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## Anisotropic flow, entropy, and $L^{p-M i n k o w s k i ~ p r o b l e m ~}$

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Abstract. We provide a natural simple argument using anistropic flows to prove the existence of weak solutions to Lutwak's $L^{p}$-Minkowski problem on $S^{n}$ which were obtained by other methods.

## 1 Introduction

For $\alpha>0$ and nonnegative $f \in L^{1}\left(\mathbb{S}^{n}\right)$ with positive integral, we are interested in finding a weak solution to the Monge-Ampére equation

$$
\begin{equation*}
u^{\frac{1}{\alpha}} \operatorname{det}\left(\bar{\nabla}_{i j}^{2} u+u \bar{g}_{i j}\right)=f \tag{1.1}
\end{equation*}
$$

or in other words, a weak solution to Lutwak's $L^{p}$-Minkowski problem on $S^{n}$ when $-n-1<p<1$ for $p=1-\frac{1}{\alpha}$ where $\bar{\nabla}$ is the Levi-Civita connection of $\mathbb{S}^{n}, \bar{g}_{i j}$, with $\bar{g}$ being the induced round metric on the unit sphere. By a weak (Alexandrov) solution, we mean the following: Given a nontrivial finite Borel measure $\mu$ on $\mathbb{S}^{n}$ (for example, $d \mu=f d \theta$ for the Lebesgue measure $\theta$ on $S^{n}$ and the $f$ in (1.1)), find a convex body $\Omega \subset \mathbb{R}^{n+1}$ with $o \in \Omega$ such that

$$
\begin{equation*}
d \mu=u^{\frac{1}{\alpha}} d S_{\Omega} \tag{1.2}
\end{equation*}
$$

where $u(x)=\max _{z \in \Omega}\langle x, z\rangle$ is the support function and $S_{\Omega}$ is the surface area measure of $\Omega$ (see [45]). If $\partial \Omega$ is $C_{+}^{2}$, then

$$
d S_{\Omega}=\operatorname{det}\left(\bar{\nabla}_{i j}^{2} u+u \bar{g}_{i j}\right) d \theta=K^{-1} d \theta,
$$

where $K(x)$ is the Gaussian curvature at the point of $\partial \Omega$ where $x \in S^{n}$ is the exterior unit normal (see [45]). Concerning the regularity of the solution of (1.1), if $f \in$ $C^{0, \beta}\left(S^{n}\right)$ and $u$ are positive, then $u$ is $C^{2, \beta}$ according to Caffarelli's regularity theory in $[15,16]$. On the other hand, even if $f$ is positive and continuous for $\alpha>\frac{1}{n}$, there might exist weak solution where $u(x)=0$ for some $x \in S^{n}$ and $u$ is not even $C^{1}$ according to Example 4.2 in [7]. Moreover, even if $f \in C^{0, \beta}\left(S^{n}\right)$ is positive, it is possible that $u(x)=0$ for some $x \in S^{n}$ for $\alpha>\frac{1}{n}$, but Choi, Kim, and Lee [19] still managed to obtain some regularity results in this case.

[^0]The case $\alpha=\frac{1}{n+2}$ of the Monge-Ampére equation (1.1) is the critical case when the left-hand side of (1.1) is invariant under linear transformations of $\Omega$, and the case $\alpha=1$ is the so-called logarithmic Minkowski problem posed by Firey [23]. Setting $p=1-\frac{1}{\alpha}<1$, the Monge-Ampére equation (1.1) is Lutwak's $L^{p}$-Minkowski problem

$$
\begin{equation*}
u^{1-p} \operatorname{det}\left(\bar{\nabla}_{i j}^{2} u+u \bar{g}_{i j}\right)=f . \tag{1.3}
\end{equation*}
$$

In this notation, (1.2) reads as

$$
\begin{equation*}
d \mu=u^{1-p} d S_{\Omega} ; \tag{1.4}
\end{equation*}
$$

that equation makes sense for any $p \in \mathbb{R}$. Within the rapidly developing $L^{p}$-BrunnMinkowski theory (where $p=1$ is the classical case originating from Minkowski's oeuvre) initiated by Lutwak [39-41], if $p>1$ and $p \neq n+1$, then Hug, Lutwak, Yang, and Zhang [30] (improving on Chou and Wang [20]) prove that (1.4) has an Alexandrov solution if and only if the $\mu$ is not concentrated onto any closed hemisphere, and the solution is unique. We note that there are examples in [25] (see also [30]) and show that if $1<p<n+1$, then it may happen that the density function $f$ is a positive continuous in (1.3) and $o \in \partial K$ holds for the unique Alexandrov solution, and actually Bianchi, Böröczky, and Colesanti [7] exhibit an example that $o \in \partial K$ even if the density function $f$ is a positive continuous in (1.3) assuming $-n-1<p<1$.

In the case $p \in(0,1)$ (or equivalently, $\alpha>1$ ), if the measure $\mu$ is not concentrated onto any great subsphere of $S^{n}$, then Chen, Li, and Zhu [17] prove that there exists an Alexandrov solution $K \in \mathcal{K}_{o}^{n}$ of (1.4) using a variational argument (see also [8]). We note that for $p \in(0,1)$ and $n \geq 2$, no complete characterization of $L^{p}$-surface area measures is known (see [12] for the case $n=1$, and [8,43] for partial results about the case when $n \geq 2$ and the support of $\mu$ is contained in a great subsphere of $S^{n}$ ).

Concerning the case $p=0$ (or equivalently, $\alpha=1$ ), the still open logarithmic Minkowski problem (1.3) or (1.4) was posed by Firey [23] in 1974. The paper [11] characterized even measures $\mu$ such that (1.4) has an even solution for $p=0$ by the socalled subspace concentration condition (see (a) and (b) in Theorem 1.1). In general, Chen, Li, and Zhu [18] proved that if a nontrivial finite Borel measure $\mu$ on $S^{n-1}$ satisfies the same subspace concentration condition, then (1.4) has a solution for $p=0$. On the other hand, Böröczky and Hegedus [10] provide conditions on the restriction of the $\mu$ in (1.4) to a pair of antipodal points.

If $-n-1<p<0$ (or equivalently, $\frac{1}{n+2}<\alpha<1$ ) and $f \in L_{\frac{n+1}{n+1+p}}\left(S^{n}\right)$ in (1.3), then (1.3) has a solution according to [8]. For a rather special discrete measure $\mu$ satisfying that $\mu$ is not concentrated on any closed hemisphere and any $n$ unit vectors in the support of $\mu$ are independent, Zhu [47] solves the $L^{p}$-Minkowski problem (1.4) for $p<0$. The $p=-n-1$ (or equivalently, $\alpha=\frac{1}{n+2}$ ) case of the $L^{p}$-Minkowski problem is the critical case because its link with the $\operatorname{SL}(n)$ invariant centro-affine curvature whose reciprocal is $u^{n+2} \operatorname{det}\left(\bar{\nabla}_{i j}^{2} u+u \bar{g}_{i j}\right)$ (see [29] or [38]). For positive results concerning the critical case $p=-n-1$, see, for example, $[28,34]$, and for obstructions for a solution, see, for example, [20, 22].

In the super-critical case $p<-n-1$ (or equivalently, $\alpha<\frac{1}{n+2}$ ), there is a recent important work by Li, Guang, and Wang [27] proving that for any positive $C^{2}$ function $f$, there exists a $C^{4}$ solution of (1.3). See also [22] for non-existence examples.

The main contribution of this paper is to provide a very natural argument based on anisotropic flows developed by Andrews [4] to handle the case $-n-1<p<1$, or equivalently, the case $\frac{1}{n+2}<\alpha<\infty$.

Entropy functional. For any convex body $\Omega$, a fixed positive function $f$ on $\mathbb{S}^{n}$ and $\alpha \in(0, \infty)$, define

$$
\begin{equation*}
\mathcal{E}_{\alpha, f}(\Omega):=\sup _{z \in \Omega} \mathcal{E}_{\alpha, f}(\Omega, z), \tag{1.5}
\end{equation*}
$$

where

$$
\mathcal{E}_{\alpha, f}(\Omega, z):= \begin{cases}\frac{\alpha}{\alpha-1} \log \left(\int_{\mathbb{S}^{n}} u_{z}(x)^{1-\frac{1}{\alpha}} f(x) d \theta(x)\right), & \alpha \neq 1  \tag{1.6}\\ \int_{\mathbb{S}^{n}} \log \left(u_{z}(x)\right) f(x) d \theta(x), & \alpha=1\end{cases}
$$

Here, $u_{z}(x):=\sup _{y \in \Omega}\langle y-z, x\rangle$ is the support function of $\Omega$ in direction $x$ with respect to $z_{0}$ and $\int_{\mathbb{\mathbb { S }}^{n}} h(x) d \theta(x)=\frac{1}{\omega_{n}} \int_{\mathbb{S}^{n}} h(x)$ with $\omega_{n}$ being the surface area of $\mathbb{S}^{n}$ and $\theta$ is the Lebesgue measure on $S^{n}$. When $\alpha=1$ and $f(x) \equiv 1$, then the above quantity agrees with the entropy in [26], first introduced by Firey [23] for the centrally symmetric $\Omega$. General integral quantities were studied by Andrews in [2, 4]. Here, we shall assume that $\int_{\mathbb{\mathbb { S }}^{n}} f(x) d \theta(x)=1$, namely, $\frac{1}{\omega_{n}} f(x) d \theta(x)$ is a probability measure. For the special case $f \equiv 1, \mathcal{E}_{\alpha, f}(\Omega)$ becomes the entropy $\mathcal{E}_{\alpha}(\Omega)$ in [6].

For positive $f \in C^{\infty}\left(\mathbb{S}^{n}\right)$, consider the anisotropic flow for convex hypersurfaces $\tilde{X}(\cdot, \tau): M_{\tau} \rightarrow \mathbb{R}^{n+1}:$

$$
\begin{equation*}
\frac{\partial}{\partial \tau} \tilde{X}(x, \tau)=-f^{\alpha}(v) \tilde{K}^{\alpha}(x, \tau) v(x, \tau) \tag{1.7}
\end{equation*}
$$

where $v(x, \tau)$ is the unit exterior normal at $\tilde{X}(x, \tau)$ of $\tilde{M}_{\tau}=\tilde{X}(M, \tau)$, and $\tilde{K}(x, \tau)$ is the Gauss curvature of $\tilde{M}_{\tau}$ at $\tilde{X}(x, \tau)$. Andrews [4] proved that flow (1.7) contracts to a point under finite time if the initial hypersurface $M_{0}$ is strictly convex. Under a proper normalization, the normalized anisotropy flow of (1.7) is

$$
\begin{equation*}
\frac{\partial}{\partial t} X(x, t)=-\frac{f^{\alpha}(v) K^{\alpha}(x, t)}{\int_{\mathbb{S}^{n}} f^{\alpha} K^{\alpha-1}} v(x, t)+X(x, t) . \tag{1.8}
\end{equation*}
$$

The basic observation is that a critical point for entropy $\mathcal{E}_{\alpha, f}(\Omega)$ defined in (1.5) under volume normalization is a solution to equation (1.1). The entropy is monotone along flow (1.8). One may view (1.1) is an "optimal solution" to this variational problem as the flow (1.8) provides a natural path to reach it. This approach was devised in [5] with the aim to obtain convergence of the normalized flow (1.8). The main arguments in [5] follows those in [6,26] where convergence of isotropic flows by power of Gauss curvature (i.e., $f=1$ ) was established. Unfortunately, the entropy point estimate in [ 6, 26] fails for general anisotropic flows except $\frac{1}{n+2}<\alpha \leq \frac{1}{n}$ [4]. The convergence was obtained in [5] assuming $M_{0}$ and $f$ are invariant under a subgroup $G$ of $O(n+1)$ which has no fixed point. We note that an inverse Gauss curvature flow argument was considered by Bryan, Ivaki, and Scheuer [14] to produce a origin-symmetric solution to (1.1).

