



#### RESEARCH ARTICLE

# Cusps, Kleinian groups, and Eisenstein series

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#### Abstract

We study the Eisenstein series associated to the full rank cusps in a complete hyperbolic manifold. We show that given a Kleinian group  $\Gamma < \mathrm{Isom}^+(\mathbb{H}^{n+1})$ , each full rank cusp corresponds to a cohomology class in  $H^n(\Gamma, V)$ , where V is either the trivial coefficient or the adjoint representation. Moreover, by computing the intertwining operator, we show that different cusps give rise to linearly independent classes.

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# 1. Introduction

We say  $\Gamma$  is a Kleinian group if it is a discrete isometry subgroup of  $G = \operatorname{Isom}^+(\mathbb{H}^n)$ , the orientation preserving isometry group of  $\mathbb{H}^n$ . One of the main themes in hyperbolic geometry is to study the number of cusps in the associated quotient manifold  $\Gamma \backslash \mathbb{H}^n$ . When  $\Gamma$  contains parabolic elements, every cusp corresponds to a  $\Gamma$ -conjugacy class of maximal parabolic subgroups in  $\Gamma$ . In dimension 3, the celebrated work of Sullivan [Sul81] shows that finitely generated Kleinian groups always have finitely many cusps, and the number of cusps is bounded by 5N - 4, where N denotes the number of generators in  $\Gamma$  (see also the work of Kra [Kra84]). However, starting in dimension 4, cusp finiteness theorem fails. The first example was due to Kapovich [Kap95], where he constructed a finitely generated free Kleinian group

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 $\Gamma < \text{Isom}^+(\mathbb{H}^4)$  that has infinitely many rank one cusps. In a recent paper [IMM22], Italiano et al. constructed a finitely generated Kleinian group  $\Gamma < \mathrm{Isom}^+(\mathbb{H}^n)$  which has infinitely many full rank cusps in dimensions  $5 \le n \le 8$ . Moreover,  $\Gamma$  can be made finitely presented in dimensions 7 and 8. On the other hand, it is proved in [LW20] that the number of cusps is bounded by the first Betti number, provided the critical exponent is smaller than 1. One general approach to show a cusp finiteness theorem is to first associate each maximal parabolic subgroup  $\Gamma_i < \Gamma$  with a cohomology class  $\alpha_i \in H^*(\Gamma, V)$ after choosing a suitable coefficient module V, then to show the corresponding cohomology classes for different cusps are linearly independent. Finally, if we know the overall dimension of  $H^*(\Gamma, V)$  is finite, then the number of cusps must also be finite. For example, in Sullivan's proof he chose V to be the polynomial space of degree at most 4, and constructed a cross homomorphism from  $\Gamma$  to V (thus representing a class in  $H^1(\Gamma, V)$  via the Borel series associated to each cusp. Then he showed these representing classes are linearly independent, hence, the number of cusps is bounded in terms of the first Betti number (number of generators of  $\Gamma$ ). Analogous to the Borel series, we can use the Eisenstein series to associate a cusp with a cohomology class in  $H^*(\Gamma, V)$ . Much work was done by Harder [Har75, Har87], Schwermer [Sch94, Sch83], and many others when  $\Gamma$  is an arithmetic lattice in a semisimple Lie group. Using the Borel-Serre compactification, the Eisenstein cohomology naturally arises from the cohomology of the boundary, which has deep relations to the arithmetic aspects of  $\Gamma$ , such as the special values of L-functions. The main purpose of this note is to extend the Eisenstein construction to the context of general Kleinian groups with full rank cusps. We make use of the Poincaré series to obtain absolute convergence of the Eisenstein series. Thus, each full rank cusp corresponds to a cohomology class on the quotient manifold. In order to distinguish these cohomology classes arising from different cusps, we compute the intertwining operators and use them to show that these cohomology classes are indeed linearly independent. The computation of the intertwining operators is very difficult in general. We follow the general approach of Harish-Chandra [HC68] but instead use the Lie group decompositions over the reals, including the Bruhat and Langlands decompositions. In particular, our proof does not rely on the finite volume property or arithmeticity of  $\Gamma$ . In the case of trivial coefficient, we prove,

**Theorem 1.1.** Let  $\Gamma < \text{Isom}^+(\mathbb{H}^{n+1})$  be a torsion free discrete subgroup. If the critical exponent  $\delta(\Gamma) < n$  or  $\Gamma$  is of convergence type, then for any parabolic subgroup  $\Gamma_i < \Gamma$  of rank n, and any generating cohomology class  $\alpha_i \in H^n(\Gamma_i, \mathbb{R}) \cong \mathbb{R}$ , there is a harmonic form  $E(\alpha_i)$  on  $\Gamma \backslash \mathbb{H}^{n+1}$  constructed via the Eisenstein series, such that

- 1. The restriction homomorphism  $H^n(\Gamma, \mathbb{R}) \to H^n(\Gamma_i, \mathbb{R})$  sends  $[E(\alpha_i)]$  to  $\alpha_i$ .
- 2. If  $\Gamma_i, \Gamma_j$  are not  $\Gamma$ -conjugate, then the restriction homomorphism  $H^n(\Gamma, \mathbb{R}) \to H^n(\Gamma_j, \mathbb{R})$  sends  $[E(\alpha_i)]$  to 0.

