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# DISCRETE TRACY–WIDOM OPERATORS

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*Abstract* Integrable operators arise in random matrix theory, where they describe the asymptotic eigenvalue distribution of large self-adjoint random matrices from the generalized unitary ensembles. We consider discrete Tracy–Widom operators and give sufficient conditions for a discrete integrable operator to be the square of a Hankel matrix. Examples include the discrete Bessel kernel and kernels arising from the almost Mathieu equation and the Fourier transform of Mathieu's equation.

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#### 1. Introduction

We consider Tracy–Widom operators arising from first-order recurrence relations

$$a(j+1) = T(j)a(j), \quad j = 1, 2, \dots,$$
 (1.1)

where a(j) is a real  $2 \times 1$  vector and T(x) is a  $2 \times 2$  real matrix with entries that are rational functions of x, and such that det T(j) = 1. Then with the skew-symmetric matrix

$$J = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix},$$

the Tracy–Widom operator  $K: \ell^2(\mathbb{N}) \to \ell^2(\mathbb{N})$  has matrix

$$K(m,n) = \frac{\langle Ja(m), a(n) \rangle}{m-n}, \quad m \neq n,$$
(1.2)

with respect to the usual orthonormal basis. In specific examples, there are natural ways of defining the diagonal K(n, n), as we discuss below. We recall that a real sequence  $(\phi_j) \in \ell^2(\mathbb{N})$  gives a Hankel matrix  $\Gamma_{\phi} = [\phi_{j+k-1}]_{j,k=1}^{\infty}$ ; clearly  $\Gamma_{\phi}$  is symmetric, and  $\Gamma_{\phi}$ is Hilbert–Schmidt if and only if  $\sum_{n=1}^{\infty} n |\phi_n|^2$  converges [14].

We consider whether such a K may be expressed as  $K = \Gamma_{\phi}^2$ , where  $\Gamma_{\phi}$  is a Hankel matrix that is self-adjoint and Hilbert–Schmidt.

G. Blower and A. McCafferty

A significant example from random matrix theory is the discrete Bessel kernel

$$B(m,n) = \sqrt{\theta} \frac{J_m(2\sqrt{\theta})J_{n-1}(2\sqrt{\theta}) - J_n(2\sqrt{\theta})J_{m-1}(2\sqrt{\theta})}{m-n},$$
(1.3)

as considered by Borodin *et al.* [5] and Johansson [11]. They showed that *B* is the square of the Hilbert–Schmidt Hankel matrix  $[J_{m+k-1}(2\sqrt{\theta})]$ , and thus obtained information about the spectrum of *B* itself.

Tracy and Widom observed that many of the fundamentally important kernels in random matrix theory have the form

$$W(x,y) = \frac{f(x)g(y) - f(y)g(x)}{x - y}, \quad x \neq y,$$
(1.4)

where f, g are bounded real functions in  $L^2(0, \infty)$  such that

$$\frac{\mathrm{d}}{\mathrm{d}x} \begin{bmatrix} f(x)\\ g(x) \end{bmatrix} = \begin{bmatrix} \alpha(x) & \beta(x)\\ -\gamma(x) & -\alpha(x) \end{bmatrix} \begin{bmatrix} f(x)\\ g(x) \end{bmatrix}, \qquad (1.5)$$

and  $\alpha(x), \beta(x)$  and  $\gamma(x)$  are real rational functions [16]. Then there is a bounded linear operator  $W: L^2(0, \infty) \to L^2(0, \infty)$  given by

$$Wh(x) = \int_0^\infty W(x,y)h(y)\,\mathrm{d} y.$$

Of particular importance is the case in which W is a trace-class operator such that  $0 \leq W \leq I$ , since such a W is associated with a determinantal point process as in [15]. In order to verify this property in special cases, Tracy and Widom showed that W is the square of a Hankel operator  $\Gamma$  which is self-adjoint and Hilbert–Schmidt. The spectral theory and realization of Hankel operators is well understood [8,13], so such a factorization is valuable. In [3,4], we developed this method further and considered which differential equations lead to kernels that can be factored as squares of Hankel operators. The factorization theorems involve the notion of operator monotone functions.

Here we consider various discrete Tracy–Widom kernels and their factorization as Hankel products, using the formal analogy between differential equations and difference equations which suggests likely factorization theorems. We can write the matrix in (1.5) as  $J\Omega(x)$ , where  $\Omega(x)$  is real and symmetric, and then consider the analogous one-step transition matrix to be  $T(x) = \exp(J\Omega(x))$ . The functions  $x \mapsto x$  and  $x \mapsto -1/x$  are operator monotone increasing on  $(0, \infty)$ , and they appear in the transition matrices for the discrete Bessel kernel in §3 and the discrete analogue of the Laguerre kernel in §4.

The functions  $x \mapsto x^2$  and  $x \mapsto -1/x^2$  are not operator monotone increasing by [9], and so we cannot hope to have simple factorization theorems when they appear in the transition matrix T. In the special case of the parabolic cylinder equation  $-\phi_n''(x) + (\frac{1}{4}x^2 - \frac{1}{2})\phi_n(x) = n\phi_n(x)$ , Aubrun [1] recovered a factorization of the corresponding kernel

$$K(x,y) = \frac{\phi_n(x)\phi_{n-1}(y) - \phi_{n-1}(x)\phi_n(y)}{x - y}$$

in the form

$$K = \Gamma_{\phi} \Gamma_{\psi} + \Gamma_{\psi} \Gamma_{\phi}, \tag{1.6}$$

547

where  $\Gamma_{\phi}$  and  $\Gamma_{\psi}$  are bounded and self-adjoint Hankel operators. The parabolic cylinder function is a confluent form of Mathieu's functions associated with the elliptic cylinder, since the parabola is the limiting case of an ellipse as the eccentricity increases to 1 [17, p. 427]. Hence, it is natural to factorize kernels associated with Mathieu's equation

$$\frac{\mathrm{d}^2 u}{\mathrm{d}z^2} + (\alpha + \beta \cos z)u(z) = 0 \tag{1.7}$$

in the form (1.6). In §5 we consider the first-order difference equation associated with the Fourier transform of Mathieu's equation. In §6, we consider almost Mathieu operators.