Since we are only interested in finding a weak solution to (1.2), one only needs certain "weak" convergence of the flow (1.8). The key steps are to control diameter
with entropy under appropriate conditions on measure $\mu=f d \theta$ on $\mathbb{S}^{n}$ and use monotonicity of entropy to produce a solution to (1.2). The following is our main result.

Theorem 1.1 For $\alpha>\frac{1}{n+2}$ and finite nontrivial Borel measure $\mu$ on $\mathbb{S}^{n}, n \geq 1$, there exists a weak solution of (1.2) provided the following holds:
(i) If $\alpha>1$ and $\mu$ is not concentrated onto any great subsphere $x^{\perp} \cap \mathbb{S}^{n}, x \in \mathbb{S}^{n}$.
(ii) If $\alpha=1$ and $\mu$ satisfies that for any linear $\ell$-subspace $L \subset \mathbb{R}^{n+1}$ with $1 \leq \ell \leq n$, we have
(a) $\mu\left(L \cap \mathbb{S}^{n}\right) \leq \frac{\ell}{n+1} \cdot \mu\left(\mathbb{S}^{n}\right)$;
(b) equality in (a) for a linear $\ell$-subspace $L \subset \mathbb{R}^{n+1}$ with $1 \leq d \leq n$ implies the existence of a complementary linear $(n+1-\ell)$-subspace $\widetilde{L} \subset \mathbb{R}^{n+1}$ such that $\operatorname{supp} \mu \subset L \cup \widetilde{L}$.
(iii) If $\frac{1}{n+2}<\alpha<1$ and $d \mu=f d \theta$ for nonnegative $f \in L^{\frac{n+1}{n+2-\frac{1}{\alpha}}}\left(\mathbb{S}^{n}\right)$ with $\int_{\mathbb{S}^{n}} f>0$.

Let us briefly discuss what is known about uniqueness of the solution of the $L^{p_{-}}$ Minkowski problem (1.4). If $p>1$ and $p \neq n$, then Hug, Lutwak, Yang, and Zhang [30] proved that the Alexandrov solution of the $L^{p}$-Minkowski problem (1.4) is unique. However, if $p<1$, then the solution of the $L^{p}$-Minkowski problem (1.3) may not be unique even if $f$ is positive and continuous. Examples are provided by Chen, Li, and Zhu [17, 18] if $p \in[0,1)$, and Milman [42] shows that for any $C \in \mathcal{K}_{(0)}$, one finds $q \in$ $(-n, 1)$ such that if $p<q$, then there exist multiple solutions to the $L^{p}$-Minkowski problem (1.4) with $\mu=S_{C, p}$; or in other words, there exists $K \in \mathcal{K}_{(0)}$ with $K \neq C$ and $S_{K, p}=S_{C, p}$. In addition, Jian, Lu , and Wang [33] and Li, Liu, and Lu [37] prove that for any $p<0$, there exists positive even $C^{\infty}$ function $f$ with rotational symmetry such that the $L^{p}$-Minkowski problem (1.3) has multiple positive even $C^{\infty}$ solutions. We note that in the case of the centro-affine Minkowski problem $p=-n, \mathrm{Li}$ [36] even verified the possibility of existence of infinitely many solutions without affine equivalence, and Stancu [46] related unique solution in the cases $p=0$ and $p=-n$.

The case when $f$ is a constant function in the $L^{p}$-Minkowski problem (1.3) has received a special attention since [23]. When $p=-(n+1),(1.3)$ is self-similar solution of affine curvature flow. It is proved by Andrews that all solutions are centered ellipsoids. If $n=2$ and $p=2$, the uniqueness was proved by Andrews [3]. For general $n$ and $p>-(n+1)$, through the work of Lutwak [40], Guan-Ni [26], and Andrews, Guan, and Ni [6], Brendle, Choi, and Daskalopoulos [13] finally classified that the only solutions are centered balls. See also [21, 32, 44] for other approaches. Stability versions of these results have been obtained by Ivaki [31], but still no stability version is known in the case $p \in[0,1)$ if we allow any solutions of (1.3) not only even ones.

Concerning recent versions of the $L^{p}$-Minkowski problem, see [9].
The paper is structured as follows: The required diameter bounds are discussed in Section 2. Section 3 verifies the main properties of the Entropy, Section 4 proves our main result (Theorem 4.1) about flows, and finally Theorem 1.1 is proved in Section 5 via weak approximation.

## 2 Entropy and diameter estimates

For $\delta \in[0,1)$ and linear $i$-subspace $L$ of $\mathbb{R}^{n+1}$ with $1 \leq \operatorname{dim} L \leq n$, we consider the collar

$$
\Psi\left(L \cap \mathbb{S}^{n}, \delta\right)=\left\{x \in \mathbb{S}^{n}:\langle x, y\rangle \leq \delta \text { for } y \in L^{\perp} \cap \mathbb{S}^{n}\right\} .
$$

Let $B(1) \subset \mathbb{R}^{n+1}$ be the unit ball centered at the origin.
Theorem 2.1 Let $\alpha>\frac{1}{n+2}$, let $\int_{\mathbb{S}^{n}} f=1$ for a bounded measurable function $f$ on $\mathbb{S}^{n}$ with $\inf f>0$, and let $\Omega \subset \mathbb{R}^{n+1}$ be a convex body such that $|\Omega|=|B(1)|$ and $\operatorname{diam} \Omega=D$. For any $\delta, \tau \in(0,1)$, we have
(i) if $\alpha>1$, and $\int_{\bar{\Psi}\left(z^{\wedge} \cap \mathbb{S}^{n}, \delta\right)} f \leq 1-\tau$ for any $z \in S^{n}$, then

$$
\exp \left(\frac{\alpha-1}{\alpha} \mathcal{E}_{\alpha, f}(\Omega)\right) \geq \gamma_{1} \tau \delta^{1-\frac{1}{\alpha}} D^{1-\frac{1}{\alpha}},
$$

where $\gamma_{1}>0$ depends on $n$ and $\alpha$;
(ii) if $\alpha=1$, and

$$
f_{\Psi\left(L \cap \mathbb{S}^{n}, \delta\right)} f<\frac{(1-\tau) i}{n+1}
$$

for any linear $i$-subspace $L$ of $\mathbb{R}^{n+1}, i=1, \ldots, n$, then

$$
\mathcal{E}_{1, f}(\Omega) \geq \tau \log D+\log \delta-4 \log (n+1) ;
$$

(iii) if $\frac{1}{n+2}<\alpha<1, p=1-\frac{1}{\alpha}($ where $-n-1<p<0), \tau \leq \frac{1}{2} \int_{\mathbb{S}^{n}} f \cdot u^{1-\frac{1}{\alpha}}$ and

$$
\begin{equation*}
f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, \delta\right)} f^{\frac{n+1}{n+1+p}} \leq \tau^{\frac{n+1}{n+1+p}}, \tag{2.1}
\end{equation*}
$$

for any $z \in S^{n-1}$, then

$$
\text { either } D \leq 16 n^{2} / \delta^{2}, \text { or } D \leq\left(\frac{1}{2} f_{\mathbb{S}^{n}} f \cdot u^{1-\frac{1}{\alpha}}\right)^{\frac{2}{p}} \text {. }
$$

Moreover, if $\tau \leq \frac{1}{2} \exp \left(\frac{\alpha-1}{\alpha} \mathcal{E}_{\alpha, f}(\Omega)\right)$, then

$$
\text { either } D \leq 16 n^{2} / \delta^{2} \text {, or } D \leq\left(\frac{1}{2} \exp \left(\frac{\alpha-1}{\alpha} \varepsilon_{\alpha, f}(\Omega)\right)\right)^{\frac{2}{p}} \text {. }
$$

Remark 2.2 We note that for any $\alpha \geq 1$, bounded $f$ with $\inf f>0$ and $\int_{\mathbb{S}^{n}} f=1$, and $\tau \in(0,1)$, there exists $\delta \in(0,1)$ such that conditions in (i) and (ii) hold. In the case of $1>\alpha>\frac{1}{n+2}$, (iii) holds if in addition that $\tau \leq \frac{1}{2} \exp \left(\frac{1-\alpha}{\alpha} \varepsilon_{\alpha, f}(\Omega)\right)$ for the convex body $\Omega \subset \mathbb{R}^{n+1}$.

Proof Given $\alpha>\frac{1}{n+2}$, bounded $f$ with inf $f>0$ and $\int_{\mathbb{S}^{n}} f=1$, and $\tau \in(0,1)$, the existence of suitable $\delta \in(0,1)$ follows from the fact that the Lebesgue measure is a Borel measure.

Now, we assume that the conditions in (i)-(iii) hold. We may assume that the centroid of $\Omega$ is the origin; thus, Kannan, Lovász, and Simonovics [35] yield the existence of an $o$-symmetric ellipsoid such that

$$
\begin{equation*}
E \subset \Omega \subset(n+1) E \text {, and hence }-\Omega \subset(n+1) \Omega . \tag{2.2}
\end{equation*}
$$

Let $u$ be the support function of $\Omega$, and let $R=\max \{\|y\|: y \in \Omega\} \geq D / 2$ and $z_{0} \in \mathbb{S}^{n}$ such that $R z_{0} \in \partial \Omega$. We observe that the definition of the entropy yields

$$
\begin{aligned}
& f_{\mathbb{S}^{n}} f u^{1-\frac{1}{\alpha}} \leq \exp \left(\frac{1-\alpha}{\alpha} \mathcal{E}_{\alpha, f}(\Omega)\right) \text { if } \alpha>1 ; \\
& f_{\mathbb{S}^{n}} f \log u \leq \varepsilon_{0, f}(\Omega) ; \\
& f_{\mathbb{S}^{n}} f u^{1-\frac{1}{\alpha}} \geq \exp \left(\frac{1-\alpha}{\alpha} \varepsilon_{\alpha, f}(\Omega)\right) \text { if } \frac{1}{n+2}<\alpha<1 .
\end{aligned}
$$

Case 1: $\alpha>1$.
According to the condition in (i), we may choose $\zeta \in\{+1,-1\}$ such that

$$
f_{\Phi} f \geq \frac{\tau}{2} \text { for } \Phi=\left\{x \in \mathbb{S}^{n}:\left\langle x, \zeta z_{0}\right\rangle>\delta\right\}
$$

and hence $\frac{R \zeta z_{0}}{n+1} \in \Omega$ by (2.2). Since $u_{\sigma}(x) \geq\left\langle\frac{R \zeta z_{0}}{n+1}, x\right\rangle \geq \frac{R \delta}{n+1}$ for $x \in \Phi$, we have

$$
f_{\mathbb{S}^{n}} f u^{1-\frac{1}{\alpha}} \geq f_{\Phi} f\left(\frac{R \delta}{n+1}\right)^{1-\frac{1}{\alpha}} \geq \frac{\tau}{2} \cdot\left(\frac{D \delta}{2(n+1)}\right)^{1-\frac{1}{\alpha}}
$$

