# SYMMETRIC MULTIPARAMETER PROBLEMS AND DEFICIENCY INDEX THEORY

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## 1. Introduction

In this article we study the multiparameter generalization of standard deficiency index theory. A classical result in this area states that if T is a symmetric operator in a Hilbert space then the dimension of the null space of  $T^* - \lambda I$ ,  $\lambda \in \mathbb{C}$ , is constant for  $\lambda$ belonging to the upper (or lower) half-plane and further, when these two constants are equal, T admits a self-adjoint extension.

The multiparameter problem to be discussed can be posed as follows. Let  $H_1, \ldots, H_k$  be Hilbert spaces and consider operators

- (i)  $T_r: D(T_r) \subset H_r \to H_r, \ \overline{D(T_r)} = H_r, \ T_r \subset T_r^*, \ T_r \text{ closed},$
- (ii)  $V_{rs}: H_r \rightarrow H_r, V_{rs} = V_{rs}^*, V_{rs}$  bounded,  $1 \leq r, s \leq k$ .

It is customary to assume some definiteness condition on the array of operators  $[V_{rs}]$  and here we shall impose what is known as *uniform right definiteness* (URD). This can be described as follows. Let

$$V_{rs}^+ = I \otimes \cdots \otimes V_{rs} \otimes \cdots \otimes I : H \to H$$
 where  $H = \bigotimes_{r=1}^k H_r$ .

The operator  $\Delta_0$  is then defined as

$$\Delta_0 = \det[V_{rs}^+]$$

where the determinant is expanded formally. This construction is standard in multiparameter theory—see the survey paper [3], the monograph [9] or the lecture notes [7] for compendia of recent results in the area. Our definiteness condition, then, is:

URD: 
$$\Delta_0 \geq cI$$
 on H for some  $c > 0$ .

For  $\lambda = (\lambda_1, \dots, \lambda_k) \in \mathbb{C}^k$ , we put

$$W_r(\lambda) = T_r - \sum_{s=1}^k \lambda_s V_{rs}: D(T_r) \subset H_r \to H_r, \qquad 1 \le r \le k.$$

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Our first task is to investigate

$$W_r(\bar{\lambda})^* = T_r^* - \sum_{s=1}^k \lambda_s V_{rs} : D(T_r^*) \subset H_r \to H_r, \qquad 1 \le r \le k,$$

and to find regions in  $\mathbb{C}^k$  where the dimensions of the null spaces of these operators are constant. We evaluate these constants in terms of the deficiency indices of  $T_r$ . These results make up Section 2.

Secondly we shall consider the operators  $\Phi_1, \ldots, \Phi_k$  defined on the algebraic tensor product  $\bigotimes_{r=1}^k D(T_r^*)$  by

$$\Phi_1 = \det \begin{vmatrix} T_1^* & V_{12} \dots & V_{1k} \\ \vdots & & \\ T_k^* & V_{k2} \dots & V_{kk} \end{vmatrix}, \dots, \Phi_k = \det \begin{vmatrix} V_{11} & \dots & T_1^* \\ \vdots & & \\ V_{k1} & \dots & T_k^* \end{vmatrix},$$

where again the determinants are to be expanded formally. We put  $\Omega_r = \Delta_0^{-1} \Phi_r$ ,  $1 \le r \le k$ , and we draw a connection between solutions of  $(\Omega_r - \lambda_r I)x = 0$  and of  $W_r(\bar{\lambda})^* x_r = 0$ ,  $1 \le r \le k$ . Specific contributions to this multiparameter problem for systems of ordinary differential equations can be found in [5], [10].

We close with some open questions in the area.