For a compact and self-adjoint operator W, the spectrum consists of real eigenvalues  $\lambda_j$ , which may be ordered so that the sequence of singular numbers  $s_j = |\lambda_j|$  satisfies  $s_1 \ge s_2 \ge \cdots$ . While the factorization  $W = \Gamma^2$  immediately determines the spectrum of W from the spectrum of  $\Gamma$ , a factorization (1.6) imposes bounds upon the singular numbers of K in terms of the eigenvalues of  $\Gamma_{\phi}$  and  $\Gamma_{\psi}$ .

# 2. Factoring discrete Tracy–Widom operators as squares of Hankel matrices

**Theorem 2.1.** Let T(j) and B(j) be  $2 \times 2$  real matrices, and let (a(j)) be a sequence of real  $2 \times 1$  vectors such that

$$a(j+1) = T(j)a(j), \quad j \in \mathbb{N},$$
(2.1)

$$a(j) \to 0, \quad j \to \infty,$$
 (2.2)

$$\sum_{j=1}^{\infty} \|B(j)a(j)\|^2 < \infty.$$
(2.3)

Suppose further that there exists a real symmetric matrix C with eigenvalues 0 and  $\lambda$ , where  $\lambda < 0$ , such that

$$\frac{T(n)^{\mathrm{T}}JT(m) - J}{m - n} = B(n)^{\mathrm{T}}CB(m), \quad m \neq n, \ m, n \in \mathbb{N}.$$
(2.4)

Let  $\phi(j) = |\lambda|^{1/2} \langle v_{\lambda}, B(j)a(j) \rangle$ , where  $v_{\lambda}$  is a real unit eigenvector corresponding to  $\lambda$ . Then  $\Gamma_{\phi}$  is compact and  $K = \Gamma_{\phi}^2$  has entries

$$K(m,n) = \frac{\langle Ja(m), a(n) \rangle}{m-n}, \quad m \neq n, \ m, n \in \mathbb{N}.$$
(2.5)

**Proof.** Let K(m, n) be as in (2.5) and let

$$G(m,n) = K(m,n) - \sum_{k=1}^{\infty} \phi(m+k-1)\phi(n+k-1), \quad m \neq n, \ m,n \in \mathbb{N},$$
(2.6)

where the infinite sum converges because of the condition (2.3). Observe that  $a(j) \to 0$  implies that the first term in G(m, n) tends to 0 as m or  $n \to \infty$ , and that the same is true of the Hankel sum:

$$\left|\sum_{k=1}^{\infty} \phi(m+k-1)\phi(n+k-1)\right| \leq \left(\sum_{k=1}^{\infty} |\phi(m+k-1)|^2\right)^{1/2} \left(\sum_{k=1}^{\infty} |\phi(n+k-1)|^2\right)^{1/2} \to 0 \quad (2.7)$$

as m or  $n \to \infty$ . Hence,  $G(m, n) \to 0$  as m or  $n \to \infty$ . Now let U be the real orthogonal matrix with  $v_{\lambda}$  in the first column, and the eigenvector of C corresponding to the eigenvalue 0 in the second column. Then  $U^{\mathrm{T}}CU = \mathrm{diag}(\lambda, 0)$ , and we have

$$K(m+1, n+1) - K(m, n) = \frac{1}{m-n} (\langle JT(m)a(m), T(n)a(n) \rangle - \langle Ja(m), a(n) \rangle)$$
  
$$= \frac{1}{m-n} \langle (T(n)^{\mathrm{T}} JT(m) - J)a(m), a(n) \rangle$$
  
$$= \langle B(n)^{\mathrm{T}} CB(m)a(m), a(n) \rangle$$
  
$$= \langle U \operatorname{diag}(\lambda, 0) U^{\mathrm{T}} B(m)a(m), B(n)a(n) \rangle$$
  
$$= \lambda \langle \operatorname{diag}(1, 0) U^{\mathrm{T}} B(m)a(m), \operatorname{diag}(1, 0) U^{\mathrm{T}} B(n)a(n) \rangle$$
  
$$= -\phi(m)\phi(n).$$
(2.8)

The above calculation and the equality

$$\sum_{k=1}^{\infty} \phi(m+k)\phi(n+k) - \phi(m+k-1)\phi(n+k-1) = -\phi(m)\phi(n)$$
(2.9)

together imply that G(m + 1, n + 1) = G(m, n), and so in fact G(m, n) = 0 for all  $m, n \in \mathbb{N}$ , which gives

$$K(m,n) = \sum_{k=1}^{\infty} \phi(m+k-1)\phi(n+k-1), \quad m \neq n, \ m, n \in \mathbb{N},$$
(2.10)

where the right-hand side is the (m, n)th entry of the square of a Hankel matrix  $\Gamma_{\phi}$ . Furthermore, we observe that K is the composition of the discrete Hilbert transform with the compact operators  $\ell^2(\mathbb{N}; \mathbb{C}) \to \ell^2(\mathbb{N}; \mathbb{C}^2)$  given by  $(x_n) \mapsto (Ja(n)x_n)$  and the adjoint of  $(x_n) \mapsto (a(n)x_n)$ ; so K is compact. We deduce that  $\Gamma_{\phi}$  is also compact.  $\Box$ 

**Proposition 2.2.** Let a(j), T(j) and B(j) satisfy conditions (2.1) and (2.2), and suppose further that

$$\sum_{j=1}^{\infty} j \|B(j)a(j)\|^2 < \infty.$$
(2.11)

Now let  $K = \Gamma_{\phi}^2$  be as in Theorem 2.1. Then

- (i) K is a positive semidefinite and trace-class operator,
- (ii) for each  $n \in \mathbb{N}$ , there exist self-adjoint Hankel operators  $\Gamma_n$ , where  $\Gamma_n$  has rank at most n, such that

$$s_n(K) = \|\Gamma_\phi - \Gamma_n\|^2,$$
 (2.12)

so  $\Gamma_n^2 \to K$  as  $n \to \infty$ .