Case 2: $\alpha=1$.
To simplify notation, we consider the Borel probability measure $\mu(A)=\int_{\bar{A}} f$ on $S^{n}$. Let $e_{1}, \ldots, e_{n+1} \in \mathbb{S}^{n}$ be the principal directions associated with the ellipsoid $E$ in (2.2), and let $r_{1}, \ldots, r_{n+1}>0$ be the half axes of $E$ with $r_{i} e_{i} \in \partial E$ where we may assume that $r_{1} \leq \cdots \leq r_{n+1}$. In particular, (2.2) yields that

$$
\begin{equation*}
(n+1)^{n+1} \prod_{i=1}^{n+1} r_{i}=\frac{|(n+1) E|}{|B(1)|} \geq \frac{|\Omega|}{|B(1)|}=1 \tag{2.3}
\end{equation*}
$$

We observe that for any $v \in \mathbb{S}^{n}$, there exists $e_{i}$ such that $\left|\left\langle v, e_{i}\right\rangle\right| \geq \frac{1}{\sqrt{n+1}}>\frac{\delta}{n+1}$. For $i=1, \ldots, n+1$, we define

$$
B_{i}=\left\{v \in \mathbb{S}^{n}:\left|\left\langle v, e_{i}\right\rangle\right| \geq \frac{\delta}{n+1} \text { and }\left|\left\langle v, e_{j}\right\rangle\right|<\frac{\delta}{n+1} \text { for } j>i\right\} .
$$

In particular, $B_{i} \subset \Psi\left(L_{i} \cap \mathbb{S}^{n}, \delta\right)$ for $i=1, \ldots, n$ and $L_{i}=\operatorname{lin}\left\{e_{1}, \ldots, e_{i}\right\}$.
It follows that $\mathbb{S}^{n}$ is partitioned into the Borel sets $B_{1}, \ldots, B_{n+1}$, and as $B_{i} \subset \Psi\left(L_{i} \cap\right.$ $\left.\mathbb{S}^{n}, \delta\right)$ for $i=1, \ldots, n$, we have

$$
\begin{equation*}
\mu\left(B_{1}\right)+\cdots+\mu\left(B_{i}\right) \leq \frac{i(1-\tau)}{n+1} \text { for } i=1, \ldots, n, \tag{2.4}
\end{equation*}
$$

$$
\begin{equation*}
\mu\left(B_{1}\right)+\cdots+\mu\left(B_{n+1}\right)=1 . \tag{2.5}
\end{equation*}
$$

For $\zeta=\frac{1-\tau}{n+1}$, we have $0<\zeta<\frac{1}{n+1}$, and define

$$
\begin{gather*}
\beta_{i}=\mu\left(B_{i}\right)-\zeta \text { for } i=1, \ldots, n  \tag{2.6}\\
\beta_{n+1}=\mu\left(B_{n+1}\right)-\zeta-\tau \tag{2.7}
\end{gather*}
$$

where (2.4) and (2.5) yield

$$
\begin{gather*}
\beta_{1}+\cdots+\beta_{i} \leq 0 \text { for } i=1, \ldots, m-1,  \tag{2.8}\\
\beta_{1}+\cdots+\beta_{n+1}=0 . \tag{2.9}
\end{gather*}
$$

As $r_{i} e_{i} \in \Omega$, it follows from the definition of $B_{i}$ that $u(x) \geq\left\langle x, r_{i} e_{i}\right\rangle \geq r_{i} \cdot \frac{\delta}{n+1}$ for $x \in$ $B_{i}, i=1, \ldots, n+1$. We deduce from applying (2.3), (2.5)-(2.9), $r_{1} \leq \cdots \leq r_{n+1}$, and $\zeta<\frac{1}{n+1}$ that

$$
\begin{aligned}
\int_{\mathbb{S}^{n}} \log u d \mu & =\sum_{i=1}^{n+1} \int_{B_{i}} \log u d \mu \\
& \geq \sum_{i=1}^{n+1} \mu\left(B_{i}\right) \log r_{i}+\sum_{i=1}^{n+1} \mu\left(B_{i}\right) \log \frac{\delta}{n+1}=\sum_{i=1}^{n+1} \mu\left(B_{i}\right) \log r_{i}+\log \frac{\delta}{n+1} \\
& =\sum_{i=1}^{n+1} \beta_{i} \log r_{i}+\sum_{i=1}^{n+1} \zeta \log r_{i}+\tau \log r_{n+1}+\log \frac{\delta}{n+1} \\
& \geq \sum_{i=1}^{n+1} \beta_{i} \log r_{i}+\zeta \log \frac{1}{(n+1)^{n+1}}+\tau \log r_{n+1}+\log \frac{\delta}{n+1} \\
& =\left(\beta_{1}+\cdots+\beta_{n+1}\right) \log r_{n+1}+\sum_{i=1}^{n}\left(\beta_{1}+\cdots+\beta_{i}\right)\left(\log r_{i}-\log r_{i+1}\right) \\
& -(n+1) \zeta \log (n+1)+\tau \log r_{n+1}+\log \frac{\delta}{n+1} \\
& \geq-\log (n+1)+\tau \log r_{n+1}+\log \frac{\delta}{n+1} .
\end{aligned}
$$

Now, $D \leq(n+1) \operatorname{diam} E=2(n+1) r_{n+1} \leq(n+1)^{2} r_{n+1}$ and $\tau<1$, and hence

$$
\begin{aligned}
-\log (n+1)+\tau \log r_{n+1}+\log \frac{\delta}{n+1} & \geq-\log (n+1)+\tau \log \frac{D}{(n+1)^{2}}+\log \frac{\delta}{n+1} \\
& =\log \left(\delta D^{\tau}\right)-(2+2 \tau) \log (n+1) \\
& \geq \tau \log D+\log \delta-4 \log (n+1) .
\end{aligned}
$$

In particular, we conclude that

$$
\mathcal{E}_{1, f}(\Omega) \geq f_{\mathbb{S}^{n}} f \log u=\int_{\mathbb{S}^{n}} \log u d \mu \geq \tau \log D+\log \delta-4 \log (n+1) .
$$

Case 3: $\frac{1}{n+2}<\alpha<1$.
In this case, $-(n+1)<1-\frac{1}{\alpha}<0$. We may assume that

$$
D \geq 16 n^{2} / \delta^{2}
$$

and we consider

$$
\begin{aligned}
& \Phi_{0}=\left\{x \in \mathbb{S}^{n}: u(x)>\sqrt{2 R}\right\}, \\
& \Phi_{1}=\left\{x \in \mathbb{S}^{n}: u(x) \leq \sqrt{2 R}\right\} .
\end{aligned}
$$

Concerning $\Phi_{0}$, we have

$$
\begin{equation*}
f_{\Phi_{0}} f \cdot u^{1-\frac{1}{\alpha}} \leq(2 R)^{\frac{1}{2}\left(1-\frac{1}{\alpha}\right)} \int_{\Phi_{0}} f \leq D^{\frac{1}{2}\left(1-\frac{1}{\alpha}\right)}=D^{\frac{p}{2}} . \tag{2.10}
\end{equation*}
$$

On the other hand, we have $\pm \frac{R}{(n+1)} z_{0} \in \Omega$ by (2.2), thus any $x \in \Phi_{1}$ satisfies

$$
\sqrt{2 R} \geq u(x) \geq\left|\left\langle x, \frac{R}{n+1} z_{0}\right\rangle\right|
$$

and hence $\left|\left\langle x, z_{0}\right\rangle\right| \leq(n+1) \sqrt{\frac{2}{R}} \leq \frac{4 n}{\sqrt{D}} \leq \delta$; or in other words,

$$
\Phi_{1} \subset \Psi\left(z_{0}^{\perp} \cap \mathbb{S}^{n}, \delta\right)
$$

It follows from $|\Omega|=|B(1)|$ and the Blaschke-Santaló inequality (cf. [45]) that

$$
\int_{\mathbb{S}^{n}} u^{-(n+1)} \leq(n+1)|B(1)|=\omega_{n}, \text { and hence } \int_{\mathbb{S}^{n}} u^{-(n+1)} \leq 1 .
$$

For $p=1-\frac{1}{\alpha} \in(-n-1,0)$, Hölder's inequality and $\int_{\Phi_{1}} f^{\frac{n+1}{n+1+p}}<\tau^{\frac{n+1}{n+1+p}}$ yield

$$
f_{\Phi_{1}} f \cdot u^{1-\frac{1}{\alpha}} \leq\left(f_{\Phi_{1}} f^{\frac{n+1}{n+1+p}}\right)^{\frac{n+1+p}{n+1}}\left(f_{\Phi_{1}} u_{\sigma}^{-(n+1)}\right)^{\frac{|p|}{n+1}} \leq\left(f_{\Phi_{1}} f^{\frac{n+1}{n+1+p}}\right)^{\frac{n+1+p}{n+1}} \leq \tau .
$$

Finally, adding the last estimate to (2.10) yields

$$
\exp \left(\frac{\alpha-1}{\alpha} \mathcal{E}_{\alpha, f}(\Omega)\right) \leq f_{\mathbb{S}^{n}} f \cdot u^{1-\frac{1}{\alpha}} \leq D^{\frac{p}{2}}+\tau
$$

and hence the conditions either $\tau \leq \frac{1}{2} \int_{\mathbb{S}^{n}} f \cdot u^{1-\frac{1}{\alpha}}$ or $\tau \leq \frac{1}{2} \exp \left(\frac{1-\alpha}{\alpha} \mathcal{E}_{\alpha, f}(\Omega)\right)$ on $\tau$ implies (iii).

## 3 Anisotropic flows and monotonicity of entropies

The following theorem was proved by Andrews in [4] (see also for a discussion of contracting of non-homogeneous fully nonlinear anisotropic curvature flows in [24]).
Theorem 3.1 [4] For any $\alpha>0$ and positive $f \in C^{\infty}\left(\mathbb{S}^{n}\right)$ and any initial smooth, strictly convex hypersurface $\tilde{M}_{0} \subset \mathbb{R}^{n+1}$, the hypersurfaces $\tilde{M}_{\tau}$ given by the solution of (1.7) exist for a finite time $T$ and converge in Hausdorff distance to a point $p \in \mathbb{R}^{n+1}$ as $\tau$ approaches $T$.

Assuming

$$
f_{\mathbb{S}^{n}} f=1, \quad\left|\Omega_{0}\right|=|B(1)|,
$$

solution (1.7) yields a smooth convex solution to the normalized flow (1.8) with volume preserved.

Set

$$
\begin{equation*}
h_{z}(x, t) \doteqdot f(x) u_{z}^{-\frac{1}{\alpha}}(x, t) K(x, t), \quad d \sigma_{t}(x)=\frac{u_{z}(x, t)}{K(x, t)} d \theta(x) . \tag{3.1}
\end{equation*}
$$

Note that $\int_{\mathbb{S}^{n}} d \sigma_{t}(x)=\int_{\mathbb{S}^{n}} d \theta(x)=1$.

Since the un-normalized flow (1.7) shrinks to a point in finite time, we may assume that it is the origin. Then the support function $u(x, t)$ is positive for the normalized flow (1.8).

