

ICE MOVEMENT AND TEMPERATURE DISTRIBUTION IN GLACIERS AND ICE SHEETS*

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ABSTRACT. The distribution of temperature throughout an ice sheet has been considered, taking into account the influence of ice movement as well as other items previously considered, such as conduction, the geothermal outflow of heat and heat generated by ice movement. By making certain simplifying assumptions, a quantitative method of estimating the temperature distribution near the centre of an ice sheet has been put forward.

It is shown that even a small mean annual accumulation will have considerable effect on the temperature distribution in a large ice sheet. For a moderate rate of accumulation a substantial fraction of the total thickness of ice at the centre of a large ice sheet may be isothermal at the prevailing surface ice temperature. Under these conditions at some distance from the centre, the change in the surface ice temperature with elevation may produce a temperature gradient opposite to normal, that is the temperature falls with increasing depth below the surface, due to the outward movement of the ice. Observed temperature gradients on ice sheets fit the proposed hypotheses roughly, but it appears that climatic change should also be taken into account.

It is suggested that a rise from temperatures below melting point at the base of ice sheets may provide an explanation of the occasional catastrophic advances of certain glaciers.

ZUSAMMENFASSUNG. Es wird die Temperaturverteilung durch eine ganze Eiskappe hindurch erwogen. Der Einfluss der Eisbewegung sowohl wie anderer schon früher erwogener Punkte wurde dabei in Betracht gezogen, wie z. B. Wärmeleitung, geothermischer Wärmeauslauf und durch Eisbewegung erzeugte Wärme. Es wurden gewisse vereinfachende Voraussetzungen getroffen, die eine quantitative Methode zur Bestimmung der Temperaturverteilung nahe dem Zentrum einer Eiskappe ermöglichten.

Es wird gezeigt, dass selbst eine im Jahresdurchschnitt kleine Akkumulation eine beträchtliche Einwirkung auf die Temperaturverteilung in einer grossen Eiskappe hat. Bei mässiger Akkumulation kann ein beträchtlicher Teil der gesamten Eisdichte im Zentrum einer grossen Eiskappe bei der vorherrschenden Oberflächen-Eistemperatur isothermisch sein. Unter diesen Umständen kann die Änderung der Oberflächen-Eistemperatur mit Ansteigen in einiger Entfernung vom Zentrum ein dem Normalen entgegengesetztes Temperaturgefälle produzieren, das heisst die Temperatur fällt mit zunehmender Tiefe unter der Oberfläche auf Grund der nach aussen gerichteten Bewegung des Eises. Beobachtete Temperaturgefälle in Eisflächen stimmen ungefähr mit der vorgeschlagenen Hypothese überein, aber es scheint, dass klimatische Wechsel auch in Betracht gezogen werden sollten.

Es wird erwogen, dass Erhöhung von Temperaturen unter dem Schmelzpunkt am Grunde von Eiskappen eine Erklärung für das gelegentlich katastrophale Vorrücken gewisser Gletscher aufweisen könnte.

INTRODUCTION

Previous discussion of this subject by Brockamp^{1,2} has considered the effect of the outflow of heat from the Earth, and of internal friction on the temperature of the ice sheet. The effect of accumulation in modifying the winter "cold wave" in a snowfield has been considered in detail by Benfield.³ A large accumulation was shown to affect the form of the cold wave appreciably, but the effect with a small accumulation was not significant.

The object of this paper is to show for a polar ice sheet (that is one where the layers of ice or firn near the surface remain below freezing point throughout the year) that even a small accumulation will produce a considerable effect on the temperature gradient in the ice. A quantitative expression for the variation of this gradient with accumulation, ice thickness and depth has been derived which should be applicable near the centre of a stable ice sheet. Under certain circumstances, which probably apply in the southern half of the Greenland Ice Cap, it is shown that a negative temperature gradient (*i.e.* temperature decreasing downwards) may be formed as a stable feature.

