Let  $Q$  be the stabilizer in  $M^1$  of the line spanned by  $v_0$ . Then  $Q$  is a maximal parabolic subgroup of  $M^1$ , and is the intersection of  $M^1$  with the maximal parabolic subgroup of  $H$  corresponding to the vertex  $\alpha_2$ . It has an abelian unipotent radical  $U \cong J$ , and the derived group of its Levi factor is isomorphic to the group  $L$  introduced earlier in (2.3).

Now, we have [MS, Lemma 7.5]:

**Lemma 2.7** *Q has 4 orbits on  $\Omega$ , which are given by:*

$$\begin{aligned} \mathcal{O}_0 &= \{(0, 0, 0, d) : d \in F^\times\}, \\ \mathcal{O}_1 &= \{(0, 0, Y, d) : \text{rank}(Y) = 1 \text{ and } d \in F\}, \\ \mathcal{O}_2 &= \{(0, Y, 2B \times Y, (B, B, Y)) : \text{rank}(Y) = 1 \text{ and } B \in J\}, \\ \mathcal{O}_3 &= \{a(1, Z, Z \times Z, \det(Z)) : a \in F^\times \text{ and } Z \in J\}. \end{aligned}$$

Now note that the characters of the compact group  $V \backslash V(\mathbb{A})$  can be parametrized by  $\bar{V}$  as follows. Fix a non-trivial character  $\psi = \prod_v \psi_v$  of  $F \backslash \mathbb{A}$ . The Killing form  $\langle \cdot, \cdot \rangle$  induces a non-degenerate pairing of  $V$  with  $\bar{V}$ . Then, for  $x \in \bar{V}$ , the corresponding character  $\psi_x$  is given by:

$$\psi_x(n) = \psi(\langle x, n \rangle).$$

Similarly, the characters of  $N_v$  can be parametrized by  $\bar{V}_v$  using the Killing form and  $\psi_v$ . Henceforth, we shall regard the elements of  $\Omega$  as characters of  $V \backslash V(\mathbb{A})$ .

Finally, the modulus character of  $P$  is unramified, and is given by:

$$(2.8) \quad \delta_P = (2 + d)\alpha_0,$$

when regarded as an element of  $X^\bullet(A) \otimes_{\mathbb{Z}} \mathbb{C}$ . Similarly, the modulus character of  $P_0$ , the minimal parabolic, is given by:

$$(2.9) \quad \delta_0 = (4 + 2d)\alpha_1 + (6 + 4d)\alpha_2 + \left(16 + \frac{16}{3}d\right)\alpha_3 + \left(12 + \frac{8}{3}d\right)\alpha_4.$$

### 3 Local Minimal Representation

We now summarize some important facts about the local minimal representation  $\Pi_v$  of  $H_v$ , and refer the reader to [GrW], [S], [R2] and [T] for more details.

For any  $s \in \mathbb{C}$ , consider the degenerate principal series representation:

$$(3.1) \quad I_v(s) = \text{Ind}_{P_v}^{H_v} \delta_P^{\frac{1}{2}+s}$$

where for  $v$  real,  $I_v(s)$  denotes the Harish-Chandra module of the corresponding smooth representation  $I_v(s)^\infty$ . Also, let

$$s_0 = \begin{cases} \frac{11}{34}, & \text{if } D = \mathbb{H}; \\ \frac{19}{58}, & \text{if } D = \mathbb{O}. \end{cases}$$

Now we have:

**Proposition 3.2**

- (i) If  $v$  is the real place, then  $\Pi_v$  is a quotient of  $I_v(s_0)$ , and occurs with multiplicity one in  $I_v(s_0)$ .
- (ii) If  $v \notin S$ , then  $\Pi_v$  is spherical and is the unique irreducible quotient of  $I_v(s_0)$ . Moreover, if  $\psi$  is a character of  $N_v$ , then ( $v \neq 2$  if  $D = \mathbb{O}$ )

$$\dim(\Pi_v)_{N_v, \psi} = \begin{cases} 1, & \text{if } \psi \in \Omega_v; \\ 0, & \text{otherwise.} \end{cases}$$

- (iii) If  $v \in S$  is finite, then  $I_v(s_0)$  has a unique irreducible quotient. If, further,  $v \neq 2$ , then  $\dim(\Pi_v)_{N_v, \psi}$  is as given in (ii) above.

**Remarks** In (i), we suspect, but do not know, that  $\Pi_\infty$  is the unique irreducible quotient of  $I_\infty(s_0)$ . We shall see later that in (iii), the unique irreducible quotient of  $I_v(s_0)$  is exactly  $\Pi_v$ .

#### 4 Eisenstein Series and Intertwining Operators

Now we can begin the construction of the automorphic theta module, *i.e.*, an embedding of  $\Pi = \hat{\otimes}_v \Pi_v$  into the space of automorphic forms  $\mathcal{A}(H)$  on  $H$ . The above results about the local minimal representations suggest that we should consider the global induced representation:

$$(4.1) \quad I(s) = \text{Ind}_{P(\mathbb{A})}^{H(\mathbb{A})} \delta_P^{\frac{1}{2}+s} = \hat{\otimes}_v I_v(s),$$

where the restricted tensor product is formed using the unique  $K_v$ -spherical vector  $\Gamma_{v,s}$ , normalized by  $\Gamma_{v,s}(1) = 1$ . Indeed, as a corollary of Proposition 3.2, we have:

**Proposition 4.2** *The global induced representation  $I(s_0) = \hat{\otimes}_v I_v(s_0)$  has a unique irreducible quotient  $\Pi'$ , with  $\Pi'_\infty \cong \Pi_\infty$ .*

The following lemma is straightforward:

**Lemma 4.3**

- (1) Let  $\chi_s: P_0(\mathbb{A}) \rightarrow \mathbb{C}^\times$  be the unramified character given by:

$$\chi_s = \delta_P^{\frac{1}{2}+s} \cdot \delta_0^{-\frac{1}{2}}.$$

Then

$$I(s) \subset \text{Ind}_{P_0(\mathbb{A})}^{H(\mathbb{A})} \chi_s \cdot \delta_0^{\frac{1}{2}} := I_0(\chi_s).$$

(2) If  $H_v$  is split, let

$$\chi_s^0 = \delta_P^{\frac{1}{2}+s} \cdot \delta_{B_v}^{-\frac{1}{2}}.$$

Then

$$I_v(s) \subset \text{Ind}_{B_v}^{H_v} \chi_s^0 \cdot \delta_{B_v}^{\frac{1}{2}}.$$