Lemma 3.2 (a) The entropy $\mathcal{E}_{\alpha, f}\left(\Omega_{t}\right)$ defined in (1.5) is monotonically decreasing,

$$
\begin{equation*}
\mathcal{E}_{\alpha, f}\left(\Omega_{t_{2}}\right) \leq \mathcal{E}_{\alpha, f}\left(\Omega_{t_{1}}\right), \quad \forall t_{1} \leq t_{2} \in[0, \infty) . \tag{3.2}
\end{equation*}
$$

(b) There is $D>0$ depending only on $\inf f, \sup f, \alpha, \Omega_{0}$ such that

$$
\begin{equation*}
\operatorname{diam} \Omega_{t}=D(t) \leq D, \forall t \geq 0 . \tag{3.3}
\end{equation*}
$$

(c) $\forall t_{0} \in[0, \infty)$,

$$
\begin{equation*}
\mathcal{E}_{\alpha, f}\left(\Omega_{t_{0}}, 0\right) \geq \mathcal{E}_{\alpha, f, \infty}+\int_{t_{0}}^{\infty}\left(\frac{\int_{\overline{\mathbb{S}}^{n}} h^{\alpha+1}(x, t) d \sigma_{t}}{\mathbb{\mathbb { S }}^{n} h(x, t) d \sigma_{t} \cdot \int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha}(x, t) d \sigma_{t}}-1\right) d t \tag{3.4}
\end{equation*}
$$

where

$$
h(x, t)=h_{0}(x, t), \mathcal{E}_{\alpha, f, \infty} \doteqdot \lim _{t \rightarrow \infty} \mathcal{E}_{\alpha, f}\left(\Omega_{t}\right)
$$

Proof (a) We follow argument in [26]. For each $T_{0}>$ fixed, pick $T>T_{0}$. Let $a^{T}=$ $\left(a_{1}^{T}, \ldots, a_{n+1}^{T}\right)$ be an interior point of $\Omega_{T}$. Set $u^{T}=u-e^{t-T} \sum_{i=1}^{n+1} a_{i}^{T} x_{i}$; it satisfies equation

$$
\begin{equation*}
\frac{\partial}{\partial t} u^{T}(x, t)=-\frac{f^{\alpha}(x) K^{\alpha}(x, t)}{\int_{\mathbb{\mathbb { S }}^{n}} f^{\alpha} K^{\alpha-1}}+u^{T}(x, t) . \tag{3.5}
\end{equation*}
$$

Note that since $a^{T}$ is an interior point of $\Omega_{T}$ and $u(x, T)$ is the support function of $\Omega_{T}$ with respect to $a^{T}, u^{T}(x, T)>0, \forall x \in \mathbb{S}^{n}$. We claim

$$
u^{T}(x, t)>0, \forall t \in[0, T) .
$$

Suppose $u^{T}\left(x_{0}, t^{\prime}\right) \leq 0$ for some $0<t^{\prime}<T, x_{0} \in \mathbb{S}^{n}$, and equation (3.5) implies $u^{T}\left(x_{0}, t\right)<0$ for all $t>t^{\prime}$, which contradicts to $u^{T}(x, T)>0$.

Set $a^{T}(t)=e^{t-T} a^{T}$. By the claim, $a^{T}(t)$ is in the interior of $\Omega_{t}, \forall t \leq T$. Denote

$$
d \sigma_{T, t}=u^{T}(x, t) K^{-1}(x, t) d \theta
$$

we rewrite equation (3.3) as

$$
\begin{equation*}
\frac{\partial}{\partial t} u_{a^{T}(t)}(x, t)=-\frac{f^{\alpha}(x) K^{\alpha}(x, t)}{\int_{\mathbb{S}^{n}} h_{a^{T}(t)}^{\alpha}(x, t) d \sigma_{T, t}}+u_{a^{T}(t)}(x, t) . \tag{3.6}
\end{equation*}
$$

We have

$$
\frac{\partial}{\partial t} \varepsilon_{\alpha, f}\left(\Omega_{t}, a^{T}(t)\right)=\frac{\int_{\mathbb{\mathbb { S }}^{n}} h_{a^{T}(t)}^{\alpha+1}(x, t) d \sigma_{T, t}}{\int_{\mathbb{\mathbb { S }}^{n}} h_{a^{T}(t)}(x, t) d \sigma_{T, t} \cdot \int_{\mathbb{\mathbb { S }}^{n}} h_{a^{T}(t)}^{\alpha}(x, t) d \sigma_{T, t}}+1
$$

Thus, $\forall t<T$,

$$
\begin{align*}
& \mathcal{E}_{\alpha, f}\left(\Omega_{t}, a^{T}(t)\right)-\mathcal{E}_{\alpha, f}\left(\Omega_{T}, a^{T}\right)  \tag{3.7}\\
& =\int_{t}^{T} f_{\mathbb{S}^{n}}\left(\frac{\int_{\mathbb{S}^{n}} h_{a^{T}(t)}^{\alpha+1}(x, t) d \sigma_{T, t}}{\int_{\mathbb{S}^{n}} h_{a^{T}(t)}(x, t) d \sigma_{T, t} \cdot \int_{\mathbb{S}^{n}} h_{a^{T}(t)}^{\alpha}(x, t) d \sigma_{T, t}}-1\right) d t \geq 0 .
\end{align*}
$$

Therefore,

$$
\mathcal{E}_{\alpha, f}\left(\Omega_{t}\right) \geq \mathcal{E}_{\alpha, f}\left(\Omega_{T}, a^{T}\right), \forall t<T
$$

Since $a^{T}$ is arbitrary, (3.2) is proved.
(b) The boundedness of $D(t)$ follows from Theorem 2.1 combined with the estimate $\mathcal{E}_{\alpha, 1}\left(\Omega_{t}\right) \leq \mathcal{E}_{\alpha, 1}(B(1))$ from (a) (see also [6,26]). The only nontrivial case is when $\frac{1}{n+2}<\alpha<1$ because we have to choose a $\tau$ independent of $t$. However, we may choose any $\tau \in(0,1)$ with $\tau \leq \frac{1}{2} \exp \left(\frac{1-\alpha}{\alpha} \varepsilon_{\alpha, f}(B(1))\right)$ according to $\mathcal{E}_{\alpha, 1}\left(\Omega_{t}\right) \leq$ $\mathcal{E}_{\alpha, 1}(B(1))$.
(c) $\forall \varepsilon>0, \forall t_{0}$ fixed, pick $T>T_{0}>t_{0}$. As $\mathcal{E}_{\alpha, f}\left(\Omega_{T}\right)$ is bounded by (a), $\exists a^{T}$ inside $\Omega_{T}$ such that $\mathcal{E}_{\alpha, f}\left(\Omega_{T}\right) \leq \mathcal{E}_{\alpha, f}\left(\Omega_{T}, a^{T}\right)+\varepsilon$. By (3.7),

$$
\begin{aligned}
& \mathcal{E}_{\alpha, f}\left(\Omega_{t_{0}}, a^{T}\left(t_{0}\right)\right)-\mathcal{E}_{\alpha, f}\left(\Omega_{T}\right) \\
& \geq \int_{t_{0}}^{T_{0}} f_{\mathbb{S}^{n}}\left(\frac{\int_{\mathbb{\mathbb { S }}^{n}} h_{a^{T}(t)}^{\alpha+1}(x, t) d \sigma_{T, t}}{\int_{\mathbb{S}^{n}} h_{a^{T}(t)}(x, t) d \sigma_{T, t} \cdot \int_{\mathbb{\mathbb { S }}^{n}} h_{a^{T}(t)}^{\alpha}(x, t) d \sigma_{T, t}}-1\right) d t-\varepsilon .
\end{aligned}
$$

As $\left|a^{T}\right| \leq D, \forall T$, let $T \rightarrow \infty$,

$$
a^{T}(t) \rightarrow 0, u^{T}(x, t) \rightarrow u(x, t), \text { uniformly for } 0 \leq t \leq T_{0}, x \in \mathbb{S}^{n} .
$$

We obtain $\forall t_{0}<T_{0}$,

$$
\mathcal{E}_{\alpha, f}\left(\Omega_{t_{0}}, 0\right)-\mathcal{E}_{\alpha, f, \infty} \geq \int_{t_{0}}^{T_{0}} f_{\mathbb{S}^{n}}\left(\frac{\int_{\overline{\mathbb{S}}^{n}} h^{\alpha+1}(x, t) d \sigma_{t}}{\widehat{\mathbb{S}}^{n} h(x, t) d \sigma_{t} \cdot \int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha}(x, t) d \sigma_{t}}-1\right) d t-\varepsilon .
$$

Then let $T_{0} \rightarrow \infty$, as $\varepsilon>0$ is arbitrary, we obtain (3.4).

## 4 Weak convergence

The goal of this section is to prove the following statement.
Theorem 4.1 For a $C^{\infty}$ function $f: \mathbb{S}^{n} \rightarrow(0, \infty)$ and $\alpha>\frac{1}{n+2}$ with $_{\mathbb{\mathbb { S }}^{n}} f=1$, there exist $\lambda>0$ and a convex body $\Omega \subset \mathbb{R}^{n+1}$ with $o \in \Omega$ whose support function $u$ is a (possibly weak) solution of the Monge-Ampère equation

$$
\begin{equation*}
u^{\frac{1}{\alpha}} \operatorname{det}\left(\bar{\nabla}_{i j}^{2} u+u \bar{g}_{i j}\right)=f \tag{4.1}
\end{equation*}
$$

and $\Omega$ satisfies that

$$
\begin{equation*}
\mathcal{E}_{\alpha, f}(\lambda \Omega) \leq \mathcal{E}_{\alpha, f}(B(1)), \quad|\lambda \Omega|=|B(1)|, \tag{4.2}
\end{equation*}
$$

where $C^{-1}<\lambda<C$ for a $C>1$ depending only on the $\alpha, \tau, \delta$ in Theorem 2.1 such that $f$ satisfies the conditions in Theorem 2.1.

From now on, we will assume that the $f$ in Theorem 4.1 satisfies the corresponding condition in Theorem 2.1 and $\Omega_{0}=B(1)$ in (1.8). We note that for any $z \in B(1), v_{z} \leq 2$ for the support function $v_{z}$ of $B(1)$ at $z$, and hence if $\alpha>\frac{1}{n+2}$, then

$$
\mathcal{E}_{\alpha, f_{k}}(B(1)) \leq\left\{\begin{align*}
\frac{\alpha}{\alpha-1} \cdot \log 2^{1-\frac{1}{\alpha}}, & \text { if } \alpha \neq 1,  \tag{4.3}\\
\log 2, & \text { if } \alpha=1 .
\end{align*}\right.
$$

The following is a consequence of Theorem 2.1 and Lemma 3.2.
Lemma 4.2 There exist $C_{\alpha, \tau, \delta}>0, D_{\alpha, \tau, \delta}>0$, and $c_{\alpha, \tau, \delta} \in \mathbb{R}$ depending only on constants $\alpha, \tau, \delta$ in Theorem 2.1 such that, along (1.8), we have

$$
\begin{equation*}
D(t) \leq D_{\alpha, \tau, \delta}, \mathcal{E}_{\alpha, f}\left(\Omega_{t}, 0\right) \geq c_{\alpha, \tau, \delta}, \frac{1}{C_{\alpha, \tau, \delta}} \leq f_{\mathbb{S}^{n}} h(x, t) d \sigma_{t} \leq C_{\alpha, \tau, \delta} . \tag{4.4}
\end{equation*}
$$

Proof For each $\alpha>\frac{1}{n+2}$ fixed with condition on $f$ as in Theorem 2.1, $\mathcal{E}_{\alpha, f}\left(\Omega_{t}\right)$ is bounded from below in terms of the diameter $D(t)$. Since $\left|\Omega_{t}\right|=|B(1)|$, we have $D(t) \geq 2$ by the Isodiametric Inequality (cf. [45]). By Theorem 2.1, $\mathcal{E}_{\alpha, f}\left(\Omega_{t}\right)$ is bounded from below by a constant $c_{\alpha, \tau, \delta}>0$, and hence $\mathcal{E}_{\alpha, f, \infty} \geq c_{\alpha, \tau, \delta}$. It follows from Lemma 3.2 that $\mathcal{E}_{\alpha, f}\left(\Omega_{t}\right) \leq \mathcal{E}_{\alpha, f}(B(1))$, and this estimate combined with (4.3) and Theorem 2.1 yields $D(t) \leq D_{\alpha, \tau, \delta}$ where $D_{\alpha, \tau, \delta}$ depends only on constants in condition on $f$ in Theorem 2.1. Finally, the inequalities follow from Lemma 3.2.