**Proof.** (i) The Hilbert–Schmidt norm of  $\Gamma_{\phi}$  satisfies

$$\|\Gamma_{\phi}\|_{HS}^2 = \sum_{k=1}^{\infty} k\phi(k)^2 < \infty.$$

Hence,  $K = \Gamma_{\phi}^2$  is of trace class.

(ii) Since  $\Gamma_{\phi}$  is self-adjoint, the singular numbers satisfy  $s_n(K) = s_n(\Gamma_{\phi}^2) = s_n(\Gamma_{\phi})^2$ . By the Adamyan–Arov–Krein theorem [14], there exists a unique Hankel operator  $\Gamma_n$ with rank at most n such that  $s_n(\Gamma_{\phi}) = \|\Gamma_{\phi} - \Gamma_n\|$ . Evidently,  $\Gamma_n^*$  is also a Hankel operator of rank at most n such that  $s_n(\Gamma_{\phi}) = \|\Gamma_{\phi} - \Gamma_n^*\|$ , so, by uniqueness,  $\Gamma_n = \Gamma_n^*$ . We have  $\|\Gamma_n - \Gamma_{\phi}\| \to 0$  as  $n \to \infty$ , so  $\Gamma_n^2 \to \Gamma_{\phi}^2$  as  $n \to \infty$ .

**Definition 2.3.** For a compact and self-adjoint operator W on a Hilbert space H, the

spectral multiplicity function  $\nu_W : \mathbb{R} \to \{0, 1, ...\} \cup \{\infty\}$  is given by

$$\nu_W(\lambda) = \dim\{x \in H : Wx = \lambda x\}, \quad \lambda \in \mathbb{R}.$$
(2.13)

**Proposition 2.4.** Let K be as in Proposition 2.2. Then the following hold:

- (i)  $\nu_K(0) = 0 \text{ or } \nu_K(0) = \infty;$
- (ii)  $\nu_K(\lambda) < \infty$  and  $\nu_K(\lambda) = \nu_{\Gamma_\phi}(\sqrt{\lambda}) + \nu_{\Gamma_\phi}(-\sqrt{\lambda})$  for all  $\lambda > 0$ ;
- (iii) if  $\nu_K(\lambda)$  is even, then  $\nu_{\Gamma_{\phi}}(\sqrt{\lambda}) = \nu_{\Gamma_{\phi}}(-\sqrt{\lambda});$
- (iv) if  $\nu_K(\lambda)$  is odd, then  $|\nu_{\Gamma_{\phi}}(\sqrt{\lambda}) \nu_{\Gamma_{\phi}}(-\sqrt{\lambda})| = 1$ .

**Proof.** Part (i) follows from Beurling's theorem (see [14, p. 15]), while (ii) is elementary. Peller *et al.* show in [13] that, for any compact and self-adjoint Hankel operator  $\Gamma_{\phi}$ , the spectral multiplicity function satisfies  $|\nu_{\Gamma_{\phi}}(\lambda) - \nu_{\Gamma_{\phi}}(-\lambda)| \leq 1$ . Using this, and (ii), statements (iii) and (iv) follow immediately.

#### 3. The discrete Bessel kernel

We show how Theorem 2.1 can be applied to the discrete Bessel kernel to recover a result from [5,11].

G. Blower and A. McCafferty

**Proposition 3.1.** Let  $J_n(z)$  be the Bessel function of the first kind of order n, let  $J_n = J_n(2\sqrt{\theta})$ , where  $\theta > 0$ ; let  $\phi(n) = J_{n+1}$  and  $a(n) = [\sqrt{\theta}J_n, J_{n+1}]^{\mathrm{T}}$ . Then the Hankel operator  $\Gamma_{\phi}$  is Hilbert–Schmidt, and  $B = \Gamma_{\phi}^2$  has entries

$$B(m,n) = \frac{\langle Ja(m), a(n) \rangle}{m-n}, \quad m \neq n, \ m, n \in \mathbb{N}.$$
(3.1)

**Proof.** It is clear that (2.1) holds, since we have the recurrence relation

$$J_{n+2}(2z) = \frac{n+1}{z} J_{n+1}(2z) - J_n(2z), \qquad (3.2)$$

giving a(n+1) = T(n)a(n), where

$$T(n) = \begin{bmatrix} 0 & \sqrt{\theta} \\ \frac{-1}{\sqrt{\theta}} & \frac{n+1}{\sqrt{\theta}} \end{bmatrix}.$$
 (3.3)

Note that

$$\frac{T(n)^{\rm T} J T(m) - J}{m - n} = C,$$
(3.4)

where C = diag(0, -1), which is clearly of rank 1. The non-zero eigenvalue of C is  $\lambda = -1$ , and a corresponding unit eigenvector is  $v_{\lambda} = [0, 1]^{\text{T}}$ , so

$$|\lambda|^{1/2} \langle v_{\lambda}, a(n) \rangle = J_{n+1} = \phi(n).$$

We now verify condition (2.11), and thus (2.2). Note that

$$\frac{1}{\theta} \sum_{n=1}^{\infty} n J_{n+1}^2 < \frac{1}{\theta} \sum_{n=1}^{\infty} (n+1)^2 J_{n+1}^2 = \sum_{n=1}^{\infty} (J_{n+2} + J_n)^2 \leqslant 4 \sum_{n=1}^{\infty} J_n^2.$$
(3.5)

The standard formula [17, p. 379]

$$e^{i2\sqrt{\theta}\sin\psi} = J_0(2\sqrt{\theta}) + 2\sum_{m=1}^{\infty} J_{2m}(2\sqrt{\theta})\cos 2m\psi + 2i\sum_{m=1}^{\infty} J_{2m-1}(2\sqrt{\theta})\sin(2m-1)\psi$$
(3.6)

and Parseval's identity can be used to show that

$$J_0(2\sqrt{\theta})^2 + 2\sum_{m=1}^{\infty} J_m(2\sqrt{\theta})^2 = 1 \quad \text{for all } \theta > 0,$$

and hence that the sum on the right-hand side of (3.5) is finite.