THEORY

It is necessary to make certain assumptions regarding an ice sheet as a starting point from which to tackle this problem. We consider only the case of a stable ice sheet, that is one which is of constant size, and in which the temperature distribution does not vary with time. We will regard the ice sheet as having a two-dimensional form of a type outlined by Nye.⁴ Fig. 1 (p. 525) shows a cross-section of such an ice sheet, and serves to explain the various symbols.

* Based on a paper read at the tenth General Assembly of the I.U.G.G., Rome, 1954 and expanded.

At the centre of the ice sheet H is the total thickness and h the height above the base of the ice sheet, H' and h' are the corresponding quantities at distance r from the centre, R being the half width. Other symbols used are:

θ =temperature

t =time

q =rate of outflow of geothermal heat per unit area

K =thermal conductivity of ice (5.3×10^{-3} cal./cm./sec./degree C.)

k =thermal diffusivity of ice (1.18×10^{-2} cm.²/sec.)

$\dot{A} = \frac{dA}{dt}$ = rate of accumulation of ice on upper surface of ice sheet. This is assumed constant over the ice sheet.

The upper surface of the ice sheet is assumed to fit the equation⁴

$$H' = \sqrt{2h_0(R-r)} \quad \dots \dots \dots (1)$$

where h_0 is a constant.

The rate of outward movement of a column distant r from the centre is

$$\frac{dr}{dt} = \frac{\dot{A}r}{H'} \quad \dots \dots \dots (2)$$

We also assume that any vertical column remains vertical in spite of movement over the subglacial floor. We consider that the effect of accumulation, or of decreasing thickness due to outward movement, increases the cross-section of a vertical column uniformly at all depths. Thus at the centre of the ice sheet, if we neglect effects of compression, which are only significant close to the surface, the rate of vertical movement of the ice at any level h due to surface accumulation is given by

$$-\frac{dh}{dt} = \frac{h}{H} \dot{A} \quad \dots \dots \dots (3)$$

The above mechanism approximately fits that proposed by Nye⁴ and some such simplified picture is necessary to facilitate analysis. Alternative mechanisms can be used which require modification of (3), as is done in the case of ice shelves later.

TEMPERATURE GRADIENT AT THE CENTRE OF AN ICE SHEET

Consider a vertical prism near the centre of an ice sheet. The general equation of heat flow is

$$\frac{\partial \theta}{\partial t} = k \nabla^2 \theta \quad \dots \dots \dots (4)$$

We will consider the horizontal temperature gradient in the ice sheet as negligible, as in the case of central Greenland, its mean value cannot exceed 1° C. per 10 km. We can then simplify (4) to

$$\frac{d\theta}{dh} \cdot \frac{dh}{dt} = k \frac{d^2 \theta}{dh^2} \quad \dots \dots \dots (5)$$

By substituting (3) in (5) and integrating we can obtain a solution for the temperature gradient at any level as

$$\frac{d\theta}{dh} = \left(\frac{d\theta}{dh} \right)_{\text{bottom}} e^{-\frac{\dot{A}}{2Hk} h^2} \quad \dots \dots \dots (6)$$

e being the exponential function and $\left(\frac{d\theta}{dh} \right)_{\text{bottom}}$ the temperature gradient at the base of the ice sheet. This can be estimated approximately as it is also given by

$$\left(\frac{d\theta}{dh} \right)_{\text{bottom}} = \frac{q}{K} \quad \dots \dots \dots (7)$$

Tables (1) and (2) show how these temperature gradients vary near the centre of two ice sheets, of thickness 3000 m. and 400 m. respectively.

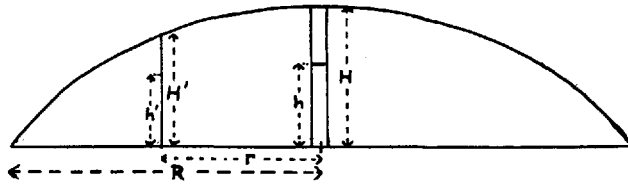


Fig. 1. Symbols used in discussion of an ice sheet

TABLES I AND II

Fraction of the upward flow of heat through the base at the centre of an ice sheet passing through a horizontal surface at a height *h* above the base.

