## Tensor Tomography

This chapter solves the tensor tomography problem for simple surfaces following Paternain et al. (2013). We shall in fact prove a stronger result in which the absence of conjugate points is replaced by the assumption that $I_{0}^{*}$ is surjective. In order to do this we introduce the notion of holomorphic integrating factors and prove their existence, which will be important in later chapters.

### 10.1 Holomorphic Integrating Factors

Let $(M, g)$ be a compact non-trapping surface having strictly convex boundary, and consider the geodesic X-ray transform $I_{m}$ that acts on symmetric $m$-tensor fields. Recall that the solenoidal injectivity of $I_{m}$ is equivalent with a uniqueness statement for the transport equation (see Proposition 6.4.4). We will focus on proving this uniqueness statement.

Suppose that $u \in C^{\infty}(S M)$ solves

$$
\begin{equation*}
X u=-f \text { in } S M,\left.\quad u\right|_{\partial S M}=0, \tag{10.1}
\end{equation*}
$$

where $f$ has degree $m$. For simplicity, assume that $f \in \Omega_{m}$. By Lemma 6.1.3, in the special coordinates $(x, \theta)$ on $S M$ we may write

$$
f(x, \theta)=\tilde{f}(x) e^{i m \theta}
$$

Recall that we already know how to deal with the case where $m=0$ (this is the injectivity of $I_{0}$ proved in Theorem 4.4.1). Let us try to reduce to this case simply by multiplying the equation (10.1) by $e^{-i m \theta}$. This gives a new transport equation for $e^{-i m \theta} u$ :

$$
\begin{equation*}
(X+a)\left(e^{-i m \theta} u\right)=-\tilde{f}(x),\left.\quad e^{-i m \theta} u\right|_{\partial S M}=0 \tag{10.2}
\end{equation*}
$$

where $a:=-e^{i m \theta} X\left(e^{-i m \theta}\right)$. Note that $a \in \Omega_{-1} \oplus \Omega_{1}$, since $X=\eta_{+}+\eta_{-}$ and

$$
e^{i m \theta} \eta_{ \pm}\left(e^{-i m \theta}\right) \in \Omega_{ \pm 1}
$$

We have now reduced the equation (10.1), where the right-hand side has degree $m$, to a new transport equation (10.2) where the right-hand side has degree 0 . However, the price to pay is that the new equation has a nontrivial attenuation factor $a$. One could ask if there is another reduction that would remove this factor. The next example gives such a reduction in elementary ODE theory.

Example 10.1.1 (Integrating factor) Consider the ODE

$$
u^{\prime}(t)+a(t) u(t)=f(t), \quad u(0)=0
$$

The standard method for solving this ODE is to introduce the integrating factor $w(t)=\int_{0}^{t} a(s) d s$, so that the equation is equivalent with

$$
\left(e^{w} u\right)^{\prime}(t)=\left(e^{w} f\right)(t), \quad\left(e^{w} u\right)(0)=0
$$

Using an integrating factor has removed the zero-order term from the equation, which can now be solved just by integration. The solution is

$$
u(t)=e^{-w(t)} \int_{0}^{t}\left(e^{w} f\right)(s) d s
$$

In geodesic X-ray transform problems, we are often dealing with equations such as

$$
X u+a u=-f \text { in } S M,\left.\quad u\right|_{\partial S M}=0,
$$

where $a \in C^{\infty}(S M)$ is an attenuation factor and $f \in C^{\infty}(S M)$. We would like to use an integrating factor $w \in C^{\infty}(S M)$ satisfying $X w=a$ in $S M$, which reduces the equation to

$$
X\left(e^{w} u\right)=-e^{w} f \text { in } S M,\left.\quad e^{w} u\right|_{\partial S M}=0
$$

This can always be done, for instance by choosing $w=u^{-a}$ (which may not be smooth at $\left.\partial_{0} S M\right)$. However, in many applications one has special structure, in particular, $f$ often has finite degree (e.g. $f=\tilde{f}(x)$ as in (10.2)). The problem with applying an arbitrary integrating factor is that multiplication by $e^{w}$ may destroy this special structure. For example, if $f=\tilde{f}(x)$, then $e^{w} f$ could have Fourier modes of all degrees.