Set

$$
\begin{equation*}
\eta(t)=\int_{\mathbb{S}^{n}} h(x, t) d \sigma_{t} . \tag{4.5}
\end{equation*}
$$

We note that $\int_{\overline{\mathbb{S}}^{n}} h(x, t) d \sigma_{t}$ is monotone and bounded from below and above by Lemma 4.2, and hence we have

$$
\begin{equation*}
C_{\alpha, \tau, \delta} \geq \lim _{t \rightarrow \infty} f_{\mathbb{S}^{n}} h(x, t)=\eta \geq \frac{1}{C_{\alpha, \tau, \delta}} \tag{4.6}
\end{equation*}
$$

By Lemma 3.2 and Corollary 4.2,

$$
\begin{equation*}
\int_{0}^{\infty}\left(\frac{\int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha+1}(x, t) d \sigma_{t}}{\sqrt[\mathbb{S}]{ } h(x, t) d \sigma_{t} \cdot \int_{\mathbb{S}^{n}} h^{\alpha}(x, t) d \sigma_{t}}-1\right) d t<\infty . \tag{4.7}
\end{equation*}
$$

Since the integrand is nonnegative, $\exists t_{k} \rightarrow \infty$ such that

$$
\begin{equation*}
\lim _{k \rightarrow \infty}\left(\frac{\int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha+1}\left(x, t_{k}\right) d \sigma_{t_{k}}}{\overline{\mathbb{S}}^{n} h\left(x, t_{k}\right) d \sigma_{t_{k}} \cdot \int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha}\left(x, t_{k}\right) d \sigma_{t_{k}}}-1\right)=0 . \tag{4.8}
\end{equation*}
$$

This implies

$$
\begin{equation*}
\lim _{k \rightarrow \infty} \frac{\left(\int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha+1}\left(x, t_{k}\right) d \sigma_{t_{k}}\right)^{\frac{1}{1+\alpha}}}{\int_{\mathbb{\mathbb { S }}^{n}} h\left(x, t_{k}\right) d \sigma_{t_{k}}}=\lim _{k \rightarrow \infty} \frac{\left(\int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha+1}\left(x, t_{k}\right) d \sigma_{t_{k}}\right)^{\frac{\alpha}{1+\alpha}}}{\int_{\mathbb{\mathbb { S }}^{n}} h^{\alpha}\left(x, t_{k}\right) d \sigma_{t_{k}}}=1 . \tag{4.9}
\end{equation*}
$$

After considering a subsequence, we may assume that

$$
\begin{equation*}
\Omega_{t_{k}} \rightarrow \Omega, \quad u\left(x, t_{k}\right) \rightarrow u(x), \tag{4.10}
\end{equation*}
$$

where $u$ is the support function of $\Omega$. In view of (4.9) and (4.6),

$$
\begin{equation*}
\lim _{k \rightarrow \infty} f_{\mathbb{S}^{n}} h^{\alpha+1}\left(x, t_{k}\right) d \sigma_{t_{k}}=\eta^{1+\alpha}, \lim _{k \rightarrow \infty} f_{\mathbb{S}^{n}} h^{\alpha}\left(x, t_{k}\right) d \sigma_{t_{k}}=\eta^{\alpha} \tag{4.11}
\end{equation*}
$$

The following lemma is crucial for the weak convergence, which is a refined form of the classical Hölder inequality. ${ }^{1}$

Lemma 4.3 Let $p, q \in \mathbb{R}^{+}$with $\frac{1}{p}+\frac{1}{q}=1$, and $\operatorname{set} \beta=\min \left\{\frac{1}{p}, \frac{1}{q}\right\}$. Let $(M, \mu)$ be a measurable space; $\forall F \in L^{p}, G \in L^{q}$,

$$
\begin{equation*}
\int_{M}|F G| d \mu \leq\|F\|_{L^{p}}\|G\|_{L^{q}}\left(1-\beta \int_{M}\left(\frac{|F|^{\frac{p}{2}}}{\left(\int_{M}|F|^{p} d \mu\right)^{\frac{1}{2}}}-\frac{|G|^{\frac{q}{2}}}{\left(\int_{M}|G|^{q} d \mu\right)^{\frac{1}{2}}}\right)^{2}\right) \tag{4.12}
\end{equation*}
$$

Proof We first prove the following Claim. $\forall s, t \in \mathbb{R}$,

$$
\begin{equation*}
e^{\frac{s}{p}+\frac{t}{q}} \leq \frac{e^{s}}{p}+\frac{e^{t}}{q}-\beta\left(e^{\frac{s}{2}}-e^{\frac{t}{2}}\right)^{2} \tag{4.13}
\end{equation*}
$$

We may assume $t \geq s$, set $\tau=t-s$, and (4.13) is equivalent to

$$
\begin{equation*}
e^{\frac{\tau}{q}} \leq \frac{1}{p}+\frac{e^{\tau}}{q}-\beta\left(1-e^{\frac{\tau}{2}}\right)^{2}, \forall \tau \geq 0 . \tag{4.14}
\end{equation*}
$$

Set

$$
\xi(\tau)=\frac{1}{p}+\frac{e^{\tau}}{q}-\beta\left(1-e^{\frac{\tau}{2}}\right)^{2}-e^{\frac{\tau}{q}} .
$$

We have $\xi(0)=0$,

$$
\xi^{\prime}(\tau)=\frac{e^{\frac{\tau}{q}}}{q} \rho, \text { where } \rho(\tau)=e^{\frac{\tau}{p}}(1-\beta q)+q \beta e^{\frac{\tau}{\tau}-\frac{\tau}{q}}-1
$$

If $\beta=\frac{1}{q}$, then $\frac{1}{q} \leq \frac{1}{2}$; since $\tau \geq 0$,

$$
\rho(\tau)=e^{\frac{\tau}{p}}(1-\beta q)+q \beta e^{\frac{\tau}{2}-\frac{\tau}{q}}-1=e^{\frac{\tau}{2}-\frac{\tau}{q}}-1 \geq 0 .
$$

If $\beta=\frac{1}{p}$, then $\frac{1}{q} \geq \frac{1}{2}$; we have

$$
\begin{aligned}
\rho^{\prime}(\tau) & =e^{\frac{\tau}{p}}\left(\frac{1-\beta q}{p}+\beta q\left(\frac{1}{2}-\frac{1}{q}\right) e^{\frac{\tau}{2}-\frac{\tau}{q}}\right) \\
& \geq e^{\frac{\tau}{p}}\left(\frac{1-\beta q}{p}+\beta q\left(\frac{1}{2}-\frac{1}{q}\right)\right) \\
& \geq e^{\frac{\tau}{p}} \beta q\left(\frac{1}{2}-\frac{1}{p}\right) \geq 0 .
\end{aligned}
$$

We conclude that

$$
\rho(\tau) \geq 0, \forall \tau \geq 0
$$

In turn,

$$
\xi^{\prime}(\tau) \geq 0, \forall \tau \geq 0
$$

This yields (4.14) and (4.13). The Claim is verified.

[^1]Back to the proof of the lemma. We may assume

$$
F \geq 0, g \geq 0, \quad \int F^{p}>0, \int G^{q}>0
$$

Set

$$
e^{s}=\frac{F^{p}}{\int F^{p}}, \quad e^{t}=\frac{G^{q}}{\int G^{q}} .
$$

Put them into (4.13) and integrate, as $\frac{1}{p}+\frac{1}{q}=1$,

$$
\frac{\int F G}{\left(\int F^{p}\right)^{\frac{1}{p}}\left(\int G^{q}\right)^{\frac{1}{q}}} \leq\left(1-\beta \int\left(\frac{F^{\frac{p}{2}}}{\left(\int F^{p}\right)^{\frac{1}{2}}}-\frac{G^{\frac{q}{2}}}{\left(\int G^{q}\right)^{\frac{1}{2}}}\right)^{2}\right) .
$$

We prove weak convergence.
Proposition $4.4 \quad \forall \alpha>\frac{1}{n+2}$, suppose that (4.10) and (4.11) hold. Denote

$$
u_{k}=u\left(x, t_{k}\right), \sigma_{n, k}=\sigma_{n}\left(u_{i j}\left(x, t_{k}\right)+u\left(x, t_{k}\right) \delta_{i j}\right)
$$

Then

$$
\begin{equation*}
\lim _{k \rightarrow \infty} f_{\mathbb{S}^{n}}\left|u_{k}^{\frac{1}{\alpha}} \sigma_{n, k}-\frac{f}{\eta}\right| d \theta=0 \tag{4.15}
\end{equation*}
$$

where $\eta$ is defined in (4.5) which is bounded from below and above in (4.6). As a consequence, there is a convex body $\Omega \subset \mathbb{R}^{n+1}$ with $o \in \Omega$,

$$
|\Omega|=|B(1)|, \quad \mathcal{E}_{\alpha, f}\left(\Omega_{t}\right) \leq \mathcal{E}_{\alpha, f}(B(1)),
$$

and its support function $u$ satisfies

$$
\begin{equation*}
u^{\frac{1}{\alpha}} S_{\Omega}=\frac{1}{\eta} f d \theta \tag{4.16}
\end{equation*}
$$

Proof We only need to verify (4.15). By (4.11), it is equivalent to prove

$$
\begin{equation*}
\lim _{k \rightarrow \infty} f_{\mathbb{S}^{n}}\left|u_{k}^{\frac{1}{\alpha}} \sigma_{n, k}-f \eta^{-1}\left(t_{k}\right)\right| d \theta=0 \tag{4.17}
\end{equation*}
$$

