

# KOI-256's Magnetic Activity Under the Influence of the White Dwarf

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#### Abstract

We present the findings about chromospheric activity nature of KOI-256 obtained from the Kepler Mission data. First, it was found that there are some sinusoidal variations out-of-eclipses due to cool spot activity. The sinusoidal variations modelled by the SPOTMODEL program indicate that the active component has two different active regions. Their longitudinal variation revealed that one of them has a migration period of 3.95 yrs, while the other has a migration period of 8.37 yrs. Second, 225 flares were detected from the short cadence data in total. The parameters, such as increase ( $T_r$ ) and decay ( $T_d$ ) times, total flare time ( $T_t$ ), equivalent durations (P), were calculated for each flare. The distribution of equivalent durations versus total flare times in logarithmic scale is modelled to find flare activity level. The *Plateau* value known as the saturation level of the active component was calculated to be 2.3121  $\pm$  0.0964 s, and the *Half-life* value, which is required flare total time to reach the saturation, was computed to be 2233.6 s. In addition, the frequency of  $N_1$ , which is the number of flares per an hour in the system, was found to be 0.05087 h<sup>-1</sup>, while the flare frequency  $N_2$  that the flare-equivalent duration emitting per an hour was found to be 0.00051. Contrary to the spot activity, it has been found that the flares are in tends to appear at specific phases due to the white dwarf component.

Keywords: binaries: eclipsing – methods: data analysis – stars: flare – stars: individual (KOI-256) – stars: low-mass – techniques: photometric

# 1 INTRODUCTION

A total of 65% of the stars in our galaxy are red dwarf stars and 75% of them show flare activity (Rodonó 1986). The vast majority of the red dwarfs found in the open star clusters and association show flare activity (Mirzoyan 1990; Pigatto 1990), while the number of UV Ceti stars in the clusters decrease from a young cluster to the older one (Skumanich 1972; Marcy & Chen 1992; Pettersen 1991; Stauffer 1991). The flare activity results in mass loss, which plays an important role in stellar evolution.

Although various studies have been carried out since the first flare was observed by R. C. Carrington and R. Hodgson on 1859 September 1 (Carrington 1859; Hodgson 1859), it is not fully understood what the flare and its process are. However, the flare activity of the dMe stars is modelled on the basis of the Solar Flare Event (Gershberg 2005; Benz & Güdel 2010). The studies, which have been continuing to understand the flare events of dMe stars have revealed that there are some differences between the flare energy levels of stars (Gershberg & Shakhovskaya 1983).

The level of highest energy releasing on the Sun are about  $10^{30}$ – $10^{31}$  erg in the flare events (Gershberg 2005; Benz 2008). It seems that this level is about  $10^{31}$  erg (Haisch,

Strong, & Rodonó 1991), if RS CVn stars, the chromospheric active binaries, are considered. The observations lasting over decades show that the energy levels of flares occurring on dMe stars could have increased from 10<sup>28</sup> to 10<sup>34</sup> erg (Gershberg 2005). Considering the stars in the Pleiades cluster and Orion association, it is seen that these values have reached 10<sup>36</sup> erg (Gershberg & Shakhovskaya 1983). There are significant differences between the energy level of flare stars from different spectral types. For example, it is well known that there are serious differences between the mass loss rates and the flare energy level, if the Sun is compared to a dMe star. Nevertheless, the flare events occurring on the dMe stars are tried to be explained with the Solar Flare processes. As a result, it is clear that the flares in the different stars should be well studied and the similarities and differences between them should be analysed. Gershberg (2005) and Hudson & Khan (1997) have suggested that magnetic reconnection processes are the source of energy for flare events. In order to determine the important points in the flare process, it is necessary that the cause of the differences in the flare frequency and also the flare energy spectra should be determined.

Besides Gershberg & Shakhovskaya (1983), there are several studies about the flare energies in the literature such

as Gershberg et al. (1972), Lacy, Moffett, & Evans (1976), Walker (1981), Pettersen, Coleman, & Evans (1984), and Mavridis & Avgoloupis (1986). However, according to Dal & Evren (2010, 2011), the results obtained in these studies are unsatisfactory for comparing the stars from different spectral types. For example, Gershberg et al. (1972) derived the flare energy spectra for several dMe stars, such as AD Leo, EV Lac, UV Cet, and YZ CMi, while Gershberg & Shakhovskaya (1983) derived the flare energy spectra for lots of stars from the galactic field to compare with the flare energy spectra of some stars from Pleiades and Orion association. They showed that these stars are located in different points in the distribution of energy. They correctly indicated that this distribution is caused due to different ages of the stars. On the other hand, Dal & Evren (2010, 2011) indicated that there is one more reason to cause this distribution. The second reason is including the luminosity parameter in calculations, of flare energies, which leads to incorrect results, as this caused the stratifying in the flare energy spectra. Indeed, the studies of Dal & Evren (2010, 2011) show that the flare equivalent durations in the logarithmic scale vary within a certain rule versus the flare total time, and it depends on the spectral type of the stars, on which the flares occur, when the relations between the flare parameters are examined. On the other hand, two different flare frequencies were defined by Ishida et al. (1991) for the flare activity. The frequency of  $N_1$  indicates the number of flares per an hour, while the frequency of  $N_2$  describes the flare-equivalent duration emitting per an hour. Leto et al. (1997) clearly show that the frequency of EV Lac's flares vary over time.

On the other hand, for the first time in the literature, Kron (1952) discovered that UV Ceti-type stars also exhibited spot activity known as BY Dra Syndrome. Kron (1952) observed sinusoidal variations out-of-eclipses in the light curves of the YY Gem, which is a binary system, and Kunkel (1975) called this phenomenon BY Dra Syndrome and explained as a fact that these variations were caused by heterogeneous temperature distribution on the surface of the star. In the case of the Sun, Berdyugina & Järvinen (2005) found two stable active longitudes separated by 180° from each other, and they indicated that these longitudes are exhibiting semi-rigid behaviour. However, according to some authors, these longitudes are not the persistent active structures, which show variation in the time (Lopez 1961; Stanek 1972; Bogart 1982). The difference between the regular activity oscillations shown by these longitudes, called the Flip-Flop, is very important for the north-south asymmetry exhibited by the star's magnetic topology. It is very important to calculate the angular velocities of these longitudes because these calculations light on the differential rotational velocities in the latitudes of the spots and spots groups.

In the case of binary or single stars, determining the parameters of stellar spot activity, such as spot latitude, radius, and longitude, is a controversial phenomenon. In the literature, there are several methods to find out these parameters (Ribárik 2002; Ribárik, Oláh, & Strassmeier 2003; Walkow-

**Table 1.** Physical parameters of KOI-256 taken from the literature.

P <sub>orb</sub>	(day)	1.3786548 (a)
а	(AU)	$0.0250  (^a)$
i	(°)	89.3 ( <sup>b</sup> )
(B-V)	( <sup>m</sup> )	1.42 (°)
$R_1$	$(R_{\odot})$	$0.540(^a)$
$R_2$	$(R_{\odot})$	0.01345 (a)
$T_1$	(K)	3 450 ( <sup>a</sup> )
$T_2$	(K)	7 100 ( <sup>a</sup> )
$M_1$	$({ m M}_{\odot})$	$0.51  (^a)$
$M_2$	$({ m M}_{\odot})$	$0.592  (^a)$
[M/H]	_	$0.31  (^a)$
d	(pc)	$155  (^d)$
J, H, K	$\binom{m}{}$	12.701-12.000-11.782 (e)
Spectral type	_	$M3 V + WD (^d)$

- (a) Muirhead et al. (2013).
- (b) Slawson et al. (2011).
- (c) Walkowicz & Basri (2013b).
- (d) Muirhead et al. (2014).
- (e) Zacharias et al. (2004).

icz, Basri, & Valenti 2013a; Jeffers & Keller 2009). For example, the SPOTMODEL program (Ribárik 2002; Ribárik et al. 2003) required two band observation and the inclination of the rotation axis to be able to model the distribution of the spots on the stellar surface. However, it must be noted that the system's maximum brightness level ever observed has also a very important role to determine the spot radius depending on spot latitude. According to Walkowicz et al. (2013a) and Jeffers & Keller (2009), this method does not work for the observations such as the data acquired in the Kepler Mission. Because the Kepler observations provide us monochromatic data, which have been detrended while combining different observation parts.

In this study, we examine the flare and spot activity exhibiting by KOI-256, and examine the parameters derived from the One Phase Exponential Association (hereafter OPEA) model, and also the spot migration considering the sinusoidal variation out-of-the flare. It is seen from the literature that KOI-256 is a system, which the primary component is a white dwarf, the secondary component is a main-sequence star from the spectral type of M3. In addition, there are some clues about third body, which could be a planet. Borucki et al. (2011) indicated that KOI-256 is a planet candidate system. There are lots of studies such as Ritter & Kolb (2003), Borucki et al. (2011), Slawson et al. (2011), Walkowicz & Basri (2013b), Muirhead et al. (2013), Muirhead et al. (2014), and Zacharias et al. (2004) in the literature. Several physical parameters of KOI-256 were computed in these studies, using some colour calibrations explained by these authors. As it can be seen from these parameters listed in Table 1, KOI-256 is very interesting system. Although one of the components is with a dwarf, Walkowicz & Basri (2013b) indicated that the system's age is 0.01 Gyr. However, using the Equation (2.3) given by Gänsicke (1997) with the stellar parameters taken from Muirhead et al. (2013), we estimated the age of the system as  $\sim$ 2 Gyr, considering the cooling of the white dwarf component. Indeed, Muirhead et al. (2013) laterly revealed that the system is a post-common envelope binary. In the paper, the stellar spot activity analyses are described in Section 2.1, while the flare models are described in Section 2.2. The orbital period variation analysis is explained in Section 2.3. The results obtained from the analyses are summarised and discussed in Section 3.

### 2 DATA AND ANALYSES

The Kepler Mission project, which was launched to explore the exoplanets, has observed more than 150.000 sources (Borucki et al. 2010; Koch et al. 2010; Caldwell, Kolodziejczak, & Van Cleve 2010). These observations are one of the highest sensitivity photometric observations ever achieved (Jenkins et al. 2010a, 2010b). With this highest sensitivity of observation, a large number of new eclipsing binaries and lots of new variable stars have been discovered besides the exoplanets (Slawson et al. 2011; Matijevič et al. 2012). Some of these variable stars are single, and some of them are binary stars, which exhibit both the stellar spot activity and flare activity (Balaji et al. 2015). In this study, observational data of KOI-256, which is one of these binary systems, was taken from Kepler database (Slawson et al. 2011; Matijevič et al. 2012). KOI-256's short cadence data obtained in the Kepler Mission covered the time range from BJD 24 55372.460219 to BJD 24 55552.55836 and from BJD 24 56419.80351 to BJD 24 56424.01160, while the long cadence data of the system covered time ranges BJD 24 54964.51238-55206.21898, BJD 24 55276.51124-55552.54849, BJD 24 55641.51645-55931.30552, BJD 24 56015.77828-56304.13695, BJD 24 56392.24699–56424.00173. All the available data presented in both long and short cadence formats are shown in Figure 1. The short cadence data were used in the analyses of the flares, while the long cadence data were used for the analyses of sinusoidal variation. The detrended data was used among the public data provided in the Kepler database. The data sets were created in appropriate formats, which were edited in the analysis processes for the flare activity, the stellar spot activity, and the orbital period variation (O-C).

The whole of KOI-256's analysed data taken from the Kepler database is shown in the upper and middle panels of Figure 1. The light curve is plotted versus the Barycentric Julian Date in the upper panel, while it is plotted versus phase in the middle panel of the figure. In the lower panel, the light curve is plotted versus phase, expanding the *y*-axis to be easily seen the sinusoidal variations out-of-the dominant flare activity.

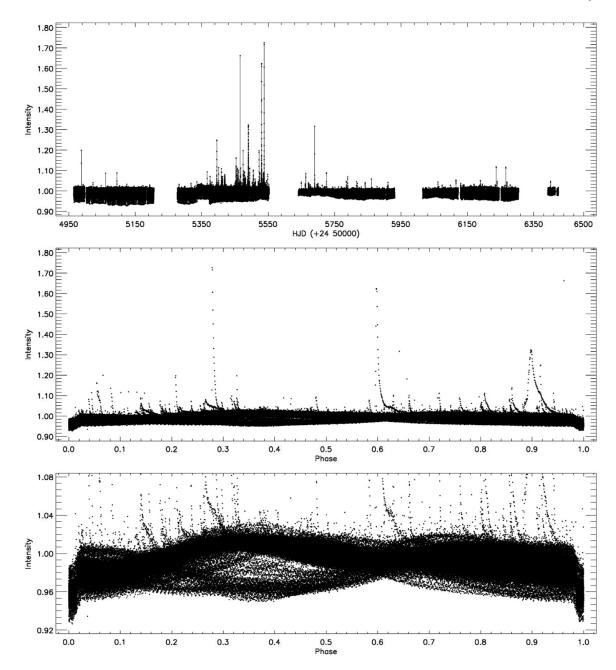
# 2.1. Stellar spot activity

Examining the out-of-eclipses light variations, it is seen that the system exhibits also sinusoidal variations. Considering the surface temperatures of the components of the system, it is understood that the variations is caused by the rotational modulation of the cool stellar spots. It is seen that both minima times and amplitudes of the light curves are varying once in a few cycles, when the consecutive cycles depend on the orbital period are examined in the light curves, in which both the flares and eclipses are removed. This situation indicates that the active regions on the components of the system evolved rapidly. Because of this, it is not possible to model the entire light curve in a single analysis, so the data of sinusoidal variation are separated into several sets. When dividing into data sets, consecutive cycles in which some characters of the light curve such as a light curve asymmetry, a spot minima phase, and the minima and maxima level were the same were collected in a single set. In this format, 138 subsets are obtained, thus, each subset is modelled separately.