# 4. A discrete analogue of the Laguerre differential equation

In this section we consider a case in which condition (2.2) is violated and we cannot hope to factor the kernel K as the square of a Hankel operator. Nevertheless, we can identify a Toeplitz operator W such that K - W factors as a product of Hankels.

**Proposition 4.1.** For  $\theta \in \mathbb{R}$ , let (a(j)) satisfy the recurrence relation a(j + 1) = T(j)a(j) with

$$T(j) = \begin{bmatrix} \theta/(j+1) & -1\\ 1 & 0 \end{bmatrix}$$
(4.1)

and  $a(1) = [\theta, 1]^{\mathrm{T}}$ . Then there exist polynomials  $p_j(\theta)$  of degree j such that

(i)  $a(j) = [p_j(\theta), p_{j-1}(\theta)]^{\mathrm{T}}.$ 

(ii) The self-adjoint Hankel matrix  $\Gamma_{\phi} = [\phi(j+k-1)]_{j,k=1}^{\infty}$  with entries

$$\phi(j) = \frac{p_j(\theta)}{j+1} \tag{4.2}$$

is a bounded linear operator such that  $\theta \Gamma_{\phi}^2 = K + W$  where K has entries

$$K(m,n) = \frac{\langle Ja(m), a(n) \rangle}{m-n}, \quad m \neq n, \ m, n \in \mathbb{N},$$
(4.3)

and W is a bounded Toeplitz operator with matrix

$$W(m-n) = \frac{\langle J^{\bar{m}-\bar{n}+1}T_{\infty}a(1), T_{\infty}a(1)\rangle}{m-n}, \quad m \neq n,$$

for some  $2 \times 2$  matrix  $T_{\infty}$ , and  $\overline{m}$  and  $\overline{n}$  are the congruence classes of m and n modulo 4.

**Proof.** When  $\theta = 0$ , the recurrence relation reduces to a(n+1) = Ja(n), with solution  $a(n) = J^{n-1}a(1)$ . This gives rise to a kernel

$$K(m,n) = \frac{\langle Ja(m), a(n) \rangle}{m-n} = \frac{\langle J^{m-n+1}a(1), a(1) \rangle}{m-n}, \quad m \neq n,$$
(4.4)

which has the shape of a Toeplitz operator, and is a variant on the discrete Hilbert transform with matrix  $[1/(m-n)]_{m\neq n}$ .

Now suppose  $\theta \neq 0$ . The matrix J satisfies  $J^4 = I$ , and so we consider the partial product of the T(j) in bunches of four, with the *j*th bunch giving

$$B(j) = T(4j)T(4j-1)T(4j-2)T(4j-3) = I - \frac{\theta}{2j}J + O(1/j^2), \quad j \in \mathbb{N}.$$
 (4.5)

Now we deduce that

$$||B(j)||^{2} = ||B(j)^{*}B(j)|| = 1 + O(1/j^{2}),$$
(4.6)

and likewise with  $B(j)^{-1}$  in place of B(j); so there exists  $C(\theta)$  such that

$$\|T(n)T(n-1)\cdots T(2)T(1)\| \leq C(\theta), \\ \|T(1)^{-1}T(2)^{-1}\cdots T(n-1)^{-1}T(n)^{-1}\| \leq C(\theta), \\ \} \qquad n = 1, 2, \dots$$

$$(4.7)$$

It follows that there exists  $\kappa(\theta) > 0$  such that  $\kappa(\theta) < ||a(n)|| < \kappa(\theta)^{-1}$  for all n, so (2.2) is violated. We introduce

$$C_k = \exp\left(\theta J \sum_{j=1}^k \frac{1}{2j}\right) B(k) B(k-1) \cdots B(1), \qquad (4.8)$$

which satisfies  $C_{k+1} - C_k = O(1/k^2)$ ; so the limit

$$T_{\infty} = \lim_{k \to \infty} C_k \tag{4.9}$$

exists. One can check that

$$\exp\left(\sum_{j=1}^{k+n} \frac{\theta J}{2j}\right)^* \exp\left(\sum_{j=1}^{k+m} \frac{\theta J}{2j}\right) \to I$$
(4.10)

as  $k \to \infty$ , and hence

$$\langle Ja(m+4k), a(n+4k) \rangle$$
  
=  $\langle JT(m+4k)T(m+4k-1)\cdots T(1)a(1), T(n+4k)T(n+4k-1)\cdots T(1)a(1) \rangle$   
 $\rightarrow \langle J^{\bar{m}+1}T_{\infty}a(1), J^{\bar{n}}T_{\infty}a(1) \rangle, \quad k \rightarrow \infty.$  (4.11)

For temporary convenience we introduce

$$\tilde{K}(m,n) = \begin{cases} \frac{\langle Ja(m), a(n) \rangle}{(m-n)} & \text{for } m \neq n, \\ 0 & \text{for } m = n. \end{cases}$$
(4.12)

The discrete Hilbert transform is bounded on  $\ell^2(\mathbb{Z}; \mathbb{C}^2)$  by [7], so  $\tilde{K}$  defines a bounded linear operator on  $\ell^2(\mathbb{N})$ , but condition (2.2) is violated.