TABLE I. CENTRAL THICKNESS OF ICE SHEET (*H*) 3000 m.

| Annual accumulation "A" cm./year | <i>h</i> — 200 m. | <i>h</i> — 500 m. | <i>h</i> — 1000 m. | <i>h</i> — 1500 m. | <i>h</i> — 2000 m. | <i>h</i> — 2500 m. | <i>h</i> — 3000 m. |
|-------------------------------------|----------------------|----------------------|-----------------------|-----------------------|-----------------------|-----------------------|-----------------------|
| 0 | 1.00 | 1.00 | 1.00 | 1.00 | 1.00 | 1.00 | 1.00 |
| 2 | .996 | .98 | .91 | .82 | .70 | .57 | .45 |
| 4 | .99 | .96 | .84 | .67 | .49 | .33 | .20 |
| 8 | .99 | .92 | .70 | .45 | .24 | .11 | .04 |
| 16 | .97 | .84 | .49 | .20 | .06 | .01 | .00 |
| 32 | .94 | .70 | .24 | .04 | .00 | .00 | .00 |
| 64 | .89 | .49 | .06 | .00 | .00 | .00 | .00 |
| 128 | .80 | .24 | .00 | .00 | .00 | .00 | .00 |
| 256 | .63 | .06 | .00 | .00 | .00 | .00 | .00 |

TABLE II. CENTRAL THICKNESS OF ICE SHEET (*M*) 400 m.

| Annual accumulation "A" cm./year | <i>h</i> — 50 m. | <i>h</i> — 100 m. | <i>h</i> — 150 m. | <i>h</i> — 200 m. | <i>h</i> — 300 m. | <i>h</i> — 400 m. |
|-------------------------------------|---------------------|----------------------|----------------------|----------------------|----------------------|----------------------|
| 0 | 1.00 | 1.00 | 1.00 | 1.00 | 1.00 | 1.00 |
| 2 | .999 | .99 | .99 | .97 | .94 | .90 |
| 4 | .997 | .99 | .97 | .95 | .89 | .81 |
| 8 | .99 | .97 | .94 | .90 | .78 | .55 |
| 16 | .99 | .95 | .89 | .81 | .62 | .47 |
| 32 | .97 | .90 | .79 | .65 | .38 | .18 |
| 64 | .95 | .81 | .62 | .42 | .15 | .03 |
| 128 | .90 | .65 | .38 | .18 | .02 | .00 |
| 256 | .81 | .42 | .15 | .03 | .00 | .00 |
| 512 | .65 | .18 | .02 | .00 | .00 | .00 |

The rate of outflow of geothermal heat (*q*) is roughly uniform over the world at about 38 cal./cm.²/year⁶ which gives a gradient of 1° C. per 44 m. at the base of the ice sheet. By integration of (6) we find that the temperature difference between any height *h* and the surface of the ice sheet is

$$\theta_H - \theta_h = \left(\frac{d\theta}{dh}\right)_{\text{bottom}} \sqrt{\frac{2Hk}{A}} \left[\operatorname{erf} \left(\sqrt{\frac{A}{2Hk}} \cdot h \right) \right]_h \dots \dots (8)$$

$$\text{where } \operatorname{erf} x = \int_0^x e^{-y^2} dy$$

Results corresponding to a number of different values of annual accumulation are shown graphically in Figs. 2 and 3 (p. 527), corresponding to values of *H* of 3000 m. and 400 m. respectively, taking *q* as 38 cal./cm.²/year.

