# A NEW PROOF OF N. J. YOUNG'S THEOREM ON THE ORBITS OF THE ACTION OF THE SYMPLECTIC GROUP 

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#### Abstract

The group of symplectic transformations acts on the unit ball of a Hilbert space. The structure of the orbits has been determined by N. J. Young in [8]. We provide a new proof of this theorem; it is slightly simpler than the original one, and does not involve Brown-Douglas-Fillmore theory. Moreover, the steps followed hopefully throw some additional light on the subject. We rely heavily on previous work of Khatskevich, Shmulyan and Shulman ( $[5,6,7]$ ); the proofs of the results used are included for completeness.


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## 1. Young's theorem

Suppose $\mathcal{H}, \mathcal{K}$ are Hilbert spaces, and let $\mathcal{B}$ denote the closed unit ball of $\mathcal{L}(\mathcal{H}, \mathcal{K})$. The problem considered by Young in [8] is that of characterizing the orbit of the action of the symplectic group on $\mathcal{B}$. The symplectic group $G$ is defined as the group of all $2 \times 2 J$-unitary operator matrices; that is, all invertible matrices $g=\left(\begin{array}{cc}A & B \\ C & D\end{array}\right)$, $A \in \mathcal{L}(\mathcal{K}), B \in \mathcal{L}(\mathcal{H}, \mathcal{K}), C \in \mathcal{L}(\mathcal{K}, \mathcal{H}), D \in \mathcal{L}(\mathcal{H})$ which satisfy the relation

$$
\left(\begin{array}{ll}
A & B \\
C & D
\end{array}\right)^{-1}=\left(\begin{array}{cc}
A^{*} & -C^{*} \\
-B^{*} & D^{*}
\end{array}\right) .
$$

The action of $G$ on $\mathcal{B}$ is given by the maps

$$
\Phi_{g}(X)=(A X+B)(C X+D)^{-1}
$$

On the open unit ball (the set of all strict contractions) $G$ acts transitively. This is no more the case on $\mathcal{B}$; the orbits of the action have a rather complicated structure, which is determined by Young's result.

Theorem 1. $X$ and $Y$ lie in the same orbit if and only if the following conditions are satisfied:
(i) there exist invertible operators $P$ and $Q$, such that

$$
\begin{equation*}
I-Y^{*} Y=P^{*}\left(I-X^{*} X\right) P, I-Y Y^{*}=Q^{*}\left(I-X X^{*}\right) Q \tag{1}
\end{equation*}
$$

(ii) if $X$ is essentially unitary (and thus also Y by (1)), then

$$
\begin{equation*}
\operatorname{ind} X=\operatorname{ind} Y \tag{2}
\end{equation*}
$$

Conditions (i) are simple consequences of the identities:

$$
\begin{align*}
& I-\Phi_{g}(X)^{*} \Phi_{g}(X)=\left(X^{*} C^{*}+D^{*}\right)^{-1}\left(I-X^{*} X\right)(C X+D)^{-1}  \tag{3}\\
& I-\Phi_{g}(X) \Phi_{g}(X)^{*}=\left(X B^{*}+A^{*}\right)^{-1}\left(I-X X^{*}\right)\left(B X^{*}+A\right)^{-1} \tag{4}
\end{align*}
$$

The necessity of (ii) is less apparent and will be shown later.
Remark. It should be mentioned that Theorem 1 has been independently obtained, in the different context of Krein spaces, by Azizov ([2]). A readable exposition can be found in [3].

## 2. Preliminary facts

We will use repeatedly the following simple lemma (see, for instance, [4]).
Lemma 1 (a) For $A, B$ bounded operators, the following statements are equivalent.
(i) There exists $X$ bounded, such that $A=X B$.
(ii) There exists $Y$ bounded, such that $A^{*} A=B^{*} Y B$.
(iii) range $A^{*} \subset$ range $B^{*}$.
(iv) $A^{*} A \leq c B^{*} B$ for some constant $c>0$.
(b) If $Y$ is invertible in (ii), then we also have $B=X^{\prime} A$ for some bounded $X^{\prime}$.