The  $p_i(\theta)$  satisfy the recurrence relation

$$p_{n+1}(\theta) + p_{n-1}(\theta) = \frac{\theta}{n+1} p_n(\theta)$$
(4.13)

with  $p_0(\theta) = 1$  and  $p_1(\theta) = \theta$ , so clearly  $p_j(\theta)$  is a polynomial of degree j such that  $|p_n(\theta)| \leq \kappa(\theta)^{-1}$  for all  $\theta$  and n. Furthermore,

$$\frac{T(n)^{\mathrm{T}}JT(m) - J}{m - n} = \begin{bmatrix} \frac{-\theta}{(m+1)(n+1)} & 0\\ 0 & 0 \end{bmatrix},$$
(4.14)

so that, by the calculation in the proof of Theorem 2.1,

$$\tilde{K}(m+1, n+1) - \tilde{K}(m, n) = -\theta \frac{p_m(\theta)p_n(\theta)}{(m+1)(n+1)}.$$
(4.15)

We can write

$$\tilde{K}(m,n) - \tilde{K}(m+N+1,n+N+1) = \theta \sum_{k=0}^{N} \frac{p_{m+k}(\theta)p_{n+k}(\theta)}{(m+k+1)(n+k+1)}, \quad m \neq n, \quad (4.16)$$

where the limit

$$\lim_{N \to \infty} \tilde{K}(m+N+1, n+N+1) = W(m-n), \quad m \neq n,$$
(4.17)

exists and is finite since the sequence  $\phi(j) = p_j(\theta)/(j+1)$  is square summable, so

$$\tilde{K}(m,n) = \theta \sum_{k=0}^{\infty} \frac{p_{m+k}(\theta)p_{n+k}(\theta)}{(m+k+1)(n+k+1)} + W(m-n), \quad m \neq n.$$
(4.18)

We now define W(0) = 0, and let

$$K(m,n) = \theta \sum_{k=0}^{\infty} \frac{p_{m+k}(\theta)p_{n+k}(\theta)}{(m+k+1)(n+k+1)} + W(m-n),$$
(4.19)

so that the matrix of K equals the matrix of  $\tilde{K}$ , except on the principal diagonal, and the principal diagonal of K is a bounded sequence; hence, K is a bounded linear operator and also satisfies the preceding identities for m = n. Let S be the shift operator on  $\ell^2(\mathbb{N})$ . Now  $\theta \Gamma_{\phi}^2$  equals the limit in the weak operator topology of the sequence  $K - S^{*n}KS^n$ as  $n \to \infty$ , so  $\Gamma_{\phi}$  is bounded and hence  $W = K - \theta \Gamma_{\phi}^2$  is also bounded. We recognize the matrix of W as

$$W(m-n) = \frac{\langle J^{\bar{m}-\bar{n}+1}T_{\infty}a(1), T_{\infty}a(1)\rangle}{m-n}, \quad m \neq n.$$

**Remark 4.2.** The generating function  $f(z) = \sum_{j=0}^{\infty} p_j(\theta) z^j$  satisfies the differential equation

$$(1+z^2)f'(z) + (2z-\theta)f(z) = 0$$

with initial condition f(0) = 1, and hence

$$f(z) = \left(\frac{1-iz}{1+iz}\right)^{i\theta/2} \frac{1}{1+z^2}.$$
(4.20)

For comparison, Laguerre's equation [16] may be expressed as

$$\frac{\mathrm{d}}{\mathrm{d}x} \begin{bmatrix} u(x)\\ u'(x) \end{bmatrix} = \begin{bmatrix} 0 & 1\\ \frac{1}{4} - \frac{(n+1)}{x} & 0 \end{bmatrix} \begin{bmatrix} u(x)\\ u'(x) \end{bmatrix}, \qquad (4.21)$$

with solution  $u(x) = xe^{-x/2}L_n^{(1)}(x)$ , where

$$L_n^{(1)}(x) = \frac{x^{-1} \mathrm{e}^x}{n!} \frac{\mathrm{d}^n}{\mathrm{d}x^n} (x^{n+1} \mathrm{e}^{-x}), \quad x > 0,$$
(4.22)

is the Laguerre polynomial of degree n and parameter 1. The Laplace transform of u is the rational function

$$\mathcal{L}(u;\lambda) = (n+1)\frac{(\lambda - \frac{1}{2})^n}{(\lambda + \frac{1}{2})^{n+2}}, \quad \text{Re}\,\lambda > -\frac{1}{2}.$$
(4.23)

## 5. The Fourier transform of Mathieu's equation

Let  $\beta \neq 0$  be a real number; then there exists a sequence of real values of  $\alpha$  such that Mathieu's equation

$$\frac{\mathrm{d}^2 u}{\mathrm{d}\theta^2} + (\alpha + \beta \cos \theta) u(\theta) = 0 \tag{5.1}$$

has a real periodic solution with period  $2\pi$  or  $4\pi$ . The odd or even periodic solutions are known as Mathieu functions, and various determinants describe the dependence of the eigenvalues  $\alpha$  on  $\beta$ , as in [12, 17]. Here we are concerned with some matrices that arise from the Fourier transform of the differential equation.

