

MASS TRANSFER BETWEEN BINARY STARS

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The transfer of mass from one component of a binary system to its companion can involve ejection of the mass by radiation pressure, by Roche lobe overflow, by centrifugal action through the coupling of the upper atmosphere to the stellar rotation through magnetic fields, and by the dynamics of the atmospheric fluid under the action of an energy source at the photosphere or lower chromosphere. We will primarily address the latter mechanism, which commonly goes under the title of stellar winds. The major determinant of a stellar wind is the energy source. In the case of a binary system, the dynamic effects of the system on the flow pattern, particularly for close binaries, is a second major determinant, once an energy source is available to drive the wind. Roche lobe overflow might be thought of as a limiting case in which the system dynamics predominate.

In addition to ejection, the transfer of mass requires capture of the ejected matter by the companion. The first issue is free particle vs. continuum flow. We will only consider continuum flow, as the applicability of the free particle concept to any binary system with significant mass transfer is doubtful, because of the scale height and because of the likely presence of magnetic fields in the interstellar plasma. The need to consider continuum flow in the capture process complicates the analysis considerably, particularly for close binaries where the coriolis force enters strongly with the collective motion of the plasma. We have not solved this problem, which requires multi-dimensional flow analysis with only one simplifying symmetry. However, for more distant binaries, rather simple considerations show that the capture of matter for continuum flow can be enhanced over that for free particle flow.

Now for the principal topic, mass ejection through fluid dynamics. As shown by Parker years ago, a static stellar atmosphere is unstable, provided the temperature is high enough. The corona of the sun, $>1,000,000^{\circ}\text{K}$, flows outward in a continuous, though highly variable wind as predicted by Parker and discovered by the Mariner II Spacecraft. It seems clear from the work of Parker and others that any star with a hot corona will have a stellar wind.

The existence of the hot corona itself once appeared paradoxical, since its energy presumably comes from the solar interior through the relatively cold photosphere and chromosphere to the corona. The spontaneous flow of heat from the cold photosphere to the hot corona directly violates the Second Law of Thermodynamics. The resolution of the paradox, of course, is that the energy does not flow as heat, but rather as mechanical energy that is dissipated into heat in the chromosphere and corona. In the solar corona only a fraction of the mechanical energy appears as heat--most shows up in the kinetic energy of bulk motion of the solar wind. The stagnation temperature of the relatively cold solar wind is $\approx 30,000,000^{\circ}\text{K}$.

For the sun, the mechanical energy flow is in the form of acoustic waves generated by the deep convection zone. As the waves move outward their amplitude increases until dissipation occurs. This phenomenon of the increase in amplitude of waves traveling from heavier to lighter media is widespread. It accounts for the breaking of waves on a beach, the heating of the ionosphere by thunder, and the cracking of a whip. In each case the wave, traveling with constant energy, is propagated by decreasing mass, which can only contain the energy by increasing the amplitude of its motion. A sufficient amplitude leads to dissipation so that small disturbances in a dense medium can produce great heating in a contiguous medium of lesser density.

For single main sequence stars, acoustic waves from deep convection zones have had great success in explaining heating of the upper atmosphere. However, stellar winds can only be observed for the sun, so that the accepted model, based on the systematic decrease in chromospheric emissions and increase in stellar rotation from late to early type stars is inferential.

In studying binaries we have a richer field. It is possible to observe the effects of winds, if they exist, as they impinge on companions. There are sources of disturbance in the lower atmosphere other than acoustic noise associated with deep convection zones that can drive stellar winds, for example, g-mode oscillations and x-ray heating. We will conclude by describing some results of model calculations for binary systems showing how relatively small disturbances in the lower chromosphere or photosphere can cause substantial mass ejection in binary systems.

The model is a straightforward numerical integration of the equations of motion, including the energy equation, with an initial static atmosphere and with various temperature fluctuations imposed at the base of the corona. The initial static atmosphere is isothermal in the corona, with a linear temperature profile in the chromosphere. The flow is calculated along the line of centers between the two binary components, in the rotating frame of the system, so that the centrifugal force is included. The coriolis force is not included, so that the present model is not valid for close binaries with large orbital angular velocities. Energy flows by convection, by thermal conduction, and by propagation of longitudinal pressure waves.

There are two ad hoc assumptions: The flow is one dimensional in a diverging/converging flux tube configured to produce radial flow at

each star, and dissipation of waves is included as a coefficient multiplied by the velocity gradient. The flux tube parameters can be adjusted to correspond to the capture cross-section, once that problem is solved. Meanwhile, it permits the examination of the effect of constricting the flow on the ejection of mass. The dissipation coefficient is set so that waves from the base of the atmosphere are dissipated totally in the lower corona, thereby providing the energy to heat the corona and drive the stellar wind.

The model has been applied to several situations: using a linear temperature pulse it has been shown that a substantial rise in the temperature at the base of the corona can account for the mass ejections needed to cause the x-ray turn-ons observed on HZ-Herculis. It is hypothesized that G-mode oscillations may account for the temperature fluctuations through the amplitude amplification mechanism described above. By starting the disturbance further down in the atmosphere the model shows that rather small temperature fluctuations lead to large mass ejections.

An application that tends to validate the model itself is the quiet sun. Beginning with a static chromosphere and corona, a periodic oscillation in the lower corona of about 10% of the static value initiates mass ejection that grows into a typical quiet-time solar wind in a few hours.

If the oscillations were initiated as temperature fluctuations, the fluctuation amplitude in the photosphere would be $\pm 500^{\circ}\text{K}$. The corresponding velocity field is ± 4 km/sec., about what is observed in the photosphere. The fluctuating quantity used to initiate the oscillations is not very important, only introducing a phase shift in the overall motion.

For sources of disturbance other than convection zones (say, x-ray heating), the altitude of the initiating disturbance may vary. This altitude is very important in determining the disturbance amplitude required for a given velocity of ejection, since for disturbances higher in the chromosphere, the amplification with outward propagation is reduced.

There is also a considerable effect of altitude on mass loss, both because of the amplification and because there is less mass in the atmosphere at high altitudes. For the HZ-Herculis x-ray turn-ons, a disturbance at the base of the corona must be of the order of the coronal temperature in amplitude, whereas at lower altitude 10% disturbances will suffice.

The dissipation also has a significant effect on mass loss. For disturbances originating at low altitudes, a small dissipation coefficient permits waves to travel upwards until conduction in the corona provides damping of the waves. The result will be a thin chromosphere and low mass loss. Dissipation in the chromosphere heats and expands it, thereby increasing the mass loss.

The principal result of the modeling study is that any modest disturbance at low altitudes can cause mass loss through a stellar wind mechanism, due to the amplification of the disturbance as it propagates into the thinner atmosphere. The possible sources of the disturbance are many, so that this approach is capable of explaining many mass

ejection phenomena. Binary stars offer a good application to such an approach, since it is possible to observe the effects of mass ejection. For a better model, the ad hoc assumptions of dissipation coefficients and flux tubes need replacing by careful analysis of decay mechanisms and the fluid dynamics capture problem.