To be able to model the sinusoidal variations, the prewhitened light curves were obtained. In this step, we first removed data parts, in which all instant light increase due to the flares are seen, from whole data. In addition, the data parts, in which the primary minima are seen due to eclipses, were also removed from whole data.

The pre-whitened light curves are modelled with the SPOT-MODEL program (Ribárik 2002; Ribárik et al. 2003) to find out the spot distribution parameters, such as the spot radius, latitude, and especially longitudinal distribution, on the stars. To model the spots, the SPOTMODEL program requires two-band observations or a temperature factor (kw = $[T_{\rm spot}/T_{\rm stellar surface}]^2)$  for the stellar spot. However, the data analysed in this study consist of monochromatic observations that are presented publicly in the Kepler database. In this point, considering both the study of Botsula (1978), which first revealed the spot temperature factor for the stellar activity, and also the light curve analysis obtained from analogue systems (Clausen, Helt, & Olsen 2001; Thomas & Weiss 2008), it was assumed that the secondary component exhibits magnetic activity. The inclination of rotation axis is taken 89.1° as it was given by Muirhead et al. (2013). Then, taking the different values from kw = 0.60 to kw = 0.90 for the spot temperature factor in agreement with the values found by these studies for analogue systems, the first few sets were tried to be modelled. As a result, it was seen that the best solution can be obtained by taking the spot temperature factor of kw =0.65 for both spots. For this reason, the spot temperature factors were taken as kw = 0.65 for both spots in the all sets. Taking the spot temperature factor as constant value in the models for each set, the longitudinal, latitudinal distributions, and radius variations of the spot area on the active component were determined.

In the analysis, the parameters such as the longitudes (l), latitudes (b), and radii of the spots (g) parameters were taken as free parameters. The parameters obtained from the models for each set are listed in Table 2, while six examples for the derived models are shown in Figure 2. In the left panel of the figure, the observations and models are plotted versus Heliocentric Julian Day as a time, while they were plotted



**Figure 1.** All of the long cadence and short cadence data taken from the Kepler database for the KOI-256 are shown. The light variation was plotted in the plane of intensity taken from database as detrended form versus the phase computed by using orbital period.

versus phase computed using by orbital period. As seen from the figure, the synthetic models absolutely fit the observations. The variations of latitude (b) and radius (g) for both spots are shown in Figures 3. In addition, the most important parameter of the models, longitude (l), is plotted versus time in Figure 4. The longitudinal variations were fitted by a linear function for both spots. Using these linear fits, the migration periods were computed. The migration period was found to be 3.95 yrs for first spotted area and 8.37 yrs for second area.

To test whether the findings about longitudinal variations are close to the real nature of the stellar surface, it was tested

by another method. Using the Fourier Transform, the minima times of sinusoidal variations, where the amplitude is larger, were computed. Then, the orbital phases called as  $\theta_{\rm min}$  were computed for these sinusoidal minima times. In Figure 5, the variation of  $\theta_{\rm min}$  are plotted versus time. This variation was fitted by a linear function similar to the longitudinal variations. In the same process, the migration period was computed from the phase shift of the  $\theta_{\rm min}$  using this linear function. In a result, the period of the spotted area migration was found to be 9.126 yrs. As it is expected, this value is in agreement with one of the migration periods found from the longitudinal variations.

 Table 2. The Spot parameters obtained with SPOTMODEL.

JD	b <sub>1</sub> (Latitude)	b <sub>2</sub> (Latitude)	g <sub>1</sub> (Radius)	g <sub>2</sub> (Radius)	l <sub>1</sub> (Longitude)	l <sub>2</sub> (Longitude)
54966.749927	$63.277 \pm 3.430$	$79.015 \pm 0.252$	$15.917 \pm 1.808$	$31.179 \pm 0.410$	$77.848 \pm 1.494$	$378.418 \pm 2.961$
54971.061541	$57.628 \pm 4.704$	$78.572 \pm 0.234$	$14.244 \pm 1.643$	$30.481 \pm 0.313$	$79.038 \pm 1.083$	$342.094 \pm 2.181$
54975.230111	$64.687 \pm 2.303$	$77.990 \pm 0.323$	$17.743 \pm 1.498$	$29.433 \pm 0.492$	$75.469 \pm 1.578$	$336.431 \pm 3.112$
54979.960614	$64.074 \pm 2.669$	$78.130 \pm 0.250$	$16.693 \pm 1.533$	$30.294 \pm 0.401$	$73.857 \pm 1.447$	$336.859 \pm 2.638$
54984.742199	$56.901 \pm 5.588$	$79.318 \pm 0.190$	$13.122 \pm 1.723$	$31.410 \pm 0.259$	$80.089 \pm 1.018$	$343.844 \pm 1.955$
54988.992489	$56.368 \pm 4.173$	$78.341 \pm 0.173$	$13.458 \pm 1.272$	$31.008 \pm 0.204$	$82.514 \pm 0.806$	$340.931 \pm 1.367$
54993.058867	$62.702 \pm 3.014$	$78.436 \pm 0.211$	$15.332 \pm 1.407$	$30.875 \pm 0.289$	$78.478 \pm 1.145$	$338.160 \pm 1.944$
54996.491784	$56.142 \pm 6.666$	$78.138 \pm 0.259$	$13.270 \pm 2.007$	$30.898 \pm 0.325$	$77.494 \pm 1.324$	$337.974 \pm 2.173$
55004.481471	$65.054 \pm 3.030$	$79.779 \pm 0.218$	$16.620 \pm 1.904$	$32.811 \pm 0.426$	$72.950 \pm 1.555$	$338.632 \pm 3.309$
55007.985892	$64.527 \pm 2.464$ $61.244 \pm 5.076$	$78.929 \pm 0.234$ $78.986 \pm 0.349$	$17.507 \pm 1.559$	$31.759 \pm 0.409$	$69.841 \pm 1.448$	$332.064 \pm 2.888$
55011.061195 55013.492827	$59.668 \pm 4.704$	$77.560 \pm 0.349$	$16.060 \pm 2.370$ $16.324 \pm 1.498$	$32.337 \pm 0.489$ $32.693 \pm 0.289$	$72.068 \pm 1.765$ $73.878 \pm 3.926$	$331.181 \pm 3.486$ $331.057 \pm 2.571$
55017.640899	$66.308 \pm 2.303$	$77.999 \pm 0.759$	$17.122 \pm 1.533$	$31.309 \pm 0.325$	$69.324 \pm 9.060$	$332.951 \pm 3.291$
55019.949921	$65.718 \pm 2.988$	$80.074 \pm 0.234$	$17.122 \pm 1.933$ $17.455 \pm 2.071$	$33.626 \pm 0.482$	$66.448 \pm 1.757$	$330.176 \pm 3.845$
55022.422411	$61.803 \pm 6.938$	$79.715 \pm 0.316$	$15.211 \pm 3.222$	$33.209 \pm 0.581$	$67.316 \pm 2.231$	$335.756 \pm 4.538$
55024.496438	$64.030 \pm 4.381$	$79.073 \pm 0.383$	$16.711 \pm 2.566$	$32.067 \pm 0.613$	$67.204 \pm 2.127$	$328.939 \pm 4.392$
55027.520633	$62.317 \pm 3.422$	$78.843 \pm 0.207$	$15.440 \pm 1.623$	$31.722 \pm 0.325$	$68.190 \pm 1.271$	$331.344 \pm 2.327$
55032.475806	$74.406 \pm 2.185$	$76.669 \pm 0.207$	$26.567 \pm 3.707$	$26.866 \pm 3.598$	$47.323 \pm 1.324$	$301.919 \pm 2.630$
55036.991646	$64.178 \pm 3.007$	$78.693 \pm 0.272$	$16.326 \pm 1.647$	$32.266 \pm 0.363$	$68.294 \pm 1.406$	$324.269 \pm 2.377$
55036.991646	$64.124 \pm 3.020$	$78.697 \pm 0.271$	$16.296 \pm 1.645$	$32.273 \pm 0.361$	$68.314 \pm 1.399$	$324.309 \pm 2.366$
55044.051454	$64.961 \pm 2.458$	$78.258 \pm 0.285$	$17.635 \pm 1.564$	$32.375 \pm 0.403$	$65.672 \pm 1.476$	$321.194 \pm 2.559$
55047.525157	$61.982 \pm 3.970$	$78.775 \pm 0.290$	$15.990 \pm 1.888$	$32.788 \pm 0.375$	$68.351 \pm 1.422$	$323.413 \pm 2.575$
55050.304115	$61.227 \pm 4.059$	$79.092 \pm 0.294$	$15.909 \pm 1.874$	$32.386 \pm 0.386$	$64.700 \pm 1.411$	$320.649 \pm 2.736$
55053.767589	$62.461 \pm 2.430$	$78.973 \pm 0.210$	$16.465 \pm 1.245$	$32.507 \pm 0.276$	$66.362 \pm 0.981$	$321.805 \pm 1.907$
55057.874710	$59.698 \pm 3.869$	$79.232 \pm 0.199$	$15.005 \pm 1.542$	$33.013 \pm 0.265$	$66.069 \pm 1.045$	$324.048 \pm 1.938$
55061.634456	$63.015 \pm 2.974$	$78.326 \pm 0.269$	$17.080 \pm 1.633$	$32.120 \pm 0.393$	$62.205 \pm 1.452$	$321.125 \pm 2.605$
55065.230727	$35.995 \pm 2.974$	$77.523 \pm 0.798$	$9.311 \pm 1.633$	$30.464 \pm 0.826$	$61.544 \pm 3.690$	$328.098 \pm 4.977$
55067.182111	$75.691 \pm 1.294$	$69.690 \pm 5.531$	$31.217 \pm 1.982$	$19.782 \pm 5.433$	$24.124 \pm 9.473$	$279.457 \pm 2.929$
55070.277756	$63.922 \pm 3.508$	$78.844 \pm 0.169$	$15.607 \pm 1.909$	$33.771 \pm 0.352$	$56.672 \pm 1.575$	$325.965 \pm 2.571$
55074.425711	$66.321 \pm 3.445$	$78.750 \pm 0.187$	$16.152 \pm 2.283$	$33.518 \pm 0.458$	$53.919 \pm 2.210$	$324.187 \pm 3.291$
55077.848279	$67.671 \pm 4.066$	$77.458 \pm 0.357$	$17.269 \pm 3.311$	$32.052 \pm 0.865$	$50.804 \pm 1.324$	$321.033 \pm 5.382$
55083.824984 55084.734264	$68.052 \pm 2.600$ $66.361 \pm 5.317$	$77.793 \pm 0.203$ $77.680 \pm 0.478$	$17.142 \pm 2.122$ $17.179 \pm 3.827$	$33.066 \pm 0.505$ $32.598 \pm 0.952$	$53.435 \pm 2.540$ $52.764 \pm 1.324$	$322.217 \pm 3.232$ $318.453 \pm 2.377$
55087.278192	$71.806 \pm 3.214$	$76.954 \pm 0.719$	$21.685 \pm 4.859$	$31.254 \pm 2.171$	$46.173 \pm 1.324$	$313.208 \pm 2.366$
55094.102831	$61.631 \pm 6.295$	$75.903 \pm 0.504$	$15.308 \pm 2.932$	$31.042 \pm 0.669$	$55.341 \pm 2.659$	$320.091 \pm 3.583$
55096.421985	$63.947 \pm 4.856$	$76.563 \pm 0.390$	$15.843 \pm 2.604$	$32.215 \pm 0.557$	$57.951 \pm 2.617$	$318.779 \pm 3.213$
55100.978555	$57.876 \pm 5.887$	$77.649 \pm 0.170$	$12.960 \pm 1.872$	$34.039 \pm 0.236$	$62.181 \pm 1.227$	$324.958 \pm 1.532$
55105.453390	$58.804 \pm 5.542$	$76.923 \pm 0.338$	$15.505 \pm 2.264$	$32.774 \pm 0.450$	$53.787 \pm 1.688$	$314.621 \pm 2.377$
55108.089248	$73.020 \pm 2.975$	$74.902 \pm 0.241$	$27.145 \pm 5.195$	$27.606 \pm 4.971$	$32.067 \pm 1.324$	$291.603 \pm 2.366$
55111.430042	$58.736 \pm 3.466$	$75.829 \pm 0.241$	$15.596 \pm 1.416$	$31.835 \pm 0.305$	$51.917 \pm 1.171$	$313.049 \pm 1.653$
55116.763048	$63.599 \pm 2.688$	$76.973 \pm 0.264$	$17.157 \pm 1.592$	$32.509 \pm 0.396$	$52.047 \pm 1.535$	$312.718 \pm 2.365$
55120.236651	$67.100 \pm 2.534$	$77.592 \pm 0.382$	$19.611 \pm 2.222$	$32.561 \pm 0.719$	$47.943 \pm 2.692$	$308.686 \pm 4.336$
55122.484275	$65.708 \pm 2.534$	$78.011 \pm 0.594$	$17.748 \pm 4.261$	$33.933 \pm 1.028$	$54.446 \pm 4.318$	$314.813 \pm 2.782$
55125.610515	$68.529 \pm 2.427$	$78.667 \pm 0.360$	$20.282 \pm 2.677$	$33.143 \pm 0.921$	$47.553 \pm 3.319$	$312.377 \pm 2.022$
55128.225931	$67.447 \pm 3.104$	$78.347 \pm 0.393$	$19.391 \pm 2.933$	$32.574 \pm 0.945$	$48.032 \pm 3.430$	$313.894 \pm 2.197$
55130.514420	$61.196 \pm 6.095$	$75.843 \pm 0.615$	$16.127 \pm 2.980$	$30.903 \pm 0.782$	$59.660 \pm 2.665$	$322.497 \pm 4.008$
55132.710959	$67.228 \pm 2.467$	$78.319 \pm 0.351$	$18.628 \pm 2.028$	$32.079 \pm 0.614$	$49.971 \pm 2.332$	$309.434 \pm 3.950$
55140.403958	$68.362 \pm 1.572$	$78.515 \pm 0.253$	$19.713 \pm 1.621$	$31.138 \pm 0.598$	$42.595 \pm 2.093$	$307.516 \pm 3.901$
55150.589779	$68.603 \pm 2.413$	$79.115 \pm 0.203$	$18.488 \pm 2.404$	$31.868 \pm 0.727$	$42.700 \pm 2.881$	$314.182 \pm 2.219$
55158.415607	$67.166 \pm 3.511$	$77.979 \pm 0.500$	$19.066 \pm 3.177$	$30.642 \pm 1.132$	$41.285 \pm 3.827$	$308.543 \pm 2.145$
55161.858566	$73.557 \pm 1.875$	$74.930 \pm 0.500$	$28.440 \pm 2.855$	$24.497 \pm 0.289$	$22.522 \pm 2.332$	$273.844 \pm 2.366$
55165.352611	$70.567 \pm 3.037$	$78.492 \pm 0.562$	$22.987 \pm 4.487$	$30.529 \pm 0.325$	$32.892 \pm 2.093$	$295.508 \pm 1.653$
55169.888745	$68.733 \pm 3.314$	$78.484 \pm 0.562$	$20.493 \pm 3.744$	$31.984 \pm 1.419$	$36.448 \pm 1.324$	$301.355 \pm 2.571$
55173.862975	$68.021 \pm 4.607$	$78.484 \pm 0.715$	$19.730 \pm 4.602$	$32.198 \pm 1.587$ $32.093 \pm 0.767$	$35.569 \pm 1.324$ $38.022 \pm 2.685$	$299.308 \pm 3.291$
55176.887073 55180.493518	$65.779 \pm 3.268$ $63.531 \pm 4.174$	$77.793 \pm 0.463$ $77.176 \pm 0.476$	$18.602 \pm 2.495$ $17.253 \pm 2.451$	$32.093 \pm 0.767$ $31.224 \pm 0.674$	$38.022 \pm 2.085$ $37.329 \pm 2.391$	$298.566 \pm 4.758$ $297.495 \pm 3.965$
55187.440787	$63.331 \pm 4.174$ $68.403 \pm 3.204$	$77.176 \pm 0.476$ $78.177 \pm 0.544$	$17.253 \pm 2.451$ $19.051 \pm 2.973$	$31.224 \pm 0.674$ $32.592 \pm 0.932$	$37.329 \pm 2.391$ $37.112 \pm 1.147$	$297.493 \pm 3.963$ $295.397 \pm 2.693$
55191.159632	$69.524 \pm 3.215$	$78.177 \pm 0.344$ $78.457 \pm 0.610$	$20.404 \pm 3.621$	$32.092 \pm 0.932$ $32.022 \pm 1.342$	$37.112 \pm 1.147$ $29.811 \pm 1.449$	$293.397 \pm 2.093$ $290.715 \pm 2.285$
55194.786534	$65.594 \pm 4.560$	$78.938 \pm 0.273$	$15.563 \pm 2.677$	$32.022 \pm 1.342$ $33.146 \pm 0.509$	$37.822 \pm 2.447$	$302.929 \pm 3.696$
55198.556476	$67.491 \pm 3.688$	$79.212 \pm 0.268$	$16.620 \pm 2.751$	$32.992 \pm 0.616$	$36.101 \pm 2.754$	$302.929 \pm 3.090$ $303.617 \pm 2.496$
55203.368529	$66.201 \pm 3.404$	$79.536 \pm 0.149$	$15.008 \pm 1.965$	$32.592 \pm 0.010$ $33.587 \pm 0.326$	$38.443 \pm 1.768$	$305.854 \pm 2.577$
55278.564869	$69.191 \pm 2.174$	$77.595 \pm 0.736$	$21.329 \pm 2.214$	$30.962 \pm 0.977$	$31.706 \pm 3.075$	$279.055 \pm 2.056$
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Table 2. Continued.