Equation (8) gives an estimate of ice temperature only if the calculated value of the bottom temperature ($\theta_{b,e}$) is equal to or below the melting point of the ice at the appropriate pressure

DOWNWARD HEAT FLOW FROM SURFACE

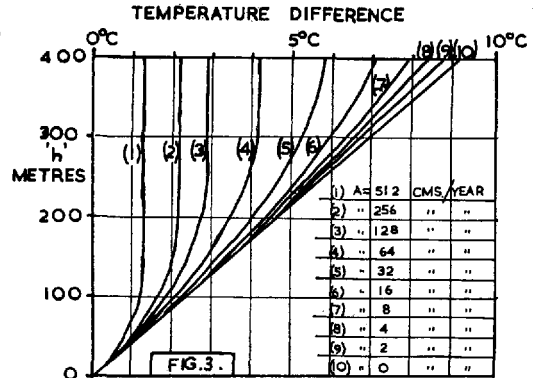
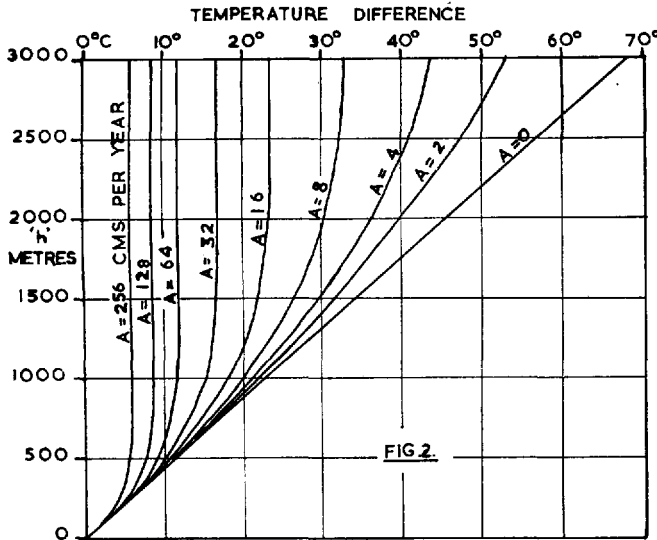
In addition to the factors mentioned so far, the temperature of the surface layers of ice and firn varies with its height above sea level. Brockamp¹ gives 0.9° C. per 100 m. elevation as the change in ice temperature due to height which agrees approximately with the Antarctic results of Schytt¹⁰. Due to outward and downward movement of ice away from the centre of the ice sheet, the mean temperature at which each annual layer of firn accumulates will be slightly warmer than that of the layer of firn of the previous year. In the simplest case, where the velocity of ice movement is V , the surface slope α , the net rate of accumulation \dot{A} , and λ equals 0.9° C. per 100 m., the negative temperature gradient produced by this effect, if we neglect conduction of heat, will be

$$\frac{d\theta}{dh'} = \frac{V \sin \alpha}{\dot{A}} \cdot \lambda \approx \frac{\alpha V \lambda}{\dot{A}} \dots \dots \dots (11)$$

provided α is small.

We can estimate the magnitude of this effect for our simplified ice sheet as follows :
Differentiation of (1) gives

$$\alpha \approx \frac{dH'}{dr} = \frac{-h_0}{\sqrt{2h_0(R-r)}} \dots \dots \dots (12)$$



Figs. 2 and 3. Temperature difference between the base and other levels near the centre of an ice sheet, Fig. 2 central thickness 3000 m., Fig. 3 central thickness 400 m.

Combining (1) and (2) we get,

$$V = \frac{dr}{dt} = \frac{\dot{A}r}{\sqrt{2h_0(R-r)}} \dots \dots \dots (13)$$

from (11), (12) and (13)

$$\frac{d\theta}{dh'} = \frac{r}{2(R-r)} \cdot \lambda \dots \dots \dots (14)$$

This simple expression indicates that one should get an increase in the negative temperature gradient close to the surface as one moves out from the centre of an ice sheet, provided the upward flow of heat from the base is negligible. If the upward flow of heat is appreciable, one would anticipate a decrease in the negative gradient shown by (11) or (14), a positive gradient being formed when the upward flow of heat outweighs this tendency to form a negative gradient. While (14)

indicates the order of magnitude of such a negative gradient close to the surface, it does not take account of thermal conduction. This has been done however for one particular case later in this paper.

