# COMPARISON THEOREMS FOR LINEAR ELLIPTIC EQUATIONS 

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#### Abstract

Two comparison theorems, one of pointwise type and one of integral type, will be obtained for linear elliptic equations of order $2 m$ on an exterior domain in $R^{n}$.


1. Introduction. Using Gårding's inequality, we will obtain comparison theorems for the linear elliptic differential equations

$$
\begin{equation*}
L u:=\sum_{|\alpha|,|\beta|=0}^{m}(-1)^{|\alpha|} D^{\alpha}\left[A_{\alpha \beta}(x) D^{\beta} u\right]=0 \quad\left(x \in \Omega \subseteq R^{n}\right) \tag{1}
\end{equation*}
$$

and

$$
\begin{equation*}
\ell u:=\sum_{\alpha|,|\beta|=0}^{m}(-1)^{|\alpha|} D^{\alpha}\left[a_{\alpha \beta}(x) D^{\beta} u\right]=0 . \tag{2}
\end{equation*}
$$

Here, $\Omega$ is an unbounded open set, the coefficient functions are sufficiently smooth, and we make use of the multi-index notation employed in [1]. Our results generalize two known comparison theorems:
(i) work of the author [5] in which $\ell u=(-1)^{m} \Delta^{m} u+h(x) u$;
(ii) work of Butler and Erbe [2] on the ordinary differential equations

$$
\begin{equation*}
L_{N} v+p v=0 \tag{3}
\end{equation*}
$$

and

$$
\begin{equation*}
L_{N} v+q v=0 \tag{4}
\end{equation*}
$$

where $L_{N}$ is a linear, disconjugate, ordinary differential operator of order $N$.
We remind the reader that an $N$-th order, ordinary, linear differential operator $L_{N}$ is said to be disconjugate on an interval $J$ iff the equation $L_{N} y=0$ has no nontrivial solution with $N$ zeros, counting multiplicities, on $J$.

[^0]Remark 1.1. According to the Pólya-Levin disconjugacy criterion [9,7], an $N$-th order, ordinary, linear differential operator $L_{N}$ with real-valued, locally integrable coefficients is disconjugate on a nondegenerate interval $J$ if, and only if, there exist sufficiently smooth, nonvanishing functions $\rho_{0}, \rho_{1}, \ldots, \rho_{N}$ such that, in the interior of $J$, we have

$$
\begin{equation*}
L_{N} v=\rho_{N} \frac{d}{d t}\left[\rho_{N-1} \frac{d}{d t}\left[\cdots \rho_{1} \frac{d}{d t}\left[\rho_{0} v\right] \cdots\right]\right] \tag{5}
\end{equation*}
$$

2. Definitions and preliminary results. Let $G$ be a nonempty, open (possibly unbounded) subset of $\Omega$. If $k$ is a nonnegative integer, we define the seminorm $|\cdot|_{k, G}$, the weighted seminorm $|\cdot|_{k, G, w}$ and the norm $\|\cdot\|_{k, G}$ as follows:

$$
\begin{gather*}
|u|_{k, G}=\left[\sum_{|\alpha|=k} \int_{G}\left|D^{\alpha} u\right|^{2} d x\right]^{1 / 2}  \tag{6}\\
|u|_{k, G, w}=\left[\sum_{|\alpha|=k} \int_{G}(k!/ \alpha!)\left|D^{\alpha} u\right|^{2} d x\right]^{1 / 2},  \tag{7}\\
\|u\|_{k, G}=\left[\sum_{j=0}^{k}|u|_{j, G}^{2}\right]^{1 / 2} \tag{8}
\end{gather*}
$$

The definition of $|u|_{k, G, w}$ is motivated by the following formula, which is valid for any real-valued $\phi$ in $C_{0}^{\infty}(G)$ :

$$
\begin{equation*}
(-1)^{k} \int_{G} \phi \Delta^{k} \phi d x=\sum_{|\alpha|=k} \int_{G}(k!/ \alpha!)\left|D^{\alpha} \phi\right|^{2} d x . \tag{9}
\end{equation*}
$$

To compare the seminorms $|\cdot|_{m, G}$ and $|\cdot|_{m, G, w}$, we let

$$
\begin{equation*}
c_{0}=\max \{m!/ \alpha!:|\alpha|=m\} \tag{10}
\end{equation*}
$$