JD	$b_1$ (Latitude)	b <sub>2</sub> (Latitude)	g <sub>1</sub> (Radius)	g <sub>2</sub> (Radius)	$l_1$ (Longitude)	$l_2$ (Longitude)
55286.115289	$68.067 \pm 2.204$	$78.100 \pm 0.702$	$21.246 \pm 2.111$	$31.863 \pm 0.875$	$33.457 \pm 2.514$	$278.883 \pm 2.760$
55290.477987	$67.725 \pm 2.826$	$77.905 \pm 0.809$	$21.153 \pm 2.580$	$32.151 \pm 1.016$	$31.765 \pm 3.020$	$277.276 \pm 2.495$
55294.299182	$66.081 \pm 2.682$	$77.905 \pm 0.629$	$20.212 \pm 2.098$	$32.634 \pm 0.728$	$29.837 \pm 2.170$	$274.620 \pm 4.009$
55297.916038	$64.709 \pm 3.003$	$78.589 \pm 0.350$	$17.892 \pm 1.877$	$33.916 \pm 0.445$	$34.050 \pm 1.583$	$283.230 \pm 2.813$
55301.829192	$59.677 \pm 3.514$	$78.421 \pm 0.271$	$16.143 \pm 1.513$	$33.575 \pm 0.300$	$33.796 \pm 1.050$	$282.969 \pm 1.885$
55305.967125	$63.953 \pm 2.074$	$78.028 \pm 0.303$	$18.464 \pm 1.283$	$33.333 \pm 0.344$	$30.741 \pm 1.147$	$277.405 \pm 1.998$
55311.729577	$64.647 \pm 2.103$	$77.422 \pm 0.488$	$20.267 \pm 1.502$	$32.553 \pm 0.528$	$29.781 \pm 1.449$	$272.577 \pm 2.740$
55314.723194	$65.233 \pm 1.989$	$77.395 \pm 0.588$	$21.063 \pm 1.526$	$31.947 \pm 0.610$	$28.626 \pm 1.609$	$268.837 \pm 3.056$
55318.248101	$64.094 \pm 2.395$	$78.449 \pm 0.471$	$19.721 \pm 1.567$	$32.508 \pm 0.499$	$29.410 \pm 1.307$	$270.048 \pm 2.801$
55321.507363	$67.747 \pm 2.158$	$77.049 \pm 0.830$	$22.326 \pm 1.986$	$31.277 \pm 0.930$	$23.642 \pm 2.585$	$263.668 \pm 2.343$
55324.327288	$63.231 \pm 1.828$	$77.392 \pm 0.361$	$19.259 \pm 1.119$	$32.241 \pm 0.360$	$30.961 \pm 1.027$	$270.101 \pm 1.825$
55327.893061	$66.778 \pm 2.175$	$78.474 \pm 0.570$	$20.473 \pm 1.753$	$32.267 \pm 0.632$	$28.538 \pm 1.718$	$269.419 \pm 3.508$
55331.346445	$65.131 \pm 1.995$	$77.933 \pm 0.491$	$19.913 \pm 1.409$	$31.583 \pm 0.504$	$27.701 \pm 1.431$	$267.882 \pm 2.688$
55334.850912	$61.023 \pm 3.036$	$77.925 \pm 0.411$	$17.550 \pm 1.484$	$31.577 \pm 0.394$	$27.811 \pm 1.149$	$269.490 \pm 2.184$
55339.489476	$66.995 \pm 2.168$	$76.507 \pm 0.859$	$21.517 \pm 1.808$	$30.605 \pm 0.845$	$26.610 \pm 2.198$	$263.648 \pm 3.641$
55343.126760	$66.924 \pm 1.603$	$76.246 \pm 0.673$	$22.089 \pm 1.338$	$30.500 \pm 0.656$	$24.992 \pm 1.652$	$260.579 \pm 2.749$
55347.601842	$66.501 \pm 1.506$	$75.809 \pm 0.605$	$22.068 \pm 1.207$	$30.153 \pm 0.593$	$23.722 \pm 1.523$	$260.257 \pm 2.451$
55351.157383	$61.440 \pm 2.749$	$77.325 \pm 0.496$	$18.424 \pm 1.453$	$31.620 \pm 0.446$	$28.428 \pm 1.200$	$266.097 \pm 2.191$
55353.486873	$68.538 \pm 1.992$	$77.246 \pm 0.897$	$22.360 \pm 1.850$	$30.903 \pm 0.913$	$25.707 \pm 2.319$	$258.643 \pm 3.919$
55356.490687	$68.026 \pm 2.091$	$78.397 \pm 0.784$	$21.722 \pm 1.892$	$31.396 \pm 0.835$	$24.741 \pm 2.067$	$260.950 \pm 4.201$
55359.494498	$69.289 \pm 2.484$	$79.019 \pm 0.361$	$22.689 \pm 2.538$	$31.424 \pm 0.289$	$25.281 \pm 2.819$	$258.597 \pm 2.118$
55362.978506	$70.732 \pm 1.760$	$78.142 \pm 0.570$	$24.305 \pm 1.956$	$30.223 \pm 0.325$	$23.033 \pm 2.644$	$255.656 \pm 2.010$
55366.636200	$67.624 \pm 1.958$	$78.918 \pm 0.847$	$22.111 \pm 1.787$	$30.882 \pm 0.858$	$25.284 \pm 1.788$	$259.436 \pm 4.395$
55369.783038	$65.885 \pm 2.325$	$79.667 \pm 0.717$	$20.771 \pm 1.808$	$31.341 \pm 0.721$	$25.963 \pm 1.545$	$261.395 \pm 4.223$
55373.723379	$69.688 \pm 0.734$	$78.148 \pm 0.519$	$24.727 \pm 0.741$	$29.172 \pm 0.517$	$20.718 \pm 0.884$	$248.008 \pm 1.984$
55376.393757	$67.823 \pm 0.835$	$78.289 \pm 0.447$	$22.743 \pm 0.775$	$28.870 \pm 0.459$	$20.538 \pm 0.897$	$256.363 \pm 2.215$
55379.845729	$66.688 \pm 0.657$	$79.422 \pm 0.266$	$21.339 \pm 0.554$	$28.900 \pm 0.282$	$19.259 \pm 0.559$	$257.882 \pm 1.646$
55383.285777 55387.551685	$66.523 \pm 0.764$ $66.708 \pm 0.595$	$79.697 \pm 0.297$ $80.077 \pm 0.213$	$21.045 \pm 0.628$ $20.520 \pm 0.486$	$28.839 \pm 0.311$ $28.623 \pm 0.231$	$19.361 \pm 0.608$ $18.404 \pm 0.466$	$257.566 \pm 1.866$ $258.571 \pm 1.482$
55392.343077	$66.536 \pm 0.711$	$80.077 \pm 0.213$ $80.006 \pm 0.264$	$20.320 \pm 0.480$ $20.452 \pm 0.566$	$27.872 \pm 0.283$	$14.864 \pm 0.460$ $14.864 \pm 0.561$	$258.371 \pm 1.482$ $254.923 \pm 1.804$
55396.930800	$71.601 \pm 0.462$	$76.197 \pm 0.556$	$25.365 \pm 0.441$	$27.872 \pm 0.283$ $23.531 \pm 0.519$	$6.999 \pm 0.799$	$234.923 \pm 1.804$ $234.162 \pm 1.513$
55401.464365	$74.165 \pm 0.778$	$75.475 \pm 1.243$	$26.908 \pm 0.738$	$21.316 \pm 1.208$	$3.240 \pm 1.728$	$230.587 \pm 2.751$
55404.659853	$69.298 \pm 0.678$	$80.712 \pm 0.441$	$20.368 \pm 0.738$ $22.353 \pm 0.661$	$27.070 \pm 0.433$	$12.554 \pm 0.654$	$241.520 \pm 2.359$
55408.234381	$71.552 \pm 0.702$	$79.246 \pm 0.707$	$24.458 \pm 0.692$	$25.333 \pm 0.641$	$10.011 \pm 0.757$	$226.661 \pm 2.048$
55410.950691	$70.428 \pm 0.765$	$78.487 \pm 0.739$	$24.069 \pm 0.675$	$24.809 \pm 0.628$	$9.584 \pm 0.743$	$225.428 \pm 1.920$
55414.314060	$67.766 \pm 0.758$	$82.063 \pm 0.363$	$20.547 \pm 0.635$	$27.598 \pm 0.327$	$13.545 \pm 0.460$	$238.909 \pm 2.124$
55417.755753	$67.843 \pm 0.951$	$82.471 \pm 0.418$	$20.333 \pm 0.817$	$27.539 \pm 0.413$	$11.979 \pm 0.619$	$243.091 \pm 3.137$
55421.889452	$68.636 \pm 0.774$	$83.132 \pm 0.284$	$19.923 \pm 0.717$	$27.831 \pm 0.338$	$14.122 \pm 0.559$	$253.879 \pm 3.132$
55426.614352	$67.900 \pm 0.822$	$83.328 \pm 0.236$	$19.255 \pm 0.702$	$28.520 \pm 0.289$	$14.666 \pm 0.522$	$256.507 \pm 2.886$
55429.912649	$69.051 \pm 1.108$	$83.615 \pm 0.400$	$20.039 \pm 1.077$	$29.060 \pm 0.470$	$15.990 \pm 0.770$	$253.600 \pm 4.560$
55433.240206	$72.787 \pm 2.094$	$84.692 \pm 1.007$	$21.453 \pm 2.940$	$27.567 \pm 1.691$	$10.110 \pm 2.408$	$255.676 \pm 18.000$
55435.678588	$72.312 \pm 0.851$	$81.306 \pm 0.802$	$24.269 \pm 0.965$	$26.261 \pm 0.811$	$11.136 \pm 0.960$	$231.213 \pm 3.429$
55438.429596	$75.418 \pm 0.567$	$75.819 \pm 1.028$	$27.761 \pm 0.512$	$21.045 \pm 0.940$	$6.928 \pm 0.997$	$220.547 \pm 1.531$
55441.877717	$72.414 \pm 0.732$	$81.364 \pm 0.679$	$23.981 \pm 0.821$	$26.171 \pm 0.677$	$11.250 \pm 0.800$	$230.361 \pm 2.831$
55445.513138	$65.529 \pm 1.033$	$84.504 \pm 0.178$	$17.364 \pm 0.663$	$29.201 \pm 0.202$	$14.208 \pm 0.373$	$254.044 \pm 2.528$
55448.922765	$67.560 \pm 0.842$	$83.195 \pm 0.306$	$19.168 \pm 0.660$	$27.787 \pm 0.284$	$12.294 \pm 0.436$	$240.545 \pm 2.369$
55452.903836	$71.702 \pm 0.642$	$79.618 \pm 0.654$	$23.844 \pm 0.611$	$24.590 \pm 0.570$	$9.030 \pm 0.615$	$221.961 \pm 1.739$
55458.651004	$67.454 \pm 0.602$	$83.186 \pm 0.298$	$20.039 \pm 0.475$	$27.345 \pm 0.234$	$13.438 \pm 0.248$	$225.998 \pm 1.374$
55464.327988	$64.347 \pm 1.018$	$83.193 \pm 0.313$	$18.492 \pm 0.611$	$27.115 \pm 0.251$	$13.667 \pm 0.346$	$230.084 \pm 1.759$
55467.113359	$75.280 \pm 0.418$	$71.792 \pm 0.986$	$27.533 \pm 0.285$	$18.648 \pm 0.678$	$5.116 \pm 0.399$	$203.328 \pm 0.491$
55470.866573	$71.397 \pm 0.516$	$78.158 \pm 0.634$	$24.780 \pm 0.436$	$23.353 \pm 0.497$	$9.349 \pm 0.344$	$210.027 \pm 0.862$
55474.258800	$67.772 \pm 0.884$	$82.891 \pm 0.587$	$21.131 \pm 0.732$	$26.546 \pm 0.438$	$11.610 \pm 0.334$	$215.610 \pm 1.840$
55477.719474	$67.483 \pm 0.709$	$81.740 \pm 0.445$	$20.901 \pm 0.544$	$26.305 \pm 0.323$	$14.101 \pm 0.265$	$215.318 \pm 1.092$
55482.134025	$70.552 \pm 0.486$	$76.821 \pm 0.567$	$23.782 \pm 0.370$	$22.756 \pm 0.410$	$10.330 \pm 0.302$	$209.804 \pm 0.649$
55486.917386	$71.526 \pm 0.541$	$78.525 \pm 0.583$	$23.993 \pm 0.459$	$23.859 \pm 0.461$	$8.352 \pm 0.355$	$210.133 \pm 0.868$
55491.373477	$73.059 \pm 0.638$	$75.799 \pm 0.844$	$24.945 \pm 0.517$	$22.085 \pm 0.669$	$8.792 \pm 0.687$	$214.372 \pm 1.077$
55495.871451	$73.925 \pm 0.403$	$70.787 \pm 0.776$	$25.801 \pm 0.260$	$19.198 \pm 0.499$	$7.578 \pm 0.319$	$204.576 \pm 0.374$
55499.097473	$72.097 \pm 0.543$	$74.413 \pm 0.689$	$24.582 \pm 0.388$	$22.083 \pm 0.490$	$10.643 \pm 0.318$	$205.162 \pm 0.476$
55501.742858	$73.291 \pm 0.506$	$71.403 \pm 0.825$	$25.372 \pm 0.339$	$20.259 \pm 0.550$	$7.233 \pm 0.426$	$204.146 \pm 0.470$
55504.467590	$72.524 \pm 0.462$	$71.336 \pm 0.772$	$25.018 \pm 0.297$	$20.040 \pm 0.491$	$9.101 \pm 0.280$	$203.230 \pm 0.385$
55507.360894	$67.462 \pm 0.627$	$76.505 \pm 0.487$	$21.467 \pm 0.401$	$23.736 \pm 0.317$	$13.054 \pm 0.234$	$208.283 \pm 0.462$
55510.117978	$71.784 \pm 0.618$	$73.560 \pm 0.827$	$23.813 \pm 0.418$	$21.027 \pm 0.554$	$11.875 \pm 0.375$	$207.317\pm0.535$