In the above theorem, the *critical exponent*  $\delta(\Gamma)$  of  $\Gamma$  is defined as

$$\delta(\Gamma) = \inf\{s \ : \ \sum_{\gamma \in \Gamma} e^{-sd(O,\gamma O)} < \infty\}.$$

Note that if  $\Gamma \subset \operatorname{Isom}(\mathbb{H}^{n+1})$ , then  $0 \le \delta(\Gamma) \le n$ . For simplicity, we sometimes write  $\delta$  for  $\delta(\Gamma)$  if the context is clear. The group  $\Gamma$  is said to be of convergence type if the above infimum is achieved. The additional assumption on the critical exponent or on the convergence type of  $\Gamma$  is to assure the absolute convergence of the Eisenstein series. This is necessary for our theorem to hold because in the case  $\Gamma$  is a nonuniform lattice (where  $\delta = n$  and  $\Gamma$  is of divergence type), the degree n-homology classes coming from the cusps form a linearly dependent system, thus by Stokes's theorem, the result in our Theorem 1.1 will never hold. To our surprise, by examining the entire argument in our proof, the nonconvergence of the Eisenstein series is the only place where it fails. However, if we choose the coefficient module to be the Lie algebra  $\mathfrak g$  of G, equipped with the natural adjoint action of  $\Gamma$  inherited from G, then the absolute convergence issue will be resolved. This does not contradict the example of nonuniform lattices, since we do not have Stokes's theorem for  $\mathfrak g$  coefficient. More precisely, we prove,

**Theorem 1.2.** Let  $\Gamma < \operatorname{Isom}^+(\mathbb{H}^{n+1})$  be a torsion free discrete subgroup. Then for any parabolic subgroup  $\Gamma_i < \Gamma$  of rank n, and any cohomology class  $\alpha_i \in H^n(\Gamma_i, \operatorname{Ad})$ , there is a closed differential form  $E(\alpha_i)$  on  $\Gamma \backslash \mathbb{H}^{n+1}$  constructed via the Eisenstein series, such that

1. The restriction homomorphism  $H^n(\Gamma, Ad) \to H^n(\Gamma_i, Ad)$  sends  $[E(\alpha_i)]$  to  $\alpha_i$ . In particular, there is a surjective homomorphism

$$H^n(\Gamma, Ad) \to H^n(\Gamma_i, Ad).$$

2. If  $\Gamma_i$ ,  $\Gamma_j$  are not  $\Gamma$ -conjugate, then the restriction homomorphism  $H^n(\Gamma, Ad) \to H^n(\Gamma_j, Ad)$  sends  $[E(\alpha_i)]$  to 0.

Both Theorems 1.1 and 1.2 give a way to control the number of full rank cusps N on the quotient manifold. The case of trivial coefficient implies  $N \leq \beta_n(\Gamma)$  (given that  $\delta < n$  or that  $\Gamma$  is of convergence type), where  $\beta_n(\Gamma)$  denotes the n-th Betti number of  $\Gamma$ , but this is clear since a full rank cusp is always a topological end and  $H^{n+1}(\Gamma, \Gamma_i, \mathbb{R}) = 0$  unless  $\Gamma$  is a cocompact lattice. So the surjectivity of the restriction homomorphism  $H^n(\Gamma, \mathbb{R}) \to H^n(\Gamma_i, \mathbb{R})$  follows immediately from the long exact sequence for the pair  $(\Gamma, \Gamma_i)$ . In the case of adjoint representation, we obtain a similar bound.

**Corollary 1.3.** Let  $\Gamma < \text{Isom}^+(\mathbb{H}^{n+1})$  be a torsion free discrete subgroup. Then the number of full rank toric cusps of  $\Gamma \backslash \mathbb{H}^{n+1}$  is bounded by

$$N \leq \frac{1}{n} \dim(H^n(\Gamma, \mathrm{Ad})).$$

**Remark 1.4.** The reason why we need to add the toric cusp condition is that in general  $H^n(\Gamma_i, Ad)$  could be trivial (see Proposition 3.5 and Remark 3.6). If the group  $\Gamma$  is LERF, that is any finitely generated subgroup is closed in the profinite topology, then we can always pass onto a finite cover of  $\Gamma \backslash \mathbb{H}^{n+1}$  to assure a given full rank cusp is toric.

### Organization of the paper

In Section 2, we review vector-valued differential forms, their equivalent perspective as functions on Lie groups, and the Lie algebra cohomology. In Section 3, we construct a cohomology class for each full rank parabolic fixed point with the coefficient either  $\mathbb{R}$  or the Lie algebra of Isom<sup>+</sup>( $\mathbb{H}^{n+1}$ ) with the adjoint representation. In Section 4, we construct the Eisenstein series and discuss its closeness and convergence. In Section 5, we investigate the restriction of the Eisenstein series to the horosphere corresponding to any given cusp. In Section 6, we prove Theorems 1.1, 1.2, and Corollary 1.3.