Now let  $f_s = \otimes_v f_{v,s}$  be a standard section of  $I(s)$ . For  $\text{Re}(s)$  sufficiently large, we form the Eisenstein series:

$$(4.4) \quad E(g, f_s) := \sum_{\gamma \in P \backslash H} f_s(\gamma g), \quad g \in H(\mathbb{A}),$$

which admits a meromorphic continuation to the whole complex plane, and defines an automorphic form of  $H$  at a point  $s$  of holomorphy. We are interested in the analytic behaviour of  $E(g, f_s)$  at  $s = s_0$ , which is the same as that of its constant term  $E_{P_0}(g, f_s)$  along  $P_0$ . By a standard computation [MW, p. 92],

$$(4.5) \quad E_{P_0}(g, f_s) = \sum_{w \in \Psi} M(w, \chi_s)(f_s)(g)$$

where

$$(4.6) \quad \Psi = \{w \in W : w\Delta_M \subset \Phi^+\}$$

is the set of distinguished coset representatives for  $P_0 \backslash H/P$ , and for any  $w \in W$ ,

$$(4.7) \quad M(w, \chi_s) : I_0(\chi_s) \longrightarrow I_0(w\chi_s)$$

is the standard intertwining operator, which, for  $\text{Re}(s)$  sufficiently large, is given by:

$$(4.8) \quad \begin{aligned} M(w, \chi_s)(f_s)(g) &= \prod_v M_v(w, \chi_s)(f_{v,s})(g_v) \\ &= \prod_v \int_{U_{w,v}} f_{v,s}(w^{-1}u_v g_v) du_v \end{aligned}$$

with

$$(4.9) \quad U_w := \prod_{\alpha > 0, w^{-1}\alpha < 0} U_\alpha.$$

Here, we have chosen a representative of  $w \in W$  in  $H$ , and denoted it by  $w$  as well. The global operator does not depend on the choice of this representative, but the local operators  $M_v(w, \chi_s)$  do. However, the local operators corresponding to two different choices differ up to multiplication by a non-vanishing entire function of  $s$ , so that their analytic properties are the same. In fact, by choosing the maximal compact subgroup  $K$  suitably, we can and do normalize this choice by requiring that  $w \in K_v$  for all  $v \notin S$ .

For this normalized choice of  $w$ , there is a meromorphic function  $c_v(w, s)$  such that

$$(4.10) \quad M_v(w, \chi_s)(\Gamma_{v,s}) = c_v(w, s)\Gamma_{v,w(\chi_s)}$$

where  $\Gamma_{v,w(\chi_s)}$  is the normalized spherical vector in  $I_{0,v}(w(\chi_s))$ . This function, which is called the  $c$ -function, was computed by Gindikin-Karpelevich in the real case, and by Langlands [L] in the  $p$ -adic case. See also [R, Lemmas 6 and 7]. We have:

**Proposition 4.11** *Suppose that  $H_v$  is split. Then*

$$c_v(w, s) = \prod_{\alpha > 0, w\alpha < 0} c_v(\alpha, s),$$

where, if  $\alpha$  is a long root,

$$c_v(\alpha, s) = \frac{\zeta_v(\langle \chi_s, \alpha^\vee \rangle)}{\zeta_v(\langle \chi_s, \alpha^\vee \rangle + 1)},$$

and if  $\alpha$  is a short root,

$$\prod_{\beta \rightarrow \alpha} \frac{\zeta_v(\langle \chi_s^0, \beta^\vee \rangle)}{\zeta_v(\langle \chi_s^0, \beta^\vee \rangle + 1)}.$$

In the second case, the product is taken over all roots  $\beta \in \Phi^0$  which project onto  $\alpha \in \Phi$ , and  $\chi_s^0$  is the unramified character of  $B_v$  defined in Lemma 4.3 (2).

### 5 An Automorphic Theta Module

Now let  $f_s = \otimes_v f_{v,s}$  be a factorizable standard section, such that  $f_{v,s}$  is the normalized spherical vector  $\Gamma_{v,s}$  for all  $v \notin S$ . Then we want to understand the analytic properties of  $E(g, f_s)$  at  $s = s_0$ . By (4.5), it suffices to understand the analytic properties of  $M(w, \chi_s)(f_s)$ , for  $w \in \Psi$ . The previous proposition allows us to evaluate

$$M_S(w, \chi_s)(f_s)(1) = \prod_{v \notin S} c_v(w, s) := c_S(w, s).$$

For  $w \in \Psi$ , which has cardinality 24, let:

$$(5.1) \quad \Phi_w = \{\alpha \in \Phi^+ : w\alpha < 0\} \subset \Phi^+ \setminus \Phi_M^+.$$

$\alpha$	$\alpha^\vee$	$c_\nu(\alpha, s)$ for $D = \mathbb{H}$	$c_\nu(\alpha, s)$ for $D = \mathbb{O}$
(1,0,0,0)	(1,0,0,0)	$\zeta_\nu(17s + \frac{15}{2})/\zeta_\nu(17s + \frac{17}{2})$	$\zeta_\nu(29s + \frac{27}{2})/\zeta_\nu(29s + \frac{29}{2})$
(1,1,0,0)	(1,1,0,0)	$\zeta_\nu(17s + \frac{13}{2})/\zeta_\nu(17s + \frac{15}{2})$	$\zeta_\nu(29s + \frac{25}{2})/\zeta_\nu(29s + \frac{27}{2})$
(1,1,2,0)	(1,1,1,0)	$\zeta_\nu(17s + \frac{5}{2})/\zeta_\nu(17s + \frac{7}{2})$	$\zeta_\nu(29s + \frac{9}{2})/\zeta_\nu(29s + \frac{11}{2})$
(1,2,2,0)	(1,2,1,0)	$\zeta_\nu(17s + \frac{3}{2})/\zeta_\nu(17s + \frac{5}{2})$	$\zeta_\nu(29s + \frac{7}{2})/\zeta_\nu(29s + \frac{9}{2})$
(1,1,2,2)	(1,1,1,1)	$\zeta_\nu(17s - \frac{3}{2})/\zeta_\nu(17s - \frac{1}{2})$	$\zeta_\nu(29s - \frac{7}{2})/\zeta_\nu(29s - \frac{5}{2})$
(1,2,2,2)	(1,2,1,1)	$\zeta_\nu(17s - \frac{5}{2})/\zeta_\nu(17s - \frac{3}{2})$	$\zeta_\nu(29s - \frac{9}{2})/\zeta_\nu(29s - \frac{7}{2})$
(1,2,4,2)	(1,2,2,1)	$\zeta_\nu(17s - \frac{13}{2})/\zeta_\nu(17s - \frac{11}{2})$	$\zeta_\nu(29s - \frac{25}{2})/\zeta_\nu(29s - \frac{23}{2})$
(1,3,4,2)	(1,3,2,1)	$\zeta_\nu(17s - \frac{15}{2})/\zeta_\nu(17s - \frac{13}{2})$	$\zeta_\nu(29s - \frac{27}{2})/\zeta_\nu(29s - \frac{25}{2})$
(2,3,4,2)	(2,3,2,1)	$\zeta_\nu(34s)/\zeta_\nu(34s + 1)$	$\zeta_\nu(58s)/\zeta_\nu(58s + 1)$