Then it is easily seen that

$$
\begin{equation*}
|u|_{m, G} \leq|u|_{m, G, w} \leq c_{0}^{1 / 2}|u|_{m, G} . \tag{11}
\end{equation*}
$$

In (6) and (8), when there is no danger of confusion, we will omit the subscript $G$.
Let the Sobolev spaces $H_{k}(G)$ and $H_{k}^{0}(G)$ be defined as in [4]. If $G$ is bounded, and if there exists a nontrivial function $u$ in $H_{m}^{0}(G) \cap C^{2 m}(G)$ such that (1) holds, then $G$ is called a nodal domain for $L$ or a nodal domain for (1). If for every positive number $r$ the region $\{x \in \Omega:|x|>r\}$ contains a nodal domain for $L$, then (1) is said to be nodally oscillatory in $\Omega$.

Using integration by parts, we can easily show that if $G$ is any nonempty, open (possibly unbounded) subset of $\Omega$, then for every real-valued $\phi$ in $C_{0}^{\infty}(G)$ we have:

$$
\begin{align*}
\int_{G} \phi L \phi d x= & \sum_{|\alpha|=|\beta|=m} \int_{G} A_{\alpha \beta}(x) D^{\alpha} \phi D^{\beta} \phi d x+\int_{G} \phi^{2} A_{0,0}(x) d x \\
& +\sum_{|\alpha|+|\beta|=1}^{2 m-1} \int_{G} A_{\alpha \beta} D^{\alpha} \phi D^{\beta} \phi d x  \tag{12}\\
:= & f_{G}[L ; \phi]+\int_{G} \phi^{2} A_{0,0}(x) d x
\end{align*}
$$

and

$$
\begin{align*}
\int_{G} \phi \ell \phi d x= & \sum_{|\alpha|=|\beta|=m} \int_{G} a_{\alpha \beta}(x) D^{\alpha} \phi D^{\beta} \phi d x+\int_{G} \phi^{2} a_{0,0}(x) d x \\
& +\sum_{|\alpha|+|\beta|=1}^{2 m-1} \int_{G} a_{\alpha \beta} D^{\alpha} \phi D^{\beta} \phi d x  \tag{13}\\
:= & f_{G}[\ell ; \phi]+\int_{G} \phi^{2} a_{0,0}(x) d x .
\end{align*}
$$

Define the set $K(L ; \Omega)$ as follows:

$$
\begin{equation*}
K(L ; \Omega)=\left\{\sum_{|\alpha|=|\beta|=m} A_{\alpha \beta}(x) \xi^{\alpha+\beta}: x \in \Omega, \xi \in R^{n},|\xi|=1\right\} ; \tag{14}
\end{equation*}
$$

and define $K(\ell ; \Omega)$ by replacing $A_{\alpha, \beta}$ in (14) with $a_{\alpha \beta}$.
We will suppose that the differential operators $L$ and $\ell$ are uniformly strongly elliptic in the following sense: there exist constants $E_{0}, e_{0}$, and $e_{1}$ such that

$$
\begin{equation*}
0<E_{0}:=\inf K(L ; \Omega) \tag{15}
\end{equation*}
$$

and

$$
\begin{equation*}
0<e_{0}:=\inf K(\ell ; \Omega) \leq \sup K(\ell ; \Omega):=e_{1} \tag{16}
\end{equation*}
$$

Modifying the well-known proof of the global version of Gårding's inequality [1, Theorem 7.6], we can establish the following two results.

Lemma 2.1. Suppose that the principal coefficient-functions $A_{\alpha, \beta}(|\alpha|=|\beta|=m)$ are uniformly continuous on $\Omega$, and that the intermediate coefficient-functions $A_{\alpha \beta}(1 \leq$ $|\alpha|+|\beta| \leq 2 m-1)$ are bounded and measurable on $\Omega$. Let $G$ be any nonempty, open subset of $\Omega$ and let $f_{G}[L ; \phi]$ be as in (12). Then there exist constants $c_{1} \in(0, \infty)$ and $c_{2} \in[0, \infty)$ such that, for every real-valued $\phi$ in $C_{0}^{\infty}(G)$, we have:

$$
\begin{equation*}
f_{G}[L ; \phi] \geq c_{1} E_{0}\|\phi\|_{m, G}^{2}-c_{2}|\phi|_{0, G}^{2} \tag{17}
\end{equation*}
$$

The constant $c_{1}$ may be expressed explicitly in terms of the integers $m$ and $n$; the constant $c_{2}$ may be expressed explicitly in terms of the following quantities: $\sup \left\{\left|A_{\alpha \beta}(x)\right|: x \in \Omega\right.$; $1 \leq|\alpha|+|\beta| \leq 2 m-1\}, m, n, E_{0}$, and the modulus of continuity for the principal coefficients.