Table 2. Continued.

JD	b <sub>1</sub> (Latitude)	b <sub>2</sub> (Latitude)	g <sub>1</sub> (Radius)	g <sub>2</sub> (Radius)	l <sub>1</sub> (Longitude)	l <sub>2</sub> (Longitude)
55512.869273	$70.697 \pm 0.476$	$72.657 \pm 0.551$	$23.350 \pm 0.307$	$21.827 \pm 0.355$	$11.878 \pm 0.250$	$205.546 \pm 0.331$
55516.326187	$69.021 \pm 0.592$	$77.486 \pm 0.453$	$21.440 \pm 0.416$	$24.709 \pm 0.303$	$13.343 \pm 0.212$	$205.681 \pm 0.400$
55519.665612	$69.159 \pm 0.778$	$77.485 \pm 0.562$	$21.107 \pm 0.540$	$24.592 \pm 0.380$	$13.768 \pm 0.289$	$205.216 \pm 0.448$
55522.058312	$71.307 \pm 0.671$	$73.982 \pm 0.728$	$22.541 \pm 0.450$	$22.173 \pm 0.476$	$12.652 \pm 0.295$	$202.894 \pm 0.441$
55524.814719	$71.323 \pm 0.815$	$73.845 \pm 0.836$	$21.754 \pm 0.529$	$21.647 \pm 0.535$	$12.264 \pm 0.379$	$203.441 \pm 0.472$
55527.451254	$72.219 \pm 0.551$	$74.986 \pm 0.485$	$21.968 \pm 0.381$	$23.235 \pm 0.334$	$15.422 \pm 0.205$	$201.570 \pm 0.232$
55530.126272	$78.627 \pm 0.245$	$-1.110 \pm 98.070$	$26.256 \pm 0.134$	$8.622 \pm 0.427$	$10.499 \pm 0.479$	$186.390 \pm 0.446$
55532.172637	$73.401 \pm 0.513$	$71.201 \pm 0.652$	$22.975 \pm 0.322$	$20.574 \pm 0.411$	$15.939 \pm 0.219$	$199.223 \pm 0.223$
55534.780911	$71.429 \pm 0.603$	$73.233 \pm 0.630$	$22.334 \pm 0.388$	$21.695 \pm 0.405$	$17.332 \pm 0.239$	$198.730 \pm 0.245$
55536.878020	$69.268 \pm 0.866$	$74.579 \pm 0.801$	$21.263 \pm 0.535$	$22.150 \pm 0.498$	$20.782 \pm 0.238$	$202.037 \pm 0.371$
55538.400281	$72.296 \pm 0.638$	$70.433 \pm 1.057$	$23.209 \pm 0.381$	$18.925 \pm 0.606$	$18.506 \pm 0.324$	$207.543 \pm 0.393$
55541.482607	$72.195 \pm 0.414$	$73.048 \pm 0.475$	$22.644 \pm 0.270$	$21.419 \pm 0.307$	$15.201 \pm 0.162$	$199.068 \pm 0.194$
55546.558864	$72.005 \pm 0.557$	$74.864 \pm 0.465$	$22.111 \pm 0.394$	$24.176 \pm 0.323$	$16.019 \pm 0.208$	$202.905 \pm 0.235$
55550.658428	$73.082 \pm 0.560$	$74.177 \pm 0.524$	$22.881 \pm 0.399$	$23.505 \pm 0.373$	$14.069 \pm 0.245$	$200.617 \pm 0.240$
56421.907557	$74.336 \pm 0.319$	$112.006\pm0.807$	$26.349 \pm 0.199$	$18.653 \pm 0.451$	$102.414\pm0.212$	$111.379\pm0.257$

## 2.2. Flare parameters and models

In order to understand the nature of the flare activity and to find out the flare behaviour of the system, we tried to determine the flares occurring on the active component. For the reason, using the synthetic light curves, all the other variations apart from the flares were removed from the entire light curve. Since the system is an eclipsing binary, all minima light variations between the phases of 0.04 around the minima points, where two components are external tangent to each other, were removed from the entire light curve. However, the separated observations due to the technical reasons were also removed from the data.

In order to specify the flare parameters, such as start and end points of a flare and its energy, it should be defined the quiet level of the light curve. However, it was seen some sinusoidal variations due to the rotational modulation exhibited by one of the components in addition to the flares. For this reason, the synthetic light curves were derived for the light variations out-of-the flares using the Fourier transform given by Equation (1), and all the data were modelled with these synthetic light curves for the all phases:

$$L(\theta) = A_0 + \sum_{i=1}^{N} A_i \cos(i\theta) + \sum_{i=1}^{N} B_i \sin(i\theta), \qquad (1)$$

here  $A_0$  is the zero point,  $\theta$  is the phase, while  $A_i$  and  $B_i$  are the amplitude parameters (Scargle 1982). The synthetic light curves derived by Equation (1) were assumed as the quiescent level for each flare. Two examples for the flares detected from the observations and their quiescent levels are shown in Figure 6. After modelling light variations out of both the eclipses and flares, the flare parameters are calculated, as it was previously described by Dal (2012) in detail.

The flare rise  $(T_r)$  and decay  $(T_d)$  times, the flare amplitudes and the equivalent durations (P) were calculated after the start and end points of a flare were determined. These flare parameters calculated for each flare detected from entire data

sets are listed in Table 3. In this table, the times of flare maxima in the first column, the flare equivalent durations (s), the flare rise times (s), the flare decay times (s), the flare total times (s), and the flare amplitudes in the last column were listed, respectively.

A total of 225 flares were detected from the available observational data of the system taken from the Kepler database. The equivalent duration values for each flare are calculated by Equation (2) given by Gershberg et al. (1972):

$$P = \int [(I_{\text{flare}} - I_0)/I_0] dt, \qquad (2)$$

where P is the flare equivalent duration in seconds,  $I_{\rm flare}$  is flux at the moment of a flare, and  $I_0$  is the flux of the system in the quiescent level, which were modelled by the Fourier method for the parts out-of-eclipses. Considering the reason explained by Dal & Evren (2010, 2011), which is mentioned in Section 1, the flare energies were not used in this study. In order to derive the models, the equivalent duration parameter was used instead of the flare energy. Using the orbital period of the system, the phases for each flare was calculated, and the phase distribution of these 225 flares is shown in Figure 7. In the figure, it is plotted the flare total number computed for each phase range of 0.10.