In this section we prove an important technical result about the existence of a certain solution of the transport equation $X w=a$ when $a \in \Omega_{-1} \oplus \Omega_{1}$ (i.e. $a$ represents a 1 -form on $M$ ), where $w$ is fibrewise holomorphic in the
sense of Definition 6.1.14. This provides some control on the Fourier support of $e^{w} f$; e.g. if $f=\tilde{f}(x)$, then $e^{w} f$ is at least holomorphic. This result, which goes back to Salo and Uhlmann (2011) in the case of simple surfaces with $a \in \Omega_{0}$, will unlock the solution to several geometric inverse problems in two dimensions.

Proposition 10.1.2 (Holomorphic integrating factors, part I) Let ( $M, g$ ) be a compact non-trapping surface with strictly convex boundary. Assume that $I_{0}^{*}$ is surjective. Given $a_{-1}+a_{1} \in \Omega_{-1} \oplus \Omega_{1}$, there exists $w \in C^{\infty}(S M)$ such that $w$ is holomorphic and $X w=a_{-1}+a_{1}$. Similarly there exists $\tilde{w} \in C^{\infty}(S M)$ such that $\tilde{w}$ is anti-holomorphic and $X \tilde{w}=a_{-1}+a_{1}$.

Proof We do the proof for $w$ holomorphic; the proof for $\tilde{w}$ anti-holomorphic is analogous (or can be obtained by conjugation).

First we note that one can find $f_{0} \in C^{\infty}(M)$ satisfying $\eta_{+} f_{0}=-a_{1}$. Indeed, by Remark 3.4.17 $M$ is diffeomorphic to the closed unit disk $\overline{\mathbb{D}}$ and there are global special coordinates $(x, \theta)$ in $S M$. By Lemma 6.1.8 one has in these coordinates

$$
\eta_{+} f_{0}=e^{-\lambda} \partial_{z}\left(f_{0}\right) e^{i \theta}, \quad a_{1}=\tilde{a}_{1}\left(x_{1}, x_{2}\right) e^{i \theta}
$$

Thus it is enough to find $f_{0} \in C^{\infty}(\overline{\mathbb{D}})$ solving the equation

$$
\partial_{z}\left(f_{0}\right)=-e^{\lambda} \tilde{a}_{1} \quad \text { in } \mathbb{D}
$$

This equation can be solved for instance by extending the function on the righthand side smoothly as a function in $C_{c}^{\infty}(\mathbb{C})$, and then by applying a Cauchy transform (inverse of $\partial_{z}$ ).

Since $I_{0}^{*}$ is surjective, there exists $q \in C^{\infty}(S M)$ such that $X q=0$ and $q_{0}=f_{0}$ (see Theorem 8.2.2). Recalling that $X=\eta_{+}+\eta_{-}$and looking at Fourier coefficients of $X q$, we see that $\eta_{+} q_{k-1}+\eta_{-} q_{k+1}=0$ for all $k$. Hence

$$
\begin{equation*}
X\left(q_{2}+q_{4}+\cdots\right)=\eta_{-} q_{2}=-\eta_{+} q_{0}=a_{1} \tag{10.3}
\end{equation*}
$$

Next, we solve $\eta_{-} g_{0}=a_{-1}$ and use surjectivity of $I_{0}^{*}$ to find $p \in C^{\infty}(S M)$ such that $X p=0$ and $p_{0}=g_{0}$. Hence

$$
\begin{equation*}
X\left(p_{0}+p_{2}+\cdots\right)=\eta_{-} p_{0}=a_{-1} \tag{10.4}
\end{equation*}
$$

Combining (10.3) and (10.4) and setting $w=\sum_{k \geq 0} p_{2 k}+\sum_{k \geq 1} q_{2 k}$, we see that $w$ is holomorphic and $X w=a_{-1}+a_{1}$.