Table 1: Long Roots

To use Proposition 4.11, we need to enumerate the set  $\Phi_w$  and compute  $c_\nu(\alpha, s)$ . The values of  $c_\nu(\alpha, s)$  are given in Table 1 for the 9 long roots in  $\Phi^+ \setminus \Phi_M^+$ , and in Table 2 for the 6 short roots.

Let

$$\begin{aligned}
 (5.2) \quad c_S(\alpha, s) &= \prod_{\nu \notin S} c_\nu(\alpha, s), \\
 \zeta_S(s) &= \prod_{\nu \notin S} \zeta_\nu(s).
 \end{aligned}$$

$\alpha$	$\alpha^\vee$	$c_\nu(\alpha, s)$ for $D = \mathbb{H}$	$c_\nu(\alpha, s)$ for $D = \mathbb{O}$
(1,1,1,0)	(2,2,1,0)	$\zeta_{D_\nu}(17s + \frac{9}{2}) / \zeta_{D_\nu}(17s + \frac{13}{2})$	$\zeta_{D_\nu}(29s + \frac{17}{2}) / \zeta_{D_\nu}(29s + \frac{25}{2})$
(1,1,1,1)	(2,2,1,1)	$\zeta_{D_\nu}(17s + \frac{5}{2}) / \zeta_{D_\nu}(17s + \frac{9}{2})$	$\zeta_{D_\nu}(29s + \frac{9}{2}) / \zeta_{D_\nu}(29s + \frac{17}{2})$
(1,1,2,1)	(2,2,2,1)	$\zeta_{D_\nu}(17s + \frac{1}{2}) / \zeta_{D_\nu}(17s + \frac{5}{2})$	$\zeta_{D_\nu}(29s + \frac{1}{2}) / \zeta_{D_\nu}(29s + \frac{9}{2})$
(1,2,2,1)	(2,4,2,1)	$\zeta_{D_\nu}(17s - \frac{1}{2}) / \zeta_{D_\nu}(17s + \frac{3}{2})$	$\zeta_{D_\nu}(29s - \frac{1}{2}) / \zeta_{D_\nu}(29s + \frac{7}{2})$
(1,2,3,1)	(2,4,3,1)	$\zeta_{D_\nu}(17s - \frac{5}{2}) / \zeta_{D_\nu}(17s - \frac{1}{2})$	$\zeta_{D_\nu}(29s - \frac{9}{2}) / \zeta_{D_\nu}(29s - \frac{1}{2})$
(1,2,3,2)	(2,4,3,2)	$\zeta_{D_\nu}(17s - \frac{9}{2}) / \zeta_{D_\nu}(17s - \frac{5}{2})$	$\zeta_{D_\nu}(29s - \frac{17}{2}) / \zeta_{D_\nu}(29s - \frac{9}{2})$

Table 2: Short Roots

From the above results, one sees that when  $D = \mathbb{H}$ ,  $c_S(\alpha, s)$  is finite and non-zero at  $s = s_0$  except possibly for the following  $\alpha$ 's:

$\alpha$	$c_S(\alpha, s)$	Behaviour at $s = s_0$
$\alpha(0) = (1,3,4,2)$	$\zeta_S(-2) / \zeta_S(-1)$	Zero of order 1
$\alpha(1) = (1,2,4,2)$	$\zeta_S(-1) / \zeta_S(0)$	Pole of order $ S  - 1$
$\alpha(2) = (1,2,3,2)$	$(\zeta_S(1) \cdot \zeta_S(0)) / (\zeta_S(3) \cdot \zeta_S(2))$	Zero of order $ S  - 2$

As for the case  $D = \mathbb{O}$ ,  $c_S(\alpha, s)$  is finite and non-zero at  $s = s_0$  except for the following  $\alpha$ 's:

$\alpha$	$c_S(\alpha, s)$	Behaviour at $s = s_0$
$\alpha(0) = (1,3,4,2)$	$\zeta_S(-4) / \zeta_S(-3)$	Zero of order 1
$\alpha(1) = (1,2,4,2)$	$\zeta_S(-3) / \zeta_S(-2)$	Pole of order 1

Recall that:

$$(5.3) \quad c_\nu(w, s) = \prod_{\alpha \in \Phi_w} c_\nu(\alpha, s).$$

For those  $w \in \Psi$ , such that  $\Phi_w$  does not contain any of the roots  $\alpha(i)$ ,  $i = 0, 1, 2$ , the function  $c_S(w, s) = \prod_{\alpha \in \Phi_w} c_S(\alpha, s)$  is thus finite and non-zero at  $s = s_0$ . The only  $w$ 's not accounted for are:

$$(5.4) \quad \begin{aligned} w_0 &= (1, 2, 3, 2, 1, 4, 3, 2, 1, 3, 2, 4, 3, 2, 1), \\ w_{-1} &= (2, 3, 2, 1, 4, 3, 2, 1, 3, 2, 4, 3, 2, 1), \\ w_{-2} &= (3, 2, 1, 4, 3, 2, 1, 3, 2, 4, 3, 2, 1). \end{aligned}$$

For these, we list the set  $\Phi_w$ :

$$\begin{aligned} \Phi_{w_0} &= \Phi^+ - \Phi_M^+, \\ \Phi_{w_{-1}} &= \Phi_{w_0} - \{\alpha(0)\}, \\ \Phi_{w_{-2}} &= \Phi_{w_0} - \{\alpha(0), \alpha(1)\}. \end{aligned}$$

Now we have:

**Lemma 5.5**

- (i)  $c_S(w_{-2}, s)$  is holomorphic at  $s = s_0$ ; indeed it has a zero of order  $|S| - 2$  if  $D = \mathbb{H}$ .
- (ii)  $c_S(w_0, s)$  is finite and non-zero at  $s = s_0$ .
- (iii)  $c_S(w_{-1}, s)$  has a pole of order 1 at  $s = s_0$ .