LEMMA 2.2. Suppose that the regularity hypotheses in Lemma 2.1 hold, with $A_{\alpha \beta}$ replaced by $a_{\alpha \beta}$. Let $G$ be any nonempty, open subset of $\Omega$, and let $f_{G}[\ell ; \phi]$ be as in (13). Then there exist positive constants $c_{5}$ and $c_{6}$ such that, for every real-valued $\phi$ in $C_{0}^{\infty}(G)$, we have:

$$
\begin{equation*}
f_{G}[\ell ; \phi] \leq c_{5}|\phi|_{m, G}^{2}+c_{6}|\phi|_{0, G}^{2} . \tag{18}
\end{equation*}
$$

The constants $c_{5}$ and $c_{6}$ may be expressed explicitly in terms of the following quantities:
$m, n, e_{1}, \sup \left\{\left|a_{\alpha, \beta}(x)\right|: x \in \Omega ; 1 \leq|\alpha|+|\beta| \leq 2 m-1\right\}$ and the modulus of continuity for the principal coefficients.
3. The main results. Our first comparison theorem is a generalization of the scalar case of [5, Theorem 3.1]. Note that in [5] we considered the case where $L$ is vector-valued, and we compared the differential operator $L$ with the differential operator $(-1)^{m} \Delta^{m}+$ $h$, where $h: \Omega \rightarrow R^{N \times N}$ is a continuous matrix-valued function, $\Delta$ denotes the Laplace operator, and $\Delta^{m}:=\Delta\left(\Delta^{m-1}\right)$ whenever $m \geq 2$.

Theorem 3.1. Suppose that

$$
\begin{equation*}
0<c_{5} \leq c_{4}:=c_{0}^{-1} c_{1} E_{0} \tag{19}
\end{equation*}
$$

and that for all $x \in \Omega$ we have

$$
\begin{equation*}
A_{0,0}(x)-a_{0,0}(x) \geq c_{2}+c_{6} \tag{20}
\end{equation*}
$$

If (1) is nodally oscillatory in $\Omega$, then (2) is also nodally oscillatory in $\Omega$.
Proof. If (1) is nodally oscillatory in $\Omega$, then for every positive number $r$ the region $\{x \in \Omega:|x|>r\}$ contains a nodal domain $G$ for the differential operator $L$. Thus, there exists a nontrivial real-valued function $u$ in $C^{2 m}(G) \cap H_{m}^{0}(G)$ such that (1) holds.

Furthermore, (13), (18) (i.e., Lemma 2.2) and (11) together imply that for every realvalued $\phi$ in $C_{0}^{\infty}(G)$ we have:

$$
\begin{align*}
\int_{G} \phi \ell \phi d x & =f_{G}[\ell ; \phi]+\int_{G} \phi^{2} a_{0,0}(x) d x \\
& \leq c_{5}|\phi|_{m, G}^{2}+\int_{G}|\phi|^{2}\left[a_{0,0}(x)+c_{6}\right] d x  \tag{21}\\
& \leq c_{5}|\phi|_{m, G, w}^{2}+\int_{G}|\phi|^{2}\left[a_{0,0}(x)+c_{6}\right] d x .
\end{align*}
$$

From (21), (19) and (20) we deduce that for every $\phi \in C_{0}^{\infty}\left(G ; R^{1}\right)$ we have

$$
\begin{equation*}
\int_{G} \phi \ell \phi d x \leq c_{4}|\phi|_{m, G, w}^{2}+\int_{G}|\phi|^{2}\left[A_{0,0}(x)-c_{2}\right] d x . \tag{22}
\end{equation*}
$$