### 10.2 Tensor Tomography

Our main result gives a positive answer to the tensor tomography problem in the case of surfaces with $I_{0}^{*}$ surjective.

Theorem 10.2.1 (Tensor tomography) Let $(M, g)$ be a compact non-trapping surface with strictly convex boundary and $I_{0}^{*}$ surjective. The transform $I_{m}$ is $s$-injective for any $m \geq 0$.

We note that for the case of $(M, g)$ simple and $m=2$, solenoidal injectivity of $I_{2}$ was proved in Sharafutdinov (2007) using the solution to the boundary rigidity problem. We begin with a simple observation that holds in any dimension.

Lemma 10.2.2 Let $(M, g)$ be a compact non-trapping manifold with strictly convex boundary. If $I_{0}^{*}: C_{\alpha}^{\infty}\left(\partial_{+} S M\right) \rightarrow C^{\infty}(M)$ is surjective, then $I_{0}: C^{\infty}(M) \rightarrow C^{\infty}\left(\partial_{+} S M\right)$ is injective.

Proof Suppose that $f \in C^{\infty}(M)$ satisfies $I_{0} f=0$. If $I_{0}^{*}$ is surjective, there is $w \in C_{\alpha}^{\infty}\left(\partial_{+} S M\right)$ such that $I_{0}^{*} w=f$. Hence we can write

$$
\|f\|^{2}=\left(f, I_{0}^{*} w\right)_{L^{2}(M)}=\left(I_{0} f, w\right)_{L_{\mu}^{2}\left(\partial_{+} S M\right)}=0
$$

and thus $f=0$.
The next result is the master result from which tensor tomography is derived. It asserts, in terms of the transport equation, that $\left.I\right|_{\Omega_{m}}: \Omega_{m} \rightarrow C^{\infty}\left(\partial_{+} S M\right)$ is injective whenever $I_{0}^{*}$ is surjective.

Theorem 10.2.3 (Injectivity of $\left.I\right|_{\Omega_{m}}$ ) Let $(M, g)$ be a compact non-trapping surface with strictly convex boundary and $I_{0}^{*}$ surjective. Assume that $m \in \mathbb{Z}$, and let $u \in C^{\infty}(S M)$ be such that

$$
X u=-f \in \Omega_{m},\left.\quad u\right|_{\partial S M}=0
$$

Then $u=0$ and $f=0$.
The proof is based on another important injectivity result, where the fact that $f$ has one-sided Fourier support is used to deduce that $u$ has one-sided Fourier support. A more precise result in this direction will be given in Proposition 10.2.6.

Proposition 10.2.4 Let $(M, g)$ be a compact non-trapping surface with strictly convex boundary and $I_{0}$ injective. If $u \in C^{\infty}(S M)$ is odd and satisfies

$$
X u=-f \text { in } S M,\left.\quad u\right|_{\partial S M}=0
$$

where $f$ is holomorphic (respectively anti-holomorphic), then $u$ is holomorphic (respectively anti-holomorphic).

Proof We prove the case where $f$ is holomorphic. Write $q:=\sum_{k=-\infty}^{-1} u_{k}$. Since $f$ is holomorphic, we have $(X u)_{k}=0$ for $k \leq-1$, and using the decomposition $X=\eta_{+}+\eta_{-}$this gives that $\eta_{+} u_{k-1}+\eta_{-} u_{k+1}=0$ for $k \leq-1$. Thus we obtain that

$$
X q=\eta_{+} u_{-1},\left.\quad q\right|_{\partial S M}=0
$$

Now $\eta_{+} u_{-1}$ only depends on $x$, and hence the injectivity of $I_{0}$ implies that $\eta_{+} u_{-1}=0$. This proves that $q=0$ showing that $u$ is holomorphic.