It remains now to analyze the finitely many terms  $M_\nu(w, \chi_s)(f_{\nu,s})$  for  $\nu \in S$ . Using the functional equation for intertwining operators, we can write  $M_\nu(w, \chi_s)$  as a product of simple intertwining operators, *i.e.*, those corresponding to simple reflections. Using the well-known analytic properties of intertwining operators for rank 1 groups, we deduce that if  $w \neq w_i$ ,  $i = 0, -1$ , the integral defining  $M_\nu(w, \chi_s)(f_{\nu,s})$  converges at  $s = s_0$ , for any choice of  $f_{\nu,s}$ . Thus, for  $w \neq w_i$ ,  $i = 0, -1$ ,  $M(w, \chi_s)(f_s)$  is holomorphic at  $s = s_0$ .

Similarly, we deduce that, at  $s = s_0$ ,  $M_\nu(w_{-1}, \chi_s)$  is holomorphic for all  $\nu \in S$ , whereas  $M_\nu(w_0, \chi_s)$  is holomorphic for  $\nu$  finite, and can have a pole of order  $\leq 1$  at the real place. Now we have the following crucial lemma:

**Lemma 5.6** For  $\nu \in S$ , the intertwining operator

$$M_\nu(w_{-1}, \chi_{s_0}): I_\nu(s_0) \longrightarrow I_{0,\nu}(w_{-1}(\chi_{s_0}))$$

is not identically zero.

**Proof** First, note that the double coset  $P_\nu w_0 P_\nu = P_\nu w_0 N_\nu$  is open in  $H_\nu$ , and any element  $g \in P_\nu w_0 P_\nu$  has a unique expression  $g = pw_0n$  with  $p \in P_\nu$  and  $n \in N_\nu$ .

For  $\nu$  real, we let  $\phi$  be a smooth real-valued non-negative function on  $N_\nu$  with compact support, and set

$$f_s(g) = \begin{cases} \delta(p)^{\frac{1}{2}+s}\phi(n), & \text{if } g = pw_0n; \\ 0, & \text{otherwise.} \end{cases}$$

Then  $f_s \in I_\nu(s)^\infty$  is an entire section, though not standard. By results of Vogan and Wallach [Wa, Chapter 10],  $M_\nu(w_{-1}, \chi_{s_0})$  is an intertwining operator defined on  $I_\nu(s_0)^\infty$ , and is continuous with respect to the natural Fréchet topology on the smooth representations. Now, for  $\text{Re}(s)$  sufficiently large, and  $g = w_{-1}w_0$ ,

$$\begin{aligned} M_\nu(w_{-1}, \chi_s)(f_s)(g) &= \int_{g^{-1}U_{w_{-1}}g} f_s(w_0u) \, du \\ &= \int_{g^{-1}U_{w_{-1}}g} \phi(u) \, du \end{aligned}$$

since  $g^{-1}U_{w_{-1}}g \subset N_\nu$ . For a suitable choice of  $\phi$ , we can certainly ensure that this last integral is non-zero. Since it is also independent of  $s$ , the meromorphic function

$$s \mapsto M_\nu(w_{-1}, \chi_s)(f_s)(g)$$

is constant and non-zero. Since  $I_\nu(s_0)$  is dense in  $I_\nu(s_0)^\infty$ , we deduce that  $M_\nu(w_{-1}, \chi_{s_0})$  must be non-zero on some  $K_\nu$ -finite vector. This proves the lemma for  $\nu$  real.

Now assume that  $\nu \in S$  is finite. Let  $f_s \in I(s)$  be defined by:

$$f_s(g) = \begin{cases} \delta(p)^{\frac{1}{2}+s}, & \text{if } g = pw_0n \in P_\nu w_0(N_\nu \cap K_\nu); \\ 0, & \text{otherwise.} \end{cases}$$

Then  $f_s$  is a standard section, and the same argument as above proves the lemma. ■

**Corollary 5.7**

- (i) For  $\nu \in S$  finite,  $M_\nu(w_0, \chi_s)$  is holomorphic at  $s = s_0$  and  $M_\nu(w_0, \chi_{s_0})$  is not identically zero on  $I_\nu(s_0)$ .
- (ii) For  $\nu$  real,  $M_\nu(w_0, \chi_s)$  has a pole of order at most 1 at  $s = s_0$ , and this pole is attained for some vector in  $I_\nu(s_0)$ .

**Proof** This follows from the factorization

$$M_\nu(w_0, \chi_s) = M_\nu(w_1, w_{-1}(\chi_s)) \circ M_\nu(w_{-1}, \chi_s),$$

the lemma above and the well-known analytic properties of  $SL_2$ -intertwining operators. ■

Now we have:

**Theorem 5.8** For any standard section  $f_s \in I(s)$ , the Eisenstein series  $E(g, f_s)$  has at most a simple pole at  $s = s_0$ . Moreover, this pole is actually attained for some standard section  $f_s$  with  $f_{v,s} = \Gamma_{v,s}$  for all  $v \notin S$ . Let

$$(5.9) \quad \begin{aligned} \theta: I(s_0) &\rightarrow \mathcal{A}(H) \\ f &\mapsto \text{Res}_{s=s_0} E(g, f_s) \end{aligned}$$

and let  $\Theta$  be its image. Then  $\Theta$  is an irreducible square-integrable automorphic representation isomorphic to  $\Pi$ .