For any contraction $T$, the symbol $D_{T}$ will denote the defect operator $\left(I-T^{*} T\right)^{1 / 2}$; we have thus $D_{T^{*}}=\left(I-T T^{*}\right)^{1 / 2}$. As customary, for any operator $T,|T|=\left(T^{*} T\right)^{1 / 2}$

The next result is a structure theorem for $J$-unitary $2 \times 2$ operator matrices. It seems to have been several times rediscovered; a good reference is [1].

Theorem 2. For any J-unitary operator matrix $g=\left(\begin{array}{ll}A & B \\ C & D\end{array}\right)$ there exist a strict contraction $T \in \mathcal{L}(\mathcal{H}, \mathcal{K})(\|T\|<1)$ and unitary operators $U_{1} \in \mathcal{L}(\mathcal{K}), U_{2} \in \mathcal{L}(\mathcal{H})$ such that

$$
\left(\begin{array}{ll}
A & B \\
C & D
\end{array}\right)=\left(\begin{array}{cc}
U_{1} & 0 \\
0 & U_{2}
\end{array}\right)\left(\begin{array}{cc}
-D_{T}^{-1} & T D_{T}^{-1} \\
-T^{*} D_{T}^{-1} & D_{T}^{-1}
\end{array}\right)
$$

As a consequence, we have, after a few computations

$$
\begin{equation*}
\Phi_{g}(X)=U_{1} \phi_{T}(X) U_{2}^{*} \tag{5}
\end{equation*}
$$

where

$$
\begin{equation*}
\phi_{T}(X)=T-D_{T} X\left(I-T^{*} X\right)^{-1} D_{T} \tag{6}
\end{equation*}
$$

Note also, for further reference, that relations (3) and (4) become

$$
\begin{align*}
& I-\phi_{T}(X)^{*} \phi_{T}(X)=D_{T}\left(I-X^{*} T\right)^{-1}\left(I-X^{*} X\right)\left(I-T^{*} X\right)^{-1} D_{T}  \tag{7}\\
& I-\phi_{T}(X) \phi_{T}(X)^{*}=D_{T^{*}}\left(I-X T^{*}\right)^{-1}\left(I-X X^{*}\right)\left(I-T X^{*}\right)^{-1} D_{T} \tag{8}
\end{align*}
$$

Actually, formula (6) makes sense in an even larger context, namely when $T$ and $X$ are two contractions subjected to the condition that $I-T^{*} X$ should be invertible. Obviously in this case $I-X T^{*}$ is also invertible, and equations (7) and (8) remain valid.

Formula (5) suggests the introduction of some ad-hoc terminology. We will call two contractions $X$ and $Y$ orthogonally equivalent, and write $X \simeq Y$, iff there exist unitary operators $U \in \mathcal{L}(\mathcal{K})$ and $V \in \mathcal{L}(\mathcal{H})$ such that $Y=U X V$. Note that this is indeed an equivalence relation on $\mathcal{B}$ (in [6] it is called a congruence). According to formula (5), the orbit of $X$ coincides with the set of all contractions orthogonally equivalent to an operator of the form $\phi_{T}(X)$, with $\|T\|<1$. The following simple lemma gives a useful criterion.

Lemma ([6]). If $|X|$ is unitary equivalent to $|Y|$ and $\left|X^{*}\right|$ is unitarily equivalent to $\left|Y^{*}\right|$, then $X \simeq Y$.

Proof. If $|X|=U^{*}|Y| U$, then the map $X h \mapsto Y U h$ is an isometric map from range $X$ to range $Y$. The second condition insures that it can be extended to a unitary $V$, since the orthogonal complements of the ranges (which coincide with ker $\left|X^{*}\right|$ and ker $\left.\left|Y^{*}\right|\right)$ have the same dimensions. Then $X=V^{*} Y U$.

## 3. A symmetry formula

There is a hidden symmetry between $T$ and $X$ which is not apparent in the definition of $\phi_{T}(X)$. This is shown by the following result from [6].

Theorem 3. If $\|T\|<1$, and $\|X\| \leq 1$, then $\phi_{T}(X) \simeq \phi_{X}(T)$.