### 2. Preliminary

# 2.1. Vector-valued differential forms

Let V be a finite dimensional real vector space and  $\rho: G \to \operatorname{Aut}(V)$  be any continuous (and hence smooth) representation. Since  $\rho$  restricts to any discrete subgroup  $\Gamma$ , and  $\Gamma$  naturally acts from the left on X = G/K, where K is the stabilizer of a point in G, it follows that  $\Gamma$  also left acts on the trivial bundle  $V \times X$  via

$$\gamma \cdot (v, x) = (\rho(\gamma)(v), \gamma x).$$

Endowed with the trivial connection on  $V \times X$ , it induces a flat bundle structure on the quotient manifold  $M = \Gamma \setminus X$ , which we denote by  $V_\rho$ . It is known that the cohomology with the associated local system  $H^*(\Gamma, V_\rho)$  can be computed using the De Rham complex  $\Omega^*(X, V)^\Gamma$ , where the codifferential operator  $d: \Omega^k(X, V) \to \Omega^{k+1}(X, V)$  is defined by

$$\begin{split} d\omega(X_1,..,X_{k+1}) := \sum_i (-1)^{i+1} X_i \omega(X_1,...,\widehat{X}_i,...,X_{k+1}) \\ + \sum_{i < j} (-1)^{i+j} \omega([X_i,X_j],X_1,...,\widehat{X}_i,...,\widehat{X}_j,...,X_{k+1}). \end{split}$$

To this end, any cohomology class in  $H^*(\Gamma, V_\rho)$  can be represented by a  $\Gamma$ -invariant V-valued closed differential form on X.

### 2.2. Matsushima-Murakami formalism

For the purpose of computations, it is convenient to view alternatively the abovementioned vector valued differential forms as a smooth function from G to  $\text{Hom}(\Lambda^*\mathfrak{g}, V)$ . We follow the original treatment in [MM63, Section 4].

Given any  $\eta \in \Omega^*(X, V)^{\Gamma}$ , we can pullback the differential form to G under the projection  $\pi : G \to X = G/K$ , followed by a twist of a G-action. Define a differential form  $\widetilde{\eta} \in \Omega^*(G, V)$  by

$$\widetilde{\eta}_s := \rho(s^{-1})(\eta \circ \pi)_s, \quad \forall s \in G.$$

Then one can check  $\eta \in \Omega^*(X, V)^{\Gamma}$  if and only if  $\widetilde{\eta}$  satisfies

- 1.  $\widetilde{\eta} \circ L_{\gamma} = \widetilde{\eta}$  for any  $\gamma \in \Gamma$ ,
- 2.  $\widetilde{\eta} \circ R_k = \rho(k^{-1})\widetilde{\eta}$  for any  $k \in K$ , and
- 3.  $i(Y)\widetilde{\eta} = 0$  for any  $Y \in \mathfrak{k}$ ,

where  $\mathfrak{k}$  denotes the Lie algebra of K, and  $L_{\gamma}$ ,  $R_k$  denote the left and right multiplications. Furthermore, we can view  $\widetilde{\eta}$  as a function on G whose values are in  $\operatorname{Hom}(\Lambda^*\mathfrak{g},V)$  by identifying  $T_sG$  with  $T_eG \cong \mathfrak{g}$  via left translation, and the above constraints then turn into

- 1.  $\widetilde{\eta}(\gamma g) = \widetilde{\eta}(g)$  for any  $\gamma \in \Gamma$  and  $g \in G$ ,
- 2.  $\widetilde{\eta}(gk) = \operatorname{Ad}_{\mathfrak{q}}^*(k^{-1}) \otimes \rho(k^{-1})(\widetilde{\eta}(g))$  for any  $k \in K$  and  $g \in G$ , and
- 3.  $i(Y)\widetilde{\eta} = 0$  for any  $Y \in \mathfrak{k}$ ,

where  $\operatorname{Ad}^*_{\mathfrak{g}}$  is the dual adjoint representation of G on  $\Lambda^*(\mathfrak{g})$ . Fix a basepoint on X, we write  $\mathfrak{g}=\mathfrak{k}\oplus\mathfrak{p}$ , the Cartan decomposition. Then there is a natural identification between  $\Omega^*(X,V)^\Gamma$  and functions  $\varphi\in C^\infty(G/\Gamma,\operatorname{Hom}(\Lambda^*\mathfrak{p},V))$  which satisfies

$$\varphi(gk) = \operatorname{Ad}_{\mathfrak{p}}^*(k^{-1}) \otimes \rho(k^{-1})(\varphi(g)), \quad \forall k \in K, g \in G.$$