Proof of Theorem 10.2.3 We follow the approach described at the beginning of Section 10.1. Let $r:=e^{-i m \theta}$ and observe that $r^{-1} X r \in \Omega_{-1} \oplus \Omega_{1}$ since

$$
e^{i m \theta} \eta_{ \pm}\left(e^{-i m \theta}\right) \in \Omega_{ \pm 1}
$$

By Proposition 10.1.2, there is a holomorphic $w \in C^{\infty}(S M)$ and antiholomorphic $\tilde{w} \in C^{\infty}(S M)$ such that $X w=X \tilde{w}=-r^{-1} X r$. Since $r^{-1} X r$ is odd, without loss of generality we may replace $w$ and $\tilde{w}$ by their even parts so that $w$ and $\tilde{w}$ are even. A simple calculation shows that

$$
\begin{equation*}
X\left(e^{w} r u\right)=e^{w}\left(X-r^{-1} X r\right)(r u)=-e^{w} r f \tag{10.5}
\end{equation*}
$$

with a similar equation for $\tilde{w}$. Since $r f \in \Omega_{0}, e^{w} r f$ is holomorphic and $e^{\tilde{w}} r f$ is anti-holomorphic.

Assume now that $m$ is even, the proof for $m$ odd being very similar. Then we may assume that $u$ is odd and thus $e^{w} r u$ and $e^{\tilde{w}} r u$ are odd. By Proposition 10.2.4, since we have

$$
X\left(e^{w} r u\right)=-e^{w} r f,\left.\quad e^{w} r u\right|_{\partial S M}=0,
$$

we see that $e^{w} r u$ is holomorphic and thus $r u=e^{-w}\left(e^{w} r u\right)$ is holomorphic. Arguing with $\tilde{w}$ we deduce that $r u$ is also anti-holomorphic. Thus one must have $r u \in \Omega_{0}$. This implies that $u \in \Omega_{m}$, and using that $X u \in \Omega_{m}$ we see that $X u=0$ and finally $u=f=0$ as desired.

One can explicitly compute $r^{-1} \mathrm{Xr}$ in the proof above using isothermal coordinates in which the metric is $e^{2 \lambda}\left(d x_{1}^{2}+d x_{2}^{2}\right)$ :

Exercise 10.2.5 Show that

$$
r^{-1} X r=m \eta_{+}(\lambda)-m \eta_{-}(\lambda)
$$

By inspecting the proof of Proposition 10.1.2 show that the conclusion of Theorem 10.2.3 still holds if we assume that $I_{0}$ is injective and there is a
smooth $q$ such that $X q=0$ with $q_{0}=\lambda$. Hence surjectivity of $I_{0}^{*}$ is only needed for the function $\lambda$ !

We will give two corollaries of Theorem 10.2.4.
Proposition 10.2.6 Let $(M, g)$ be a compact non-trapping surface with strictly convex boundary and $I_{0}^{*}$ surjective. Let $u \in C^{\infty}(S M)$ be such that

$$
X u=-f,\left.\quad u\right|_{\partial S M}=0
$$

Suppose $f_{k}=0$ for $k \geq m+1$ for some $m \in \mathbb{Z}$. Then $u_{k}=0$ for $k \geq m$. Similarly, if $f_{k}=0$ for $k \leq m-1$ for some $m \in \mathbb{Z}$, then $u_{k}=0$ for $k \leq m$.