When the relationship between some flare parameters is examined, it can be seen that the flare equivalent duration varies versus the flare total time within a certain rule. In fact, as it had been demonstrated by Dal & Evren (2010, 2011), the regression calculations processed with the SPSS V17.0 (Green, Salkind, & Akey 1999) and GraphPad Prism V5.02 (Dawson & Trapp 2004) programs in this study show that the OPEA is the best function to model the distribution of the flare equivalent durations in the logarithmic scale versus the flare total time. The OPEA function is a special function that has a term 'Plateau' and is represented by Equation (3) (Motulsky 2007; Spanier & Oldham 1987):

$$y = y_0 + (\text{Plateau} - y_0) \times (1 - e^{-k \times x}),$$
 (3)

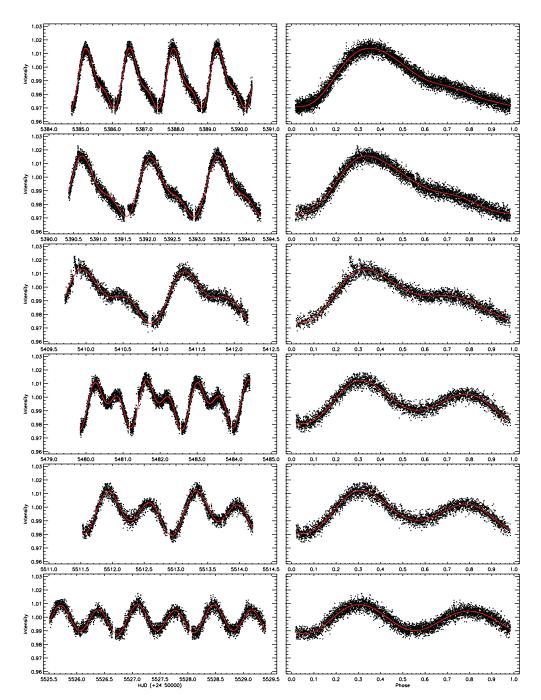
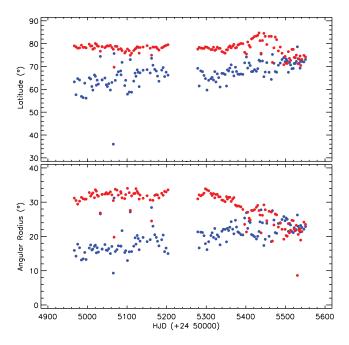


Figure 2. The observed light curve samples and their synthetic model fits derived by the SPOTMODEL analyses. In the figure, the filled circles represent the observations as intensity in detrended form, while the lines (red) represent the synthetic fits. In the left panels, the data and synthetic model are plotted versus time as Heliocentric Julian Day, while they are plotted versus phase computed using epoch and orbital period given in Equation (6) in the right panels.

where the parameter y is equivalent duration in the logarithmic scale, x is the total time of a flare, and  $y_0$  is the theoretical equivalent duration in the logarithmic scale for the minimum flare total time, as it had been previously defined by Dal & Evren (2010). In other words, the parameter  $y_0$  defines the minimum equivalent duration that someone can obtain for a flare occurring on the active component. Therefore,  $y_0$  value depends on the brightness of the observed target

and on the sensitivity of the used optical system. The Plateau value defines the upper limit of the equivalent durations for the flares observed on a particular star. According to Dal & Evren (2011), this parameter is defined as the saturation level for the flare activity in the observed wavelength range for the observed target.

Using the least-squares method, the distribution of the equivalent duration versus the flare total time was modelled



**Figure 3.** The variations of both spot latitude and spot radius parameters obtained with the SPOTMODEL program versus time. The filled blue circles represent the first spot and the filled red circles represent the second spot.

by the OPEA function for the flares. The obtained model is shown in Figure 8 with a 95% confidence interval. The computed parameters from the OPEA model are listed in Table 4. The span value listed in the table is the difference between Plateau and  $y_0$  values. The half-life value is half of the first x value, which the flare equivalent durations in the logarithmic scale reach the maximum level defined as the Plateau value. In other words, it is half value of the first total time, where the highest flare energy is seen for the flare.

In the Kepler Mission database, KOI-256's short cadence data are available for the time range from BJD 24 55372.460219 to BJD 24 55552.55836 and from BJD 24 56419.80351 to BJD 24 56424.01160. In total, the KOI-256 has been observed in short cadence format during 184.30624 d (4423.34976 h). From these data, 225 flares were obtained and their total equivalent durations were computed of 8169.834 s. In the literature, two different flare frequencies were defined by Gershberg et al. (1972). These frequencies are given by Equations (4) and (5):

$$N_1 = \Sigma n_{\rm f} / \Sigma T_{\rm t}, \tag{4}$$

$$N_2 = \Sigma P / \Sigma T_{\rm t},\tag{5}$$

where  $\Sigma n_{\rm f}$  is the total number of the obtained flare, while  $\Sigma T_{\rm t}$  defines the total observing time of the star.  $\Sigma P_{\rm u}$  is the sum of the equivalent durations of all detected flares. According to these definitions, the  $N_1$  frequency was found to be 0.05087 h<sup>-1</sup>, while the  $N_2$  frequency was found to be 0.00051.

Gershberg et al. (1972) described the flare frequency distribution separately calculated for the different energy limits of

the flares detected from a star, which defines the flare energy character for that star. However, using the flare equivalent duration instead of the flare energy parameter in this study, since the flare energy parameter depends on the quiet intensity level of the star in the observing band, the flare frequencies were calculated for different flare equivalent duration limits for the 225 flares. The obtained cumulative flare frequency distribution is shown in Figure 9.

# 2.3. Orbital period variation

The minima times were computed from the KOI-256s short cadence data from the first quarter to the quarter 17, which were taken from the publicly available Kepler database, without any corrections. The minima times were computed with a script depending on the method described by Kwee & van Woerden (1956). In this method, taking symmetrically increasing and decreasing parts of minima, the minima times were computed by fitting these parts with polynomial function. Before computing the minima times, the flares as the activity exhibited by the system were removed from the light curves. In the second step, (O - C)I residuals were determined for the obtained minima times. However, some of them include very large errors. It is seen that the flare activity occurred during these minima, when the light curves were examined for these minima times with large errors. Thus, these minima times were removed from the (O - C) data. As a result, 125 minima times were determined in the analyses. The obtained (O - C)I residuals were adjusted by the linear correction given in Equation (6):

$$JD(Hel.) = 2454965.55513(4) + 1^{d}.3786503(1) \times E.$$
 (6)

The computed minima times (O-C)I and the (O-C)II residuals obtained by applying a linear correction are listed in Table 5. In the table, the minima times, cycles, the minima type, (O-C)I, and (O-C)II residuals are listed, respectively. An interesting variation is seen in the variation of obtained (O-C)II residuals versus time in Figure 10.

First, checking the minima types given in the third column of Table 5, it was practically examined whether there is a marker about the existence of secondary minimum, which is the subject of discussion in the literature. Second, it was examined whether there was any separation in the (O-C)II residuals plotted in Figure 10. If there was any secondary minima, it would be expected a separation between the primary and secondary minima time residuals due to stellar spot activity, as it was demonstrated by Tran et al. (2013) and Balaji et al. (2015) for the first time. However, as it is seen from Figure 10, there is no decomposition in the minima time residuals as the primary and secondary minima.

### 3 RESULTS AND DISCUSSION

The analysis of KOI-256s data taken from the Kepler database (Slawson et al. 2011; Matijevič et al. 2012) indicates that the

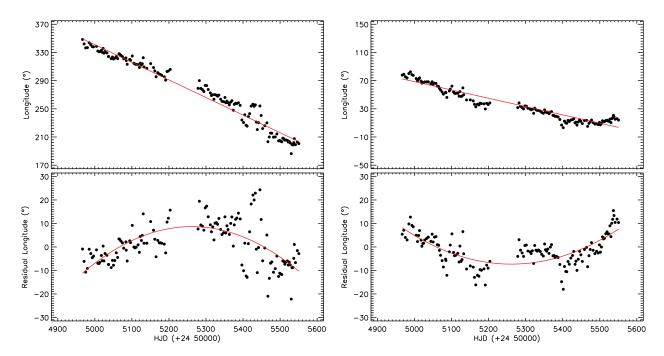
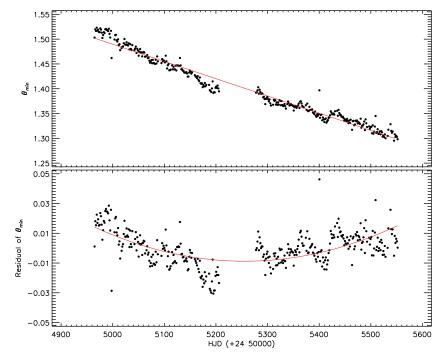


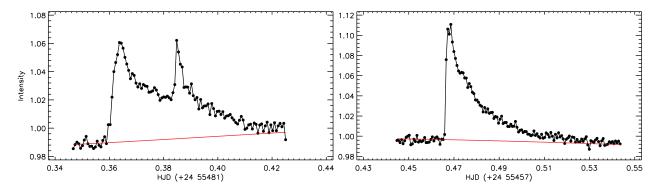
Figure 4. The variations of the first spot longitude as degree (left) and the second spot (right) are shown in the upper, while the residuals after the linear correction to the longitude variations are shown in the lower panel.



**Figure 5.** O-C diagram for the observed minima times of sinusoidal variation versus time in each data subset shown with  $\theta_{\min}$  term and its linear fit are shown in the upper panel. The filled circles represent the  $\theta_{\min}$  variation, while the line (red) show the linear fit. In the bottom panel, the residuals of  $\theta_{\min}$  is shown with its parabolic fit, which is plotted to show the trends clearly for the readers.

system has high chromospheric activity. However, it is necessary to compare the activity level between similar stars in order to reach a more definite result about the activity nature. In the literature, KOI-256 is an eclipsing binary system, with a white dwarf as a primary component and a main-sequence

star from M3 spectral type as a secondary component (Muirhead et al. 2014). The temperatures of the components were determined as  $T_1 = 7\,100\,\mathrm{K}$  and  $T_2 = 3\,450\,\mathrm{K}$  in the analyses of the spectral data together with the photometrical data of the system (Muirhead et al. 2014).



**Figure 6.** Two flare examples detected from the system. The light variation was plotted in the plane of intensity taken from database as detrended form versus time. The filled black circles represent the observations, while the red lines represent the quiescent level modelled by the Fourier method.

The eclipsing binary system KOI-256 was observed in the short cadence format from HJD 24 54964.51238 to HJD 24 56424.011602 in total of 184.30624 d (4423.34976 h). Within this study, 225 flares were determined and the parameters of each flare were calculated. Using the frequency descriptions defined by Gershberg et al. (1972), the flare frequencies for KOI-256 were calculated as  $N_1 = 0.05087 \text{ h}^{-1}$  and  $N_2 =$ 0.00051, respectively. The flare frequencies of KOI-256 were compared with the flare frequencies found from the young main-sequence dMe dwarfs, known as UV Ceti-type, exhibiting flare activity. The flare frequencies of KIC 9761199 were found to be  $N_1 = 0.01351 \text{ h}^{-1}$  and  $N_2 = 0.00006$  (Yoldaş & Dal 2016), while they were found to be  $N_1 = 0.01735 \,h^{-1}$  and  $N_2 = 0.00001$  for KIC 9641031 (Yoldaş & Dal 2017). KOI-256, when compared to these two systems, seems to have the highest flare frequency values among similar systems. On the other hand, the  $N_2$  frequencies were computed to be 0.088 for EQ Peg and 0.086 in the case of AD Leo, while the  $N_1$  frequencies were determined as 1.331 h<sup>-1</sup> for AD Leo, a UV Ceti type single star, and 1.056 h<sup>-1</sup> for EV Lac (Dal & Evren 2011). Compared to these single stars of the UV Ceti type, the flare frequencies of KOI-256 are found to be quite low. However, it is known in the literature that the flare frequency of EV Lac varies with time (Leto et al. 1997). As it is seen in Figure 11, the monthly frequencies vary with time like EV Lac, when the flare frequencies are computed for each month in the case of KOI-256.

The Plateau value of the OPEA model depending on the flare equivalent duration distribution in the logarithmic scale versus the flare total time is found to be  $2.3121 \pm 0.0964$  s over detected 225 flares. This value is 3.014 s for EV Lac ( $B - V = 1^m.554$ ), 2.935 s for EQ Peg ( $B - V = 1^m.574$ ) and 2.637 s for V1005 Ori ( $B - V = 1^m.307$ ) (Dal & Evren 2011). The Plateau value was found to be  $1.232 \pm 0.069$  s for KIC 9641031 ( $B - V = 0^m.74$ ) (Yoldaş & Dal 2016) and 1.951  $\pm 0.069$  for KIC 9761199 ( $B - V = 1^m.303$ ) (Yoldaş & Dal 2017). If it is considered that one of the components is a white dwarf in the case of KOI-256, it is clear in this study that the chromospherical active component is the main-sequence star from the spectral type of M3. The B - V colour index for this component is given as ( $1^m.42$ ) in the literature (Walkowicz

& Basri 2013b). Therefore, the flares with high energy exhibited by KOI-256 are strong enough to be compared with the flares exhibited by the single UV Ceti-type stars. Dal & Evren (2011) found that the Plateau value is constant for a star, but varies from one star to another depending on their B-V colour indexes. The authors have defined this value as the saturation level of flares on a star, which is maximum energy limit a flare can reach.

