## Tensor Tomography

This chapter solves the tensor tomography problem for simple surfaces following Paternain et al. (2013). We shall in fact prove a stronger result in which the absence of conjugate points is replaced by the assumption that  $I_0^*$  is surjective. In order to do this we introduce the notion of holomorphic integrating factors and prove their existence, which will be important in later chapters.

## **10.1 Holomorphic Integrating Factors**

Let (M, g) be a compact non-trapping surface having strictly convex boundary, and consider the geodesic X-ray transform  $I_m$  that acts on symmetric *m*-tensor fields. Recall that the solenoidal injectivity of  $I_m$  is equivalent with a uniqueness statement for the transport equation (see Proposition 6.4.4). We will focus on proving this uniqueness statement.

Suppose that  $u \in C^{\infty}(SM)$  solves

$$Xu = -f \text{ in } SM, \qquad u|_{\partial SM} = 0, \tag{10.1}$$

where *f* has degree *m*. For simplicity, assume that  $f \in \Omega_m$ . By Lemma 6.1.3, in the special coordinates  $(x, \theta)$  on *SM* we may write

$$f(x,\theta) = \tilde{f}(x)e^{im\theta}.$$

Recall that we already know how to deal with the case where m = 0 (this is the injectivity of  $I_0$  proved in Theorem 4.4.1). Let us try to reduce to this case simply by multiplying the equation (10.1) by  $e^{-im\theta}$ . This gives a new transport equation for  $e^{-im\theta}u$ :

$$(X+a)(e^{-im\theta}u) = -\tilde{f}(x), \qquad e^{-im\theta}u|_{\partial SM} = 0, \tag{10.2}$$

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where  $a := -e^{im\theta}X(e^{-im\theta})$ . Note that  $a \in \Omega_{-1} \oplus \Omega_1$ , since  $X = \eta_+ + \eta_$ and

$$e^{im\theta}\eta_{\pm}(e^{-im\theta})\in\Omega_{\pm 1}$$

We have now reduced the equation (10.1), where the right-hand side has degree m, to a new transport equation (10.2) where the right-hand side has degree 0. However, the price to pay is that the new equation has a nontrivial attenuation factor a. One could ask if there is another reduction that would remove this factor. The next example gives such a reduction in elementary ODE theory.

Example 10.1.1 (Integrating factor) Consider the ODE

$$u'(t) + a(t)u(t) = f(t), \qquad u(0) = 0.$$

The standard method for solving this ODE is to introduce the *integrating factor*  $w(t) = \int_0^t a(s) ds$ , so that the equation is equivalent with

$$(e^{w}u)'(t) = (e^{w}f)(t), \qquad (e^{w}u)(0) = 0.$$

Using an integrating factor has removed the zero-order term from the equation, which can now be solved just by integration. The solution is

$$u(t) = e^{-w(t)} \int_0^t \left( e^w f \right)(s) \, ds$$

In geodesic X-ray transform problems, we are often dealing with equations such as

$$Xu + au = -f$$
 in  $SM$ ,  $u|_{\partial SM} = 0$ ,

where  $a \in C^{\infty}(SM)$  is an attenuation factor and  $f \in C^{\infty}(SM)$ . We would like to use an integrating factor  $w \in C^{\infty}(SM)$  satisfying Xw = a in SM, which reduces the equation to

$$X(e^w u) = -e^w f$$
 in  $SM$ ,  $e^w u|_{\partial SM} = 0$ .

This can always be done, for instance by choosing  $w = u^{-a}$  (which may not be smooth at  $\partial_0 SM$ ). However, in many applications one has special structure, in particular, f often has finite degree (e.g.  $f = \tilde{f}(x)$  as in (10.2)). The problem with applying an arbitrary integrating factor is that multiplication by  $e^w$  may destroy this special structure. For example, if  $f = \tilde{f}(x)$ , then  $e^w f$  could have Fourier modes of all degrees.

In this section we prove an important technical result about the existence of a certain solution of the transport equation Xw = a when  $a \in \Omega_{-1} \oplus \Omega_1$ (i.e. *a* represents a 1-form on *M*), where *w* is fibrewise *holomorphic* in the sense of Definition 6.1.14. This provides some control on the Fourier support of  $e^w f$ ; e.g. if  $f = \tilde{f}(x)$ , then  $e^w f$  is at least holomorphic. This result, which goes back to Salo and Uhlmann (2011) in the case of simple surfaces with  $a \in \Omega_0$ , will unlock the solution to several geometric inverse problems in two dimensions.