Proof Suppose $f_{k}=0$ for $k \geq m+1$. Let $w:=\sum_{m}^{\infty} u_{k}$. Using the equation $X u=-f$ and the hypothesis on $f$, we see that

$$
X w=\eta_{-} u_{m}+\eta_{-} u_{m+1} \in \Omega_{m-1} \oplus \Omega_{m} .
$$

Applying Theorem 10.2 .3 to the even and odd parts of $w$, we deduce that $w=0$ and thus $u_{k}=0$ for $k \geq m$. Similarly, arguing with $\sum_{-\infty}^{m} u_{k}$ we deduce that $u_{k}=0$ for $k \leq m$ if $f_{k}=0$ for $k \leq m-1$.

The next corollary is an obvious consequence of the previous proposition.
Corollary 10.2.7 (Tensor tomography, transport version) Let $(M, g)$ be a non-trapping surface with strictly convex boundary and $I_{0}^{*}$ surjective. Let $u \in C^{\infty}(S M)$ be such that

$$
X u=f,\left.\quad u\right|_{\partial S M}=0 .
$$

Suppose $f_{k}=0$ for $|k| \geq m+1$ for some $m \geq 0$. Then $u_{k}=0$ for $|k| \geq m$ (when $m=0$, this means $u=f=0$ ).

By Proposition 6.4.4, the previous result also proves Theorem 10.2.1.

### 10.3 Range for Tensors

In this section we explain how some of the ideas of the previous section can be employed to give a description of the range for the X-ray transform acting on symmetric tensors of any rank, pretty much in the spirit of Theorem 9.6.2.

Let $(M, g)$ be a non-trapping surface with strictly convex boundary. Pick a function $h: S M \rightarrow S^{1} \subset \mathbb{C}$ such that $h \in \Omega_{1}$. Such a function always exists: for instance, in global isothermal coordinates we may simply take $h=e^{i \theta}$. Our description of the range will be based on this choice of $h$. Define the 1 -form

$$
A:=-h^{-1} X h
$$

Observe that since $h \in \Omega_{1}$, then $h^{-1}=\bar{h} \in \Omega_{-1}$. Also $X h=\eta_{+} h+\eta_{-} h \in$ $\Omega_{2} \oplus \Omega_{0}$, which implies that $A \in \Omega_{1} \oplus \Omega_{-1}$. It follows that $A$ is the restriction to $S M$ of a purely imaginary 1 -form on $M$.

First we will describe the range of the geodesic ray transform $I$ restricted to $\Omega_{m}$ :

$$
\mathbf{I}_{m}:=\left.I\right|_{\Omega_{m}}: \Omega_{m} \rightarrow C^{\infty}\left(\partial_{+} S M, \mathbb{C}\right) .
$$

Observe that if $u$ solves the transport equation $X u=-f$ where $f \in \Omega_{m}$ and $\left.u\right|_{\partial_{-} S M}=0$, then $h^{-m} u$ solves $(X-m A)\left(h^{-m} u\right)=-h^{-m} f$ and $\left.h^{-m} u\right|_{\partial_{-} S M}=0$. Also note that $h^{-m} f \in \Omega_{0}$. Thus

$$
\begin{equation*}
I_{-m A}\left(h^{-m} f\right)=\left(\left.h^{-m}\right|_{\partial_{+} S M}\right) \mathbf{I}_{m}(f), \tag{10.6}
\end{equation*}
$$

where the left-hand side is an attenuated X-ray transform with attenuation $-m A$ as given in Definition 5.3.3. The relation in (10.6) is telling us that if we know how to describe the range of $I_{A}$ acting on $C^{\infty}(M)$, where $A$ is a purely imaginary 1-form, then we would know how to describe the range of $\mathbf{I}_{m}$. It turns out that this is possible to do even in much greater generality, namely when $A$ is a connection (cf. Theorem 14.5.5). We will return to this topic in later chapters; for the time being we content ourselves with a description of the results.