Proof. Denote $S=D_{X}\left(I-T^{*} X\right)^{-1} D_{T}$. A direct computation shows that

$$
\begin{aligned}
\phi_{T}(X)^{*} \phi_{T}(X) & =I-S^{*} S \\
\phi_{X}(T)^{*} \phi_{X}(T) & =I-S S^{*} .
\end{aligned}
$$

Since $D_{T}$ is invertible, $\operatorname{dim} \operatorname{ker} S=\operatorname{dim} \operatorname{ker} S^{*}=\operatorname{dim} \operatorname{ker} D_{X}$. The polar decomposition of $S$ implies then that $S^{*} S$ and $S S^{*}$ are unitarily equivalent, and thus also $\phi_{T}(X)^{*} \phi_{T}(X)$ and $\phi_{X}(T)^{*} \phi_{X}(T)$. A similar argument shows that $\phi_{T}(X) \phi_{T}(X)^{*}$ and $\phi_{X}(T) \phi_{X}(T)^{*}$ are unitarily equivalent; the theorem follows by Lemma 2.

Note that in [6] an example is given of two contractions (neither one strict), for which $I-T^{*} X$ is invertible, consequently $\phi_{T}(X)$ and $\phi_{X}(T)$ both can be defined, but are not orthogonally equivalent.

In this moment we may already prove the remaining assertion in the proof of the necessity of Young's conditions. Indeed, if $D_{X}^{2}$ is compact, then $D_{X}$ is also compact. If $Y=\phi_{X}(T)=X-D_{X} \cdot T\left(I-X^{*} T\right)^{-1} D_{X}$, it follows that $Y-X$ is compact and therefore ind $X=$ ind $Y$. Furthermore, the index is invariant with respect to orthogonal equivalence; therefore formula (5) and Theorem 3 yield (2).

## 4. Shmulyan's equivalence

In [7] Shmulyan introduces the following equivalence relation on $\mathcal{B}: X$ and $Y$ are said to be S-equivalent if and only if there exist bounded operators $A$ and $B$ such that

$$
\begin{equation*}
Y-X=D_{X} \cdot A D_{X}=D_{Y} \cdot B D_{Y} \tag{9}
\end{equation*}
$$

We will write $X \cong Y$ in this case. Obviously $X \cong Y$ implies $X^{*} \cong Y^{*}$. Also, note that the operators $A$ and $B$ can be chosen such that $\operatorname{ker} D_{X} \subset \operatorname{ker} A$, $\operatorname{ker} D_{X} \subset \subset \operatorname{ker} A^{*}$, $\operatorname{ker} D_{Y} \subset \operatorname{ker} B, \operatorname{ker} D_{Y^{*}} \subset \operatorname{ker} B^{*} ;$ we will always suppose that this is achieved.

If $Y=X+D_{X} . A D_{X}$, a simple computation shows that

$$
\begin{equation*}
D_{Y}^{2}=D_{X}\left[\left(I-A^{*} X\right)\left(I-X^{*} A\right)-A^{*} A\right] D_{X} \tag{10}
\end{equation*}
$$

By Lemma 1 , range $D_{Y} \subset$ range $D_{X}$. By symmetry considerations we have that, if $X \cong Y$, then

$$
\begin{equation*}
\text { range } D_{X}=\text { range } D_{Y}, \quad \text { range } D_{X}=\text { range } D_{Y} . \tag{11}
\end{equation*}
$$

It is obvious that all strict contractions are S-equivalent, since their defect operators are invertible. Also, the following lemma is immediate.

Lemma 3. If $X \cong Y$ and $X^{\prime} \cong Y^{\prime}$, then $X \oplus X^{\prime} \cong Y \oplus Y^{\prime}$.

More interesting is the following theorem, proved in [7]:
Theorem 4. The set $\left\{\phi_{X}(T) \mid\|T\|<1\right\}$ coincides with the $S$-equivalence class of $X$.
Proof. Suppose first that $\|T\|<1$. By the definition of $Y=\phi_{X}(T)$ we have already obtained one of the equalities in (9). On the other hand, since $D_{T}$ and $D_{T}^{2}$. are invertible, relations (7) and (8) (with $T$ and $X$ interchanged) and Lemma 1(b) show that there exist bounded operators $A$ and $B$, such that $D_{X}=A D_{Y}, D_{X^{*}}=D_{Y} . B$. Then

$$
X=Y-D_{Y} \cdot B T\left(I-X^{*} T\right)^{-1} A D_{Y}
$$

and the second equality in (9) is obtained.
Suppose now that $Y \cong X$. We have first $Y=X+D_{X} . A D_{X}$ for some bounded operator $A$; note that $\operatorname{ker} D_{X} \subset \operatorname{ker} A$ and $\operatorname{ker} D_{X^{*}} \subset \operatorname{ker} A^{*}$.