Under such identification, the coboundary operator  $d: \Omega^k(X, V) \to \Omega^{k+1}(X, V)$  as described gives rise to the coboundary operator

$$d: C^{\infty}(G, \operatorname{Hom}(\Lambda^k \mathfrak{g}, V)) \to C^{\infty}(G, \operatorname{Hom}(\Lambda^{k+1} \mathfrak{g}, V))$$

given by ([MM63, Proof of Proposition 4.1])

$$d\varphi(X_1, ..., X_{k+1}) = \sum_{i} (-1)^{i+1} (X_i + \rho(X_i)) \varphi(X_1, ..., \widehat{X}_i, ..., X_{k+1})$$

$$+ \sum_{i < j} (-1)^{i+j} \varphi([X_i, X_j], X_1, ..., \widehat{X}_i, ..., \widehat{X}_j, ..., X_{k+1}).$$
(1)

Here, we abuse notation and still use  $\rho$  to denote the induced Lie algebra representation  $\rho: \mathfrak{g} \to \operatorname{End}(V)$ . Note that in the original statement of [MM63, Proposition 4.1], the second term of the above equation (1) vanishes. This is because their function  $\varphi$  is valued on  $\operatorname{Hom}(\Lambda^*\mathfrak{p},V)$  and that  $[X_i,X_j] \in \mathfrak{k}$  for any pair  $X_i,X_j \in \mathfrak{p}$ . However, for the purpose of computation, besides the usual coset model

 $\mathbb{H}^n = G/K = \mathrm{SO}^+(n,1)/\mathrm{SO}(n)$ , we will also use  $\mathbb{H}^n = P_\xi/K_\xi$ , where  $P_\xi$  is a maximal parabolic group for the parabolic fixed point  $\xi$ , and  $K_\xi = P_\xi \cap K$ . Let us be a little more verbose here as this description is essential to the computations of  $\phi_{\xi,s}$  in Sections 3, 5, and 6.

Under the Langlands decomposition, we have  $P_{\xi} = N_{\xi} A_{\xi} K_{\xi}$ , and accordingly, the Lie algebra splits as  $\mathfrak{p}_{\xi} = \mathfrak{n}_{\xi} \oplus \mathfrak{a}_{\xi} \oplus \mathfrak{m}_{\xi}$ . Thus, from the above discussions, any differential form  $\eta \in \Omega^{k}(X, V)$  can be viewed as a function  $\widetilde{\eta} \in C^{\infty}(P_{\xi}, \operatorname{Hom}(\Lambda^{k}\mathfrak{p}_{\xi}, V))$  which satisfies

1. 
$$\widetilde{\eta}(pm) = \operatorname{Ad}_{\mathfrak{p}_{\xi}}^*(m^{-1}) \otimes \rho(m^{-1})(\widetilde{\eta}(p))$$
 for any  $m \in K_{\xi}$  and  $p \in P_{\xi}$ , and 2.  $i(Y)\widetilde{\eta} = 0$  for any  $Y \in \mathfrak{m}_{\xi}$ ,

where the second property shows that we can further view  $\widetilde{\eta}$  as in  $C^{\infty}(P_{\xi}, \operatorname{Hom}(\Lambda^{k}(\mathfrak{a}_{\xi} \oplus \mathfrak{n}_{\xi}), V))$ . If  $H < P_{\xi}$ , then  $\eta$  is H-invariant if and only if  $\widetilde{\eta}$  is H-left invariant as a function. Note that the Lie bracket  $[\mathfrak{a}_{\xi}, \mathfrak{n}_{\xi}]$  stays in  $\mathfrak{n}_{\xi}$ , so, in particular, the second term in (1) will possibly be nonzero (see Proposition 3.3 and compare the proof of [Har75, Lemma 3.1]).

**Remark 2.1.** Our convention uses left action of G on X = G/K, which is different from that in [Har75]. So there are sign differences in the expression of the coboundary operators.

**Lemma 2.2.** For any  $g \in G$ , if  $L_g$  denotes the left action on  $C^{\infty}(G, \operatorname{Hom}(\Lambda^*\mathfrak{g}, V))$ , that is  $(L_g \varphi)|_a = \varphi|_{ga}$  for any  $\varphi \in C^{\infty}(G, \operatorname{Hom}(\Lambda^k\mathfrak{g}, V))$  and any  $a \in G$ , then

$$L_{g} \circ d = d \circ L_{g}$$
.

In particular,  $\varphi$  is closed if and only if  $L_g \varphi$  is closed, and  $\varphi$  is a coboundary if and only if  $L_g \varphi$  is a coboundary.