On the other hand, the half-life parameter of the OPEA model was found to be 2233.6 s for KOI-256. This value is almost 10 times higher than the value obtained on single flare stars. This value is 433.10 s for DO Cep  $(B - V = 1^{m}.604)$ , 334.30 s for EQ Peg, and 226.30 s for V1005 Ori (Dal & Evren 2011). In the same way, the half-life parameter is 2291.7 s for KIC 9641031, while it is 1014 s for KIC 9761199, which are binary systems including a dMe-type component (Yoldaş & Dal 2016, 2017). As a result, when the flare total time for the stars such as EQ Peg, V1005 Ori, and DO Cep reaches a few minutes, the their flare energies can easily reach the Plateau level. However, in the case of KOI-256, a flare event must last at least 37 min, about a few tens of minutes, in order to reach the maximum flare energy described as the Plateau level for this system. Similar durations are also observed in the case of both flare rise and total times. The maximum flare rise time  $(T_{\rm r})$  observed in the single UV Ceti-type stars is about 2042 s for V1005 Ori, 1967 s for CR Dra (Dal & Evren 2011), while the longest flare rise time obtained in KOI-256 is 3942.749 s. The longest flare rise time obtained flares on KIC 9641031 is 5178.87 s (Yoldaş & Dal 2016), while it is 1118.099 s for KIC 9761199 (Yoldaş & Dal 2017). Similarly, the maximum flare total time for the flares obtained on the V1005 Ori is 5 236 s, while it is 4 955 s for CR Dra flares (Dal & Evren 2011). However, the longest flare obtained on KOI-256 lasts along 22185.361 s.

Considering the stellar spot activity together with the flare activity detected from KOI-256, it is obvious that the system exhibits a high level of chromospheric activity. Using very sensitive observations provided by Kepler Mission (Jenkins et al. 2010a, 2010b), some variations with small amplitudes, which are impossible to observe with ground-based classic telescopes, can be easily detected. When the system light

**Table 3.** The parameters calculated for each flare detected with analysis of the short cadence data obtained by Kepler Mission for KOI-256 are listed.

JD (+2 400 000)	P	$T_{\rm r}$	$T_{\rm d}$	$T_{\rm t}$	Amplitude
(+2 400 000)	(s)	(s)	(s)	(s)	(intensity)
55375.342090	44.20740	235.39162	2353.99133	2589.38294	0.07436
55375.947616	16.56043	176.54371	1118.15597	1294.69968	0.02621
55376.350846	15.26603	411.94483	765.05904	1177.00387	0.04393
55376.476174	22.31022	58.84790	1176.99523	1235.84314	0.05763
55377.033340	12.30394	176.55322	1176.98573	1353.53894	0.02639
55377.296257	27.67483	588.49718	882.75485	1471.25203	0.03511
55378.175597	20.63149	176.55235	2236.29206	2412.84442	0.03071
55378.328852	1.99895	58.85654	235.39248	294.24902	0.01837
55380.097067	34.60774	235.40026	2824.77888	3060.17914	0.02929
55380.590205	7.63004	58.84790	765.05731	823.90522	0.02825
55380.741416	8.25863 19.27064	294.24816	823.90522 1176.99264	1118.15338 1412.39290	0.01060
55381.267930 55381.637784	7.80086	235.40026 529.63978	647.34336	1176.98314	0.04652 0.01822
55383.117879	1.83830	58.83926	58.84704	117.68630	0.01822
55383.318131	1.72242	58.84790	58.85654	117.70445	0.02840
55385.586972	6.76596	1176.99005	235.39939	1412.38944	0.00939
55385.605363	15.97846	1176.99005	2118.58157	3295.57162	0.01176
55386.161845	1.38046	176.54285	117.70445	294.24730	0.01294
55386.294665	9.55332	176.56099	1353.53290	1530.09389	0.02204
55386.812322	10.15885	1059.29424	823.90349	1883.19773	0.01717
55389.135650	15.95644	1412.38771	1176.98832	2589.37603	0.01390
55390.912711	8.01176	470.79878	823.90176	1294.70054	0.02250
55391.021692	4.32847	353.08570	353.09520	706.18090	0.01673
55391.457613	9.42058	117.70358	706.19731	823.90090	0.02878
55392.580110	4.03619	117.70358	470.79792	588.50150	0.01606
55392.744261	9.80095	294.24643	1471.23302	1765.47946	0.01100
55393.700563	23.57782	235.39853	2118.56429	2353.96282	0.03602
55394.247507	4.93193	117.70358	529.64496	647.34854	0.02670
55395.160216	13.90038	117.70358	706.19558	823.89917	0.04806
55395.578427	404.82960	1118.12832	13417.62970	14535.75802	0.27309
55397.006747	38.35936	765.04262	2118.56947	2883.61210	0.03148
55401.269904	42.89049	117.69408	1412.37994	1530.07402	0.13944
55404.761341	7.68292	176.54198	765.04003	941.58202	0.03032
55405.599122	3.58347	117.70272	353.09174	470.79446	0.01701
55406.776100	7.62807	176.54976	529.63373	706.18349	0.03081
55407.746018	9.85367	176.54198	823.88534	1000.42733	0.02354
55408.136300	58.03656	470.79360	3060.13853	3530.93213	0.05706
55409.594580	75.45501	235.38902	2824.75555	3060.14458	0.13119
55409.832291	34.40761	176.54976	3001.29667	3177.84643	0.04241
55409.890187	26.17318	117.70272	1824.31181	1942.01453	0.05453
55412.681417 55412.876899	50.01182 6.91584	294.23434 176.54976	2765.89642 588.47904	3060.13075 765.02880	0.05973 0.02533
55413.391825	6.37731	58.84704	529.63978	588.48682	0.02509
55414.867810	5.48103	117.70186	588.49546	706.19731	0.02086
55414.925024	12.11842	235.39680	1294.67549	1530.07229	0.01774
55416.699338	12.14917	176.54112	1647.76464	1824.30576	0.02079
55417.073953	4.53704	176.54112	647.33299	823.87411	0.01461
55417.940336	2.68543	58.84704	294.24298	353.09002	0.01401
55418.479781	7.96692	294.24298	1000.42214	1294.66512	0.01591
55418.700463	40.81725	1353.52080	2589.33802	3942.85882	0.03181
55418.793776	4.98763	176.54026	588.48595	765.02621	0.01339
55419.427897	82.95145	823.87238	4237.09920	5060.97158	0.04302
55419.497371	185.47791	1118.12400	8827.28755	9945.41155	0.07175
55421.411309	8.11929	235.38730	941.58288	1176.97018	0.02561
55423.293915	68.96291	353.08915	5237.51011	5590.59926	0.05211
55425.451690	5.81840	58.85568	588.48336	647.33904	0.02426
55425.581783	11.64399	941.58029	647.33040	1588.91069	0.02157
55427.109523	10.32274	235.38643	1000.42646	1235.81290	0.02436
55428.162527	6.31033	117.69322	470.78150	588.47472	0.02500
55429.939553	14.85190	294.24125	1294.66598	1588.90723	0.03114

 Table 3. Continued.

JD	P	$T_{ m r}$	$T_{\rm d}$	$T_{\mathrm{t}}$	Amplitude
(+2 400 000)	(s)	(s)	(s)	(s)	(intensity)
55431.149892	6.49379	117.70186	882.71424	1000.41610	0.02000
55431.203700	14.69386	411.95088	1294.65648	1706.60736	0.02178
55433.127822	1.64001	176.53939	117.70099	294.24038	0.01724
55433.193209	5.77048	58.84704	1176.96240	1235.80944	0.01492
55434.151534	1.12874	176.55667	58.83754	235.39421	0.01448
55434.157664	2.97310	117.69235	411.94224	529.63459	0.01664
55434.353825	6.63105	117.68371	1588.90378	1706.58749	0.01544
55435.784160	4.06328 0.73911	117.70099	529.63459 58.85482	647.33558 117.70099	0.02182
55435.835924 55435.926512	3.61430	58.84618 235.40198	706.17312	941.57510	0.01707 0.01044
55436.045026	7.41432	58.84618	1294.65389	1353.50006	0.01044
55436.068183	3.51315	117.68371	529.64237	647.32608	0.01205
55436.081125	2.71089	353.08656	588.48077	941.56733	0.00666
55436.558584	23.75906	353.08656	2236.22035	2589.30691	0.01960
55436.636912	7.82704	353.08570	823.87411	1176.95981	0.00867
55436.794249	11.70227	176.54717	882.71165	1059.25882	0.03533
55438.701360	21.11367	117.69235	3413.18621	3530.87856	0.03614
55438.913866	5.61630	353.08570	1118.11363	1471.19933	0.01361
55441.226917	32.58644	647.32522	2412.76147	3060.08669	0.04902
55441.943445	3.33076	235.38470	58.85482	294.23952	0.02369
55443.211671	6.27441	235.39335	941.57251	1176.96586	0.01589
55443.229380	2.03953	470.78582	470.77805	941.56387	0.00682
55443.527706	0.95940	58.84618	117.70099	176.54717	0.01647
55443.531792	11.15569	235.39248	1294.64957	1530.04205	0.02500
55445.415740 55446.473501	11.63756 11.65998	235.39248 941.56214	1059.26400 2000.82614	1294.65648 2942.38829	0.02883 0.01389
55446.954364	23.94391	529.63200	1883.12429	2412.75629	0.01389
55447.452254	37.09812	1176.94598	3942.80266	5119.74864	0.03134
55450.550613	26.19646	235.39248	2353.90579	2589.29827	0.04492
55450.614637	15.81569	1118.10672	1294.65389	2412.76061	0.02139
55451.372709	5.64617	117.70099	353.08397	470.78496	0.03361
55453.086373	20.35148	176.55494	765.01325	941.56819	0.07455
55453.672806	149.01479	235.38384	3118.92422	3354.30806	0.18088
55456.080515	72.95417	529.63805	3942.77933	4472.41738	0.07085
55456.193579	12.16574	176.54544	941.55869	1118.10413	0.03073
55456.308004	15.39212	117.69149	1530.02477	1647.71626	0.02762
55456.611096	5.30449	117.69149	470.77546	588.46694	0.02077
55456.622675	12.84198	235.38298	2118.51763	2353.90061	0.01174
55457.468608	193.53680	411.92842	8827.13117	9239.05958	0.11971
55458.638065	6.67818	235.38298	882.72029	1118.10326	0.01673
55460.915679	9.38679	58.84531	765.02016	823.86547	0.02893
55460.929301 55461.137038	3.05990 10.42624	117.70013 58.84531	529.62854 1000.41091	647.32867 1059.25622	0.00933 0.03052
55461.401988	84.63194	176.54630	5002.01914	5178.56544	0.10579
55462.177765	11.80962	353.08224	1118.11104	1471.19328	0.02530
55463.505896	16.78388	117.70013	823.85597	941.55610	0.05735
55464.891260	15.02389	1471.18291	1353.49142	2824.67434	0.01315
55465.276084	10.70152	117.68198	1294.65475	1412.33674	0.02691
55465.425926	14.17575	176.54544	1176.93734	1353.48278	0.04345
55465.482458	18.31911	235.39939	1530.01958	1765.41898	0.05152
55465.952419	41.43259	58.84618	58.84531	117.69149	0.68383
55467.303046	68.23197	117.70013	3060.06336	3177.76349	0.13012
55469.090940	10.93777	176.53680	941.55350	1118.09030	0.03609
55469.613346	18.50133	823.85424	2530.43395	3354.28819	0.01302
55470.019964	6.44893	117.70013	823.85338	941.55350	0.02676
55470.607755	4.84475	117.69926	823.85424	941.55350	0.01444
55471.355605	15.53370	235.39075	1176.93475	1412.32550	0.03245
55472.217880	19.35750	176.54458	3883.92019	4060.46477	0.01227
55472.319364 55472.813844	5.77485 16.21783	411.93533 117.70013	882.69869 882.70733	1294.63402 1000.40746	0.01891 0.05861
1 1 + 1 / . O L 1044	10.41/03	117.70013	004.70733	1000.40740	0.05001

Table 3. Continued.