COMPARISON WITH OBSERVED TEMPERATURES

Bottom Temperatures

Near the centre of the Greenland Ice Cap we can take $H=3000$ m., $A=30$ cm. per year and the surface temperature $\theta_H=20^\circ$ C. Use of Fig. 2 or equation (8) then gives a bottom temperature of about -12° C. Both Haefeli¹¹ and Holtzscherer¹² conclude that the base of at least part of the ice sheet is below freezing point. Holtzscherer estimates the basal temperature near the centre at around -10° C. from seismic shooting information. Away from the centre Haefeli believes the bottom temperature may rise to freezing point, but Holtzscherer's refraction shooting results indicate a layer which may consist of frozen moraine or similar material at Camp IV and VI near the edge of the ice sheet.

Qualitatively this paper agrees with the picture presented by Haefeli, as it suggests that the basal temperature will increase as we move out from the centre of the ice sheet. However, Holtzscherer's observations suggest that this increase is not sufficient to raise the basal temperature to melting point.

In the case of a small ice hill¹³ near the base of the Norwegian British Swedish Antarctic Expedition 1949-52, with an annual accumulation around 40 cm. of ice, one would expect a basal temperature around -14° C. as the data in Fig. 3 fit this case. Calculations of the shear stress at the base of the ice hill suggest that the ice there is a few degrees below melting point.¹³

Temperature Gradients near the Surface

Measurements of temperature gradients close to the surface are available. Ahlmann¹⁵ has summarized earlier measurements, and works presented by Bauer¹⁴ and Heuberger,¹⁷ Schytt,¹⁰ and Fisher^{7,8,9} give more recent data. Temperature gradients obtained by the German Greenland expeditions¹⁵ close to the edge of the ice sheet at Borg in 1912-13 and West Station in 1931 were found to be four or five times the anticipated gradient due to the geothermal supply alone. This appears to fit the concept of a heat supply due to the rapid shear in the lowest layers of ice, but it will also be due in part at least to the increase of temperature gradient caused by the relatively rapid thinning of the ice during it's outward motion. The negative gradient first suggested by the German results at Eismitte has now been confirmed and amplified by the results of Expéditions Polaires Françaises^{14,17}. Comparison of their results with those given by equation (14) are shown below :

| Location | r/R | Calculated -ve gradient | Observed -ve gradient | |
|------------------------|-------|-------------------------|-----------------------|---------------------|
| | | | (1) 20-40 m. depth | (2) 20-120 m. depth |
| Station Centrale | 1/4 | 0.15° C. per 100 m. | 1.5° C. per 100 m. | 0.8° C. per 100 m. |
| Camp VI | 6/7 | 2.7° C. per 100 m. | 5° C. per 100 m. | 3.4° C. per 100 m. |

Qualitatively the increase in negative gradient with distance from the centre of the ice sheet is shown, but the quantitative agreement is not good.

In the case of Camp VI, a more detailed analysis than that given by (14) has been made, which also takes into account the term due to thermal conduction.

We make the simplifying assumption, that as the ice moves outwards over a limited region the temperature gradient near the surface remains constant, although the temperature as a whole changes in proportion to the change of surface elevation, then we can express the heat balance by the equation as follows :

$$\frac{d^2\theta}{dh'^2} - \frac{A}{kH'} \cdot k' \frac{d\theta}{dh'} = \frac{\alpha V\lambda}{k} \dots \dots \dots (15)$$