We also note that (12), (17) (i.e., Lemma 2.1), (8) and (11) imply that for every $\phi \in$ $C_{0}^{\infty}\left(G ; R^{1}\right)$ we have

$$
\begin{align*}
\int_{G} \phi L \phi d x & =f_{G}[L ; G]+\int_{G} \phi^{2} A_{0,0}(x) d x \\
& \geq c_{1} E_{0}\|\phi\|_{m, G}^{2}+\int_{G}|\phi|^{2}\left[A_{0,0}(x)-c_{2}\right] d x \\
& \geq c_{1} E_{0}|\phi|_{m, G}^{2}+\int_{G}|\phi|^{2}\left[A_{0,0}(x)-c_{2}\right] d x  \tag{23}\\
& \geq c_{0}^{-1} c_{1} E_{0}|\phi|_{m, G, w}^{2}+\int_{G}|\phi|^{2}\left[A_{0,0}(x)-c_{2}\right] d x \\
& =c_{4}|\phi|_{m, G, w}^{2}+\int_{G}|\phi|^{2}\left[A_{0,0}(x)-c_{2}\right] d x .
\end{align*}
$$

It follows, from (22) and (23), that for every real-valued $\phi$ in $C_{0}^{\infty}(G)$ we have

$$
\begin{equation*}
\int_{G} \phi \ell \phi d x \leq \int_{G} \phi L \phi d x \tag{24}
\end{equation*}
$$

Since $u$ is in $H_{m}^{0}(G) \cap C^{2 m}(G)$ and satisfies (1), and since $C_{0}^{\infty}(G)$ is dense in $H_{m}^{0}(G)$, it follows from (24) that

$$
\begin{equation*}
\int_{G} u \ell u d x \leq \int_{G} u L u d x=0 . \tag{25}
\end{equation*}
$$

A standard variational argument [6,3] may now be employed to find a nonempty open set $G^{\prime} \subseteq G$ such that zero is the smallest eigenvalue of the boundary-value problem

$$
\begin{equation*}
\ell y=\lambda y, \quad y \in H_{m}^{0}\left(G^{\prime}\right) \cap C^{2 m}\left(G^{\prime}\right) \tag{26}
\end{equation*}
$$

Thus, we have shown that for every positive number $r$ the equation $\ell y=0$ has a nontrivial solution $y$, with a nodal domain $G^{\prime} \subseteq G \subseteq\{x \in \Omega:|x|>r\}$. The proof of Theorem 3.1 is now complete.

Before formulating and proving our next comparison theorem, we recall some ideas and results from [2].

Suppose that the functions $\rho_{1}, \ldots, \rho_{N}$ introduced in Remark 1.1 have the property that for each $j \in\{1, \ldots, N\}$ we have $\rho_{j} \in C^{N-j}(J ;(0, \infty))$. Define the quasiderivatives $L_{0} v, \ldots, L_{N} v$ in the usual way:

$$
\begin{equation*}
L_{0} v=\rho_{0} v, \quad L_{j} v=\rho_{j} \frac{d}{d t}\left(L_{j-1} v\right) \quad(1 \leq j \leq N) \tag{27}
\end{equation*}
$$

Furthermore, let $\Gamma:=\left\{i_{1}, \ldots, i_{k}\right\}$ and $\Lambda:=\left\{j_{1}, \ldots, j_{N-k}\right\}$ be subsets of $\{0,1, \ldots, N-1\}$ such that $0<i_{1}<i_{2}<\cdots<i_{k} \leq N-1$ and $0 \leq j_{1}<j_{2}<\cdots<j_{n-k} \leq N-1$.

For any point $a$ in the interval $J$, the first right extremal point $\theta_{1}(\Gamma, \Lambda ; a)$ for (3) is defined to be the first point $s \in J \cap(a, \infty)$ for which there exists a nontrivial solution of (3) satisfying the boundary conditions

$$
\begin{cases}L_{i} v(a)=0, & (i \in \Gamma)  \tag{28}\\ L_{j} v(s)=0, & (j \in \Lambda)\end{cases}
$$

Similarly, we can define $\tilde{\theta}_{1}(\Gamma, \Lambda ; a)$, the first right extremal point for (4).
The differential equation (3) is said to be ( $\Gamma, \Lambda$ )-disconjugate on $J$ iff for every $a \in J$ the first right extremal point $\theta_{1}(\Gamma, \Lambda ; a)$ is nonexistent.

The pair $(\Gamma, \Lambda)$ is said to be admissible iff for every integer $b \in\{1, \ldots, N-1\}$, at least $b$ members of the sequence $\left.\left(i_{1}, \ldots, i_{k}, \ldots, j_{1}, \ldots, k\right\}\right)$ are less than $b$.