**Theorem 5.1.** Suppose that *u* has Fourier expansion

$$u(\theta) = \sum_{n=-\infty}^{\infty} b_n \mathrm{e}^{\mathrm{i}n\theta}.$$

Let  $\Gamma_u$  be the Hankel matrix  $[b_{j+k-1}]_{j,k=1}^{\infty}$ , let  $\Gamma_v$  be the Hankel matrix  $[(j+k-1)b_{j+k-1}]_{j,k=1}^{\infty}$  and let  $K = (-2/\beta)(\Gamma_u\Gamma_v + \Gamma_v\Gamma_u)$ . Then K is a trace class operator on  $\ell^2(\mathbb{N})$  such that

$$\operatorname{tr}(K) = \frac{1}{\beta \pi} \int_0^{2\pi} \left| \frac{\mathrm{d}u}{\mathrm{d}\theta} \right|^2 \mathrm{d}\theta$$

and

$$K(j,k) = \frac{b_{j-1}b_k - b_j b_{k-1}}{j-k}, \quad j,k \in \mathbb{N}, \ j \neq k.$$
(5.2)

**Proof.** Since u is real, we have  $b_m = \overline{b}_{-m}$ . The recurrence relation for the Fourier coefficients

$$2(-n^2 + \alpha)b_n + \beta b_{n+1} + \beta b_{n-1} = 0$$
(5.3)

may be expressed as the first-order recurrence relation

$$\begin{bmatrix} b_n \\ b_{n+1} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & \left(\frac{2}{\beta}\right)(n^2 - \alpha) \end{bmatrix} \begin{bmatrix} b_{n-1} \\ b_n \end{bmatrix},$$
(5.4)

or in the obvious shorthand a(n + 1) = T(n)a(n). Then we have

$$\frac{T(n)^{\mathrm{T}}JT(m) - J}{m - n} = \begin{bmatrix} 0 & 0\\ 0 & \left(-\frac{2}{\beta}\right)(m + n) \end{bmatrix}.$$
(5.5)

We introduce the kernel  $\tilde{K}$  by the formula

$$\tilde{K}(m,n) = \frac{\langle Ja(m), a(n) \rangle}{m-n} = \frac{b_n b_{m-1} - b_{n-1} b_m}{m-n},$$

which therefore satisfies

$$\tilde{K}(m+1,n+1) - \tilde{K}(m,n) = \left(-\frac{2}{\beta}\right)(mb_m b_n + nb_n b_m),\tag{5.6}$$

and  $K(m,n) \to 0$  as  $m, n \to \infty$  in any way such that  $m \neq n$ . We deduce that

$$\tilde{K}(m,n) = \frac{2}{\beta} \sum_{k=0}^{\infty} (m+k)b_{m+k}b_{n+k} + (n+k)b_{n+k}b_{m+k}, \quad m,n \in \mathbb{N}, \ m \neq n$$

Hence,  $\tilde{K}(m,n)$  is the (m,n) entry of the matrix of  $K = (2/\beta)(\Gamma_u \Gamma_v + \Gamma_v \Gamma_u)$  for all  $m \neq n$ .

Since u and u'' are square integrable, the series  $\sum_{n=-\infty}^{\infty} n^4 |b_n|^2$  converges; so  $\Gamma_u$  and  $\Gamma_v$  are Hilbert–Schmidt, and K is trace class. Furthermore, we have

$$tr(K) = \sum_{m=1}^{\infty} K(m,m)$$
  
=  $\frac{4}{\beta} \sum_{m,k=1}^{\infty} (m+k-1)b_{m+k-1}^2$   
=  $\frac{4}{\beta} \sum_{m=1}^{\infty} m^2 b_m^2$   
=  $\frac{2}{\beta} \int_0^{2\pi} |u'(\theta)|^2 \frac{d\theta}{2\pi}.$  (5.7)

#### 6. Almost Mathieu operators

We introduce the almost Mathieu operator  $H: \ell^2(\mathbb{Z}) \to \ell^2(\mathbb{Z})$  by

$$(Hu)_n = u_{n+1} + u_{n-1} + \lambda \cos 2\pi (n\theta + \alpha)u_n \tag{6.1}$$

for  $u = (u_n)_{n=-\infty}^{\infty} \in \ell^2(\mathbb{Z})$ , where  $(Hu)_n$  denotes the *n*th term in the sequence  $Hu \in \ell^2(\mathbb{Z})$ . For all real  $\lambda$ ,  $\theta$ ,  $\omega$  and  $\alpha$ , the operator H is bounded and self-adjoint, with its spectrum contained in  $[-2-|\lambda|, 2+|\lambda|]$ . According to the precise values of the parameters, as we discuss below, the spectrum can consist of a mixture of the point spectrum, the continuous spectrum and the singular continuous spectrum.

**Definition 6.1.** Let *E* be an eigenvalue of *H* with the corresponding eigenvector  $(u_n)$ . Say that  $(u_n)$  decays exponentially if there exist  $C, \delta > 0$  and  $n_0 \in \mathbb{Z}$  such that

$$|u_n| \leqslant C e^{-\delta |n-n_0|}, \quad n \in \mathbb{Z}.$$
(6.2)

G. Blower and A. McCafferty

(Typically,  $\delta$  depends on  $\theta$ ,  $\alpha$  and  $\lambda$  in a complicated fashion.) Say that *H* exhibits *Anderson localization* if its spectrum is pure point and all eigenvectors decay exponentially.

**Definition 6.2.** Say that  $\theta \in \mathbb{R}$  is *Diophantine* if there exist  $c(\theta) > 0$  and  $r(\theta) > 0$  such that

$$|\sin 2\pi j\theta| \ge c(\theta)|j|^{-r(\theta)}, \quad j \in \mathbb{Z} \setminus \{0\}.$$
(6.3)

With respect to Lebesgue measure, almost all real numbers are Diophantine [2, p. 373]. Clearly, rational numbers are not Diophantine, and Liouville numbers are not Diophantine [10].

Following the work of several mathematicians, as summarized in [6], Jitomirskaya [10] obtained a satisfactory description of the spectrum of the almost Mathieu operator.

Lemma 6.3 (Jitomirskaya [10]). Suppose that  $\theta$  is Diophantine. Then there exists a set  $S_{\theta}$  such that  $\mathbb{R} \setminus S_{\theta}$  has Lebesgue measure zero, and such that, for all  $\lambda > 2$  and all  $\alpha \in S_{\theta}$ , the Mathieu operator H has an eigenvalue E such that the corresponding eigenvector  $(u_n)$  decays exponentially as  $n \to \pm \infty$ .