The coefficients in (15) can be determined for a given temperature depth curve by fitting the polynomial

$$\theta_h' = \theta_{H'} + xh' + yh'^2 + zh'^3 \dots \dots \dots (16)$$

to a series of points on the observed curve. From these points the coefficients in (15) can be determined. This has been done for the smoothed temperature-depth curve published by Bauer¹⁴ for Camp VI. Two sets of results were obtained by including terms of the third degree in (16) as follows :

| | (1) | (2) | Mean |
|------------------------------|------------------------|------------------------|------------------------|
| $\frac{\dot{A}}{kH}$ | 5.05×10^{-10} | 5.63×10^{-10} | 5.34×10^{-10} |
| $\frac{\alpha V \lambda}{k}$ | 2.58×10^{-8} | 2.86×10^{-8} | 2.72×10^{-8} |

using C.G.S. units.

Taking $k = 1.18 \times 10^{-2}$ cm.²/sec. as the value for ice,
and $H = 1370$ m.—as measured by seismic shooting,
we get $\dot{A} = 27.4$ cm./year.

This agrees reasonably well with the observed accumulation.

If we take $\alpha = 0.01$ radians (as indicated approximately on Holtzscherer's profile)
and $\lambda = 9 \times 10^{-5}$ ° C. per cm.
we get $V = 113$ m. per year.

For comparison equation (2) gives $V = 100$ m. per year if we take $\dot{A} = 30$ cm. per year and $r = 450$ km. This figure is probably too high, as at Camp IV the observed surface movement¹⁴ was 30 m. per year compared with a calculated movement of 196 m. per year.¹⁴ However, in view of the presence of streams of faster flowing ice in the ice sheet,¹⁴ we cannot base conclusions on the surface velocity at Camp VI on observations at Camp IV. A more accurate calculation should also allow for the effect of climatic change on the shape of the temperature-depth curve.

No attempt has been made to use observations at Station Centrale to determine probable accumulation and velocity, as the form of the temperature-depth curve there would not give consistent results over different sections of the curve. Instead it appears that climatic changes have more effect on the shape of the curve than the slight effect given by (14). The upper portion of this curve would certainly agree with the hypothesis of a warming of this region in recent decades, and could perhaps be used to make a quantitative estimate. In addition, the slight temperature-depth change shown on this curve agrees within limits such as may be imposed by the above factors with the appropriate temperature-depth curve of Fig. 2.

Fisher^{7,8,9} has provided figures for the temperature-depth variation of the cold glacier on the north face of Monte Rosa. The temperatures observed in his 1952-53 tunnel fit (11) qualitatively at least, but the picture for the 1954 tunnel is too complicated to agree with this analysis.

The curves for the U.S. West Base in 1940,¹⁵ and for Maudheim in 1951¹⁰ are those for floating ice shelves. The problem is somewhat different, and instead of assuming a constant geothermal supply of heat, we can assume a constant temperature of the base of the floating ice shelf equal to the freezing point of sea water. In Fig. 4 (p. 531) two curves have been calculated for Maudheim. For curve (1), it is assumed no bottom melting of the ice shelf occurred, and the ice shelf was assumed to spread out so that equation (1) applied. For curve (2) the bottom melting of the ice has been assumed equal to the surface accumulation so that equation (1) now becomes

$$\frac{dh}{dt} = \dot{A}$$

Neither curve (1) nor curve (2) fits the observed value at Maudheim, but (2) seems to fit the

observed values at the Bay of Whales (U.S. West Base) if we assume that accumulation and ice thicknesses are similar, which is very approximately true.

It appears the ice movement and accumulation at Maudheim are not entirely in a state of equilibrium as defined by these hypotheses. Other factors which we should also consider are changing surface temperatures of ice shelves due to their northward motion and increasing proximity to the sea, or due to climatic change. Also a considerable time would be needed for inland ice pushed out into an ice shelf to attain thermal equilibrium under the new conditions.