Denote $W=I-X^{*} A$. By (10), we have $D_{Y}^{2}=D_{X}\left(W^{*} W-A^{*} A\right) D_{X}$. Obviously then $W^{*} W-A^{*} A$ is positive on the range of $D_{X}$. If $h \in \operatorname{ker} D_{X}, A h=0, W h=h$, while $W^{*} h=\left(I-A^{*} X\right) h$. But $X$ maps $\operatorname{ker} D_{X}$ into ker $D_{X^{*}} ;$ since $\operatorname{ker} D_{X^{*}} \subset \operatorname{ker} A^{*}, W^{*} h=h$, and thus $\left(W^{*} W-A^{*} A\right) h=h$. Thus, $W^{*} W-A^{*} A \geq 0$.

Now, by (11) and Lemma 1, it follows that $D_{Y}^{2} \geq c D_{X}^{2}$ for some $c>0$. Then

$$
\left\langle\left(W^{*} W-A^{*} A\right) D_{X} h, D_{X} h\right\rangle=\left\langle D_{Y} h, D_{Y} h\right\rangle \geq c\left\langle D_{X} h, D_{X} h\right\rangle .
$$

Since we have shown above that $\left(W^{*} W-A^{*} A\right) h=h$ for $h \in \operatorname{ker} D_{X}$, it follows that

$$
\begin{equation*}
W^{*} W-A^{*} A \geq c I . \tag{12}
\end{equation*}
$$

Define then an operator $T$ by $T W h=-A h$ on the range of $W$ and $T=0$ on the rest. Relation (12) shows that $T$ is a strict contraction. We have $-A=T W=T-T X^{*} A \quad$ and thus $\quad-\left(I-T X^{*}\right) A=T, \quad A=-\left(I-T X^{*}\right)^{-1} T$. Then $Y=X-D_{X}\left(I-T X^{*}\right)^{-1} T D_{X}=\phi_{X}(T)$, and the theorem is completely proved.

Remark. In [5] it is proved that Shmulyan's equivalence relation has an interesting geometrical interpretation as the facial relation of equivalence on the unit ball of the Hilbert space $\mathcal{H}$.

## 5. Completing the proof

The last result suggests that we introduce a notation for the equivalence relation obtained by "composing" orthogonal equivalence with Shmulyan's equivalence. Thus, $X \sim Y$ will mean that there exists $Z$ such that $X \simeq Z$ and $Z \cong Y$. It can be seen directly that $\sim$ is an equivalence relation; however, in our case this follows from the fact that the set $\{Y \mid Y \sim X\}$ is just the orbit $X$ under the action of the symplectic group. This fact (first proved in [6]) follows from formula (5) and Theorems 3 and 4.

Thus, to complete the proof of Young's theorem, we have to show that conditions (i) and (ii) imply that $Y \sim X$.

We begin with the following lemma (compare Lemma 2.17 of [5]).

Lemma 4. If (1) is satisfied, then $|X|$ is unitarily equivalent to an operator S-equivalent to $|Y|$.

Proof. Suppose $D_{Y}^{2}=P^{*} D_{X}^{2} P$. Thus, $\left\|D_{Y} \xi\right\|=\left\|D_{X} P \xi\right\|$ for any $\xi \in \mathcal{H}$; since $P$ is invertible, we may define an isometry $\Omega: \mathcal{D}_{Y} \rightarrow \mathcal{D}_{X}$, such that

$$
\begin{equation*}
D_{X}=\Omega D_{Y} P^{-1} \tag{13}
\end{equation*}
$$

But $\mathcal{H} \ominus \mathcal{D}_{Y}=\operatorname{ker} D_{Y}=\left\{\xi \mid P \xi \in \operatorname{ker} D_{X}\right\}$; using again the invertibility of $P$, it follows that $\operatorname{dim} \operatorname{ker} D_{Y}=\operatorname{dim} \operatorname{ker} D_{X}$, and thus $\Omega$ can be extended to a unitary that we will denote by the same letter, and still satisfying (13). Then