**Proof** If the standard section  $f_s$  is such that  $f_{v,s} = \Gamma_{v,s}$  for all  $v \notin S$ , then we have seen that  $E(g, f_s)$  has at most a simple pole at  $s = s_0$ . Now we claim that this is true for any other standard sections as well. Suppose not; then for any decomposable  $f_s$  such that  $E(g, f_s)$  has a pole at  $s = s_0$  of order greater than 1, let

$$S_f := \{v \notin S : f_{v,s} \text{ is not spherical}\}.$$

Hence,  $S_f$  is non-empty. Choose  $f_s$  such that  $S_f$  is minimal, and suppose that  $v_0 \in S_f$ . Then consider:

$$\begin{aligned} I_{v_0}(s_0) &\rightarrow \mathcal{A}(H) \\ \phi &\mapsto \lim_{s \rightarrow s_0} (s - s_0)^k E(g, \phi_s) \end{aligned}$$

where  $\phi_s$  is the unique standard section satisfying:

$$\phi_{s_0} = \phi \otimes \left( \bigotimes_{v \neq v_0} f_{v,s_0} \right)$$

and  $k > 1$  is the highest order of pole attained by  $E(g, \phi_s)$  at  $s = s_0$  for all such  $\phi_s$ . Such a  $k$  exists since  $I_{v_0}(s_0)$  has finite length. By the definition of  $k$ , this map is a non-zero  $H_{v_0}$ -intertwining map, and by the minimality of  $S_f$ , it vanishes on the spherical vector. But by Proposition 3.2(ii),  $I_{v_0}(s_0)$  is generated by the spherical vector as a representation of  $H_{v_0}$ . With this contradiction, the claim is proved.

To see that the pole is actually attained for some sections, it remains to show that the poles of  $M(w_{-1}, \chi_s)$  and  $M(w_0, \chi_s)$  at  $s = s_0$  do not cancel. For this, note that the residue of  $M(w_i, \chi_s)(f_s)$  at  $s = s_0$ , when regarded as a function of the maximal split torus  $A$ , is the unramified character  $\delta_0^{\frac{1}{2}} \cdot w_i(\chi_{s_0})$ . Hence it suffices to see that these two unramified characters are different. One checks that if  $D = \mathbb{H}$ ,

$$(5.10) \quad \begin{aligned} w_0(\chi_{s_0}) &= -11\alpha_1 - 24\alpha_2 - 36\alpha_3 - 20\alpha_4, \\ w_{-1}(\chi_{s_0}) &= -13\alpha_1 - 24\alpha_2 - 36\alpha_3 - 20\alpha_4, \end{aligned}$$

whereas if  $D = \mathbb{O}$ ,

$$(5.11) \quad \begin{aligned} w_0(\chi_{s_0}) &= -19\alpha_1 - 42\alpha_2 - 60\alpha_3 - 32\alpha_4, \\ w_{-1}(\chi_{s_0}) &= -23\alpha_1 - 42\alpha_2 - 60\alpha_3 - 32\alpha_4. \end{aligned}$$

Hence, we see that the pole at  $s = s_0$  is actually attained. Moreover, the above shows that all the cuspidal exponents of  $\Theta$  have strictly negative coefficients. This implies, by Jacquet’s criterion [MW, p. 74] that  $\Theta$  is contained in  $L^2(H \setminus H(\mathbb{A})) \cap \mathcal{A}(H)$ .

Since  $\Theta$  is square-integrable, it is semi-simple. Suppose that  $\Pi_1 \subset \Theta$  is an irreducible summand. Then, for  $v$  finite,  $(\Pi_1)_v$  is the unique irreducible quotient of  $I_v(s_0)$ , by Proposition 3.2(ii) and (iii). Moreover, for  $v \notin S$ ,  $(\Pi_1)_v \cong \Pi_v$ . Now by a rigidity result of Kazhdan [R2, Proposition 57], this implies that  $(\Pi_1)_v$  is also the minimal representation for all  $v \in S$ , in particular for  $v$  the real place. Thus,  $\Pi_1 \cong \Pi$ , and in view of Proposition 4.2, we deduce that  $\Theta$  must be irreducible. This proves the theorem. ■

**Corollary 5.12** *For  $v \in S$  finite, the minimal representation  $\Pi_v$  is the unique irreducible quotient of  $I_v(s_0)$ .*

**Corollary 5.13** *For any non-zero  $f \in \Pi \subset I(-s_0)$ , let  $f_s$  be the unique standard section extending  $f$ . Then  $E(g, f_s)$  is holomorphic at  $s = -s_0$ , and  $E(g, f_{-s_0})$  generates  $\Theta$ .*

**Proof** This follows from the functional equation of Eisenstein series. ■

**Remarks** (i) It seems likely that there is exactly one automorphic realization for  $\Pi$ . As in [GRS1], this uniqueness statement would follow if the multiplicity one result for Jacquet modules in Proposition 3.2(ii) holds for all places  $v$ .

(ii) When  $D$  is an indefinite quaternion algebra, the data in this section allows one to conclude that the minimal representation for the corresponding group is automorphic as well.

## 6 Fourier Coefficients

In this section, we consider the Fourier coefficients of  $\theta = \theta(\otimes_v f_v) \in \Theta$  along the unipotent radical  $N$  of the Heisenberg parabolic subgroup  $P = M \cdot N$ . Recall that  $N$  is a Heisenberg group with center  $Z$ , and  $V = N/Z$ .

Consider the constant term of  $\theta$  along  $Z$ :

$$(6.1) \quad \theta_Z(g) = \int_{Z \backslash Z(\mathbb{A})} \theta(zg) dz.$$

Note that  $\theta_Z$  is non-zero since the constant term of  $\theta$  with respect to  $P_0$  is non-zero. We consider its Fourier expansion along the compact group  $V \setminus V(\mathbb{A})$ . For a character  $\psi$  of  $V \setminus V(\mathbb{A})$ , the  $\psi$ -Fourier coefficient of  $\theta_Z$  is:

$$(6.2) \quad \theta_\psi(g) = \int_{V \setminus V(\mathbb{A})} \theta_Z(ng)\psi(n)^{-1} dn.$$

Let  $C_\psi$  be the stabilizer of  $\psi$  in  $M^1$ . If  $\psi \in \Omega$ , then  $C_\psi$  is a conjugate of the derived group of the maximal parabolic subgroup  $Q$  of  $M^1$ . As in [GRS1] and [GrS], we have the following two important properties of  $\theta_\psi$ :

**Proposition 6.3** *Suppose that  $\psi$  is non-trivial. For any non-zero  $\theta \in \Theta$ ,  $\theta_\psi$  is non-zero if and only if  $\psi \in \Omega$ .*

**Proposition 6.4** *Suppose that  $\psi \in \Omega$ . Then for all  $c \in C_\psi(\mathbb{A})$ , we have:*

$$\theta_\psi(cg) = \theta_\psi(g).$$

Finally, we consider the constant term  $\theta_N$ , regarded as an automorphic form on  $M^1$ . By computing the constant term  $E_P(g, f_s)$  of  $E(g, f_s)$  along  $P$ , we have:

**Proposition 6.5** *As an automorphic form on  $M^1$ ,*

$$\theta_N = c + \theta'$$

where  $c$  is a constant function, and  $\theta'$  is contained in an automorphic realization of the global minimal representation of  $M^1$ .