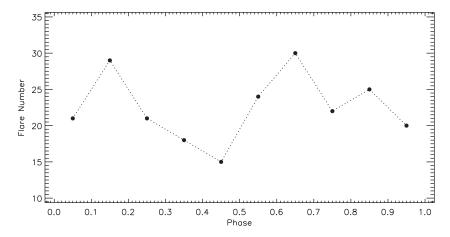
JD	P	$T_{ m r}$	$T_{ m d}$	$T_{t}$	Amplitude
(+2 400 000)	(s)	(s)	(s)	(s)	(intensity)
55474.398085	12.41232	294.22742	1294.64179	1588.86922	0.02210
55474.727057	44.88099	176.52730	1941.95923	2118.48653	0.18575
55474.900057	9.91234	176.54458	941.55264	1118.09722	0.03113
55475.423824	5.64696	117.69926	588.47126	706.17053	0.02536
55475.562087	6.45265	58.83667	706.17053	765.00720	0.03652
55476.091984	29.14921	235.38989	1353.48710	1588.87699	0.08591
55476.294952	7.44200	58.84531	941.56042	1000.40573	0.02985
55476.567392	314.03759	706.16189	21479.19898	22185.36086	0.07990
55476.696120	21.90776	529.61731	1883.11219	2412.72950	0.02629
55476.823486	86.57508	117.69926	8650.51920	8768.21846	0.03453
55478.800721	21.97714	765.00634	1647.72144	2412.72778	0.02521
55479.075204	48.89291	235.39853	2000.80195	2236.20048	0.08979
55479.596246	75.63241	411.92582	4413.52886	4825.45469	0.06286
55480.878077	3.63981	235.38125	706.16966	941.55091	0.00691
55480.897829	53.97274	176.54458	2530.42618	2706.97075	0.07388
55480.934608	46.92281	294.23520	4236.99120	4531.22640	0.02365
55481.363701	147.57825	470.77978	5001.99754	5472.77731	0.07048
55481.384815	160.06194	2295.02765	4295.83565	6590.86330	0.07118
55482.448012 55484.817558	25.10970 7.04579	235.38989	1588.86576	1824.25565	0.07721
		117.69062 176.52730	1000.40400	1118.09462	0.03365 0.02701
55486.206319 55486.736215	8.58246 5.60687	58.84531	941.55869 588.47818	1118.08598 647.32349	0.01998
55487.396200	39.15742	1941.95318	5119.68989	7061.64307	0.01998
55488.308191	75.29180	353.07965	3530.81722	3883.89686	0.01433
55488.373577	9.90490	529.62422	1883.10701	2412.73123	0.01054
55489.807290	140.09412	3942.74909	9650.89814	13593.64723	0.02690
55490.681140	1049.43937	2471.56790	15947.53056	18419.09846	0.33241
55493.194394	14.87307	176.53594	647.32262	823.85856	0.04763
55494.570213	54.05090	353.07965	2059.64986	2412.72950	0.11057
55494.645815	28.36075	941.54832	2471.57482	3413.12314	0.01765
55498.618661	9.03055	176.54371	1471.17168	1647.71539	0.01734
55498.702436	3.62009	176.54371	470.77891	647.32262	0.01493
55498.718101	13.40445	588.46867	2295.02074	2883.48941	0.01675
55499.032768	5.29147	176.54458	823.84906	1000.39363	0.00952
55500.394283	15.19193	117.68198	1706.55984	1824.24182	0.03291
55501.954679	6.79482	58.85395	1118.09117	1176.94512	0.02542
55502.011891	18.73455	588.46781	2236.18406	2824.65187	0.01551
55502.410333	11.46807	176.53507	765.01238	941.54746	0.03713
55504.495855	9.53926	117.69840	1118.08339	1235.78179	0.02625
55505.411932	55.53928	176.54371	2883.49718	3060.04090	0.08309
55505.476636	60.13398	117.69840	5001.98198	5119.68038	0.10543
55508.166292	1.79399	58.84531	58.85395	117.69926	0.03583
55508.891661	14.28894	235.38902	1000.40141	1235.79043	0.04915
55509.218588	21.62582	706.16794	2706.95347	3413.12141	0.01718
55509.391587	31.00352	294.24298	2824.65274	3118.89571	0.03682
55510.390758	36.78295	176.55235	4236.96960	4413.52195	0.05801
55510.739480	11.95360	882.70214	941.54832	1824.25046	0.02207
55511.376307	51.67114	588.46867	3177.72288	3766.19155	0.05542
55513.722008	9.63831	176.53507	1883.10442	2059.63949	0.01544
55514.407193	8.84904	117.69062	1471.17946	1588.87008	0.02198 0.01942
55515.838861 55516.291791	13.31654	235.38902	1941.94195	2177.33098 1706.56070	
55516.392594	11.92080 39.99328	353.07965 3118.89658	1353.48106 2177.32234	5296.21891	0.02661 0.01812
55516.430735	23.22504	117.68976	2706.95434	2824.64410	0.01812
55517.760920	98.70816	3177.73325	6826.23677	10003.97002	0.02622
55520.294607	67.24819	176.54371	6708.53923	6885.08294	0.02022
55522.737707	39.92723	176.53594	2236.18666	2412.72259	0.03204
55522.816034	75.88575	176.53507	2177.34134	2353.87642	0.19217
55524.700634	23.03246	58.84531	2353.87728	2412.72259	0.03093
55525.479130	50.79022	353.08051	2765.81174	3118.89226	0.09610
55526.711920	9.11512	117.69926	588.47818	706.17744	0.03393

Table 3. Continued.

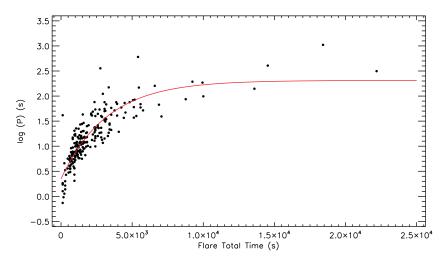
JD	P	$T_{ m r}$	$T_{ m d}$	$T_{\mathrm{t}}$	Amplitude
(+2 400 000)	(s)	(s)	(s)	(s)	(intensity)
55527.096741	8.51295	117.69840	1235.77661	1353.47501	0.01772
55527.229556	7.92308	353.07965	1176.93907	1530.01872	0.01985
55527.828923	17.13460	176.53507	2059.64467	2236.17974	0.03384
55528.409901	6.92449	176.54458	647.31485	823.85942	0.01744
55530.246145	602.45858	353.08051	5060.84976	5413.93027	0.62493
55530.609853	110.84606	529.62509	2412.72605	2942.35114	0.13490
55530.638459	22.64046	58.84531	1000.40400	1059.24931	0.05504
55531.639676	20.11803	823.86029	765.01498	1588.87526	0.03575
55532.322138	5.99543	58.84531	941.55091	1000.39622	0.02514
55532.377988	17.43303	117.69926	882.69696	1000.39622	0.04528
55532.540771	19.95533	117.69926	1412.32205	1530.02131	0.05189
55534.100491	11.44830	117.69926	1235.78698	1353.48624	0.01969
55534.874221	7.51594	117.70013	765.00634	882.70646	0.03108
55535.983734	69.75887	235.38125	3707.37907	3942.76032	0.07953
55536.041628	40.06649	117.69926	2706.97334	2824.67261	0.04786
55536.117911	27.10988	117.69149	1235.78698	1353.47846	0.07168
55536.827618	59.17967	882.70646	3354.28214	4236.98861	0.03548
55538.078801	357.68647	176.54458	2589.26717	2765.81174	0.71509
55538.144186	40.92859	117.69926	2059.64986	2177.34912	0.07322
55539.366763	7.80770	117.69062	706.17139	823.86202	0.02386
55540.609774	23.02603	235.38125	2412.73296	2648.11421	0.04814
55540.952368	5.51462	117.68285	765.02534	882.70819	0.01524
55540.963266	18.93492	117.69062	2236.19702	2353.88765	0.02537
55541.343321	3.94267	117.70013	765.00806	882.70819	0.01674
55541.395084	14.23128	411.93533	2648.12371	3060.05904	0.01220
55542.043493	1.26695	58.85395	58.84618	117.70013	0.02280
55542.082316	7.81970	176.53594	1294.64438	1471.18032	0.02129
55542.610851	26.84851	529.62682	2883.52397	3413.15078	0.03662
55542.691222	7.77636	117.69926	882.70819	1000.40746	0.02239
55542.725277	4.57177	58.84531	176.53680	235.38211	0.04254
55543.163907	14.31275	411.92669	1059.25363	1471.18032	0.03769
55544.722950	7.26750	176.52816	588.48077	765.00893	0.03342
55545.570242	5.29618	176.54544	823.86374	1000.40918	0.01745
55545.611108	16.07877	235.39075	1176.94598	1412.33674	0.05000
55547.527733	35.88220	588.46435	2530.43914	3118.90349	0.02779
55548.411805	6.62462	176.53680	706.17398	882.71078	0.01821
55548.777558	42.66308	117.69149	2000.81146	2118.50294	0.09381
55549.246839	14.48484	117.70013	2295.04061	2412.74074	0.01299
55550.629480	18.35971	470.78237	2236.19616	2706.97853	0.01681
55552.452797	10.20249	353.08310	882.72029	1235.80339	0.03406
56419.872987	42.42996	235.39507	3295.63123	3531.02630	0.07203
56421.178735	18.79582	176.54630	1647.81994	1824.36624	0.02747
56422.598915	59.26194	2707.13491	2883.67258	5590.80749	0.02899

variations out-of-eclipses without any flares is carefully examined, it is seen that the shape of eliminated light curve can vary per two or three cycles at most. This variation is caused due to both rapid evolutions and migrations of the cool stellar spots on the surface of the M3 star. It needs about 4–5 d for the evolution of the stellar spots on the active component in the case of KOI-256, while it takes a few weeks for the cool spot groups on the solar surface (Gershberg 2005). In this study, the stellar spot distribution on the surface of the active component was modelled by the SPOTMODEL program (Ribárik 2002; Ribárik et al. 2003), though there are some deficiency in the method of this program. As it was mentioned in Section 1,

the method using in the SPOTMODEL program do consider neither the spot evolution nor spot migration on the surface during long time. However, separating the data into suitable subsets by considering the light curve shape change point, we solved out this problem. The whole data were separated into 138 subsets. The cycle shapes of sinusoidal light variation are different from one subset data to the next, while the shapes of the cycles are the same among themselves in each subset data. Thus, each subset data was individually modelled. On the other hand, maximum level of the obtained brightness is not suitable for modelling by the SPOTMODEL program due to detrended data. However, the spot longitude parameters is



**Figure 7.** The distribution of flare total number in phase range of 0.10 is plotted versus the phase computed by using orbital period for 225 flares.



**Figure 8.** The distribution of the equivalent duration in the logarithmic scale (log*P*) are plotted versus the flare total time, which were sum of flare rise and decay times. The OPEA model obtained over 225 flare determined in the analyses. The fill circles represent the observed flares, while the red line represents the OPEA model.

calculated in this study, because it is an important parameter to reveal the spot migration instead of the spot latitude or radius.

The result parameters of the modelling spot variation are plotted in Figures 2–4. As it is seen from Figure 2, the models of 138 subsets indicate that two cool spots (or spotted areas) are enough to absolutely fit the observed light variation out-of-eclipses. Although both latitude and radius parameters of the spots do not externalise the realistic nature of the spots on the stellar surface as discussed earlier, their variations are also plotted in Figure 3 to note as an initial approach. The main goal of these models is longitudinal variation and it is shown in Figure 4. Observations of the cool spot groups on solar surface reveal the existence of two permanently active longitudes separated by 180° from each other. These active longitudes known as the Carrington coordinates are constant structure according to some authors, while some authors indicate that the rotation speeds of these active longitudes are not

**Table 4.** Parameters derived from the OPEA model by least squares method.

Parameter	Values	95% confidence intervals
y <sub>0</sub>	0.34741	0.2440 to 0.4508
Plateau	2.3121	2.1232 to 2.5010
K	0.00031033	0.00024245 to 0.00037821
Tau	3222.4	2644.0 to 4124.6
Half-time	2233.6	1832.7 to 2859.0
Span	1.9647	1.7914 to 2.1380
Span $R^2$	_	0.75

constant, but slowly change (Lopez 1961; Stanek 1972; Bogart 1982). Similarly, in this study, it was found that there are two stellar cool spots on the active component of KOI-256, which show the migration behaviour with different speeds.

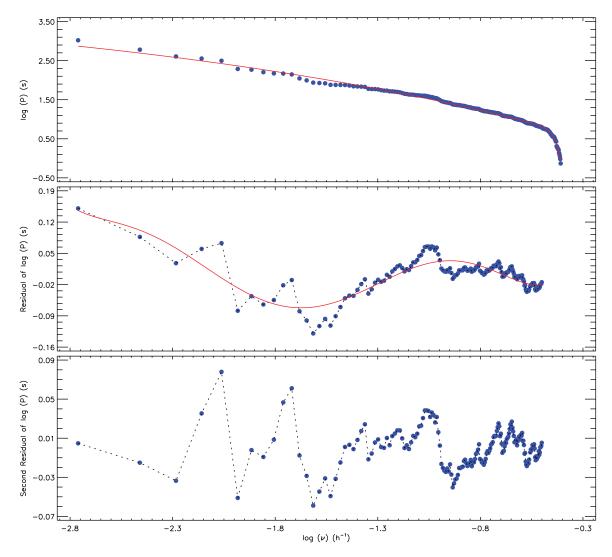


Figure 9. The cumulative flare frequencies ( $\log \nu$ ) and model computed for 225 flares obtained from KOI-256. In the upper panel, it is seen the variation of the flare equivalent durations in logarithmic scale ( $\log P$ ) versus the cumulative flare frequency, which is called the flare energy spectrum (Gershberg 2005), while the residuals obtained from the model are shown in the middle and bottom panels.

Modelling the migration movements shown in Figure 4 by the linear fits, it is estimated that the migration of the first stellar spot is 3.95 yrs and the migration period of the second one is 8.37 yrs.

To test whether these results are real, it was tested by another method. Computing  $\theta_{\min}$  parameter, which is the phases of the observed sinusoidal variation out-of-eclipses, the migration behaviour of the dominant spotted area was tried to find out. Indeed, as it is seen from Figure 5, the  $\theta_{\min}$  values obtained for all cycles were plotted versus Heliocentric Julian Day, and then, its trend was fitted by a linear function. The migration period of dominant spotted area is found to be 9.13 yrs. In addition, the most impressive finding was came from the residual  $\theta_{\min}$  according to the linear fit. The residuals show the same trend with the residuals of the longitudinal variation obtained from the SPOTMODEL analyses. This situation absolutely reveals that the main goal of the SPOTMODEL analyses have been obtained, and it does

figure out the migration of active regions close to its real nature.