*Proof.* For any  $a \in G$ ,  $\varphi \in C^{\infty}(G, \text{Hom}(\Lambda^k \mathfrak{g}, V))$ , and any  $X_1, ..., X_{k+1} \in \mathfrak{g}$ . We do the following direct computations:

$$\begin{split} d(L_g\varphi)|_a(X_1,...,X_{k+1}) &= \sum_i (-1)^{i+1} (X_i + \rho(X_i)) (L_g\varphi)|_a(X_1,...,\widehat{X}_i,...,X_{k+1}) \\ &+ \sum_{i < j} (-1)^{i+j} (L_g\varphi)|_a([X_i,X_j],X_1,...,\widehat{X}_i,...,\widehat{X}_j,...,X_{k+1}) \\ &= \sum_i (-1)^{i+1} (X_i + \rho(X_i)) \varphi|_{ga}(X_1,...,\widehat{X}_i,...,\widehat{X}_i,...,X_{k+1}) \\ &+ \sum_{i < j} (-1)^{i+j} \varphi|_{ga}([X_i,X_j],X_1,...,\widehat{X}_i,...,\widehat{X}_j,...,X_{k+1}) \\ &= d\varphi|_{ga}(X_1,...,X_{k+1}) \\ &= L_g(d\varphi)|_a(X_1,...,X_{k+1}). \end{split}$$

Thus,  $L_g \circ d = d \circ L_g$ .

# 2.3. Lie algebra cohomology

Let  $\mathfrak{g}$  be a Lie algebra and  $\rho: \mathfrak{g} \to \operatorname{End}(V)$  be a Lie algebra representation. We define the Chevalley–Eilenberg complex by

$$\cdots \to \operatorname{Hom}(\Lambda^k \mathfrak{g}, V) \xrightarrow{d} \operatorname{Hom}(\Lambda^{k+1} \mathfrak{g}, V) \to \cdots$$

and the coboundary operator is given by

$$\begin{split} d\varphi(X_1,...,X_{k+1}) &= \sum_i (-1)^{i+1} \rho(X_i) \varphi(X_1,...,\widehat{X}_i,...,X_{k+1}) \\ &+ \sum_{i < j} (-1)^{i+j} \varphi([X_i,X_j],X_1,...,\widehat{X}_i,...,\widehat{X}_j,...,X_{k+1}). \end{split}$$

The cohomology induced by the above cochain complex is called the Lie algebra cohomology with V-coefficient, denoted by  $H^*(\mathfrak{g}, V_o)$ .

Since we only work with specific Lie algebras and representations, we will make simplifications by setting  $G = \operatorname{Isom}^+(\mathbb{H}^{n+1}) \cong \operatorname{SO}^+(n+1,1)$ , and setting either  $V = \mathfrak{g}$  and  $\rho : G \to \operatorname{End}(\mathfrak{g})$  the adjoint representation, or  $V = \mathbb{R}$  and  $\rho$  the trivial representation. Let U < G be a maximal unipotent subgroup associated to some chosen maximal abelian subgroup A < G, such that U is expanding. In the case of adjoint representation, we denote  $\mathfrak{u}, \mathfrak{g}$  the Lie algebra of U, G, and  $\rho : \mathfrak{u} \to \operatorname{End}(\mathfrak{g})$  the restriction of the adjoint representation  $\rho$ . Diagonalized by the adjoint action of A, the vector space V (under the restricted root space decomposition) decomposes as  $V = V_{-2} \oplus V_0 \oplus V_2$  and that the Lie algebra  $\mathfrak{u} = V_2$ .

**Lemma 2.3.** Following the same notations above, if  $\{u_1, ..., u_n\}$  is a basis of  $\mathfrak{u}$ , and  $v \in V_{-2}$ , then there is a natural isomorphism

$$J: V_{-2} \cong H^n(\mathfrak{u}, V_{\rho}),$$

given by

$$v \mapsto (u_1^* \wedge ... \wedge u_n^*) \otimes v,$$

where  $\{u_1^*, ..., u_n^*\}$  represents the dual basis on  $\mathfrak{u}^*$ .

*Proof.* Since  $n = \dim \mathfrak{u}$ , the cochain complex stops in dimension n, so  $(u_1^* \wedge ... \wedge u_n^*) \otimes v$  is automatically closed. For the injectivity of J, it suffices to show that for any nonzero  $v \in V_{-2}$ , the closed form  $(u_1^* \wedge ... \wedge u_n^*) \otimes v$  does not lie in the image of

$$d: \operatorname{Hom}(\Lambda^{n-1}\mathfrak{u}, V) \to \operatorname{Hom}(\Lambda^n\mathfrak{u}, V).$$

We write an arbitrary element in  $\operatorname{Hom}(\Lambda^{n-1}\mathfrak{u}, V)$  as

$$\varphi = \sum_{i=1}^{n} \left( u_1^* \wedge \ldots \wedge \widehat{u_i^*} \wedge \ldots \wedge u_n^* \right) \otimes A_i,$$

for some  $A_i \in V$ . Then we compute

$$d\varphi = \left(u_1^* \wedge \dots \wedge u_n^*\right) \otimes \left(\sum_{i=1}^n (-1)^{i+1} \rho(u_i) A_i\right).$$

Since  $[V, V_2] = V_0 \oplus V_2$ , we see that  $\left(\sum_{i=1}^k (-1)^{i+1} \rho(u_i) A_i\right) \in V_0 \oplus V_2$ , and, in particular, it does not lie in  $V_{-2}$ . Thus,  $(u_1^* \wedge ... \wedge u_n^*) \otimes v$  is not a coboundary, and it represents a nontrivial cohomology class. This shows the map J is injective.