The following known criterion for $(\Gamma, \Lambda)$ to be admissible will be needed in the proof of our next comparison theorem.

PROPOSITION 3.2 (SEE [2, P. 216]). The pair $(\Gamma, \Lambda)$ is admissible if, and only if, for every pair of points $a$ and $s$ in $J$ satisfying $a<s$, there exists no nontrivial solution of the differential equation $L_{N} y=0$ satisfying (28).

REmARK 3.3. We now prove a comparison theorem which extends [2, Theorem 2.5]. To facilitate the statement of the theorem, we introduce some additional notation. We define the differential operators $M_{0}$ and $M_{1}$ as follows:

$$
\begin{align*}
M_{0} u & :=(-1)^{m} c_{4} \Delta^{m} u+\left[A_{0,0}(x)-c_{2}\right] u  \tag{29}\\
M_{1} u & :=(-1)^{m} c_{5} \Delta^{m} u+\left[a_{0,0}(x)+c_{6}\right] u . \tag{30}
\end{align*}
$$

Suppose that there exists $r_{0} \geq 0$ such that the interior of $J$ is the open interval $\left(r_{0}, \infty\right)$, and suppose that there exists $x^{0} \in R^{n}$ such that $\Omega \supset\left\{x \in R^{n}:\left|x-x^{0}\right| \geq r_{0}\right\}$. For any positive $r$, let $S_{r}=\left\{x \in R^{n}:|x|=r\right\}$. Define the real-valued functions $h_{j}(2 \leq j \leq 7)$ as follows:

$$
\begin{gather*}
h_{2}(r)=\min \left\{A_{0,0}(x)-c_{2}: x \in S_{r}\right\} \text { whenever } r \in J,  \tag{31}\\
h_{3}(x)=h_{2}(|x|) \text { whenever } x \in \Omega,  \tag{32}\\
h_{4}(r)=\max \left\{a_{0,0}(x)+c_{6}: x \in S_{r}\right\} \text { whenever } r \in J,  \tag{33}\\
h_{5}(x)=h_{4}(|x|) \text { whenever } x \in \Omega,  \tag{34}\\
h_{6}(r)=\min \left\{0, h_{2}(r)\right\} \text { whenever } r \in J,  \tag{35}\\
h_{7}(x)=h_{6}(|x|) \text { whenever } x \in \Omega . \tag{36}
\end{gather*}
$$

Following [2, p. 216], we will suppose that $h_{6}(r)$ is not identically zero and that $h_{4}(r)<0$.
Define the differential operators $M_{3}$ and $M_{4}$ as follows:

$$
\begin{equation*}
M_{3} u:=(-1)^{m} c_{5} \Delta^{m} u+h_{5}(x) u \tag{37}
\end{equation*}
$$

and

$$
\begin{equation*}
M_{4} u:=(-1)^{m} c_{4} \Delta^{m} u+h_{7}(x) u . \tag{38}
\end{equation*}
$$

Let $\Delta_{|x|}$ denote the radially symmetric form of the Laplace operator. In other words, if $|x|=r$, then

$$
\begin{equation*}
\Delta_{r}=r^{n-1} \frac{d}{d r}\left(r^{1-n} \frac{d}{d r}\right) \tag{39}
\end{equation*}
$$

Furthermore, let $M_{5}$ and $M_{6}$ denote the radially symmetric forms of the differential operators $M_{3}$ and $M_{4}$, respectively. In other words, let

$$
\begin{align*}
& M_{5} v=(-1)^{m} c_{5} \Delta_{r}^{m} v+h_{4}(r) v  \tag{40}\\
& M_{6} v=(-1)^{m} c_{4} \Delta_{r}^{m} v+h_{6}(r) v, \tag{41}
\end{align*}
$$

where $h_{4}$ and $h_{6}$ are defined in (33) and (35), respectively.

