## **ON THE NON-MINIMAL MARTIN BOUNDARY POINTS**

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Dedicated to Professor KIYOSHI NOSHIRO on his 60th birthday

1. In a Green space<sup>1)</sup>  $\hat{\mathcal{Q}}$  we can introduce Martin's topology and make it the Martin space<sup>2)</sup>  $\hat{\mathcal{Q}}$ .  $\hat{\mathcal{Q}}$  is a dense open subset of  $\hat{\mathcal{Q}}$  and the kernel

$$K(p, x) = \begin{cases} \frac{G(p, x)}{G(p, y_0)} & p \neq y_0 \\ 0 & p = y_0 \neq x \\ 1 & p = y_0 = x \end{cases}$$

can be extended continuously to  $(p, x) \in \hat{\mathcal{Q}} \times \mathcal{Q}$ , where G(p, x) is a Green function in  $\mathcal{Q}$  and  $y_0$  the fixed point of  $\mathcal{Q}$ .  $\hat{\mathcal{Q}}$  is a metric space.  $\mathcal{A} = \hat{\mathcal{Q}} - \mathcal{Q}$  is divided into two disjoint subsets  $\mathcal{A}_0$ ,  $\mathcal{A}_1$  and  $s \in \mathcal{A}_1$  is characterized by the fact that K(s, x) is a minimal positive harmonic function<sup>3)</sup> in  $x \in \mathcal{Q}$ .

2. We shall show the following theorem:

THEOREM. No point of  $\Delta_0$  is an isolated point.

**Proof.** Let  $\omega$  be an open subset of  $\mathcal{Q}$ ,  $\langle x_n \rangle$  (n = 1, 2, ...) be a sequence of points in  $\omega$  such that  $x_n \to x_0 \in \mathcal{A}$ . If we denote by  $\mathscr{H}$  the family of positive superharmonic functions in  $\mathcal{Q}$ , each of which dominates  $K(x_n, y)$  on  $\mathcal{Q} - \omega$ , then  $\inf_{v \in \mathscr{K}} v(y)$  is equal to the positive superharmonic function except a polar set. We shall write this superharmonic function  $\mathscr{C}_{K_n}^{\omega}(y)$ .

In this case

Received June 17, 1966.

<sup>&</sup>lt;sup>1)</sup> M. Brelot, G. Choquet, Espaces et lignes de Green. Annales Inst. Fourier 3 (1951), pp. 199-263.

<sup>&</sup>lt;sup>2)</sup> M. Brelot, Le problème de Dirichlet. Axiomatique et frontière de Martin. Journal de Math. 35 (1956), pp. 297-335 (pp. 329-330). Cf. also R. S. Martin, Minimal positive harmonic functions, Trans. Amer. Math. Soc., 49 (1941), pp. 137-172. M. Parreau, Sur les moyennes des fonctions harmoniques et analytiques et la classification des surfaces de Riemann, Annales Inst. Fourier 3 (1952), pp. 103-197. L. Naïm, Sur le rôle de la frontière de R. S. Martin dans la théorie du potentiel, Annales Inst. Fourier 7 (1957), pp. 183-281.
<sup>3)</sup> R. S. Martin, loc. cit., p. 137.

$$\mathscr{C}^{\omega}_{K_n}(y) = \int K(x, y) \, d\mu_n(x)$$

where  $\mu_n$  is a positive mass-distribution on  $\overset{*}{\omega} \cap \mathcal{Q}$  and the total mass of  $\mu_n$  does not exceed 1,  $\overset{*}{\omega}$  being the boundary of  $\omega$  in  $\hat{\mathcal{Q}}$ . By the theorem of choice, we can extract from  $\{\mu_n\}$  the subsequence  $\{\mu'_n\}$  such that  $\mu'_n$  converges vaguely to  $\mu$  and the carrier of  $\mu$  is contained in  $\overline{\mathfrak{D} \cap \mathcal{Q}}$ .

$$v(y) = \int K(x, y) \, d\mu(x)$$

is a positive superharmonic function in 2, and we have

$$\mathscr{C}^{\omega}_{K_{x_0}}(y) \leq v(y).$$

In fact, for fixed  $y \in Q$  and r > 0 we shall denote by  $\varepsilon_y''$  the mass-distribution which can be obtained after sweeping out the unit mass on y to the exterior of the sphere (circle) of radius r and with center y. Then

$$U^{r}(x) = \int K(x, z) d\varepsilon_{y}^{\prime r}(z)$$

is bounded and continuous on  $\hat{Q}$ . Therefore

$$\lim_{n\to\infty}\int U^r(x)\,d\mu'_n(x)=\int U^r(x)\,d\mu(x).$$

By reciprocal law

$$\lim_{n\to\infty}\int \mathscr{C}^{\omega}_{K_{\infty n'}}(z)\,d\varepsilon_{y}^{\prime r}(z)=\int v(z)\,d\varepsilon_{y}^{\prime r}(z)$$

and by Fatou's lemma

$$\int \mathscr{C}^{\omega}_{K_{x_0}}(z) \, d\varepsilon_{y}^{\prime \prime}(z) \leq \underline{\lim_{n \to \infty}} \int \mathscr{C}^{\omega}_{K_{x_n}}(z) \, d\varepsilon_{y}^{\prime \prime}(z) = \int v(z) \, d\varepsilon_{y}^{\prime \prime}(z).$$

By making  $r \to 0$  we can get for each  $y \in \Omega$ 

$$\mathscr{C}^{\omega}_{K_{x_{b}}}(y) \leq v(y).$$

If  $\mu_{l'}$  denotes the restriction of  $\mu$  to  $\Delta$  and  $\mu_2$  the restriction of  $\mu$  to Q, then

$$v(y) = \int K(x, y) d\mu_1(x) + \int K(x, y) d\mu_2(x)$$
  
=  $u(y) + w(y)$ 

where u is harmonic and w is a potential, and this is just the Riesz decomposi-

tion.