To show J is surjective, we need to show any element  $(u_1^* \wedge ... \wedge u_n^*) \otimes v$  where  $v \in V_0 \oplus V_2$  is a coboundary. Since  $[V, V_2] = V_0 \oplus V_2$ , we can write  $v = \sum_{i=1}^n [u_i, v_i]$  for some  $v_i \in V$ . Now if we set

$$\varphi = \sum_{i=1}^{n} \left( u_1^* \wedge \ldots \wedge \widehat{u_i^*} \wedge \ldots \wedge u_n^* \right) \otimes (-1)^{i+1} v_i \in \operatorname{Hom}(\Lambda^{n-1}\mathfrak{u}, V),$$

then  $d\varphi = (u_1^* \wedge ... \wedge u_n^*) \otimes v$ . This proves that  $(u_1^* \wedge ... \wedge u_n^*) \otimes v$  is a coboundary. Thus, J is surjective.  $\Box$ 

One important aspect of the Lie algebra cohomology is that it sometimes relates to the group cohomology of a Lie group G, and those classes can be identified with certain G-invariant differential forms which are harmonic. We will use the following Van Est isomorphism theorem in the context of abelian Lie groups.

**Theorem 2.4.** [VE58] Let U be an n-dimensional Lie group isomorphic to  $\mathbb{R}^n$ , and Z < U be a torsion free cocompact lattice. Let  $\rho : U \to \operatorname{Aut}(V)$  be a representation which induces the Lie algebra representation  $\rho' : \mathfrak{u} \to \operatorname{End}(V)$ . Then there is a natural isomorphism

$$\Phi: H^*(Z, V_{\rho}) \cong H^*(\mathfrak{u}, V_{\rho'}),$$

explicitly given by the following, for any Z-invariant closed differential form  $\omega \in \Omega^*(U, V_o)$ , set

$$\Phi(\omega) = \int_{U/Z} \omega(x) d\mu(x),$$

where  $d\mu$  is the Haar measure on U/Z. Thus,  $\Phi(\omega)$  is an U-invariant differential form that can be identified with an element in  $\text{Hom}(\Lambda^*(\mathfrak{u}), V_{\rho'})$ .

# 3. Construction of cohomology classes from a cusp

Suppose that  $G = \operatorname{Isom}^+(\mathbb{H}^{n+1}) \cong \operatorname{SO}^+(n+1,1)$ , and  $\Gamma < G$  is a torsion free discrete subgroup. For each cusp on the quotient manifold  $\Gamma \backslash \mathbb{H}^{n+1}$ , when lifted to the universal cover, it associates to a  $\Gamma$ -orbit  $\Gamma \xi$  on  $\partial_{\infty} \mathbb{H}^{n+1}$ , for some  $\xi \in \partial_{\infty} \mathbb{H}^{n+1}$ . The fundamental group of the cusp is isomorphic to  $\Gamma_{\xi} = \Gamma \cap P_{\xi}$ , where  $P_{\xi} = \operatorname{stab}_{G}(\xi) < G$  is the real parabolic subgroup at  $\xi$ .

Fix a basepoint  $O \in \mathbb{H}^{n+1}$ , and let  $\mathcal{H}_{\xi}(1)$  be the horosphere of the parabolic fixed point  $\xi$  through O. Under the induced Riemannian metric, it is isometric to the standard Euclidean space  $\mathbb{R}^n$ . It is known that  $\Gamma_{\xi}$  preserves and acts isometrically on  $\mathcal{H}_{\xi}(1)$ , and in fact, by Bieberbach's theorem, it acts cocompactly on a k-dimensional Euclidean subspace  $\mathbb{E}^k_{\xi}$  of  $\mathbb{R}^n \cong \mathcal{H}_{\xi}(1)$ . It follows that  $\Gamma_{\xi}$  has a finite index abelian normal subgroup  $Z_{\xi}$  which acts on  $\mathbb{E}^k_{\xi}$  by translations. We call k the rank of the cusp at  $\xi$ . For the purpose of this paper, we will only consider the full rank case k=n, and from now on all cusps are assumed to be full rank.

Under the Langlands decomposition, the real parabolic subgroup decomposes as  $P_{\xi} = N_{\xi}A_{\xi}K_{\xi}$ , where  $A_{\xi}$  is the maximal abelian subgroup which acts by translation on the geodesic connecting O and  $\xi$ . Geometrically,  $K_{\xi} \cong SO(n)$  and  $N_{\xi} \cong \mathbb{R}^n$  are the rotations and translations on  $\mathcal{H}_{\xi}(1)$ , respectively. Denote  $K_{\xi}N_{\xi} = N_{\xi}K_{\xi}$  by  $P_{\xi}(1)$ . Then it is indeed the orientation preserving isometry group of  $\mathcal{H}_{\xi}(1)$ , so  $\Gamma_{\xi}$  is a discrete subgroup in  $P_{\xi}(1)$ . We say the cusp is *toric* if  $\Gamma_{\xi}$  is isomorphic to  $\mathbb{Z}^n$ , since under the quotient, the cusp is homeomorphic to  $\mathbb{T}^n \times [0, \infty)$ .