One can further compute the constant term of  $E_P(g, f_s)$  along the unipotent radical  $U$  of the maximal parabolic  $Q$  in  $M^1$ . Then one finds that the constant term  $(\theta_N)_U$  of  $\theta_N$  along  $U$  is simply the sum of two characters of the Levi factor of  $Q$ . This reflects the fact that, for  $v \in S$ ,  $L_v$  does not have a minimal representation.

## 7 Dual Pairs and Étale Cubic Algebras

We shall use the automorphic theta module  $\Theta$  and the properties of its Fourier coefficients discussed above to study theta correspondence. In this section, we briefly describe the dual pair  $G \times G'$ .

Let  $e \in J$  be the identity, and let  $G' \subset L$  be the algebraic subgroup stabilizing  $e$ . Then  $G'$  is the automorphism group of the Jordan algebra structure on  $J$ , and

$$(7.1) \quad G' \cong \begin{cases} \text{PU}_3(D), & \text{if } D = \mathbb{H}; \\ F_4^{cpt}, & \text{if } D = \mathbb{O}. \end{cases}$$

Here,  $\text{PU}_3(D)$  is the projective unitary group in three variables with coefficients in  $D$ . Note that  $G'$  acts naturally on  $J_0$ , the space of trace zero elements in  $J$ . Moreover,  $G'_\infty$  is compact, whereas for finite  $v \in S$ ,  $G'_v$  has rank 1, and for  $v \notin S$ ,  $G'_v$  is split. Let  $G$  be the split group of type  $G_2$ . Then  $G \times G'$  is a reductive dual pair in  $H$  (see [MS] and [SG]).

There is an embedding of  $G \times G'$  in  $H$  such that

$$(7.2) \quad (G \times G') \cap P = P_2 \times G'$$

where  $P_2 = L_2 \cdot U_2$  is the Heisenberg parabolic subgroup of  $G$ . Here,  $U_2$  is a Heisenberg group with center  $Z$ , and  $\bar{V} := \bar{U}_2/Z$  is the subspace  $F \oplus Fe \oplus Fe \oplus F$  of  $\bar{V}$ . Moreover,  $L_2 \cong \text{GL}_2$ , and its action on  $\bar{V}$  is isomorphic to  $\det \otimes \text{Sym}^3(F^2)^*$ . Thus we can identify  $\bar{V}$  with the space of binary cubic forms, and the non-zero  $L_2$ -orbits are then parametrized by cubic  $F$ -algebras [Wr]. We list the non-zero orbits below:

- (i)  $S_0 = L_2 \cdot (0, 0, 0, 1)$ ; this is the orbit corresponding to  $E_0 = F[\varepsilon]/\varepsilon^3$ .
- (ii)  $S_1 = L_2 \cdot (0, 0, 1, 0)$ ; this is the orbit corresponding to  $E_1 = F \oplus F[\varepsilon]/\varepsilon^2$ .
- (iii) For each étale cubic algebra  $E$ , there is a generic orbit  $S_E$ . Given  $E$ , let  $E^0$  denote the two-dimensional space of trace zero elements. Then the norm form of  $E$  restricts to give a binary cubic form on  $E^0$ . This form is an element in the orbit  $S_E$ .

Since  $\bar{V}$  can be identified with the characters of  $Z(\mathbb{A})U_2 \backslash U_2(\mathbb{A})$ , we see that the  $L_2$ -orbits of non-trivial characters are parametrized by cubic algebras over  $F$ .

For  $\psi_E \in S_E$ , a non-zero orbit, let

$$\Omega_E = \{ \psi \in \Omega : \psi|_{U_2(\mathbb{A})} = \psi_E \}.$$

Clearly,  $G'$  acts on  $\Omega_E$ , which as a  $G'$ -set, depends only on  $E$ . As an example, the  $G'$ -set  $\Omega_1 := \Omega_{E_1}$  can be identified as:

$$\Omega_1 = \{ X \in J : \text{rank}(X) = 1 = \text{Tr}(X) \}.$$

Note that  $\Omega_1$  is clearly non-empty, but this is not always the case when  $E$  is étale. Indeed, we have (see [GrG, Proposition 1]):

**Lemma 7.3**

- (1) Suppose that  $E$  is étale. Then  $\Omega_E$  is non-empty if and only if  $E$  is totally real, in which case  $G'$  acts transitively on  $\Omega_E$ . The algebraic subgroup of  $G'$  stabilizing an element of  $\Omega_E$  is:

$$C_E = \begin{cases} \text{Res}_{E/F}(\text{SL}(D \otimes_F E)) / \mu_2, & \text{if } D = \mathbb{H}; \\ \text{Spin}_8^{E, \text{cpt}}, & \text{if } D = \mathbb{O}. \end{cases}$$

- (2)  $G'$  acts transitively on  $\Omega_1$ , and the algebraic subgroup of  $G'$  stabilizing an element is:

$$C = \begin{cases} (\text{SU}_2(D) \times \text{SL}(D)) / \mu_2, & \text{if } D = \mathbb{H}; \\ \text{Spin}_9^{\text{cpt}}, & \text{if } D = \mathbb{O}. \end{cases}$$

### 8 Cuspidality of Theta Lifts

Let  $\pi = \hat{\otimes}_v \pi_v$  be an irreducible automorphic representation of  $G'$ , and let  $\alpha$  be an automorphic form in the space of  $\pi$ . Recall that the theta lift of  $\alpha$  by  $\theta \in \Theta$  is defined to be:

$$(8.1) \quad \beta(g) = \int_{G' \backslash G'(\mathbb{A})} \alpha(g') \theta(gg') dg'.$$

Note that this integral converges because  $G' \setminus G'(\mathbb{A})$  is compact. Moreover  $\beta$  is an automorphic form on  $G$ . Let  $\Theta(\pi)$  be the subspace of  $\mathcal{A}(G)$  spanned by all such  $\beta$ 's. We would like to investigate conditions under which  $\Theta(\pi)$  is non-zero and cuspidal, when  $\pi$  is a non-trivial automorphic representation.