From now on let  $x_0$  be a point of  $\Delta_0$  and  $x_0$  be isolated. Let

$$K(x_0, y) = \int_{\Delta_1} K(x, y) d\nu(x)$$

be the canonical representation<sup>4)</sup> of  $K(x_0, y)$ . Then we can find a neighbourhood  $\hat{\delta}$  of  $x_0$  such that

$$\widehat{\delta} \cap \Delta_0 = \{x_0\}$$

and

(2) 
$$\nu(\underline{J}_1 - \overline{\hat{\delta}}) > 0.$$

If we set  $\omega = \hat{\delta} \cap \Omega$ , then

$$\mathscr{C}^{\omega}_{K_{x_0}}(y) = \int_{\Delta_1} \mathscr{C}^{\omega}_{K_x}(y) d\nu(x)$$
$$= \int_{\Delta_1 - \overline{\delta}} \mathscr{C}^{\omega}_{K_x}(y) d\nu(x) + \int_{\overline{\delta}} \mathscr{C}^{\omega}_{K_x}(y) d\nu(x)$$

The first term of the last side is harmonic, because  $\omega$  is thin at each point of  $\mathcal{A}_1 - \overline{\delta}^{(5)}$  and therefore we can get  $\mathscr{C}^{\omega}_{K_x}(y) \equiv K(x, y)$ .

We note that  $\mu_1$  is the restriction of the mass-distribution  $\mu$  to  $(\overline{\hat{\delta}} \cap \Omega) \cap A$ , which is contained in  $\overline{\hat{\delta}} \cap A$  and does not contain the point  $x_0$ . By (1) we can get  $\mu_1(A_0) = 0$ , that is,  $\mu_1$  is the canonical mass-distribution of u, and by (2)

$$u_1(y) = \int_{\Delta_1 - \overline{\delta}} K(x, y) d\nu(x) > 0.$$

Since

$$v(y) = u(y) + w(y)$$
  

$$\geq \mathscr{C}^{\omega}_{K_{x_0}}(y) \geq \int_{\delta_1 - \overline{\delta}} \mathscr{C}^{\omega}_{K_x}(y) d\nu(x) = \int_{\delta_1 - \overline{\delta}} K(x, y) d\nu(x) = u_1(y)$$

and u is the greatest harmonic minorant of v, we have

 $u \ge u_1$ ,

but the canonical mass-distribution of u has the carrier in  $\Delta_1 \cap \overline{\delta}$ , whereas the canonical mass-distribution of  $u_1$  has positive mass only in  $\Delta_1 - \overline{\delta}$ . As  $u_1 > 0$ 

<sup>&</sup>lt;sup>4)</sup> R. S. Martin, loc. cit., p. 157.

<sup>5)</sup> L. Naïm, loc. cit., p. 203 (théorème 3) and p. 205 (théorème 5).

the canonical mass-distribution of u has positive mass in  $\Delta_1 - \overline{\hat{\delta}}$ ; this is a contradiction. Q.E.D.

COROLLARY. If  $\Delta_0 \neq \phi$  then  $\Delta_0$  contains at least countable points.

*Remark.* In the above consideration we rely upon the following argument: we have always  $u \ge u_1$ , and, if  $u_1 > 0$ , then  $\mu_1$  is not the canonical mass-distribution of u.

Mr. K. Matsumoto has kindly pointed out the following result:

Let  $x_0$  be a point of  $\Delta_0$  and  $\nu$  be the canonical mass-distribution of  $K_{x_0}(y)$ , then the common part of the carrier of  $\nu$  with  $\Delta_1$  is contained in  $\overline{\Delta}_0$ .

The proof follows from the preceding remark; let E be the carrier of  $\nu$ . If  $E \cap \varDelta_1 \oplus \overline{\varDelta}_0$ , then there exist a point  $z_0$  and a set A satisfying the following conditions:

- 1) A is an open neighbourhood of  $z_0$  in  $\Delta$ ,
- 2)  $\nu(A) > 0$ ,
- 3)  $A \subset \underline{A}_1$ .

We can construct an open set (in  $\hat{Q}$ )  $G : G \cap \Delta = A$ . In this case, there exist two positive numbers  $0 < \rho_1 < \rho$  such that:

- i) dist  $(z_0, x_0) > \rho$ ,
- ii)  $\overline{U_{\rho}(z_0)} \subset G^{6_1}$ ,
- iii)  $\nu(U_{\rho_1}(z_0) \cap \varDelta) > 0.$

If we set  $\omega = \Omega - \overline{U_{\rho}(z_0)}$ , under the same notations as in the proof of the above theorem, we see from i),  $x_0 \in \overline{\omega}$  and, as  $(\overline{\mathfrak{T} \cap \Omega}) \cap \Delta \subset \overline{U_{\rho}(z_0)} \cap \Delta \subset G \cap \Delta = A \subset \Delta_1$ ,  $\mu_1$  is canonical and from  $\Delta - \overline{\omega} \supset U_{\rho_1}(z_0) \cap \Delta$ ,  $\mu_1 > 0$ , this is a contradiction.

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<sup>&</sup>lt;sup>6)</sup> We denote  $U_{\rho}(z_0) = \{x \in \hat{Q}; \text{ dist } (x, z_0) < \rho\}$ , where the metric dist  $(x, z_0)$  is the Martin's metric.