Since $D\left(t_{k}\right)$ is bounded,

$$
\begin{aligned}
f_{\mathbb{S}^{n}} u_{k}^{\frac{1}{\alpha^{2}}} \sigma_{n, k} d \theta \leq & \left(D\left(t_{k}\right)\right)^{\frac{1}{\alpha^{2}}} f_{\mathbb{S}^{n}} u_{k}^{\frac{1}{\alpha^{2}}} \sigma_{n, k} d \theta \leq\left(D\left(t_{k}\right)\right)^{\frac{1}{\alpha^{2}}}\left|\partial \Omega_{t_{k}}\right| \leq C . \\
\int_{\mathbb{S}^{n}}\left|u_{k}^{\frac{1}{\alpha}} \sigma_{n, k}-f \eta^{-1}\left(t_{k}\right)\right| d \theta & =f_{\mathbb{S}^{n}}\left|\frac{f}{\eta\left(t_{k}\right) u_{k}^{\frac{1}{\alpha}} \sigma_{n, k}}-1\right| u_{k}^{\frac{1}{\alpha}} \sigma_{n, k} d \theta \\
& \leq\left(f_{\mathbb{S}^{n}}\left|\frac{f}{\eta\left(t_{k}\right) u_{k}^{\frac{1}{\alpha}} \sigma_{n, k}}-1\right|^{1+\alpha} d \sigma_{t_{k}}\right)^{\frac{1}{1+\alpha}}\left(f_{\mathbb{S}^{n}} u_{k}^{\left(\frac{1}{\alpha}-1\right) \frac{1+\alpha}{\alpha}} d \sigma_{t_{k}}\right)^{\frac{\alpha}{1+\alpha}} \\
& =\left(f_{\mathbb{S}^{n}}\left|\frac{f}{\eta\left(t_{k}\right) u_{k}^{\frac{1}{\alpha}} \sigma_{n, k}}-1\right|^{1+\alpha} d \sigma_{t_{k}}\right)^{\frac{1}{1+\alpha}}\left(f_{\mathbb{S}^{n}} u_{k}^{\frac{1}{\alpha^{2}}} \sigma_{n, k} d \theta\right)^{\frac{\alpha}{1+\alpha}} \\
& \leq C\left(f_{\mathbb{S}^{n}}\left|f \eta^{-1}\left(t_{k}\right) u_{k}^{-\frac{1}{\alpha}} \sigma_{n, k}^{-1}-1\right|^{1+\alpha} d \sigma_{t_{k}}\right)^{\frac{1}{1+\alpha}} .
\end{aligned}
$$

By (4.8), (4.11), and Lemma 4.3, with $p=\alpha+1, F^{\frac{1}{1+\alpha}}=h\left(x, t_{k}\right), G=1$,

$$
\begin{equation*}
\lim _{k \rightarrow \infty} f\left(\left(\frac{h\left(x, t_{k}\right)}{\eta\left(t_{k}\right)}\right)^{\frac{1+\alpha}{2}}-1\right)^{2} d \sigma_{t_{k}}=0 \tag{4.19}
\end{equation*}
$$

For $t_{k}$ fixed, let

$$
\gamma_{k}(x)=f \eta^{-1}\left(t_{k}\right) u_{k}^{-\frac{1}{\alpha}} \sigma_{n, k}^{-1}=h\left(x, t_{k}\right) \eta^{-1}\left(t_{k}\right)
$$

and set

$$
\Sigma_{k}=\left\{x \in \mathbb{S}^{n}| | \gamma_{k}(x)-1 \left\lvert\, \leq \frac{1}{2}\right.\right\}
$$

It is straightforward to check that $\exists A_{\alpha} \geq 1$ depending only on $\alpha$ such that

$$
\begin{gathered}
A_{\alpha}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right| \geq\left|\gamma_{k}(x)-1\right|, \forall x \in \Sigma_{k} \\
A_{\alpha}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{2} \geq\left|\gamma_{k}(x)-1\right|^{1+\alpha}, \forall x \in \Sigma_{k}^{c}
\end{gathered}
$$

Since $\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right| \leq 2^{1+\alpha}, \forall x \in \Sigma_{k}$, let $\delta=\min \{1+\alpha, 2\}$,

$$
\begin{aligned}
f_{\mathbb{S}^{n}}\left|\gamma_{k}(x)-1\right|^{1+\alpha} d \sigma_{t_{k}} & =\frac{1}{\omega_{n}}\left(\int_{\Sigma_{k}}\left|\gamma_{k}(x)-1\right|^{1+\alpha} d \sigma_{t_{k}}+\int_{\Sigma_{k}^{c}}\left|\gamma_{k}(x)-1\right|^{1+\alpha} d \sigma_{t_{k}}\right) \\
& \leq \frac{A_{\alpha}^{1+\alpha}}{\omega_{n}}\left(\int_{\Sigma_{k}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{1+\alpha} d \sigma_{t_{k}}+\int_{\Sigma_{k}^{c}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{2} d \sigma_{t_{k}}\right) \\
& \leq \frac{\left(2 A_{\alpha}\right)^{1+\alpha}}{\omega_{n}}\left(\int_{\Sigma_{k}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{\delta} d \sigma_{t_{k}}+\int_{\Sigma_{k}^{c}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{2} d \sigma_{t_{k}}\right) \\
& \leq\left(2 A_{\alpha}\right)^{1+\alpha}\left(\int_{\mathbb{S}^{n}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{\delta} d \sigma_{t_{k}}+\int_{\mathbb{S}^{n}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{2} d \sigma_{t_{k}}\right) \\
& \leq\left(2 A_{\alpha}\right)^{1+\alpha}\left(\left(f_{\mathbb{S}^{n}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{2} d \sigma_{t_{k}}\right)^{\frac{\delta}{2}}+\int_{\mathbb{S}^{n}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{2} d \sigma_{t_{k}}\right) .
\end{aligned}
$$

By (4.19),

$$
\lim _{k \rightarrow \infty} f_{\mathbb{S}^{n}}\left|\gamma_{k}^{\frac{1+\alpha}{2}}(x)-1\right|^{2} d \sigma_{t_{k}}=0
$$

Hence,

$$
\begin{equation*}
\lim _{k \rightarrow \infty} f_{\mathbb{S}^{n}}\left|\gamma_{k}(x)-1\right|^{1+\alpha} d \sigma_{t_{k}}=0 \tag{4.20}
\end{equation*}
$$

Now, (4.17) follows from (4.18)-(4.20).
Proof Proof of Theorem 4.1. It follows from Proposition 4.4 after a proper rescaling as $\eta$ satisfies (4.6) and (4.16).

## 5 The general Monge-Ampère equations - proof of Theorem 1.1

In order to prove Theorem 1.1, we need weak approximation in the following sense.

Lemma 5.1 For $\delta, \varepsilon \in\left(0, \frac{1}{2}\right)$ and a Borel probability measure $\mu$ on $\mathbb{S}^{n}, n \geq 1$, there exists a sequence $d \mu_{k}=\frac{1}{\omega_{n}} f_{k} d \theta$ of Borel probability measures whose weak limit is $\mu$ and $f_{k} \in C^{\infty}\left(\mathbb{S}^{n}\right)$ satisfies $f_{k}>0$ and the following properties:
(i) If $\mu\left(\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, 2 \delta\right)\right) \leq 1-\varepsilon$ for any $z \in S^{n-1}$, then

$$
\begin{equation*}
f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, \delta\right)} f_{k} \leq 1-\varepsilon \text { for any } z \in S^{n-1} \tag{5.1}
\end{equation*}
$$

(ii) If $\mu\left(\Psi\left(L \cap \mathbb{S}^{n}, 2 \delta\right)\right)<(1-2 \delta) \cdot \frac{\ell}{n+1}$ for any linear $\ell$-subspace $L$ of $\mathbb{R}^{n+1}, \ell=$ $1, \ldots, n$, then

$$
\begin{equation*}
\mu_{k}\left(\Psi\left(L \cap \mathbb{S}^{n}, \delta\right)\right)<(1-\delta) \cdot \frac{\ell}{n+1} . \tag{5.2}
\end{equation*}
$$

(iii) If $d \mu=\frac{1}{\omega_{n}} f d \theta$ for $f \in L^{r}\left(\mathbb{S}^{n}\right)$ where $r>1$, and

$$
\begin{equation*}
f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, 2 \delta\right)} f^{r} \leq \varepsilon \tag{5.3}
\end{equation*}
$$

for any $z \in S^{n-1}$, then

$$
\begin{equation*}
\int_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, \delta\right)} f_{k}^{r} \leq 2^{r} \varepsilon \text { for any } z \in S^{n-1} \tag{5.4}
\end{equation*}
$$

Proof For $k \geq 1$, let $\left\{B_{k, i}\right\}_{i=1, \ldots, m(k)}$ be a partition of $S^{n}$ into spherically convex Borel measurable sets $B_{k, i}$ with $\operatorname{diam} B_{k, i} \leq \frac{1}{k}$ and $\theta\left(B_{k, i}\right)>0$. For each $B_{k, i}$, we choose a $C^{\infty}$ function $h_{k, i}: \mathbb{S}^{n} \rightarrow[0, \infty)$ such that for $M_{k, i}=\max h_{k, i}$ and the probability measure $d \tilde{\theta}=\frac{1}{\omega_{n}} d \theta$, we have:

- $h_{k, i}=0$ if $x \notin B_{k, i}$;
- $M_{k, i} \leq\left(1+\frac{1}{k}\right) \cdot \frac{\mu\left(B_{k, i}\right)}{\bar{\theta}\left(B_{k, i}\right)}$;
- $\theta\left(\left\{x \in B_{k, i}: h_{k, i}(x)<M_{k, i}\right\}\right)<\frac{1}{k} \theta\left(B_{k, i}\right)$;
- $\int_{B_{k, i}} h_{k, i} d \tilde{\theta}=\mu\left(B_{k, i}\right)$.

We consider the positive $C^{\infty}$ function $\tilde{f}_{k}=\frac{1}{k}+\sum_{i=1}^{m(k)} h_{k, i}$, and hence $f_{k}=\left(\int_{\overline{\mathbb{S}}^{n}} \tilde{f}_{k}\right)^{-1} \tilde{f}$ satisfies that the probability measure $d \mu_{k}=f_{k} d \tilde{\theta}$ tends weakly to $\mu$, and for large $k \geq$ $1 / \delta, \mu_{k}$ satisfies (i), and if (ii) holds, then $\mu_{k}$ also satisfies (5.2).

Turning to (iii), we assume that $d \mu=f d \tilde{\theta}$ for $f \in L^{r}\left(\mathbb{S}^{n}\right)$ where $r>1$, and $f$ satisfies (5.3). For any large $k$ and $i=1, \ldots, m(k)$, we deduce from the Hölder inequality that

$$
\begin{aligned}
f_{B_{k, i}} \tilde{f}_{k}^{r} & =f_{B_{k, i}}\left(h_{k, i}+\frac{1}{k}\right)^{r} \leq 2^{r-1} f_{B_{k, i}} h_{k, i}^{r}+2^{r-1} f_{B_{k, i}} \frac{1}{k^{r}} \\
& \leq 2^{r-1} \tilde{\theta}\left(B_{k, i}\right) M_{k, i}^{r}+2^{r-1} f_{B_{k, i}} \frac{1}{k^{r}} \\
& \leq 2^{r-1}\left(1+\frac{1}{k}\right)^{r} \tilde{\theta}\left(B_{k, i}\right)\left(\frac{\int_{B_{k, i}} f}{\tilde{\theta}\left(B_{k, i}\right)}\right)^{r}+2^{r-1} f_{B_{k, i}} \frac{1}{k^{r}} \\
& \leq 2^{r-1}\left(1+\frac{1}{k}\right)^{r} f_{B_{k, i}} f^{r}+2^{r-1} f_{B_{k, i}} \frac{1}{k^{r}} .
\end{aligned}
$$

Summing this estimate up for large $k$ and all $B_{k, i}$ with $B_{k, i} \cap \Psi\left(z^{\perp} \cap \mathbb{S}^{n}, \delta\right) \neq \varnothing$, and using that $\int_{\mathbb{S}^{n}} \tilde{f}_{k} \geq 2^{-1 / 2}$ for large $k$, we deduce that

$$
f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, \delta\right)} f_{k}^{r} \leq \sqrt{2} f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, \delta\right)} \tilde{f}_{k}^{r} \leq \sqrt{2} \cdot 2^{r-1}\left(1+\frac{1}{k}\right)^{r} f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, 2 \delta\right)} f^{r}+\sqrt{2} \cdot \frac{2^{r-1}}{k^{r}} \leq 2^{r} \varepsilon .
$$

For $\alpha>0$ and $p=1-\frac{1}{\alpha}$, the $L^{p}$-surface area $d S_{\Omega, p}=u^{1-p} d S_{\Omega}$ was introduced in the seminal works [39-41] for a convex body $\Omega \subset \mathbb{R}^{n+1}$ with $o \in \Omega$ and support function $u$. Since the surface area measure is weakly continuous for $p<1$, and if $K \subset \mathbb{R}^{n+1}$ is an at most $n$-dimensional compact convex set, then $S_{K, p} \equiv 0$ for $p<1$, we have the following statement.
Lemma 5.2 If convex bodies $\Omega_{m} \subset \mathbb{R}^{n+1}$ tend to a compact convex set $K \subset \mathbb{R}^{n+1}$ where $o \in \Omega_{m}, K$, and $\lim \inf _{m \rightarrow \infty} S_{\Omega_{m}, p}>0$, then $\operatorname{int} K \neq \varnothing$ and $S_{\Omega_{m}, p}$ tends weakly to $S_{K, p}$.