### 2. Deficiency indices

We use the notation of the introduction. For  $1 \le r \le k$ , we define the sets

$$M_r^+ = \left\{ \lambda \in \mathbb{C}^k \, \middle| \, \sum_{s=1}^k \, (\operatorname{Im} \lambda_s) V_{rs} \gg 0 \text{ on } H_r \right\},$$
$$M_r^- = \left\{ \lambda \in \mathbb{C}^k \, \middle| \, \sum_{s=1}^k \, (\operatorname{Im} \lambda_s) V_{rs} \ll 0 \text{ on } H_r \right\}.$$

Here, for an operator A on a Hilbert space,  $A \gg 0$  means  $A \ge \alpha I$  for some  $\alpha > 0$ . Our definiteness condition URD implies that  $M_r^+ \ne \Phi$ ,  $M_r^- \ne \phi$ ,  $1 \le r \le k$ —(see [2, Theorem 2]). The following properties are easy to establish.

**Proposition 2.1.** For each  $1 \leq r \leq k$ ,

- (i)  $M_r^+, M_r^-$  are open and convex,
- (ii)  $M_r^- = -M_r^+ = (M_r^+)^* (= \{\bar{\lambda} \mid \lambda \in M_r^+).$

Our first result is:

**Theorem 2.2.** We have dim ker  $[W_r(\lambda)^*]$  is constant for  $\lambda \in M_r^+$  and for  $\lambda \in M_r^-$ .

**Proof.** Let  $\lambda^0 \in M_r^+$  and consider  $x_r \in D(T_r)$ . Then

$$(W_r(\bar{\lambda}^0)x_r, x_r) = (T_r x_r, x_r) - \left(\sum_{s=1}^k \bar{\lambda}_s^0 V_{rs} x_r, x_r\right),$$
$$(x_r, W_r(\bar{\lambda}^0)x_r) = (T_r x_r, x_r) - \left(\sum_{s=1}^k \bar{\lambda}_s^0 V_{rs} x_r, x_r\right).$$

Thus we have

$$\begin{aligned} \|W_r(\bar{\lambda}^0)x_r\| \|x_r\| &\ge |(W_r(\bar{\lambda}^0)x_r, x_r)| \\ &\ge |\operatorname{Im}(W_r(\bar{\lambda}^0)x_r, x_r)| \\ &= \left| \left( \sum_{s=1}^k (\operatorname{Im} \lambda_s^0) V_{rs} x_r, x_r \right) \right| \end{aligned}$$

 $\geq \alpha_r ||x_r||^2$ , for some  $\alpha_r > 0$ .

Hence

$$||W_r(\overline{\lambda}^0)x_r|| \ge \alpha_r ||x_r||$$
 for all  $x_r \in D(T_r)$ .

Now if  $\lambda \in M_r^+$ , then

$$W_r(\bar{\lambda}) = W_r(\bar{\lambda}^0) - \sum_{s=1}^k (\bar{\lambda}_s - \bar{\lambda}_s^0) V_{rs}$$

and if  $\lambda$  is such that  $|\lambda_s - \lambda_s^0| < \alpha_r / \sum_{s=1}^k ||V_{rs}||$ , then

$$\left\|\sum_{s=1}^{k} \left(\bar{\lambda}_{s} - \bar{\lambda}_{s}^{0}\right) V_{rs}\right\| < \alpha_{r}.$$

The following lemma from perturbation theory is an easy consequence of [4, Corollary V.1.3, p. 111].

**Lemma 2.3.** Let A be a closed operator and B a bounded operator in a Hilbert space, satisfying

$$||B|| < \alpha, ||Ax|| \ge \alpha ||x||, x \in D(A), \text{ for some } \alpha > 0.$$

Then

$$\dim \ker(A^*) = \dim \ker(A^* + B^*).$$

Returning to the proof of our theorem, we use the lemma with  $A = W_r(\bar{\lambda}^0)$  and  $B = -\sum_{s=1}^k (\bar{\lambda}_s - \bar{\lambda}_s^0) V_{rs}$ . The result is now immediate for the lemma shows dim ker  $[W_r(\bar{\lambda})^*]$ 

to be a local constant in  $M_r^+$  and the topological properties of  $M_r^+$  show this dimension to be a global constant in  $M_r^+$ . The discussion for  $M_r^-$  follows similar lines and so the theorem is established.