Moreover, Jitomirskaya proved that the Mathieu operator has pure point spectrum for  $\lambda > 2$ , and further conjectured that the same conclusion holds for all real  $\alpha$ . For comparison, for  $\lambda = 2$  there exists a purely singular continuous spectrum, whereas for  $0 < \lambda < 2$  there exists a purely absolutely continuous spectrum.

We introduce the Hankel matrices

$$\Gamma_{\rm c} = [\cos \pi (\alpha + x\theta + k\theta)u_{x+k}]_{x,k=1}^{\infty}, \qquad \Gamma_{\rm s} = [\sin \pi (\alpha + x\theta + k\theta)u_{x+k}]_{x,k=1}^{\infty}. \tag{6.4}$$

There exists a bounded and measurable function  $\Phi : \mathbb{T} \to M_2(\mathbb{C})$  such that

$$\hat{\Phi}(k) = \begin{bmatrix} \cos \pi (\alpha + k\theta)u_k & \sin \pi (\alpha + k\theta)u_k \\ \cos \pi (\alpha + k\theta)u_k & \cos \pi (\alpha + k\theta)u_k \end{bmatrix};$$
(6.5)

the block Hankel operator associated with  $\Phi$  is  $[\hat{\Phi}(j+k)]$ , which becomes, after a rearrangement of the block form, the matrix

$$\Gamma_{\varPhi} = \begin{bmatrix} \Gamma_{\rm c} & \Gamma_{\rm s} \\ \Gamma_{\rm s} & \Gamma_{\rm c} \end{bmatrix}.$$
(6.6)

The negative Fourier coefficients of  $\Phi$  are not uniquely determined by  $\Gamma_{\Phi}$ , but may be chosen advantageously. Note that all of these operators are self-adjoint.

We introduce operators K and L by

$$\begin{bmatrix} L & K \\ K & L \end{bmatrix} = \begin{bmatrix} \Gamma_{\rm c}^2 + \Gamma_{\rm s}^2 & \Gamma_{\rm c}\Gamma_{\rm s} + \Gamma_{\rm s}\Gamma_{\rm c} \\ \Gamma_{\rm c}\Gamma_{\rm s} + \Gamma_{\rm s}\Gamma_{\rm c} & \Gamma_{\rm c}^2 + \Gamma_{\rm s}^2 \end{bmatrix} = \Gamma_{\varPhi}^2.$$
(6.7)

**Theorem 6.4.** Let  $u_n$  be as in Lemma 6.3.

(i) The matrices of K and L are given by

$$K(m,n) = \frac{u_{m-1}u_n - u_{n-1}u_m}{2\lambda\sin\pi\theta(m-n)}, \qquad m \neq n, \ m,n \in \mathbb{N};$$
  
$$L(m,n) = \cos\pi\theta(m-n)\sum_{k=1}^{\infty} u_{m+k}u_{n+k}, \quad m,n \in \mathbb{N}.$$

$$(6.8)$$

- (ii) Then  $\Gamma_{\Phi}^2$  is a positive semidefinite trace-class operator.
- (iii) Further, the eigenvectors of K,  $\Gamma_{\rm c}$  and  $\Gamma_{\rm s}$  decay exponentially as  $x \to \infty$ .
- (iv) There exists a bounded and measurable function  $\Psi_n : \mathbb{T} \to M_2(\mathbb{C})$  such that the associated Hankel operator  $[\hat{\Psi}_n(j+k)]$  has rank less than or equal to n and

$$s_n(\Gamma_{\Phi}) = \|\Gamma_{\Phi} - \Gamma_{\Psi_n}\| = \|\Phi - \Psi_n\|_{\infty}.$$
(6.9)

(v) The eigenvalues of K, L,  $\Gamma_{\rm s}$  and  $\Gamma_{\rm c}$  decay exponentially.

**Proof.** (i) First we observe that the formula for K(m, n) makes sense for  $m \neq n$  since  $\theta$  is irrational. The discrete Mathieu equation

$$u_{n+1} + u_{n-1} + \lambda \cos 2\pi (n\theta + \alpha)u_n = Eu_n \tag{6.10}$$

gives the system

$$\begin{bmatrix} u_n \\ u_{n+1} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & E - \lambda \cos 2\pi (n\theta + \alpha) \end{bmatrix} \begin{bmatrix} u_{n-1} \\ u_n \end{bmatrix}.$$
 (6.11)

Writing T(n) for the one-step transition matrix, we have

$$\frac{T(n)^{\mathrm{T}}JT(m) - J}{2\lambda\sin\pi\theta(m-n)} = -\begin{bmatrix} 0 & 0\\ 0 & 1 \end{bmatrix} (\sin\pi(n\theta + \alpha)\cos\pi(m\theta + \alpha) + \cos\pi(n\theta + \alpha)\sin\pi(m\theta + \alpha));$$
(6.12)

so with  $a(m) = [u_{m-1}, u_m]^{\mathrm{T}}$  we introduce

$$\tilde{K}(m,n) = \frac{\langle Ja(m), a(n) \rangle}{2\lambda \sin \pi \theta (m-n)},$$
(6.13)

which satisfies

$$\tilde{K}(m+1, n+1) - \tilde{K}(m, n)$$
  
=  $(\sin \pi (n\theta + \alpha) \cos \pi (m\theta + \alpha) + \cos \pi (n\theta + \alpha) \sin \pi (m\theta + \alpha)) u_n u_m$ 

and  $\tilde{K}(m+k, n+k) \to 0$  as  $k \to \infty$ . Hence, by comparing entries of the products, we find that  $\tilde{K} = \Gamma_{\rm c}\Gamma_{\rm s} + \Gamma_{\rm s}\Gamma_{\rm c}$ .