For the reader's sake, let us recall Theorem 1.1.
Theorem 5.3 For $\alpha>\frac{1}{n+2}$ and finite nontrivial Borel measure $\mu$ on $\mathbb{S}^{n}, n \geq 1$, there exists a weak solution of (1.2) provided the following holds:
(i) If $\alpha>1$ and $\mu$ is not concentrated onto any great subsphere $x^{\perp} \cap \mathbb{S}^{n}, x \in \mathbb{S}^{n}$.
(ii) If $\alpha=1$ and $\mu$ satisfies that for any linear $\ell$-subspace $L \subset \mathbb{R}^{n+1}$ with $1 \leq \ell \leq n$, we have:
(a) $\mu\left(L \cap \mathbb{S}^{n}\right) \leq \frac{\ell}{n+1} \cdot \mu\left(\mathbb{S}^{n}\right)$;
(b) equality in (a) for a linear $\ell$-subspace $L \subset \mathbb{R}^{n+1}$ with $1 \leq d \leq n$ implies the existence of a complementary linear $(n+1-\ell)$-subspace $\widetilde{L} \subset \mathbb{R}^{n+1}$ such that $\operatorname{supp} \mu \subset L \cup \widetilde{L}$.
(iii) If $\frac{1}{n+2}<\alpha<1$, assume $d \mu=f d \theta$ for nonnegative $f \in L^{\frac{n+1}{n+2-\frac{1}{\alpha}}}\left(\mathbb{S}^{n}\right)$ with $\int_{\mathbb{S}^{n}} f>0$.

Proof Let $\alpha>\frac{1}{n+2}$. After rescaling, we may assume that the $\mu$ in (1.2) is a probability measure. We consider the sequence $d \mu_{k}=\frac{1}{\omega_{n}} f_{k} d \theta$ of Lemma 5.1 of Borel probability measures whose weak limit is $\mu$ and $f_{k} \in C^{\infty}\left(\mathbb{S}^{n}\right)$ satisfies $f_{k}>0$. For each $f_{k}$, let $\Omega_{k} \subset$ $\mathbb{R}^{n+1}$ be the convex body with $o \in \Omega_{k}$ provided by Theorem 4.1 whose support function $u_{k}$ is the solution of the Monge-Ampère equation

$$
\begin{equation*}
u_{k}^{\frac{1}{\alpha}} d S_{\Omega_{k}}=f_{k} d \theta ; \tag{5.5}
\end{equation*}
$$

$\exists \lambda_{k}>0$ under control, with $\left|\lambda_{k} \Omega\right|=|B(1)|, \Omega_{k}$ satisfies that

$$
\begin{equation*}
\mathcal{E}_{\alpha, f_{k}}\left(\lambda_{k} \Omega_{k}\right) \leq \mathcal{E}_{\alpha, f_{k}}(B(1)) . \tag{5.6}
\end{equation*}
$$

We also need the observations that

$$
\begin{equation*}
\left|\Omega_{k}\right|=\frac{1}{n+1} \int_{\mathbb{S}^{n}} u_{k} d S_{\Omega_{k}} \tag{5.7}
\end{equation*}
$$

and if $p=1-\frac{1}{\alpha}$, then

$$
\begin{equation*}
S_{\Omega_{k}, p}\left(\mathbb{S}^{n}\right)=\int_{\mathbb{S}^{n}} u_{k}^{1-\frac{1}{\alpha}} d S_{\Omega_{k}}=\omega_{n} f_{\mathbb{S}^{n}} f_{k}=\omega_{n} \tag{5.8}
\end{equation*}
$$

We claim that if there exists $\Delta>0$ depending on $n, \alpha$, and $\mu$ such that

$$
\begin{equation*}
\operatorname{diam} \Omega_{k} \leq \Delta, \text { then Theorem } 5.3 \text { holds. } \tag{5.9}
\end{equation*}
$$

To prove this claim, we note that (5.9) yields the existence of a subsequence of $\left\{\Omega_{k}\right\}$ tending to a compact convex set $\Omega$ with $o \in \Omega$, which is a convex body by (5.8) and Lemma 5.2. Moreover, Lemma 5.2 also yields that $\Omega$ is an Alexandrov solution of (1.2), verifying the claim (5.9).

We divide the rest of the argument verifying Theorem 5.3 into three cases.
Case 1: $\alpha>1$.
Since $\mu$ is not concentrated to any great subsphere, there exist $\delta \in\left(0, \frac{1}{2}\right)$ depending on $\mu$ such that $\mu\left(\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, 2 \delta\right)\right) \leq 1-2 \delta$ for any $z \in S^{n-1}$. It follows from Lemma 5.1 that we may assume that

$$
\begin{equation*}
f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, \delta\right)} f_{k} \leq 1-\delta \text { for any } z \in S^{n-1} \tag{5.10}
\end{equation*}
$$

Now, Theorem 4.1 implies that $\lambda_{k} \geq c$ for a constant $c>0$ depending on $n, \delta$, and $\alpha$, and in turn Theorem 4.1, (4.3), and $\frac{1}{\alpha}-1<0$ yield that

$$
\mathcal{E}_{\alpha, f}\left(\Omega_{k}\right)=\frac{\alpha}{\alpha-1} \cdot \log \lambda_{k}^{\frac{1}{\alpha}-1}+\mathcal{E}_{\alpha, f}\left(\lambda_{k} \Omega_{k}\right) \leq \frac{\alpha}{\alpha-1} \cdot \log \lambda_{k}^{\frac{1}{\alpha}-1}+\mathcal{E}_{\alpha, f}(B(1)) \leq C
$$

for a constant $C>0$ depending on $n, \delta$, and $\alpha$. Therefore, Theorem 2.1 and (5.10) imply that the sequence $\left\{\Omega_{k}\right\}$ is bounded, and in turn the claim (5.9) implies Theorem 5.3 if $\alpha>1$.

Case 2: $\alpha=1$.
The argument is by induction on $n \geq 0$ where we do not put any restriction on the probability measure $\mu$ in the case $n=0$. For the case $n=0$, we observe that any finite measure $\mu$ on $S^{0}$ can be represented in the form $d \mu=u d S_{\Omega}$ for a suitable segment $\Omega \subset \mathbb{R}^{1}$.

For the case $n \geq 1$, assuming that we have verified Theorem 5.3(ii) in smaller dimensions, we consider a Borel measure probability $\mu$ on $S^{n}$ satisfying (a) and (b).
Case 2.1: There exists a linear $\ell$-subspace $L \subset \mathbb{R}^{n+1}$ with $1 \leq \ell \leq n$ and $\mu\left(L \cap \mathbb{S}^{n}\right)=\frac{\ell}{n+1}$. $\mu\left(\mathbb{S}^{n}\right)$.

Let $\widetilde{L} \subset \mathbb{R}^{n+1}$ be the complementary linear ( $n+1-\ell$ )-subspace with supp $\mu \subset L \cup$ $\widetilde{L}$, and hence $\mu\left(\widetilde{L} \cap \mathbb{S}^{n}\right)=\frac{n+1-\ell}{n+1} \cdot \mu\left(\mathbb{S}^{n}\right)$. It follows by induction that there exist an $\ell$-dimensional compact convex set $K^{\prime} \subset L$ and an $(n+1-\ell)$-dimensional compact convex set $\widetilde{K}^{\prime} \subset \widetilde{L}$ such that $\mu\left\llcorner\left(L \cap S^{n}\right)=\ell V_{K^{\prime}}\right.$ and $\mu\left\llcorner\left(\widetilde{L} \cap S^{n}\right)=(n+1-\ell) V_{\widetilde{K}^{\prime}}\right.$. Finally, for $K=\widetilde{L}^{\perp} \cap\left(K^{\prime}+L^{\perp}\right)$ and $\widetilde{K}=L^{\perp} \cap\left(\widetilde{K}^{\prime}+\widetilde{L}^{\perp}\right)$, there exist $\alpha, \tilde{\alpha}>0$ such that

$$
\mu=(n+1) V_{\alpha K+\tilde{\alpha} \widetilde{K}} .
$$

Case 2.2: $\mu\left(L \cap \mathbb{S}^{n}\right)<\frac{\ell}{n+1} \cdot \mu\left(\mathbb{S}^{n}\right)$ for any linear $\ell$-subspace $L \subset \mathbb{R}^{n+1}$ with $1 \leq \ell \leq n$.
It follows by a compactness argument that there exists $\delta \in\left(0, \frac{1}{2}\right)$ depending on $\mu$ such that $\mu\left(\Psi\left(L \cap \mathbb{S}^{n}, 2 \delta\right)\right)<(1-2 \delta) \cdot \frac{\ell}{n+1}$ for any linear $\ell$-subspace $L$ of $\mathbb{R}^{n+1}$, $\ell=1, \ldots, n$. We consider the sequence of probability measures $d \mu_{k}=\frac{1}{\omega_{n}} f_{k} d \theta$ of

Lemma 5.1 tending weakly to $\mu$ such that $f_{k}>0, f_{k} \in C^{\infty}\left(\mathbb{S}^{n}\right)$, and

$$
\begin{equation*}
\mu_{k}\left(\Psi\left(L \cap \mathbb{S}^{n}, \delta\right)\right)<(1-\delta) \cdot \frac{\ell}{n+1} \tag{5.11}
\end{equation*}
$$

for any linear $\ell$-subspace $L$ of $\mathbb{R}^{n+1}, \ell=1, \ldots, n$.
For each $f_{k}$, let $\Omega_{k} \subset \mathbb{R}^{n+1}$ with $o \in \Omega_{k}$ be the convex body provided by Theorem 4.1 whose support function $u_{k}$ is the solution of the Monge-Ampère equation (4.1) and satisfies (4.2) with $f=f_{k}$ and $\lambda=\lambda_{k}$ where $|B(1)|=\left|\lambda_{k} \Omega_{k}\right|$ for $\lambda_{k}>0$, and

$$
\begin{aligned}
\left|\Omega_{k}\right| & =\frac{1}{n+1} \int_{\mathbb{S}^{n}} u_{k} \operatorname{det}\left(\bar{\nabla}_{i j}^{2} u_{k}+u_{k} \bar{g}_{i j}\right) d \theta=\frac{\omega_{n}}{n+1} f_{\mathbb{S}^{n}} u_{k} \operatorname{det}\left(\bar{\nabla}_{i j}^{2} u_{k}+u_{k} \bar{g}_{i j}\right) \\
& =|B(1)| f_{\mathbb{S}^{n}} f_{k}=|B(1)|
\end{aligned}
$$

and hence $\lambda_{k}=1$. In particular, (4.3) yields

$$
\mathcal{E}_{1, f_{k}}\left(\lambda_{k} \Omega_{k}\right) \leq \mathcal{E}_{1, f_{k}}(B(1)) \leq \log 2 .
$$