Our next task is to evaluate the two constants produced in the theorem above. We denote by  $n_r^+$ ,  $n_r^-$  the deficiency indices of  $T_r$ , i.e.

$$n_r^+ = \dim \ker[T_r^* - iI], n_r^- = \dim \ker[T_r^* + iI].$$

**Theorem 2.4.** If  $\lambda \in M_r^+$  (respectively,  $M_r^-$ ) then

dim ker 
$$[W_r(\bar{\lambda})^*] = n_r^+$$
 (respectively,  $n_r^-$ ).

**Proof.** We consider the (k+2) operators in  $H_r$ 

$$T_r V_{r1} \dots V_{rk} I$$

and define

$$\widehat{M}_r^+ = \left\{ \lambda \in \mathbb{C}^{k+1} \left| \sum_{s=1}^k \operatorname{Im} \lambda_s V_{rs} + (\operatorname{Im} \lambda_{k+1}) I \gg 0 \right\}.$$

We note that  $\lambda \in M_r^+$  implies  $(\lambda, 0) \in \hat{M}_r^+$  and also  $(0, \dots, 0, i) \in \hat{M}_r^+$ . Thus we take  $\lambda \in M_r^+$  and apply Theorem 2.1 to this larger system of operators to obtain

dim ker 
$$[W_r(\bar{\lambda})^*]$$
 = dim ker  $[T_r^* - iI] = n_r^+$ .

The argument for  $\lambda \in M_r^-$  is identical.

We may now appeal to the well known properties of the deficiency indices for symmetric operators—[1, Chapter 8] is a suitable reference, to claim the following

**Corollary 2.5.** (i) T, has a self-adjoint extension if, and only if, there is a point  $\lambda \in M_r^+$  such that

$$\dim \ker [W_{\ell}(\lambda)^*] = \dim \ker [W_{\ell}(\lambda)^*].$$
(2.1)

(ii) If  $T_r$  is semi-bounded and  $\lambda \in M_r^+$  then (2.1) holds.

(iii) If  $Q_r$  is bounded and symmetric on  $H_r$  then for  $\lambda \in M_r^+$ ,

$$\dim \ker [W_r(\bar{\lambda})^* + Q_r] = \dim \ker [W_r(\bar{\lambda})] = n_r^+.$$

(iv) Suppose  $T_r$  is bounded below and  $\lambda = \sigma + it \in M_r^+$ ,  $\sigma, t \in \mathbb{R}^k$ . Then there is a real

number  $\theta_0 > 0$  so that if  $\theta > \theta_0$  a self-adjoint extension  $\tilde{T}_r(\theta)$  of  $T_r$  can be found so that

$$\left(\tilde{T}_r(\theta) + \sum_{s=1}^k \theta t_s V_{rs}\right) x_r = 0 \quad \text{for some } x_r \neq 0.$$

**Proof.** Claims (i), (ii), (iii) are easy. For (iv) we note that  $\sum_{s=1}^{k} t_s V_{rs} \gg 0$  so we select  $\theta_0$  large enough to satisfy

$$T_r + \sum_{s=1}^k \theta_0 t_s V_{rs} \gg 0.$$

The result now follows from [1, Theorem 3, p. 365].

#### 3. Decoupling the spectral parameters

In this section we study the process of decoupling the spectral parameters  $\lambda_1, \ldots, \lambda_k$  in a system of simultaneous equations  $W_r(\bar{\lambda})^* x_r = 0$ ,  $1 \le r \le k$ . This is a commonly used idea in (self-adjoint) multiparameter spectral theory and, from some points of view, forms the basis for it. We begin with some preliminaries.

Note that the operators  $I \bigotimes_a \cdots \bigotimes_a T_r \bigotimes_a \cdots \bigotimes_a I$  and  $I \bigotimes_a \cdots \bigotimes_a T_r^* \bigotimes_a \cdots \bigotimes_a I$ are closeable since they are both restrictions of the closed operator  $(I \bigotimes_a \cdots \bigotimes_a T_r \bigotimes_a \cdots \bigotimes_a I)^*$ . We shall use  $T_r^+$  and  $T_r^{*+}$  to denote these closures. The operators  $W_r(\lambda)^+$  and  $W_r(\lambda)^{*+}$  are defined in like fashion.