**Proposition 10.1.2** (Holomorphic integrating factors, part I) Let (M, g) be a compact non-trapping surface with strictly convex boundary. Assume that  $I_0^*$  is surjective. Given  $a_{-1} + a_1 \in \Omega_{-1} \oplus \Omega_1$ , there exists  $w \in C^{\infty}(SM)$  such that w is holomorphic and  $Xw = a_{-1} + a_1$ . Similarly there exists  $\tilde{w} \in C^{\infty}(SM)$  such that  $\tilde{w}$  is anti-holomorphic and  $X\tilde{w} = a_{-1} + a_1$ .

*Proof* We do the proof for w holomorphic; the proof for  $\tilde{w}$  anti-holomorphic is analogous (or can be obtained by conjugation).

First we note that one can find  $f_0 \in C^{\infty}(M)$  satisfying  $\eta_+ f_0 = -a_1$ . Indeed, by Remark 3.4.17 *M* is diffeomorphic to the closed unit disk  $\overline{\mathbb{D}}$  and there are global special coordinates  $(x, \theta)$  in *SM*. By Lemma 6.1.8 one has in these coordinates

$$\eta_+ f_0 = e^{-\lambda} \partial_z(f_0) e^{i\theta}, \qquad a_1 = \tilde{a}_1(x_1, x_2) e^{i\theta}.$$

Thus it is enough to find  $f_0 \in C^{\infty}(\overline{\mathbb{D}})$  solving the equation

$$\partial_z(f_0) = -e^{\lambda} \tilde{a}_1$$
 in  $\mathbb{D}$ .

This equation can be solved for instance by extending the function on the righthand side smoothly as a function in  $C_c^{\infty}(\mathbb{C})$ , and then by applying a Cauchy transform (inverse of  $\partial_z$ ).

Since  $I_0^*$  is surjective, there exists  $q \in C^{\infty}(SM)$  such that Xq = 0 and  $q_0 = f_0$  (see Theorem 8.2.2). Recalling that  $X = \eta_+ + \eta_-$  and looking at Fourier coefficients of Xq, we see that  $\eta_+q_{k-1} + \eta_-q_{k+1} = 0$  for all k. Hence

$$X(q_2 + q_4 + \dots) = \eta_- q_2 = -\eta_+ q_0 = a_1.$$
(10.3)

Next, we solve  $\eta_{-}g_0 = a_{-1}$  and use surjectivity of  $I_0^*$  to find  $p \in C^{\infty}(SM)$  such that Xp = 0 and  $p_0 = g_0$ . Hence

$$X(p_0 + p_2 + \dots) = \eta_- p_0 = a_{-1}.$$
 (10.4)

Combining (10.3) and (10.4) and setting  $w = \sum_{k\geq 0} p_{2k} + \sum_{k\geq 1} q_{2k}$ , we see that w is holomorphic and  $Xw = a_{-1} + a_1$ .

## **10.2 Tensor Tomography**

Our main result gives a positive answer to the tensor tomography problem in the case of surfaces with  $I_0^*$  surjective.

**Theorem 10.2.1** (Tensor tomography) Let (M, g) be a compact non-trapping surface with strictly convex boundary and  $I_0^*$  surjective. The transform  $I_m$  is *s*-injective for any  $m \ge 0$ .

We note that for the case of (M, g) simple and m = 2, solenoidal injectivity of  $I_2$  was proved in Sharafutdinov (2007) using the solution to the boundary rigidity problem. We begin with a simple observation that holds in any dimension.

**Lemma 10.2.2** Let (M,g) be a compact non-trapping manifold with strictly convex boundary. If  $I_0^*: C_\alpha^\infty(\partial_+ SM) \to C^\infty(M)$  is surjective, then  $I_0: C^\infty(M) \to C^\infty(\partial_+ SM)$  is injective.