Let $Q_{m}: C\left(\partial_{+} S M, \mathbb{C}\right) \rightarrow C(\partial S M, \mathbb{C})$ be given by

$$
Q_{m} w(x, v):= \begin{cases}w(x, v) & \text { if }(x, v) \in \partial_{+} S M \\ \left(e^{-m \int_{0}^{\tau(x, v)} A\left(\varphi_{t}(x, v)\right) d t} w\right) \circ \alpha(x, v) & \text { if }(x, v) \in \partial_{-} S M\end{cases}
$$

and let $B_{m}: C(\partial S M, \mathbb{C}) \rightarrow C\left(\partial_{+} S M, \mathbb{C}\right)$ be

$$
B_{m} g:=\left.\left[g-e^{m \int_{0}^{\tau(x, v)} A\left(\varphi_{t}(x, v)\right) d t}(g \circ \alpha)\right]\right|_{\partial_{+} S M}
$$

In other words, with $I_{1}$ denoting the X-ray transform on 1-tensors, we have

$$
Q_{m} w(x, v)= \begin{cases}w(x, v) & \text { if }(x, v) \in \partial_{+} S M \\ \left(e^{-m I_{1}(A)} w\right) \circ \alpha(x, v) & \text { if }(x, v) \in \partial_{-} S M,\end{cases}
$$

and

$$
B_{m} g=\left.\left[g-e^{m I_{1}(A)}(g \circ \alpha)\right]\right|_{\partial_{+} S M}
$$

We define

$$
P_{m,-}:=B_{m} H_{-} Q_{m}
$$

The following result from Paternain et al. (2015b) describes the range of $\mathbf{I}_{m}$.

Theorem 10.3.1 Assume that $(M, g)$ is a simple surface. A function $u \in C^{\infty}\left(\partial_{+} S M, \mathbb{C}\right)$ belongs to the range of $\mathbf{I}_{m}$ if and only if $u=$ $\left(\left.h^{m}\right|_{\partial_{+} S M}\right) P_{m,-} w$ for $w \in \mathcal{S}_{m}^{\infty}\left(\partial_{+} S M, \mathbb{C}\right)$, where this last space denotes the set of all smooth $w$ such that $Q_{m} w$ is smooth.

Suppose that $F$ is a complex-valued symmetric tensor of order $m$ and denote its restriction to $S M$ by $f$. Recall from Proposition 6.3.5 that there is a one-to-one correspondence between complex-valued symmetric tensors of order $m$ and functions in $S M$ of the form $f=\sum_{k=-m}^{m} f_{k}$, where $f_{k} \in \Omega_{k}$ and $f_{k}=0$ for all $k$ odd (respectively even) if $m$ is even (respectively odd).

Since

$$
I(f)=\sum_{k=-m}^{m} \mathbf{I}_{k}\left(f_{k}\right)
$$

we deduce directly from Theorem 10.3.1 the following.
Theorem 10.3.2 Let $(M, g)$ be a simple surface. If $m=2 l$ is even, a function $u \in C^{\infty}\left(\partial_{+} S M, \mathbb{C}\right)$ belongs to the range of the $X$-ray transform acting on complex-valued symmetric $m$-tensors if and only if there are $w_{2 k} \in$ $\mathcal{S}_{2 k}^{\infty}\left(\partial_{+} S M, \mathbb{C}\right)$ such that

$$
u=\sum_{k=-l}^{l}\left(\left.h^{2 k}\right|_{\partial_{+} S M}\right) P_{2 k,-} w_{2 k} .
$$

Similarly, if $m=2 l+1$ is odd, a function $u \in C^{\infty}\left(\partial_{+} S M, \mathbb{C}\right)$ belongs to the range of the $X$-ray transform acting on complex-valued symmetric m-tensors if and only if there are $w_{2 k+1} \in \mathcal{S}_{2 k+1}^{\infty}\left(\partial_{+} S M, \mathbb{C}\right)$ such that

$$
u=\sum_{k=-l-1}^{l}\left(\left.h^{2 k+1}\right|_{\partial_{+} S M}\right) P_{2 k+1,-} w_{2 k+1} .
$$