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- (ii) This is clear, since the entries of the matrix of  $\Gamma_{\Phi}$  are real and summable.
- (iii) For any unit vector  $\varphi$ , we have

$$\Gamma_{\rm c}\varphi(x) = \sum_{k=0}^{\infty} \cos \pi (\alpha + \theta x + \theta k) u_{x+k}\varphi_k, \qquad (6.14)$$

so by the Cauchy-Schwarz inequality we have the uniform bound

$$|\Gamma_{\rm c}\varphi(x)| \leqslant \left(\sum_{k=x}^{\infty} u_k^2\right)^{1/2},\tag{6.15}$$

where the right-hand side decays exponentially as  $x \to \infty$ . A similar result applies with  $\Gamma_{\rm s}$ , so in particular the eigenvectors of  $\Gamma_{\rm c}$  and  $\Gamma_{\rm s}$  decay exponentially at infinity.

Now let  $(\varphi_j)_{j=1}^{\infty}$  be an orthonormal basis of  $\ell^2(\mathbb{Z})$  consisting of eigenvectors of  $\Gamma_s$  with corresponding eigenvalues  $\sigma_j$ . Then

$$\Gamma_{\rm c}\Gamma_{\rm s}\varphi(x) = \sum_{j=1}^{\infty} \sigma_j \langle \varphi_j, \varphi \rangle \Gamma_{\rm c}\varphi_j(x), \qquad (6.16)$$

where

558

$$\sum_{j=1}^{\infty} |\sigma_j \langle \varphi_j, \varphi \rangle| \leqslant \left(\sum_{j=1}^{\infty} \sigma_j^2\right)^{1/2} \left(\sum_{j=1}^{\infty} \langle \varphi, \varphi_j \rangle^2\right)^{1/2} = \|\Gamma_s\|_{c^2} \|\varphi\|_{\ell^2}.$$
(6.17)

This and a similar result for  $\Gamma_{\rm s}\Gamma_{\rm c}$  imply that  $|K\varphi(x)|$  decays exponentially as  $x \to \infty$ .

(iv) This is immediate from the vectorial form of the matrical AAK theorem [14].

(v) Since  $u_n$  decays exponentially as  $n \to \infty$ , we can approximate the Hankel matrices with finite matrices up to exponentially small error terms. For instance, we can approximate  $\Gamma_s$  by  $\Gamma_s^{(N)} = [\sin \pi (\alpha + x\theta + k\theta) u_{x+k} I_{\{(x,k):x+k \leq N\}}]$ , which has rank less than N + 1 and the operator norm satisfies

$$\|\Gamma_{\rm s} - \Gamma_{\rm s}^{(N)}\| \leqslant \sum_{k=N+1}^{\infty} k|u_k|.$$

$$(6.18)$$

The *s*-numbers satisfy

$$s_n(K) \leqslant s_n\left(\begin{bmatrix} K & L\\ L & K \end{bmatrix}\right) = s_n(\Gamma_{\varPhi}^2) = s_n(\Gamma_{\varPhi})^2.$$
(6.19)

A vectorial Hankel matrix  $\Gamma_{\Psi}$  has finite rank if and only if it has a rational symbol with coefficients of finite rank by [14, p. 19]. Peller provides a formula for the rank in terms of the coefficients.

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### References

- 1. G. AUBRUN, A sharp small deviation inequality for the largest eigenvalue of a random matrix, Springer Lecture Notes in Mathematics, Volume 1857 (Springer, 2005).
- 2. J. AVRON AND B. SIMON, Almost periodic Schrödinger operators, II, The integrated density of states, *Duke Math. J.* **50** (1983), 369–391.
- 3. G. BLOWER, Operators associated with the soft and hard edges from unitary ensembles, J. Math. Analysis Applic. **337** (2008), 239–265.
- 4. G. BLOWER, Integrable operators and the squares of Hankel operators, J. Math. Analysis Applic. **340** (2008), 943–953.
- 5. A. BORODIN, A. OKOUNKOV AND G. OLSHANSKI, Asymptotics of Plancherel measures for symmetric groups, J. Am. Math. Soc. **13** (2000), 481–515.
- 6. J. BOURGAIN, On the spectrum of lattice Schrödinger operators with deterministic potential, J. Analysis Math. 87 (2002), 37–75.
- 7. G. H. HARDY, J. E. LITTLEWOOD AND G. PÔLYA, *Inequalities*, 2nd edn (Cambridge University Press, 1988).
- 8. J. W. HELTON, Discrete time systems, operator models and scattering theory, J. Funct. Analysis 16 (1974), 15–38.
- 9. R. A. HORN AND C. R. JOHNSON, *Topics in matrix analysis* (Cambridge University Press, 1991).
- S. JITOMIRSKAYA, Metal-insulator transition for the almost Mathieu operator, Annals Math. 150 (1999), 1159–1175.
- K. JOHANSSON, Discrete orthogonal polynomial ensembles and the Plancherel measure, Annals Math. 153 (2001), 259–296.
- 12. W. MAGNUS AND S. WINKLER, Hill's equation (Dover, New York, 1966).
- A. N. MEGRETSKIĬ, V. V. PELLER AND S. R. TREIL, The inverse spectral problem for self-adjoint Hankel operators, Acta Math. 174 (1995), 241–309.
- 14. V. PELLER, Hankel operators and their applications (Springer, 2003).
- A. G. SOSHNIKOV, Determinantal random point fields, Russ. Math. Surv. 55 (2000), 923–975.
- C. A. TRACY AND H. WIDOM, Fredholm determinants, differential equations and matrix models, *Commun. Math. Phys.* 163 (1994), 33–72.
- 17. E. T. WHITTAKER AND G. N. WATSON, *A course of modern analysis*, 4th edn (Cambridge University Press, 1965).