For each cusp, we lift to the universal cover and choose an arbitrary parabolic fixed point  $\xi \in \partial_{\infty} \mathbb{H}^{n+1}$  representing the cusp. For the convenience, we will describe the following construction under the upper-half plane model  $\mathbb{H}^{n+1} = \{(y, x_1, ..., x_n) \in \mathbb{R}^+ \times \mathbb{R}^n\}$ . We assume O = (1, 0, ..., 0) and  $\xi$  is in the positive infinity of the y-axis, thus,  $\mathcal{H}_{\xi}(1) = \{(y, x_1, ..., x_n) : y = 1\}$ . For a different  $\xi' \in \partial_{\infty} \mathbb{H}^{n+1}$ , the construction differs by a  $k(\xi, \xi')$ -conjugate, where  $k(\xi, \xi') \in K$  is any element which sends  $\xi$  to  $\xi'$ .

### 3.1. Trivial coefficients

In the case of trivial representation  $V = \mathbb{R}$ , we define  $\phi_{\xi}$  to be the canonical volume form on  $\mathcal{H}_{\xi}(1)$ , that is,

$$(\phi_{\mathcal{E}})_x = dx_1 \wedge ... \wedge dx_n$$

for any  $x \in \mathcal{H}_{\xi}(1)$ . It is convenient (for the purpose of computation) to view  $\phi_{\xi}$  also as a function in  $C^{\infty}(P_{\xi}(1), \Lambda^n \mathfrak{n}_{\xi}^*)$ , according to the discussion in Section 2.1. We choose the orthonormal frame  $\{u_1, ..., u_n\}$  on  $\mathfrak{n}_{\xi}$  normalized so that each  $u_i$  exponentiates to the unit translation on  $\mathcal{H}_{\xi}(1) \cong \mathbb{R}^n$  along the  $x_i$ -axis. Thus,  $\phi_{\xi}$  satisfies,

$$\begin{split} &1. \ \ \phi_{\xi}(n) = \phi_{\xi}(1) = (u_1^* \wedge \ldots \wedge u_n^*), \text{ for all } n \in N_{\xi}, \\ &2. \ \ \phi_{\xi}(pm) = \operatorname{Ad}_{\mathfrak{n}_{\xi}}^*(m^{-1})(\phi_{\xi}(p)) \text{ for all } m \in K_{\xi} \text{ and } p \in P_{\xi}(1), \end{split}$$

where  $u_i^* \in \mathfrak{n}_{\xi}$  is the dual vector of  $u_i$ . We wish to extend the differential form to the entire  $\mathbb{H}^{n+1}$ , or equivalently, extend  $\phi_{\xi} \in C^{\infty}\left(P_{\xi}(1), \Lambda^n\mathfrak{n}_{\xi}^*\right)$  to a function in  $C^{\infty}\left(P_{\xi}, \Lambda^n(\mathfrak{a}_{\xi}^* \oplus \mathfrak{n}_{\xi}^*)\right)$ . Following [Har75], we introduce the following *degree* s extension  $\phi_{\xi,s}$  of  $\phi_{\xi}$ . Let  $t_{\xi}: A_{\xi} \to \mathbb{R}$  be the character that corresponds to the positive root on  $\mathfrak{a}_{\xi}$ , that is, for any  $a \in A_{\xi}$  and  $v \in \mathfrak{g}$ , we have

$$\operatorname{Ad}(a)v = \begin{cases} t_{\xi}^{2}(a)v & \text{if } v \in \mathfrak{g}_{2} \\ 0 & \text{if } v \in \mathfrak{g}_{0} \\ t_{\xi}^{-2}(a)v & \text{if } v \in \mathfrak{g}_{-2} \end{cases}$$
 (2)

where  $\mathfrak{g} = \mathfrak{g}_2 \oplus \mathfrak{g}_0 \oplus \mathfrak{g}_{-2}$  is the root space decomposition corresponding to  $A_{\xi}$  (take  $\xi$  to be the positive direction). Therefore, if  $T \in \mathfrak{a}_{\xi}$  is the vector, such that [T, u] = 2u for all  $u \in \mathfrak{n}_{\xi}$ , then we have  $(dt_{\xi})_a(T) = t_{\xi}(a)$ , or that  $dt_{\xi}/t_{\xi}$  is dual to T. We also write  $t_{\xi,a}$  for  $t_{\xi}(a)$  when convenient, and omit the subscript  $\xi$  when there is no confusion in the context. Now we define the degree s extension  $\phi_{\xi,s}$  of  $\phi_{\xi}$  by

- 1.  $\phi_{\xi,s}(n) = \phi_{\xi}(1) = (u_1^* \wedge ... \wedge u_n^*)$ , for all  $n \in N_{\xi}$ ,
- 2.  $\phi_{\xi,s}(p) = \operatorname{Ad}_{\mathfrak{t}_{\xi}}^*(m^{-1})(\phi_{\xi}(1))t_{\xi,a}^s$ , for all  $p = nam \in N_{\xi}A_{\xi}K_{\xi}$  under the unique Langlands decomposition.