As regard applying figures obtained in Greenland to our basic assumption of a stable ice sheet, it is worth noting that Bauer¹⁴ estimates that the Greenland ice sheet is losing ice equivalent to 100 Km.³ per year in a total annual budget of around 500 Km.³. The accuracy of the observations is not sufficient to definitely exclude the possibility that the ice sheet is stable, as was earlier proposed by Loewe.¹⁸ Gravitational observations over the Greenland Ice Cap made by Expéditions Polaires Françaises, and the British North Greenland Expedition should help to show whether or not the ice sheet is in isostatic equilibrium, and consequently whether the ice sheet has been relatively stable over a long period.

TEMPERATURE DISTRIBUTION IN THE ICE CAP OF CENTRAL GREENLAND

To estimate the temperature distribution throughout an entire ice sheet involves studying the integrated sum of all the factors previously mentioned on a column of ice during its journey from the centre to the boundary of the ice sheet. The picture is a complicated one, but Fig. 5 (p. 531) provides a sketch which attempts to combine the observations near central Greenland with the arguments in this paper.

CONSEQUENCES OF PROPOSED HYPOTHESES

(a) *Basal temperatures of ice sheets*

At the centre of the Greenland Ice Cap, if the accumulation dropped to 10 cm. of ice per year, the basal temperature for equilibrium would rise to 0° C., and any further drop would result in melting of the base of the ice at the centre of the ice sheet. Our estimate of the magnitude of the geothermal heat flow here is of course very approximate. However, the decrease in accumulation believed to occur as one moves further north over the ice sheet may produce a rise in basal ice temperature in spite of lower surface temperatures. This would facilitate the flow of ice and may help to explain the lower height of the ice sheet in northern Greenland on the basis of the work of Nye⁴ and Glen¹⁹ even though the base of the ice sheet remains at around sea-level according to Holtzscherer,¹² contrary to Nye's predictions.⁵ An alternative explanation for the difference between the north and south of the Greenland Ice Cap could be that the two portions are in different phases of long-period cyclic variations of thickness as outlined below in section (b). The gravitational observations previously mentioned should help settle this point.

(b) *Catastrophic advances of glaciers*

If circumstances outlined in this paper produce basal ice temperatures below melting point, a greater shear stress will be required to produce ice flow. This means that greater surface slopes and ice thickness will be formed. Now if the basal ice temperature in such an ice sheet or glacier is raised by some phenomenon—the most likely one being by an increase in ice movement over the bed—the shear stress due to the greater ice thickness formerly accumulated will be in excess of that required to make the ice flow at this higher temperature. As this is an unstable state the flow may continue to increase until the basal temperature reaches melting point, leading finally to a relatively rapid advance of the glacier front. With basal temperatures now at melting point the flow and ice thickness may remain similar to that of a temperate glacier until conduction and the downward "advection" of cold ice again start to lower the temperature of the ice rock interface with consequent thickening of the glacier. Such a chain of events may well be cyclic with a very long

period. The mechanism may be the explanation of such catastrophic glacier advances as those of Bråselvbreen in Svalbard and of the Kutiäh Glacier²⁰ in the Karakoram recently.

Such a mechanism may not be uncommon in polar glaciers. Ahlmann¹⁶ in his Bowman Memorial lecture on *Glacier Variations and Climatic Fluctuations* says the following:

"It has already been mentioned that in Spitsbergen conditions are in many respects peculiar. A summary prepared for me by Liestöl, points out that glacier snouts that have receded over a hundred years or so, suddenly begin a rapid advance, sometimes reaching as far as the outermost, recent terminal moraines. Every glacier has its own period of variation, which is, however, dependent on the more pronounced climatic fluctuations and changes."