$$
\begin{aligned}
(I-|X|)^{1 / 2} & =(I+|X|)^{-1 / 2} D_{X}=(I+|X|)^{-1 / 2}\left(P^{*}\right)^{-1}(I+|Y|)^{1 / 2}(I-|Y|)^{1 / 2} \Omega^{*} \\
& =Q(I-|Y|)^{1 / 2} \Omega^{*}
\end{aligned}
$$

where $Q$ is again invertible. Thus

$$
I-|X|=\Omega(I-|Y|)^{1 / 2} Q^{*} Q(I-|Y|)^{1 / 2} \Omega^{*}
$$

If we denote $X^{\prime}=\Omega^{*}|X| \Omega$, then

$$
\begin{aligned}
|Y|-X^{\prime} & =I-X^{\prime}-(I-|Y|)=(I-|Y|)^{1 / 2}\left(Q^{*} Q-I\right)(I-|Y|)^{1 / 2} \\
& =D_{|Y|}(I+|Y|)^{-1 / 2}\left(Q^{*} Q-I\right)(I+|Y|)^{-1 / 2} D_{|Y|} .
\end{aligned}
$$

This yields one of the relations in (9). The other may be obtained by similar computations, starting, instead of (13), with

$$
D_{Y}=\Omega^{*} D_{X} P
$$

Unitary equivalence obviously implies orthogonal equivalence. But we have to obtain results about the operators themselves, not about their moduli.

Lemma 5. If

$$
\begin{equation*}
\operatorname{dim} \operatorname{ker} X-\operatorname{dim} \operatorname{ker} X^{*}=\operatorname{dim} \operatorname{ker} Y-\operatorname{dim} \operatorname{ker} Y^{*} \tag{14}
\end{equation*}
$$

then $|X| \sim|Y|$ implies $X \sim Y$.

Proof. If $\operatorname{dim} \operatorname{ker} X=\operatorname{dim} \operatorname{ker} X^{*}$, then $X=U|X|$ for a unitary operator $U$, and thus $X \simeq|X|$. By relation (14), the same is true for $Y$, and the lemma is thus proved.

If this does not happen, suppose, for instance, that $d=\operatorname{dim} \operatorname{ker} X-\operatorname{dim} \operatorname{ker} X^{*} \geq 0$. If $Z$ denotes the zero operator with range $\{0\}$ and domain a space of dimension $d$, then there exist unitary operators $U$ and $V$, such that $X$ is unitarily equivalent to $U|X| \oplus Z$, while $Y$ is unitarily equivalent to $V|Y| \oplus Z$. Lemma 3 yields then the desired result.

Lemma 5 takes care of the case when $X$ is essentially unitary. To conclude the proof of Young's theorem, suppose that, for instance, $I-X^{*} X$ is not compact. If $E$ is the spectral measure of $|T|$, it follows that there exists $r<1$, such that $E([0, r])$ is infinite dimensional. If $E^{\prime}=I-E([0, r])$, then there exists an isometry $\Omega_{1}: E^{\prime} \mathcal{H} \rightarrow \mathcal{H}$, with infinite dimensional cokernel, such that $\Omega_{1}|X|\left|E^{\prime} \mathcal{H}=X\right| E^{\prime} \mathcal{H}$. Let then $\Omega$ be any unitary on $\mathcal{H}$ which coincides with $\Omega_{1}$ on $E^{\prime} \mathcal{H}$. We claim that $\Omega|X| \cong X$. Indeed, $\Omega|X|$ is unitarily equivalent to $X \mid E^{\prime} \mathcal{H} \oplus X_{1}$, while $X$ is unitarily equivalent to $X \mid E^{\prime} \mathcal{H} \oplus X_{2}$, with $X_{1}, X_{2}$ strict contractions. Again Lemma 3, together with the fact that all strict contractions are $S$-equivalent yields the result.

Since neither $I-Y^{*} Y$ is compact, it follows that the same is true for $Y$. Thus $X \sim|X|$ and $Y \sim|Y|$; an application of Lemma 4 ends then the proof of Young's theorem.

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