*Proof* Suppose that  $f \in C^{\infty}(M)$  satisfies  $I_0 f = 0$ . If  $I_0^*$  is surjective, there is  $w \in C_{\alpha}^{\infty}(\partial_+ SM)$  such that  $I_0^* w = f$ . Hence we can write

$$||f||^{2} = (f, I_{0}^{*}w)_{L^{2}(M)} = (I_{0}f, w)_{L^{2}_{\mu}(\partial_{+}SM)} = 0,$$

and thus f = 0.

The next result is the master result from which tensor tomography is derived. It asserts, in terms of the transport equation, that  $I|_{\Omega_m} : \Omega_m \to C^{\infty}(\partial_+ SM)$  is injective whenever  $I_0^*$  is surjective.

**Theorem 10.2.3** (Injectivity of  $I|_{\Omega_m}$ ) Let (M, g) be a compact non-trapping surface with strictly convex boundary and  $I_0^*$  surjective. Assume that  $m \in \mathbb{Z}$ , and let  $u \in C^{\infty}(SM)$  be such that

$$Xu = -f \in \Omega_m, \quad u|_{\partial SM} = 0.$$

Then u = 0 and f = 0.

The proof is based on another important injectivity result, where the fact that f has one-sided Fourier support is used to deduce that u has one-sided Fourier support. A more precise result in this direction will be given in Proposition 10.2.6.

**Proposition 10.2.4** *Let* (M, g) *be a compact non-trapping surface with strictly convex boundary and*  $I_0$  *injective. If*  $u \in C^{\infty}(SM)$  *is odd and satisfies* 

$$Xu = -f \text{ in } SM, \qquad u|_{\partial SM} = 0,$$

where *f* is holomorphic (respectively anti-holomorphic), then *u* is holomorphic (respectively anti-holomorphic).

*Proof* We prove the case where *f* is holomorphic. Write  $q := \sum_{k=-\infty}^{-1} u_k$ . Since *f* is holomorphic, we have  $(Xu)_k = 0$  for  $k \le -1$ , and using the decomposition  $X = \eta_+ + \eta_-$  this gives that  $\eta_+ u_{k-1} + \eta_- u_{k+1} = 0$  for  $k \le -1$ . Thus we obtain that

$$Xq = \eta_+ u_{-1}, \qquad q|_{\partial SM} = 0.$$

Now  $\eta_+u_{-1}$  only depends on x, and hence the injectivity of  $I_0$  implies that  $\eta_+u_{-1} = 0$ . This proves that q = 0 showing that u is holomorphic.

*Proof of Theorem 10.2.3* We follow the approach described at the beginning of Section 10.1. Let  $r := e^{-im\theta}$  and observe that  $r^{-1}Xr \in \Omega_{-1} \oplus \Omega_1$  since

$$e^{im\theta}\eta_{\pm}(e^{-im\theta})\in\Omega_{\pm 1}.$$

By Proposition 10.1.2, there is a holomorphic  $w \in C^{\infty}(SM)$  and antiholomorphic  $\tilde{w} \in C^{\infty}(SM)$  such that  $Xw = X\tilde{w} = -r^{-1}Xr$ . Since  $r^{-1}Xr$  is odd, without loss of generality we may replace w and  $\tilde{w}$  by their even parts so that w and  $\tilde{w}$  are even. A simple calculation shows that

$$X\left(e^{w}ru\right) = e^{w}\left(X - r^{-1}Xr\right)(ru) = -e^{w}rf$$
(10.5)

with a similar equation for  $\tilde{w}$ . Since  $rf \in \Omega_0$ ,  $e^w rf$  is holomorphic and  $e^{\tilde{w}} rf$  is anti-holomorphic.

Assume now that *m* is even, the proof for *m* odd being very similar. Then we may assume that *u* is odd and thus  $e^w ru$  and  $e^{\tilde{w}} ru$  are odd. By Proposition 10.2.4, since we have

$$X(e^{w}ru) = -e^{w}rf, \qquad e^{w}ru|_{\partial SM} = 0,$$

we see that  $e^w ru$  is holomorphic and thus  $ru = e^{-w}(e^w ru)$  is holomorphic. Arguing with  $\tilde{w}$  we deduce that ru is also anti-holomorphic. Thus one must have  $ru \in \Omega_0$ . This implies that  $u \in \Omega_m$ , and using that  $Xu \in \Omega_m$  we see that Xu = 0 and finally u = f = 0 as desired.