Theorem 3.4. Suppose that for all $r \in J$ we have

$$
\begin{equation*}
\int_{r}^{\infty} t^{1-n}\left|h_{4}(t)\right| d t \geq \int_{r}^{\infty} t^{1-n}\left|h_{6}(t)\right| d t . \tag{42}
\end{equation*}
$$

If (1) is nodally oscillatory in $\Omega$, then so is (2).
Proof. Let $G$ be any nonempty open subset of $\Omega$, and let $\phi$ be any real-valued function in $C_{0}^{\infty}(G)$. Then (23), (9), (7) and (29) imply that

$$
\begin{align*}
\int_{G} \phi L \phi d x & \geq c_{4}|\phi|_{m, G, w}^{2}+\int_{G}|\phi|^{2}\left[A_{0,0}(x)-c_{2}\right] d x \\
& =\int_{G} \phi M_{0} \phi d x . \tag{43}
\end{align*}
$$

If (1) is nodally oscillatory in $\Omega$, then (43) and the arguments following (25) in the proof of Theorem 3.1 together imply that the equation

$$
\begin{equation*}
M_{0} u=0 \tag{44}
\end{equation*}
$$

is nodally oscillatory in $\Omega$.
Furthermore, (29), (31), (32), (35), (36) and (38) imply that

$$
\begin{align*}
\int_{G} \phi M_{0} \phi d x & =(-1)^{m} c_{4} \int_{G} \phi \Delta^{m} \phi+\int_{G}\left[A_{0,0}(x)-c_{2}\right]|\phi|^{2} d x \\
& \geq \int_{G}\left[(-1)^{m} c_{4} \phi \Delta^{m} \phi+h_{7}(x)|\phi|^{2}\right] d x  \tag{45}\\
& =\int_{G} \phi M_{4} \phi d x
\end{align*}
$$

Since (44) is nodally oscillatory in $\Omega$, therefore (45) and the arguments following (25) in the proof of Theorem 3.1 together imply that the equation

$$
\begin{equation*}
M_{4} u=0 \tag{46}
\end{equation*}
$$

is nodally oscillatory in $\Omega$. In other words, for every positive number $r_{1}$, the region $\{x \in$ $\left.\Omega:|x|>r_{1}\right\}$ contains a nodal domain for (46). Let $J_{1}:=J \cap\left(r_{1}, \infty\right)$. Then we can employ the method of spherical means (as in the proof of [3, Theorem 4.1]) to show that the ordinary differential equation

$$
\begin{equation*}
M_{6} v=0 \tag{47}
\end{equation*}
$$

is ( $\Gamma, \Lambda$ )-nondisconjugate on $J_{1}$ in the case where $\Gamma=\Lambda=\{0,1, \ldots, m-1\}$. (See (41) for the definition of $M_{6}$.) Because of the representation (39), we can choose

$$
\left\{\begin{array}{l}
N=2 m, \quad \rho_{N}(r)=r^{n-1}, \rho_{N-1}(r)=r^{1-n}, \ldots, \rho_{1}(r)=r^{1-n}, \rho_{0}(r)=1  \tag{48}\\
L_{N}=\Delta_{r}^{m}
\end{array}\right.
$$

in (5). Since (47) is ( $\Gamma, \Lambda$ )-disconjugate on $J_{1}$ in the case where $\Gamma=\Lambda=\{0,1, \ldots, m-$ $1\}$, it follows from (42) and [2, Theorem 2.5] that either the pair $(\Gamma, \Lambda)$ is inadmissible or the ordinary differential equation

$$
\begin{equation*}
M_{5} v=0 \tag{49}
\end{equation*}
$$

is ( $\Gamma, \Lambda$ )-nondisconjugate on $J_{1}$ (See (40) for the definition of $M_{5}$ ). But if the pair ( $\Gamma, \Lambda$ ) is inadmissible, then it follows from Proposition 3.2 that there will exist two points $a$ and $s$ in $J_{1}$, with $a<s$, such that the boundary-value problem consisting of the ordinary differential equation

$$
\begin{equation*}
L_{n} v_{0}=0 \text { on } J_{1} \tag{50}
\end{equation*}
$$

and the boundary conditions (28) (with $v$ replaced by $v_{0}$ ) has at least one nontrivial solution. It follows from (48) and (39) that the boundary-value problem consisting of the partial differential equation

$$
\begin{equation*}
\Delta^{m} u_{0}=0 \text { in } \Omega_{a, s}:=\{x \in \Omega: a<|x|<s\} \tag{51}
\end{equation*}
$$

and the boundary conditions

$$
\begin{equation*}
\Delta^{k} u_{0}=0 \text { on } \partial \Omega_{a, s} \quad(0 \leq k \leq m-1) \tag{52}
\end{equation*}
$$

has at least one radially-symmetric nontrivial solution. (Here, $\Delta^{0} u_{0}:=u_{0}$.) But since $\Delta^{m} u_{0}:=\Delta\left(\Delta^{m-1} u_{0}\right)$ whenever $m \geq 1$, it follows from the maximum principle that any solution of (51) and (52) has the property that

$$
\begin{equation*}
\Delta^{m-1} u_{0}=0 \text { throughout } \Omega_{a, s} \tag{53}
\end{equation*}
$$

