The electronic structure of manganese doped ceramic material Ba_{0.6}Sr_{0.4}TiO₃

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1. Introduction

The electronic structure of Mn doped $Ba_{0.6}Sr_{0.4}TiO_3$ ceramic materials has been determined by analyzing the low loss region of the electron energy loss spectrum in a transmission electron microscope. We calculate the interband transition strength which arises from maxima in the valence-band density of states via the dielectric function by a Kramers-Kronig transformation in the bulk and at segregated grain boundaries. The matrix and segregated boundaries are different in this regard.

2. Experiment

TEM specimens of 5 mol% Mn doped $Ba_{0.6}Sr_{0.4}TiO_3$ were prepared by the conventional method of mechanical thinning followed by argon ion beam milling (GATAN PIPS). The measurements were performed with an Enfina PEELS system attached to a VG HB601UX FEG scanning transmission electron microscope (STEM) operating at 100keV. The PEELS spectra were corrected for dark current and readout noise and the channel to channel gain variation. Under these conditions, the energy resolution was 0.45eV because of the cold field emission gun, the energy dispersion was 0.05eV and the probe size was 0.8nm for the line scan measurement across the segregated boundaries. The grain boundary segregation was measured using STEM and EDX.

In order to obtain the single scattering distributions S(E), the spectra acquired with the PEELS spectrometer were Fourier-Log deconvoluted to remove plural effects and then normalized using the method described by Egerton [1] to obtain the energy loss function $Im(-1/\epsilon)$. The real and imaginary parts of the dielectric function were obtained, after removing surface loss effects, by an FFT-based Kramers-Kronig program. From the dielectric function the interband transition strength J_{cv} was also calculated [2].

3. Results and discussion

Mn segregated boundaries were randomly selected, Fig. 1. The interband transition strength from matrix and boundaries is shown in Fig. 2. Because the dielectric constant changes at the grain boundary, the grain boundary might deserve to be regarded as a thin fillet of different material associated with a surface plasmon, although the material on either side has the same dielectric contant. The correct Kramers-Kronig Transformation is given by French [3] for boundary analyses. The interband transition strength presents a steep increase at an energy loss of about 2-3eV that may be attributed to the energy gap [4]. The interband transition strength from the segregated boundary is different from that of the bulk. There is no outstanding feature except for 2 peaks. Other notable features in the bulk (A, B and C), especially for the transitions with low energy near the band gap, become weaker at segregated grain boundaries. At A and B, there are 3 obvious peaks in the interband transition strength in the matrix. They are much weaker at the boundaries. Furthermore the two small peaks at B join into one broad peak. The characteristics of the matrix and boundaries are yet to be determined, but chemistry is an obvious candidate for the

electronic structure changes. As shown in our previous results on spectral boundaries in pure SrTiO₃, geometry is another possible cause to change the electronic structure [5]. However we found it happens only on some spectral boundaries. The chemical segregation is a more universal reason, since it is common to all segregated boundaries.

4. Conclusion

We have calculated the interband transition strength of Ba_{0.6}Sr_{0.4}TiO₃ from the matrix and from segregated boundaries. It is proved that PEELS is a powerful analytical technique for the structure, bonding and electronic properties analysis of materials.

- 5. References
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- 6. Acknowledgements

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Fig.1. EDX linescan cross the manganese segregated grain boundary



Fig.2. Interband band strength for Ba_{0.6}Sr_{0.4}TiO₃ from matrix (a) and Mn segregated boundary (b)