**Lemma 3.1.** Let  $\lambda \in \mathbb{C}^k$ . Then

$$\bigotimes_{r=1}^{k} \ker [W_{r}(\bar{\lambda})^{*}] = \bigcap_{r=1}^{k} \ker [W_{r}(\bar{\lambda})^{*+}].$$

**Proof.** We should point out that the tensor product on the left hand side above is the Hilbert tensor product, i.e. it is the closure in H of the algebraic tensor product of the closed subspaces ker $[W_{\ell}(\bar{\lambda})^*]$  and, as such, it is a closed subspace of H. The right hand side is also a closed subspace of H as it is the intersection of the kernels of the closed operators  $W_{\ell}(\bar{\lambda})^{*+}$ .

It is easy to see

$$\bigotimes_{r=1^a}^k \ker [W_r(\bar{\lambda})^*] \subset \bigcap_{r=1}^k \ker [W_r(\bar{\lambda})^{*+}]$$

and so we take closures to obtain

$$\bigotimes_{r=1}^{k} \ker [W_{r}(\bar{\lambda})^{*}] \subset \bigcap_{r=1}^{k} \ker [W_{r}(\bar{\lambda})^{*+}].$$

Now let ker  $[W_r(\bar{\lambda})^*]$  have an orthonormal basis  $\{e_r^1, e_r^2, \ldots\}$  and its orthocomplement

have an orthonormal basis  $\{f_r^1, f_r^2, \ldots\}$ . We show that any tensor of the form  $q = x_1 \otimes \cdots \otimes f_r^j \otimes \cdots \otimes x_k$  belongs to  $(\ker[W_r(\lambda)^{*+}])^{\perp}$ . To see this, first note that  $f_r^j \in (\ker[W_r(\lambda)^*])^{\perp} = \overline{R[W_r(\lambda)]}$ , and so  $f_r^j = \lim_{n \to \infty} W_r(\lambda) y_r^n$ ,  $y_r^n \in D(T_r)$ . From our discussion and definitions at the start of this section it also follows that

$$W_r(\bar{\lambda})^{*+} \subset (W_r(\bar{\lambda})^+)^*.$$

Hence if  $z \in \ker[W_r(\bar{\lambda})^{*+}]$ , we have

$$(q, z) = \lim_{n \to \infty} (x_1 \otimes \cdots \otimes W_r(\bar{\lambda}) y_r^n \otimes \cdots \otimes x_k, z)$$
$$= \lim_{n \to \infty} (x_1 \otimes \cdots \otimes y_r^n \otimes \cdots \otimes x_k, W_r(\bar{\lambda})^{*+} z)$$
$$= 0,$$

establishing our claim.

Vectors of the form  $x_1 \otimes \cdots \otimes x_k$  where  $x_r \in \{e_r^1, e_r^2, \dots, f_r^1, f_r^2, \dots\}$  form an orthonormal basis for H. Those of the type in which  $x_r \in \{e_r^1, e_r^2, \dots\}$ ,  $1 \le r \le k$ , lie within  $\bigcap_{r=1}^k \ker[W_r(\bar{\lambda})^{*+}]$  by our opening remarks and the remainder lie within  $(\bigcap_{r=1}^k \ker[W_r(\bar{\lambda})^{*+}])^{\perp}$  by the argument above. From this observation, the lemma follows immediately. A more general version of this result can be found in [6].

We now proceed with the decoupling process by first noting that with  $\Phi_1, \ldots, \Phi_k$  as defined in the introduction and  $\Omega_s = \Delta_0^{-1} \Phi_s$ ,  $1 \le s \le k$ , we have

$$T_r^{*+} - \sum_{s=1}^k V_{rs}^+ \Omega_s = 0$$
 on  $\bigotimes_{r=1^a}^k D(T_r^*)$ .