In this section, we study the cuspidality question. Hence we need to compute the constant terms of  $\beta$  along the unipotent radicals of the two maximal parabolic subgroups of  $G$ . Recall that  $P_2 = L_2 \cdot U_2$  is the Heisenberg parabolic of  $G$ . Let  $P_1 = L_1 \cdot U_1$  be the other maximal parabolic of  $G$ . We introduce the following notations:

$$U_{12} = U_1 \cap U_2, \quad V_{12} = U_{12}/Z,$$

$$V_1 = U_1/U_1 \cap U_2, \quad V_2 = U_2/U_1 \cap U_2.$$

We first compute:

$$(8.2) \quad \begin{aligned} \beta_{U_{12}}(g) &= \int_{U_{12} \setminus U_{12}(\mathbb{A})} \beta(ug) \, du \\ &= \int_{V_{12} \setminus V_{12}(\mathbb{A})} \int_{G' \setminus G'(\mathbb{A})} \alpha(g') \cdot \theta_Z(vgg') \, dg' \, dv \\ &= \int_{G' \setminus G'(\mathbb{A})} \alpha(g') \left( \theta_N(gg') + \sum_{\psi \in \mathcal{O}_0} \theta_\psi(gg') \right) \, dg' \end{aligned}$$

where  $\mathcal{O}_0$  is defined in Lemma 2.7. By Proposition 6.4,  $\theta_\psi(gg') = \theta_\psi(g)$ , for  $\psi \in \mathcal{O}_0$ . Hence,

$$(8.3) \quad \beta_{U_{12}}(g) = \int_{G' \setminus G'(\mathbb{A})} \alpha(g') \theta_N(gg') \, dg'.$$

It follows that  $\beta_{U_2} = \beta_{U_{12}}$ . Moreover, by Proposition 6.5, we know that, as an automorphic form on  $M^1$ ,  $\theta_N$  is equal to  $c + \theta'$ , where  $c$  is a constant function of  $M^1$  and  $\theta'$  lies in the global minimal representation of  $M^1$ . The integral of  $\alpha$  against  $c$  is zero, since  $\alpha$  is non-constant. On the other hand, note that  $SL_2 \times G'$  is a commuting pair in  $M^1$ , where  $SL_2$  is the derived group of  $L_2 \cong GL_2$ . Hence,  $\beta_{U_{12}}$ , regarded as a function on  $SL_2$ , is nothing but the theta lift of  $\alpha$  by  $\theta'$ . Moreover  $V_1$  is nothing but the unipotent radical of a Borel subgroup of  $SL_2$ . It follows that, as functions on  $SL_2$ ,  $\beta_{U_1}$  is simply the constant term of  $\beta_{U_{12}}$  along  $V_1$ . Hence, we need to study the Fourier coefficients of  $\beta_{U_{12}}$  along  $V_1$ .

For this, we consider the Fourier expansion of  $\theta'$  along the unipotent radical  $U \cong J$  of  $Q$ . The analogues of Propositions 6.3 and 6.4 hold for  $\theta'$  [GrS]. In particular, the characters of  $U \setminus U(\mathbb{A})$  can be parametrized by  $J$ , and

$$(8.4) \quad \theta' = \theta'_U + \sum_{X \in J: \text{rank}(X)=1} \theta'_{\psi_X}.$$

Moreover, for  $X$  of rank 1,

$$(8.5) \quad \theta'_{\psi_X}(cg) = \theta'_{\psi_X}(g)$$

for  $c \in C_{\psi_X}(\mathbb{A})$ , the stabilizer of  $\psi_X$  in  $L(\mathbb{A})$ . Recall that  $L$  is the derived group of the Levi factor of  $Q$ . As we have noted before,  $\theta'_U$  is a constant function of  $L$  and so its integral against  $\alpha$  is zero. Further, since there are no rank one elements of  $J$  with trace zero, we deduce easily that:

$$\beta_{U_1}(g) = \int_{V_1 \backslash V_1(\mathbb{A})} \beta_{U_{12}}(vg) \, dv = 0.$$

On the other hand, for any non-trivial character  $\psi$  of  $V_1 \backslash V_1(\mathbb{A})$ , let

$$\Omega_\psi = \{X \in J : \text{rank}(X) = 1 \text{ and } \psi_X|_{V_1(\mathbb{A})} = \psi\}.$$

This is a  $G'$ -homogeneous space isomorphic to  $\Omega_1$ . Then the  $\psi$ -Fourier coefficient of  $\beta_{U_{12}}$  is:

$$\begin{aligned} (\beta_{U_{12}})_{V_1, \psi}(g) &= \int_{G' \backslash G'(\mathbb{A})} \sum_{X \in \Omega_\psi} \alpha(g') \theta'_{\psi_X}(gg') \, dg' \\ &= \int_{C(\mathbb{A}) \backslash G'(\mathbb{A})} \theta'_{\psi_{X_0}}(gg') \cdot \left( \int_{C \backslash C(\mathbb{A})} \alpha(cg') \, dc \right) \, dg', \end{aligned}$$

where  $X_0$  is an element of  $\Omega_\psi$ , with stabilizer  $C$ . In conclusion, we have:

**Proposition 8.6** *An element of  $\Theta(\pi)$  is either cuspidal or is concentrated along the Heisenberg parabolic of  $G$ . Moreover,  $\Theta(\pi)$  is cuspidal if and only if the linear functional*

$$P^C : \alpha \mapsto \int_{C \backslash C(\mathbb{A})} \alpha(c) \, dc$$

*is identically zero on  $\pi$ , i.e.,  $\pi$  is not  $C$ -distinguished.*

**Proof** The sufficiency for the vanishing of  $P^C$  is clear. As for the necessity, we argue as in [GS, Section 5, Proposition 4.5] that if the period  $P^C(\alpha)$  is non-zero, then we can choose  $\theta$  such that  $(\beta_{U_{12}})_{V_1, \psi}$  is non-zero. See also the proof of Proposition 9.3 below. ■

### 9 Non-Vanishing of Theta Lifts

In this section, we investigate when  $\Theta(\pi)$  is non-zero. To do this, we study the Fourier expansion of  $\beta$  along  $U_2$ . First note:

**Lemma 9.1**  *$\beta$  is zero if and only if  $\beta_Z$  is zero.*

**Proof** Suppose that  $\beta_Z = 0$ . Let  $Z_1 \supset Z$  be the center of  $U_1$ . Then certainly  $\beta_{Z_1, \psi} = 0$  for any character  $\psi$  of  $Z_1 \backslash Z_1(\mathbb{A})$  which is trivial on  $Z(\mathbb{A})$ . But  $L_1$  acts transitively on the non-trivial characters of  $Z_1 \backslash Z_1(\mathbb{A})$ . This implies that  $\beta = 0$ , as required. ■

Hence, to investigate the non-vanishing of  $\beta$ , it suffices to consider  $\beta_Z$ . Let  $\psi_E$  be a character of  $U_2(\mathbb{A})$  in the  $L_2$ -orbit  $S_E$ . Then, as in the last section, the  $\psi_E$ -Fourier coefficient of  $\beta_Z$  is given by:

$$(9.2) \quad \beta_{\psi_E}(g) = \int_{C_E(\mathbb{A}) \backslash G'(\mathbb{A})} \theta_{\tilde{\psi}_E}(gg') P^{C_E}(\alpha, g') dg'$$

where  $\tilde{\psi}_E$  is an element of  $\Omega_E$  with stabilizer  $C_E$ , and

$$P^{C_E}(\alpha, g') = \int_{C_E \backslash C_E(\mathbb{A})} \alpha(cg') dc.$$

Now we have:

**Proposition 9.3** *Assume that  $\Theta(\pi)$  is cuspidal. Then the linear functional  $\mathcal{F}_E: \beta \mapsto \beta_{\psi_E}(1)$  is identically zero on  $\Theta(\pi)$  unless  $E$  is étale and totally real, in which case it is non-zero if and only if the linear functional  $\alpha \mapsto P^{C_E}(\alpha) := P^{C_E}(\alpha, 1)$  is non-zero on  $\pi$ .*

**Proof** By Proposition 8.6, the assumption that  $\Theta(\pi)$  is cuspidal implies that  $\beta_{\psi_{E_1}} = 0$  for all  $\beta \in \Theta(\pi)$ . Moreover, the fact that  $\pi$  is non-trivial implies that  $\beta_{\psi_{E_0}} = 0$ , since  $C_{E_0} = G'$ . Also, if  $E$  is étale but not totally real,  $\Omega_E$  is empty by Lemma 7.3, so that  $\beta_{\psi_E} = 0$ .

By (9.2), it is clear that the vanishing of  $P^{C_E}$  implies that of  $\mathcal{F}_E$ . The proof of the converse is along the lines of [GrS, Section 5, Proposition 4.5]. So suppose that  $P^{C_E}(\alpha)$  is non-zero. Note that  $P^{C_E}(\alpha, g')$  descends to a smooth function  $\alpha_E$  on  $\Omega_E(\mathbb{A})$ . Choose a neighbourhood  $\mathcal{N} = \prod_{\nu} \mathcal{N}_{\nu}$  of  $\tilde{\psi}_E \in \Omega_E(\mathbb{A})$  with  $\mathcal{N}_{\nu}$  open compact for finite  $\nu$ . By shrinking  $\mathcal{N}$ , we can ensure that  $\alpha_E(x) = \alpha_E(x_{\infty})$ , for all  $x \in \mathcal{N}$ . Note that for almost all  $\nu$ ,  $\mathcal{N}_{\nu} = \Omega_E(\mathbb{Z}_{\nu}) = G'(\mathbb{Z}_{\nu}) \cdot \tilde{\psi}_E$ .

Similarly, the restriction of  $\theta_{\tilde{\psi}_E}$  to  $M^1(\mathbb{A})$  descends to a function on  $C_{\tilde{\psi}_E}(\mathbb{A}) \backslash M^1(\mathbb{A})$ . As in [GrS, Section 5, (3.11)], for suitable  $\theta$ , this function can be written as a product  $\prod_{\nu} f_{\nu}$  where  $f_{\nu}$  is some smooth function on  $\Omega_{\nu}$  which, for almost all finite places  $\nu$ , is equal to a distinguished function  $f_{\nu}^0$  (corresponding to the normalized spherical vector in  $\Pi_{\nu}$ ). Now, for almost all places  $\nu$ , the restriction of  $f_{\nu}^0$  to  $\Omega_{E,\nu}$  is the characteristic function of  $\Omega_E(\mathbb{Z}_{\nu})$  (see [GrG, Proposition 2]). On the other hand, at the other finite places, it follows from [MS, Theorem 6.1] that  $f_{\nu}$  can be any locally constant, compactly supported function on  $\Omega_{\nu}$ . Since  $\Omega_{E,\nu}$  is a closed subset of  $\Omega_{\nu}$ , we can choose  $f_{\nu}$  such that its restriction to  $\Omega_{E,\nu}$  is the characteristic function of  $\mathcal{N}_{\nu}$ . Hence, we can take  $f^{\infty} = \bigotimes_{\nu \neq \infty} f_{\nu}$  such that its restriction to  $\Omega_E(\mathbb{A}^{\infty})$  is the characteristic function of  $\mathcal{N}^{\infty} = \prod_{\nu \neq \infty} \mathcal{N}_{\nu}$ . Thus, up to multiplication by a non-zero constant, we have:

$$\mathcal{F}_E(\beta) = \int_{\Omega_E(\mathbb{R})} f_{\infty}(x) \alpha_E(x) dx.$$

Now as in [GrS], the fact that  $\Omega_E(\mathbb{R})$  is compact implies, by the Stone-Weierstrass theorem, that the restrictions of the functions  $f_{\infty}$  span a dense subspace of the space of smooth functions on  $\Omega_E(\mathbb{R})$ . Hence the above integral is non-zero for a suitable choice of  $f_{\infty}$ . This completes the proof of the proposition. ■

By Propositions 8.6 and 9.3, we have:

**Theorem 9.4** *If  $\pi$  is a non-trivial automorphic representation of  $G'$ , then  $\Theta(\pi)$  is non-zero and cuspidal if and only if the following two conditions hold:*

- (i)  $\pi$  is not  $C$ -distinguished.
- (ii)  $\pi$  is  $C_E$ -distinguished for some totally real étale cubic algebra  $E$ .

**Acknowledgments** I am indebted to Professor Benedict Gross and Professor Gordan Savin for their guidance and advice. I would also like to thank Karl Rumelhart for a number of useful discussions on Eisenstein series and intertwining operators, and for providing me with his preprint [R].

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