However, as can be seen in Figure 4, some sinusoidal variations still remain, after a linear correction is applied to the migration movement. All of these migration behaviours are an indicator of the strong differential rotation on stellar surface. However, a high-resolution spectroscopic observation is needed to confirm this case. In the literature, the system's age is given as 0.01 Gyr by Walkowicz & Basri (2013b). However, as it is described in Section 1, using the Equation (2.3) given by Gänsicke (1997) with the stellar parameters taken from Muirhead et al. (2013), the age of the system is estimated as  $\sim$ 2 Gyr in this study. Considering the law expressed by Skumanich (1972), this age is a bit old for rapid rotation, considering high-level chromospheric activity. However, KOI-256 is a binary system. In this case, tidal effects can let the components to rotate rapidly, which lets the system has a significant influence on both the flare and spots activity behaviour.

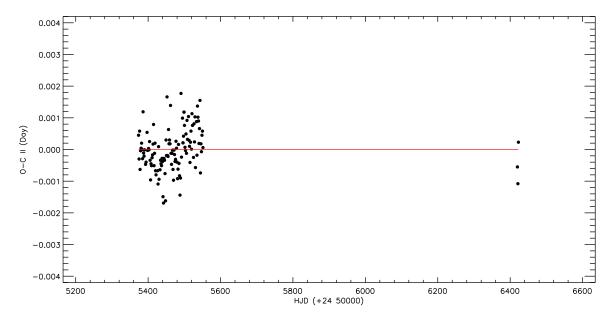


Figure 10. The variations of minima time's residuals (O - C)II obtained by applying a linear correction to minima times are shown versus time. All the residuals are shown with filled circles, while the red line represents a linear fit.

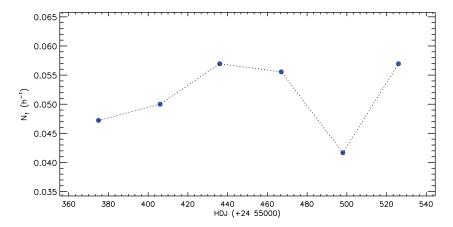


Figure 11. The monthly variation of the flare frequency of  $N_1$ , which indicates total flare number per each hour, for KOI-256 is shown for the entire observing season.

Muirhead et al. (2014) listed the semi-major axis length of the system as  $a = 4.51 \, \mathrm{R}_{\odot}$ , giving the radii of components as  $R_1 = 0.54 \, \mathrm{R}_{\odot}$  and  $R_2 = 0.01 \, \mathrm{R}_{\odot}$ . This situation indicates that the components are very close to each other, which can cause rapid activity variations due to some tidal effects on each other. On the other hand, the most interesting result in term of stellar spot and flare activity exhibited by the system is that the stellar spot areas exhibit a rapid variation in location due to the migration movement as shown in Figure 5, while it is seen interestingly that the flares are in aim to occur frequently between the phases of 0.10–0.20 and also 0.60–0.70, when the averaged phases are calculated in each 0.10 phase range for 225 flares. In the case of both spot and flare activities, the phases were computed with using the orbital period of binary system with a white dwarf and a main-sequence M3

components. This case is an unexpected situation, because the locations of flares are quite stable compared to the behaviour of the spot activity. Although flare activity is seen in all phases, it is observed that the flare activity most frequently occurs in these phase intervals. The location of the white dwarf is as close as 4.51  $R_{\odot}$  to the active component, which indicates that the components are interacting magnetically. A possible explanation is that the quickly evolving stellar spot areas under both the differential ration and tidal effects exhibit rapid migration movements, while the flares are frequently occurring in some definite longitudes on the active component surface due to the magnetic interaction of the white dwarf and active components. This situation leads KOI-256 to take an important position among its analogues for especially the future spectral studies.

**Table 5.** Minima times and their residuals.

MJD (Obs)	E	Type	O-C(I) (d)	O - C(II) (d)	Linear fit to the $O - C(I)$
55373.637352	296	I	0.00175	0.00045	0.00130
55375.015258	297	I	0.00100	-0.00030	0.00130
55376.394783	298	I	0.00188	0.00058	0.00130
55377.772232	299	I	0.00068	-0.00063	0.00130
55379.151475	300	I	0.00127	-0.00004	0.00130
55380.530204	301	I	0.00135	0.00004	0.00131
55381.909019	302	I	0.00151	0.00020	0.00131
55383.287462	303	I	0.00130	0.00000	0.00131
55384.665830	304	I	0.00102	-0.00029	0.00131
55386.045957	305 306	I I	0.00250 0.00120	0.00119 $-0.00011$	0.00131 0.00131
55387.423309 55388.801861	307	I	0.00120	-0.00011 $-0.00021$	0.00131
55390.180721	307	I	0.00110	-0.00021 $-0.00001$	0.00132
55392.937562	310	I	0.00085	- 0.00001 - 0.00047	0.00132
55394.316284	311	I	0.00092	-0.00047	0.00132
55397.074523	313	I	0.00186	0.00054	0.00132
55398.452606	314	I	0.00130	-0.00034	0.00133
55401.209976	316	I	0.00136	0.00003	0.00133
55402.588586	317	I	0.00133	-0.00001	0.00133
55403.967495	318	I	0.00158	0.00025	0.00133
55405.345561	319	I	0.00100	-0.00034	0.00134
55406.723591	320	I	0.00038	-0.00096	0.00134
55408.102754	321	I	0.00089	-0.00045	0.00134
55409.481348	322	I	0.00084	-0.00051	0.00134
55410.860243	323	I	0.00108	-0.00026	0.00134
55412.238990	324	I	0.00118	-0.00017	0.00135
55413.617981	325	I	0.00152	0.00017	0.00135
55414.997250	326	I	0.00214	0.00079	0.00135
55416.374600	327	I	0.00084	-0.00051	0.00135
55417.753642	328	I	0.00123	-0.00012	0.00135
55419.132621	329	I	0.00156	0.00020	0.00135
55420.510400	330	I	0.00069	-0.00067	0.00136
55421.888921	331	I	0.00056	-0.00080	0.00136
55426.025010	334	I	0.00069	- 0.00067	0.00136
55427.403240	335	I	0.00027	-0.00109	0.00136
55428.783068	336	I	0.00145	0.00009	0.00137
55430.160693	337 339	I I	0.00043 0.00073	- 0.00094 - 0.00064	0.00137 0.00137
55432.918294	339 340	I	0.00073	- 0.0004 - 0.00034	0.00137
55434.297248 55435.675787	341	I	0.00103	- 0.00034 - 0.00045	0.00137
55437.054380	342	I	0.00092	- 0.00043 - 0.00051	0.00137
55438.433266	343	I	0.00110	-0.00031 $-0.00028$	0.00138
55439.811817	344	I	0.00110	-0.00028	0.00138
55441.189359	345	I	-0.00013	-0.00149	0.00138
55442.567810	346	I	-0.00031	-0.00169	0.00138
55443.947869	347	I	0.00110	-0.00028	0.00139
55445.326462	348	I	0.00104	-0.00034	0.00139
55446.704701	349	I	0.00063	-0.00076	0.00139
55448.082487	350	I	-0.00023	-0.00162	0.00139
55449.463059	351	I	0.00169	0.00030	0.00139
55450.841227	352	I	0.00121	-0.00019	0.00139
55452.221730	353	I	0.00306	0.00166	0.00140
55453.598523	354	I	0.00120	-0.00019	0.00140
55454.977153	355	I	0.00118	-0.00022	0.00140
55456.356654	356	I	0.00203	0.00063	0.00140
55457.734853	357	I	0.00158	0.00018	0.00140
55459.113623	358	I	0.00170	0.00030	0.00140
55460.492162	359	I	0.00159	0.00018	0.00141
55461.872023	360	I	0.00280	0.00139	0.00141
55463.249164	361	I	0.00129	- 0.00012	0.00141
55464.627459	362 363	I	0.00094	- 0.00047	0.00141
55466.006470 55467.385217	363 364	I I	0.00130	- 0.00012 0.00002	0.00141
55467.385217	364	1	0.00139	-0.00002	0.00141

Table 5. Continued.

MJD (Obs)	E	Type	O-C(I) (d)	O - C(II) (d)	Linear fit to the $O - C(I)$
55468.763258	365	I	0.00078	- 0.00063	0.00142
55470.141569	366	I	0.00045	-0.00097	0.00142
55471.521176	367	I	0.00140	-0.00002	0.00142
55472.899686	368	I	0.00126	-0.00016	0.00142
55474.278113	369	I	0.00104	-0.00038	0.00142
55475.656836	370	I	0.00111	-0.00031	0.00142
55477.036061	371	I	0.00169	0.00026	0.00143
55478.414494	372	I	0.00147	0.00004	0.00143
55479.792703	373	I	0.00103	- 0.00040	0.00143
55481.170834	374 275	I	0.00051	- 0.00092	0.00143
55482.549788 55483.929219	375 376	I I	0.00081 0.00159	- 0.00062 0.00016	0.00143 0.00144
55485.307269	377	I	0.00099	- 0.00044	0.00144
55486.685532	378	I	0.00099	-0.00044 $-0.00083$	0.00144
55488.063580	379	I	0.00001	-0.00144	0.00144
55489.442769	380	I	0.00054	- 0.00090	0.00144
55490.824094	381	I	0.00322	0.00177	0.00144
55492.200729	382	I	0.00120	-0.00024	0.00145
55494.959266	384	I	0.00244	0.00099	0.00145
55496.337138	385	I	0.00166	0.00021	0.00145
55496.337138	385	I	0.00166	0.00021	0.00145
55497.716000	386	I	0.00187	0.00042	0.00145
55499.095410	387	I	0.00263	0.00118	0.00145
55499.095410	387	I	0.00263	0.00118	0.00145
55500.473644	388	I	0.00222	0.00076	0.00146
55501.851611	389	I	0.00153	0.00007	0.00146
55503.230158	390	I	0.00143	-0.00003	0.00146
55504.609336	391	I	0.00196	0.00049	0.00146
55505.987369	392	I	0.00134	- 0.00012	0.00146
55507.367066	393	I	0.00239	0.00092	0.00146
55508.745117	394 396	I I	0.00179 0.00251	0.00032	0.00147 0.00147
55511.503144 55512.881040	390 397	I	0.00231	0.00104 0.00029	0.00147
55514.259500	398	I	0.00170	0.00029	0.00147
55515.637647	399	Ī	0.00107	- 0.00041	0.00147
55517.016941	400	I	0.00171	0.00023	0.00148
55518.395944	401	I	0.00206	0.00058	0.00148
55519.774024	402	I	0.00149	0.00001	0.00148
55521.153793	403	I	0.00261	0.00113	0.00148
55522.532080	404	I	0.00225	0.00076	0.00148
55525.288367	406	I	0.00123	-0.00025	0.00149
55526.668079	407	I	0.00230	0.00081	0.00149
55528.046162	408	I	0.00173	0.00024	0.00149
55529.425603	409	I	0.00252	0.00103	0.00149
55530.802659	410	I	0.00092	- 0.00057	0.00149
55533.561410	412	I	0.00238	0.00088	0.00150
55534.939002	413	I	0.00132	- 0.00018	0.00150
55536.319211	414	I	0.00288	0.00137	0.00150
55537.697510	415	I	0.00252	0.00102	0.00150 0.00150
55539.076038 55540.453980	416 417	I I	0.00240 0.00169	0.00090 0.00019	0.00150
55541.833100	418	I	0.00216	0.00019	0.00151
55543.212646	419	I	0.00210	0.0005	0.00151
55544.589008	420	I	0.00300	- 0.00074	0.00151
55545.968579	421	I	0.00169	0.00074	0.00151
55547.346980	422	Ī	0.00144	- 0.00007	0.00151
55548.726156	423	I	0.00197	0.00045	0.00151
55550.104937	424	I	0.00210	0.00058	0.00152
55551.483063	425	I	0.00158	0.00006	0.00152
56420.033207	1 055	I	0.00206	-0.00055	0.00261
56421.411327	1 056	I	0.00153	-0.00108	0.00261
56422.791294	1 057	I	0.00285	0.00023	0.00261

A total of 125 minima times were determined from the all available data. It is seen that the light curve has only the primary minimum, when entire the light curves are examined. According to Tran et al. (2013) and Balaji et al. (2015), the (O-C)II residuals of the primary and secondary minima show some sinusoidal variation versus time, but in the opposed phase, if one of the components in an eclipsing binary system is a magnetically active star. However, the (O-C)II residuals obtained from the minima detected from the light curves of KOI-256 do not exhibit any separation. This situation also confirms that all the minima detected from the light curves are primary ones.

As it is seen from Figure 10, considering all 125 minima times, there is not any clear variation in the (O-C)II residuals distribution versus time obtained by improving linear correction. As shown in Figure 10, the (O - C)II residuals scatter in a range of about 50-60 s. Whereas according to Szabó et al. (2013) in the literature, KOI-256 shows an (O -C)II variation with an amplitude of 0.4464 min and a period of 41.755397 d. However, the results obtained in this study do not confirm the variation found by Szabó et al. (2013). The problem, which cause the different results can be due to the method used to compute the minima times or due to the different data formats used in these studies. In this study, the minima times were computed from the short cadence data, while the minima times were calculated from the long cadence data by Szabó et al. (2013). In addition, although the flares as the activity exhibited by the system were removed from the light curves before computing the minima times, the sinusoidal variation due to the stellar spot activity was not cleared from the data in order to test whether the primary and secondary minima were separated, as described by Tran et al. (2013) and Balaji et al. (2015). Although there is no any clue for the secondary minima, the situation can cause a scattering in a wide range for the (O - C)II residuals variation obtained in this study.

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