This follows readily from [5, Theorem 2] which states that when  $x \in \bigotimes_{r=1}^{k} D(T_{r}^{*})$ , the equations

$$T_{r}^{*+}x - \sum_{s=1}^{k} V_{rs}^{+}g^{s} = 0, \qquad 1 \le r \le k,$$
(3.1)

can be solved uniquely for  $g^1, \ldots, g^k$  with  $g^s = \Omega_s x$ . In fact we can use (3.1) to extend the domain of definition of each  $\Omega_s$  from  $\bigotimes_{r=1^a}^k D(T_r^*)$  to  $\bigcap_{r=1}^k D(T_r^{*+})$ . We continue to write  $\Omega_s$  for this extension.

**Theorem 3.2.** Let  $\lambda \in \mathbb{C}^k$ . Then

$$\bigotimes_{r=1}^{k} \ker [W_{r}(\bar{\lambda})^{*}] = \bigcap_{r=1}^{k} \ker [W_{r}(\bar{\lambda})^{*+}] = \bigcap_{r=1}^{k} \ker [\Omega_{r} - \lambda_{r}I].$$

**Proof.** In view of the previous lemma, only the second equality requires discussion. If  $x \in \bigcap_{r=1}^{k} D(T_r^{*+})$  and

$$\left(T_r^{*+}-\sum_{s=1}^k \lambda_s V_{rs}^+\right)x=0, \ 1\leq r\leq k,$$

then it follows that

$$\sum_{s=1}^{k} V_{rs}^{+}(\Omega_{s}-\lambda_{s}I)x=0, \ 1\leq r\leq k.$$

We again use [5, Theorem 2] to deduce that  $x \in \bigcap_{r=1}^{k} \ker[\Omega_s - \lambda_s I]$ . The reverse argument is similar.

**Definition 3.3.** For  $\lambda \in \mathbb{C}^k$  we define the *deficiency index for the system*  $[T_r V_{rs}]_{r,s=1}^k$  at  $\lambda$  to be

$$N(\lambda) = \prod_{r=1}^{k} \dim \ker [W_{r}(\lambda)^{*}]$$
$$= \dim \bigcap_{r=1}^{k} \ker [\Omega_{r} - \lambda_{r} I].$$

For  $\varepsilon = (\varepsilon_1, \dots, \varepsilon_k)$ ,  $\varepsilon_r = \pm 1$ ,  $1 \le r \le k$ , we put

$$M_{e} = \left\{ \lambda \in \mathbb{C}^{k} \left| \varepsilon_{r} \sum_{s=1}^{k} (\operatorname{Im} \lambda_{s}) V_{rs} \gg 0, \ 1 \leq r \leq k \right\}.$$

From [2, Theorem 2] we know that  $M_e \neq \Phi$ . In fact

$$M_e = \bigcap_{r=1}^k M_r^{e_r}$$

**Corollary 3.4.** For each  $\varepsilon$ ,  $N(\lambda)$  is constant in  $M_{\varepsilon}$ . In fact if  $\lambda \in M_{\varepsilon}$ ,

$$N(\lambda) = \prod_{r=1}^{k} n_r^{e_r}.$$

**Proof.** This is an immediate consequence of Theorem 2.4.

#### 4. Open questions

When the operators  $T_r$ ,  $1 \le r \le k$ , are self-adjoint it is known from standard multiparameter theory that  $\Omega_1, \ldots, \Omega_k$  are self-adjoint in H equipped with the inner product  $[x, y] = (\Delta_0 x, y)$  and are pairwise commuting in the sense that their spectral measures commute. If each  $T_r$  is symmetric and has equal deficiency indices then it follows that  $\Omega_1, \ldots, \Omega_k$  have restrictions which are  $[\cdot, \cdot]$ -self-adjoint and pairwise commutative. Is it possible for  $\Omega_1, \ldots, \Omega_k$  to have such restrictions when the operators  $T_r$  do not have selfadjoint extension? If  $\Gamma_1, \ldots, \Gamma_k$  are such restrictions of  $\Omega_1, \ldots, \Omega_k$ , is it possible to characterize the corresponding extensions to  $T_1, \ldots, T_k$ ?

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