Furthermore, (52) implies that

$$
\begin{equation*}
\Delta^{k} u_{0}=0 \text { on } \partial \Omega_{a, s} \quad(0 \leq k \leq m-2) \tag{54}
\end{equation*}
$$

Continuing recursively, we deduce eventually that $u_{0}=0$ throughout $\Omega_{a, s}$. This contradicts the nontrivialness of $u_{0}$, and shows that the pair $(\Gamma, \Lambda)$ cannot be inadmissible. Thus, we have proved that (49) is $(\Gamma, \Lambda)$-nondisconjugate on $J_{1}$ in the case where $\Gamma=\Lambda=\{0,1, \ldots, m-1\}$. In other words, there exist points $a$ and $s$ (in $J_{1}$ ) and a real-valued function $v \in C^{2 m}\left(J_{1}\right)$ such that (49) and (28) hold.

Introducing spherical polar coordinates in the usual way [8, p. 58], we note that, for every multi-index $\beta$, if $|x|=r$, then the expression $x^{\beta} / r^{|\beta|}$ is independent of $r$. It follows from the Chain Rule and the final statement in the last paragraph above that there exist points $a$ and $s$ (in $J_{1}$ ) and a radially-symmetric, real-valued, $C^{2 m}$ function $u: x \rightarrow v(|x|)$ such that

$$
\begin{equation*}
M_{3} u=M_{5} v=0 \tag{55}
\end{equation*}
$$

throughout the spherical shell $\Omega_{a, s}$, and

$$
\begin{equation*}
\left.D^{\alpha} u\right|_{\partial \Omega_{a, s}}=\left.\frac{x^{\alpha}}{r^{|\alpha|}}\left(\frac{\partial}{\partial r}\right)^{|\alpha|}\right|_{\partial \Omega_{a, s}}=0 \text { whenever } 0 \leq|\alpha| \leq m-1 \tag{56}
\end{equation*}
$$

But (56) implies, because of [1, Lemma 9.10], that $u \in H_{m}^{0}\left(\Omega_{a, s}\right)$. Since $r_{1}$ was chosen arbitrarily, we have therefore shown that (55) is nodally oscillatory in $\Omega$.

Furthermore, (37), (34), (33) and (30) imply that if $G$ is any nonempty open subset of $\Omega$, and if $\phi$ is any real-valued function in $C_{0}^{\infty}(G)$, then

$$
\begin{align*}
\int_{G} \phi M_{3} \phi d x & =\int_{G}\left[(-1)^{m} c_{5} \phi \Delta^{m} \phi+h_{5}(x)|\phi|^{2}\right] d x \\
& \geq \int_{G}\left[(-1)^{m} c_{5} \phi \Delta^{m} \phi+\left[a_{0,0}(x)+c_{6}\right]|\phi|^{2}\right] d x  \tag{57}\\
& =\int_{G} \phi M_{1} \phi d x .
\end{align*}
$$

Since (55) is nodally oscillatory in $\Omega$, therefore (57) and a familiar argument imply that the partial differential equation

$$
\begin{equation*}
M_{1} u=0 \tag{58}
\end{equation*}
$$

is nodally oscillatory in $\Omega$.
Finally, (30), (7), (9) and (21) imply that if $G$ is any nonempty open subset of $\Omega$, and if $\phi$ is any real-valued function in $C_{0}^{\infty}(G)$, then

$$
\begin{align*}
\int_{G} \phi M_{1} \phi d x & =\int_{G}\left[(-1)^{m} c_{5} \phi \Delta^{m} \phi+\left[a_{0,0}(x)+c_{6}\right]|\phi|^{2}\right] d x \\
& =c_{5}|\phi|_{m, G, w}^{2}+\int_{G}\left[a_{0,0}(x)+c_{6}\right]|\phi|^{2} d x  \tag{59}\\
& \geq \int_{G} \phi \ell \phi d x .
\end{align*}
$$

Since (58) is nodally oscillatory in $\Omega$, (59) implies that (2) is nodally oscillatory in $\Omega$.

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