Since $\mathcal{E}_{1, f_{k}}\left(\Omega_{k}\right)$ is bounded, (5.11) and Theorem 2.1 imply that the sequence $\Omega_{k}$ stays bounded, as well. Therefore, the claim (5.9) yields Theorem 5.3 if $\alpha=1$.
Case 3: $\frac{1}{n+2}<\alpha<1$.
We set $p=1-\frac{1}{\alpha} \in(-n-1,0)$ and $r=\frac{n+1}{n+1+p}>1$, and

$$
\begin{equation*}
\tau=\frac{1}{2} \cdot 2^{-\frac{|p|(n+1)}{|p|+n}} \tag{5.12}
\end{equation*}
$$

and choose $\delta \in\left(0, \frac{1}{2}\right)$ such that

$$
f_{\Psi\left(z^{\perp} \cap \mathbb{S}^{n}, 2 \delta\right)} f^{r} \leq \frac{\tau^{r}}{2^{r}}
$$

for any $z \in S^{n-1}$. We deduce from Lemma 5.1 that if $z \in S^{n-1}$, then

$$
\begin{equation*}
f_{\Psi\left(z^{\wedge} \cap \mathbb{S}^{n}, \delta\right)} f_{k}^{r} \leq \tau^{r} . \tag{5.13}
\end{equation*}
$$

We deduce from (5.5), (5.7), and $\left|\lambda_{k} \Omega_{k}\right|=|B(1)|=\frac{\omega_{n}}{n+1}$ that

$$
\begin{equation*}
f_{\mathbb{S}^{n}} u_{k}^{p} f_{k}=\frac{n+1}{\omega_{n}} \int_{\mathbb{S}^{n}} u_{k} d S_{\Omega_{k}}=\frac{n+1}{\omega_{n}}\left|\Omega_{k}\right|=\lambda_{k}^{-n-1} . \tag{5.14}
\end{equation*}
$$

In particular, (4.3) and the upper bound on the entropy yield that

$$
\begin{align*}
2^{p} & \leq \exp \left(p \cdot \varepsilon_{\alpha, f_{k}}(B(1))\right) \leq \exp \left(p \cdot \varepsilon_{\alpha, f}\left(\lambda_{k} \Omega_{k}\right)\right) \leq f_{\mathbb{S}^{n}}\left(\lambda_{k} u_{k}\right)^{p} f_{k} \\
& =\lambda_{k}^{p} \int_{\mathbb{S}^{n}} u_{k} d S_{\Omega_{k}}=\lambda_{k}^{p-n} \cdot \frac{n+1}{\omega_{n}} \cdot\left|\lambda_{k} \Omega_{k}\right|=\lambda_{k}^{p-n} . \tag{5.15}
\end{align*}
$$

It follows from (5.15) that $\lambda_{k} \leq 2^{\frac{|p|}{|p|+n}}$, and in turn (5.14) yields that

$$
f_{\mathbb{S}^{n}} u_{k}^{p} f_{k} \geq 2^{-\frac{\mid p(n+1)}{|p|+n}} .
$$

Therefore, $\tau \leq \frac{1}{2} \int_{\mathbb{S}^{n}} u_{k}^{p} f_{k}$ (cf. (5.12)), (5.13), and Theorem 2.1 yield that the sequence $\left\{\Omega_{k}\right\}$ is bounded, and in turn the claim (5.9) implies Theorem 5.3 if $\frac{1}{n+2}<\alpha<1$.

## References

[1] M. Aldaz, A stability version of Hölder's inequality. J. Math. Anal. Appl. 343(2008), 842-852.
[2] B. Andrews, Monotone quantities and unique limits for evolving convex hypersurfaces. Int. Math. Res. Not. IMRN 20(1997), 1001-1031.
[3] B. Andrews, Gauss curvature flow: the fate of rolling stone. Invent. Math. 138(1999), 151-161.
[4] B. Andrews, Motion of hypersurfaces by gauss curvature. Pacific J. Math. 195(2000), no. 1, 1-34.
[5] B. Andrews, K. Böröczky, P. Guan and L. Ni, Entropy and anisotropic flow by power of Gauss curvature, preprint, 2018.
[6] B. Andrews, P. Guan, and L. Ni, Flow by the power of the gauss curvature. Adv. Math. 299(2016), 174-201.
[7] G. Bianchi, K. J. Böröczky, and A. Colesanti, Smoothness in the L ${ }^{p}$ Minkowski problem for $p<1$. J. Geom. Anal. 30(2020), 680-705.
[8] G. Bianchi, K. J. Böröczky, A. Colesanti, and D. Yang, The L ${ }^{p}$-Minkowski problem for $-n<p<1$ according to Chou-Wang. Adv. Math. 341(2019), 493-535.
[9] K. J. Böröczky, The logarithmic Minkowski conjecture and the $L^{p}$-Minkowski problem. In: A. Koldobsky and A. Volberg (eds.), Harmonic analysis and convexity, DeGryuter, preprint, 2022. arxiv:2210.00194
[10] K. J. Böröczky and P. Hegedűs, The cone volume measure of antipodal points. Acta Math. Hungar. 146(2015), 449-465.
[11] K. J. Böröczky, E. Lutwak, D. Yang, and G. Zhang, The logarithmic Minkowski problem. J. Amer. Math. Soc. 26(2013), no. 3, 831-852.
[12] K. J. Böröczky and H. T. Trinh, The planar $L^{p}$-Minkowski problem for $0<p<1$. Adv. Appl. Math. 87(2017), 58-81.
[13] S. Brendle, K. Choi, and P. Daskalopoulos, Asymptotic behavior of flows by powers of the Gaussian curvature. Acta Math. 219(2017), 1-16.
[14] P. Bryan, M. Ivaki, and J. Scheuer, A unified flow approach to smooth, even Lp-Minkowski problems. Analysis \& PDE 12(2019), 259-280.
[15] L. Caffarelli, A localization property of viscosity solutions to the Monge-Ampère equation and their strict convexity. Ann. Math. 131(1990), no. 1, 129-134.
[16] L. Caffarelli, Interior $W^{2, p}$ estimates for solutions of the Monge-Ampère equation. Ann. Math. 131(1990), no. 1, 135-150.
[17] S. Chen, Q.-R. Li, and G. Zhu, On the L ${ }^{p}$ Monge-Ampère equation. J. Differential Equations 263(2017), 4997-5011.
[18] S. Chen, Q.-R. Li, and G. Zhu, The logarithmic Minkowski problem for non-symmetric measures. Trans. Amer. Math. Soc. 371(2019), 2623-2641.
[19] K. Choi, M. Kim, and T. Lee, Curvature bound for L ${ }^{p}$ Minkowski problem, preprint, 2023. arXiv:2304.11617
[20] K.-S. Chou and X.-J. Wang, The L ${ }^{p}$-Minkowski problem and the Minkowski problem in centroaffine geometry. Adv. Math. 205(2006), no. 1, 33-83.
[21] G. Crasta and I. Fragalá, Variational worn stones, preprint, 2023. arXiv:2303.11764
[22] S.-Z. Du, On the planar L ${ }^{p}$-Minkowski problem, J. Differential Equations 287 (2021), 37-77.
[23] W.-J. Firey, On the shapes of worn stones. Mathematika 21(1974), 1-11.
[24] P. Guan, J. Huang, and J. Liu, Non-homogeneous fully nonlinear contracting flows of convex hypersurfaces, to appear in Adv. Nonlinear Stud. (special issue in honour of Joel Spruck).
[25] P. Guan and C. S. Lin, On equation $\operatorname{det}\left(u_{i j}+\delta_{i j} u\right)=u^{p} f$ on $S^{n}$, preprint, 1999.
[26] P. Guan and L. Ni, Entropy and a convergence theorem for gauss curvature flow in high dimension. J. Eur. Math. Soc. 19(2017), no. 12, 3735-3761.
[27] Q. Guang, Q.-R. Li, and X.-J. Wang, The L ${ }^{p}$-Minkowski problem with super-critical exponents, preprint, 2022. arXiv:2203.05099
[28] Q. Guang, Q.-R. Li, and X.-J. Wang, Existence of convex hypersurfaces with prescribed centroaffine curvature.
https://person.zju.edu.cn/person/attachments/2022-02/01-1645171178-851572.pdf
[29] D. Hug, Contributions to affine surface area. Manuscripta Math. 91(1996), 283-301.
[30] D. Hug, E. Lutwak, D. Yang, and G. Zhang, On the $L_{p}$-Minkowski problem for polytopes. Discrete Comput. Geom. 33(2005), 699-715.
[31] M. Ivaki, On the stability of the $L^{p}$-curvature. J. Funct. Anal. 283(2022), 109684.
[32] M. Ivaki and E. Milman, Uniqueness of solutions to a class of isotropic curvature problems, preprint, 2023. arXiv:2304.12839
[33] H. Jian, J. Lu, and X.-J. Wang, Nonuniqueness of solutions to the $L^{p}$-Minkowski problem. Adv. Math. 281(2015), 845-856.
[34] H. Jian, J. Lu, and G. Zhu, Mirror symmetric solutions to the centro-affine Minkowski problem. Calc. Var. 55(2016), Article no. 41, 22 pp.
[35] R. Kannan, L. Lovász, and M. Simonovits, Isoperimetric problems for convex bodies and a localization lemma. Discrete Comput. Geom. 13(1995), 541-559.
[36] Q.-R. Li, Infinitely many solutions for the centro-affine Minkowski problem. Int. Math. Res. Not. IMRN 2019(2019), 5577-5596.
[37] Q.-R. Li, J. Liu, and J. Lu, Non-uniqueness of solutions to the dual $L^{p}$-Minkowski problem. Int. Math. Res. Not. IMRN 2022(2022), 9114-9150.
[38] M. Ludwig, General affine surface areas. Adv. Math. 224(2010), 2346-2360.
[39] E. Lutwak, Selected affine isoperimetric inequalities. In: Handbook of convex geometry, North-Holland, Amsterdam, 1993, pp. 151-176.
[40] E. Lutwak, The Brunn-Minkowski-Firey theory. I. Mixed volumes and the Minkowski problem. J. Differ. Geom. 38(1993), 131-150.
[41] E. Lutwak, The Brunn-Minkowski-Firey theory. II. Affine and geominimal surface areas. Adv. Math. 118(1996), 244-294.
[42] E. Milman, A sharp centro-affine isospectral inequality of Szegő-Weinberger type and the $L^{p}$-Minkowski problem. J. Diff. Geom., preprint, 2022. arXiv:2103.02994
[43] C. Saroglou, A non-existence result for the $L^{p}$-Minkowski problem, preprint, 2021. arXiv:2109. 06545
[44] C. Saroglou, On a non-homogeneous version of a problem of Firey. Math. Ann. 382(2022), 1059-1090.
[45] R. Schneider, Convex bodies: the Brunn-Minkowski theory, Encyclopedia of Mathematics, 44, Cambridge University Press, Cambridge, 1993.
[46] A. Stancu, Prescribing centro-affine curvature from one convex body to another. Int. Math. Res. Not. IMRN 2022(2022), 1016-1044.
[47] G. Zhu, The $L^{p}$-Minkowski problem for polytopes for $p<0$. Indiana Univ. Math. J. 66(2017), 1333-1350.

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