One can explicitly compute  $r^{-1}Xr$  in the proof above using isothermal coordinates in which the metric is  $e^{2\lambda}(dx_1^2 + dx_2^2)$ :

Exercise 10.2.5 Show that

$$r^{-1}Xr = m\eta_{+}(\lambda) - m\eta_{-}(\lambda).$$

By inspecting the proof of Proposition 10.1.2 show that the conclusion of Theorem 10.2.3 still holds if we assume that  $I_0$  is injective and there is a

smooth q such that Xq = 0 with  $q_0 = \lambda$ . Hence surjectivity of  $I_0^*$  is only needed for the function  $\lambda$ !

We will give two corollaries of Theorem 10.2.4.

**Proposition 10.2.6** *Let* (M, g) *be a compact non-trapping surface with strictly convex boundary and*  $I_0^*$  *surjective. Let*  $u \in C^{\infty}(SM)$  *be such that* 

 $Xu = -f, \quad u|_{\partial SM} = 0.$ 

Suppose  $f_k = 0$  for  $k \ge m + 1$  for some  $m \in \mathbb{Z}$ . Then  $u_k = 0$  for  $k \ge m$ . Similarly, if  $f_k = 0$  for  $k \le m - 1$  for some  $m \in \mathbb{Z}$ , then  $u_k = 0$  for  $k \le m$ .

*Proof* Suppose  $f_k = 0$  for  $k \ge m + 1$ . Let  $w := \sum_{m=1}^{\infty} u_k$ . Using the equation Xu = -f and the hypothesis on f, we see that

$$Xw = \eta_{-}u_{m} + \eta_{-}u_{m+1} \in \Omega_{m-1} \oplus \Omega_{m}$$

Applying Theorem 10.2.3 to the even and odd parts of w, we deduce that w = 0 and thus  $u_k = 0$  for  $k \ge m$ . Similarly, arguing with  $\sum_{-\infty}^m u_k$  we deduce that  $u_k = 0$  for  $k \le m$  if  $f_k = 0$  for  $k \le m - 1$ .

The next corollary is an obvious consequence of the previous proposition.

**Corollary 10.2.7** (Tensor tomography, transport version) Let (M,g) be a non-trapping surface with strictly convex boundary and  $I_0^*$  surjective. Let  $u \in C^{\infty}(SM)$  be such that

$$Xu = f, \quad u|_{\partial SM} = 0.$$

Suppose  $f_k = 0$  for  $|k| \ge m + 1$  for some  $m \ge 0$ . Then  $u_k = 0$  for  $|k| \ge m$  (when m = 0, this means u = f = 0).

By Proposition 6.4.4, the previous result also proves Theorem 10.2.1.

## **10.3 Range for Tensors**

In this section we explain how some of the ideas of the previous section can be employed to give a description of the range for the X-ray transform acting on symmetric tensors of any rank, pretty much in the spirit of Theorem 9.6.2.

Let (M,g) be a non-trapping surface with strictly convex boundary. Pick a function  $h: SM \to S^1 \subset \mathbb{C}$  such that  $h \in \Omega_1$ . Such a function always exists: for instance, in global isothermal coordinates we may simply take  $h = e^{i\theta}$ . Our description of the range will be based on this choice of h. Define the 1-form

$$A := -h^{-1}Xh.$$

Observe that since  $h \in \Omega_1$ , then  $h^{-1} = \overline{h} \in \Omega_{-1}$ . Also  $Xh = \eta_+ h + \eta_- h \in \Omega_2 \oplus \Omega_0$ , which implies that  $A \in \Omega_1 \oplus \Omega_{-1}$ . It follows that A is the restriction to *SM* of a purely imaginary 1-form on *M*.