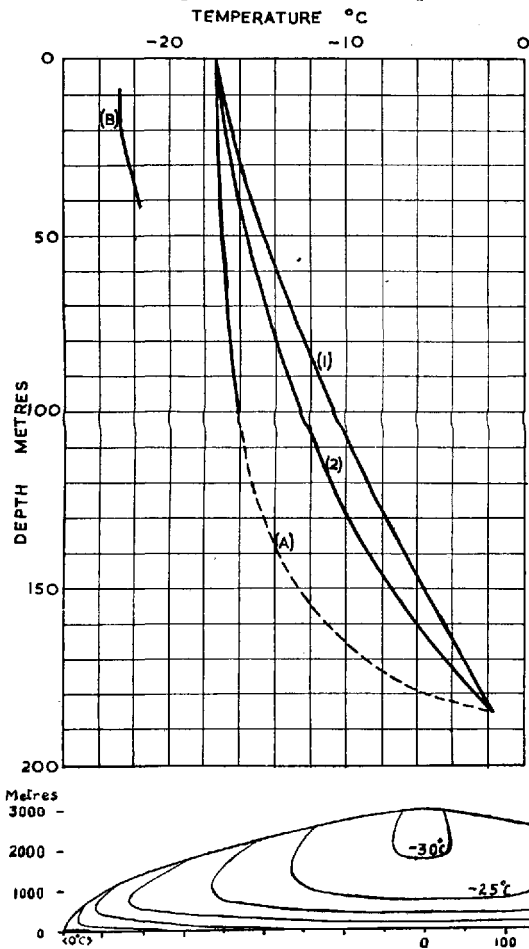


Fig. 4. Temperature-depth curves for ice shelves.

(A) Maudheim, 1951, observed.

(B) Bay of Whales, 1940, observed.

(1) Maudheim, calculated for no bottom melting.

(2) Maudheim, calculated for bottom melting equal to accumulation

Fig. 5. Sketch of the temperature distribution in a simplified ice sheet similar to central Greenland

Both Matthes and Liestöl are quoted by Ahlmann as putting forward similar explanations for such surges. He quotes Liestöl as follows:

"Over a long period of time a surplus may be built up in the accumulation area, while the lower part of the glacier concurrently wastes away until the frictional limit is exceeded, and the ice slides forward. Such a condition, he says, is most evident in polar and sub-polar glaciers and, it may also be added, is especially applicable to glaciers whose accumulation areas are situated on plateaus high above the valleys constituting their main ablation areas."

Thus the existence of such long-period surges has been observed and an explanation has been outlined. The importance of the temperature variations of the bed of the glacier has not however

been stressed, but it appears to provide the essential key to the mechanism. It is admittedly difficult to prove this by direct measurements of basal ice temperatures, but observations may be possible on the base of glaciers which have recently surged forward. It would of interest to know whether such surges also appear as intermittent ice streams, large or small, flowing out of the main Antarctic and Greenland ice sheets.

The exceptional advance of the Kutiah glacier, reported by Professor A. Desio²⁰ in this *Journal* appears to fit into the pattern of Spitsbergen glaciers as outlined by Liestöl. One concludes that the glaciers feeding it must have an interconnected accumulation zone of sufficient elevation, say 6000 m. or more, to ensure the supply of firn being appreciably below freezing point. Also one would presume that this high accumulation zone is sufficient in area, say three square kilometres or more, to provide the large volume of ice necessary for the reported advance. At present there is no evidence to show that the Kutiah glacier has previously shown such an exceptional advance, but a search for such evidence, and a knowledge of the accumulation zone of the glacier, would obviously help our understanding of this advance. In such an area with a considerable expanse of exposed rock, climatic change may trigger off an advance more quickly by conduction of heat through the rock, rather than by the slow process of modifying the firn temperatures. This would agree with Professor Desio's proposed explanation.

While this paper has not given the full solution to a complex problem, I hope it has drawn attention to the large effect of ice movement on the transfer of heat. Further field results are needed which include measurements of surface velocity in addition to accumulation and temperature gradient data.

I wish to acknowledge the help of Dr. G. Field in fitting equation (14) to field measurements, and of an Imperial Chemical Industries Research Fellowship during the tenure of which the ideas outlined in this paper were formed.

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