First we will describe the range of the geodesic ray transform *I* restricted to  $\Omega_m$ :

$$\mathbf{I}_m := I|_{\Omega_m} \colon \Omega_m \to C^{\infty}(\partial_+ SM, \mathbb{C}).$$

Observe that if *u* solves the transport equation Xu = -f where  $f \in \Omega_m$ and  $u|_{\partial_-SM} = 0$ , then  $h^{-m}u$  solves  $(X - mA)(h^{-m}u) = -h^{-m}f$  and  $h^{-m}u|_{\partial_-SM} = 0$ . Also note that  $h^{-m}f \in \Omega_0$ . Thus

$$I_{-mA}(h^{-m}f) = (h^{-m}|_{\partial_{+}SM}) \mathbf{I}_{m}(f), \qquad (10.6)$$

where the left-hand side is an attenuated X-ray transform with attenuation -mA as given in Definition 5.3.3. The relation in (10.6) is telling us that if we know how to describe the range of  $I_A$  acting on  $C^{\infty}(M)$ , where A is a purely imaginary 1-form, then we would know how to describe the range of  $\mathbf{I}_m$ . It turns out that this is possible to do even in much greater generality, namely when A is a *connection* (cf. Theorem 14.5.5). We will return to this topic in later chapters; for the time being we content ourselves with a description of the results.

Let  $Q_m: C(\partial_+ SM, \mathbb{C}) \to C(\partial SM, \mathbb{C})$  be given by

$$Q_m w(x,v) := \begin{cases} w(x,v) & \text{if } (x,v) \in \partial_+ SM, \\ (e^{-m \int_0^{\tau(x,v)} A(\varphi_t(x,v)) \, dt} w) \circ \alpha(x,v) & \text{if } (x,v) \in \partial_- SM, \end{cases}$$

and let  $B_m \colon C(\partial SM, \mathbb{C}) \to C(\partial_+ SM, \mathbb{C})$  be

$$B_mg := \left[g - e^{m\int_0^{\tau(x,v)} A(\varphi_t(x,v)) dt} (g \circ \alpha)\right]\Big|_{\partial_+ SM}.$$

In other words, with  $I_1$  denoting the X-ray transform on 1-tensors, we have

$$Q_m w(x,v) = \begin{cases} w(x,v) & \text{if } (x,v) \in \partial_+ SM, \\ (e^{-mI_1(A)}w) \circ \alpha(x,v) & \text{if } (x,v) \in \partial_- SM, \end{cases}$$

and

$$B_m g = \left[ g - e^{mI_1(A)} (g \circ \alpha) \right] \Big|_{\partial_+ SM}$$

We define

$$P_{m,-} := B_m H_- Q_m.$$

The following result from Paternain et al. (2015b) describes the range of  $I_m$ .

**Theorem 10.3.1** Assume that (M,g) is a simple surface. A function  $u \in C^{\infty}(\partial_+SM, \mathbb{C})$  belongs to the range of  $\mathbf{I}_m$  if and only if  $u = (h^m|_{\partial_+SM}) P_{m,-w}$  for  $w \in S_m^{\infty}(\partial_+SM, \mathbb{C})$ , where this last space denotes the set of all smooth w such that  $Q_m w$  is smooth.

Suppose that *F* is a complex-valued symmetric tensor of order *m* and denote its restriction to *SM* by *f*. Recall from Proposition 6.3.5 that there is a one-to-one correspondence between complex-valued symmetric tensors of order *m* and functions in *SM* of the form  $f = \sum_{k=-m}^{m} f_k$ , where  $f_k \in \Omega_k$  and  $f_k = 0$  for all *k* odd (respectively even) if *m* is even (respectively odd).

Since

$$I(f) = \sum_{k=-m}^{m} \mathbf{I}_k(f_k),$$

we deduce directly from Theorem 10.3.1 the following.

**Theorem 10.3.2** Let (M,g) be a simple surface. If m = 2l is even, a function  $u \in C^{\infty}(\partial_+SM, \mathbb{C})$  belongs to the range of the X-ray transform acting on complex-valued symmetric m-tensors if and only if there are  $w_{2k} \in S_{2k}^{\infty}(\partial_+SM, \mathbb{C})$  such that

$$u=\sum_{k=-l}^{l}\left(h^{2k}|_{\partial_{+}SM}\right)P_{2k,-}w_{2k}.$$

Similarly, if m = 2l + 1 is odd, a function  $u \in C^{\infty}(\partial_+ SM, \mathbb{C})$  belongs to the range of the X-ray transform acting on complex-valued symmetric m-tensors if and only if there are  $w_{2k+1} \in S^{\infty}_{2k+1}(\partial_+ SM, \mathbb{C})$  such that

$$u = \sum_{k=-l-1}^{l} \left( h^{2k+1} |_{\partial_{+}SM} \right) P_{2k+1,-} w_{2k+1}.$$