Thus, in view of Section 2.1,  $\phi_{\xi,s} \in C^{\infty}(P_{\xi}, \Lambda^n(\mathfrak{a}_{\xi}^* \oplus \mathfrak{n}_{\xi}^*))$  represents an  $N_{\xi}$ -invariant differential form on  $\mathbb{H}^{n+1}$ .

**Proposition 3.1.** In case of trivial representation,  $\phi_{\xi,s}$  is closed when s=2n.

The proof is deferred in Proposition 3.3.

**Remark 3.2.** It is much easier to find this unique closed  $N_{\xi}$ -invariant differential form  $\phi_{\xi,s}$  when it is viewed as in  $\Omega^n(\mathbb{H}^{n+1},\mathbb{R})^{N_{\xi}}$ , which is simply the pullback form of  $\phi_{\xi}$  under the natural projection  $\mathbb{H}^{n+1} \to \mathcal{H}_{\xi}(1)$ . In other words, it is the harmonic form

$$(\phi_{\mathcal{E},s})_x = dx_1 \wedge ... \wedge dx_n$$

for any  $x \in \mathbb{H}^{n+1}$ . However, when it comes to the computation of intertwining operators, treating  $\phi_{\xi,s}$  as a function will be more convenient for us.

# 3.2. Adjoint representation

Now, we assume  $V = \mathfrak{g}$  and  $\rho$  is the adjoint representation. We first define a harmonic form on  $\mathcal{H}_{\xi}(1)$ . By Lemma 2.3, in order to obtain nontrivial cohomology classes, we need to take  $V_{-2}$  sections. Let  $v_{\xi} \in V_{-2}$  be an arbitrary nonzero vector. We construct the unique (left)  $N_{\xi}$ -invariant n-form  $\phi_{\xi}$  on  $\mathcal{H}_{\xi}(1)$  by setting the initial value

$$(\phi_{\mathcal{E}})_O := (dx_1 \wedge ... \wedge dx_n) \otimes v_{\mathcal{E}}.$$

Since  $v_{\xi}$  does not commute with  $\mathfrak{n}_{\xi}$ , the  $N_{\xi}$ -action will twist the fiber as we vary the point on the horosphere. More precisely, the differential *n*-form  $\phi_{\xi}$  can be globally defined as

$$(\phi_{\xi})_{uO} := (dx_1 \wedge \dots \wedge dx_n) \otimes \rho(u)(v_{\xi}), \quad \forall u \in N_{\xi}.$$

Alternatively, we can view  $\phi_{\mathcal{E}}$  as a function in  $C^{\infty}(P_{\mathcal{E}}(1), \text{Hom}(\Lambda^n \mathfrak{n}_{\mathcal{E}}, V))$  which satisfies,

- 1.  $\phi_{\xi}(n) = \phi_{\xi}(1) = (u_1^* \wedge \dots \wedge u_n^*) \otimes v_{\xi}$ , for all  $n \in N_{\xi}$ ,
- 2.  $\phi_{\xi}(pm) = \operatorname{Ad}_{n_{\xi}}^{*}(m^{-1}) \otimes \rho(m^{-1})(\phi_{\xi}(p))$  for all  $m \in K_{\xi}$  and  $p \in P_{\xi}(1)$ ,

where again,  $u_i \in \mathfrak{u}_{\xi}$  is the unique vector that exponentiates the unit translation in  $x_i$ , and  $u_i^*$  is its dual vector. Now, we define the degree s extension  $\phi_{\xi,s}$  of  $\phi_{\xi}$  by

- 1.  $\phi_{\xi,s}(n) = \phi_{\xi}(1) = (u_1^* \wedge ... \wedge u_n^*) \otimes v_{\xi}$ , for all  $n \in N_{\xi}$ ,
- 2.  $\phi_{\xi,s}(p) = \phi_{\xi,s}(nam) = \mathrm{Ad}_{n_{\xi}}^*(m^{-1}) \otimes \rho(m^{-1})(\phi_{\xi}(1))t_{\xi,a}^s$ , for all  $p \in P_{\xi}$ .

Thus,  $\phi_{\xi,s} \in C^{\infty}(P_{\xi}, \operatorname{Hom}(\Lambda^n \mathfrak{p}_{\xi}, V))$  represents an  $N_{\xi}$ -invariant V-valued differential form on  $\mathbb{H}^{n+1}$ , where  $\mathfrak{p}_{\xi}$  is the Lie algebra of  $P_{\xi}$ . The following proposition computes the differential of  $\phi_{\xi,s}$  in the most generality (for later use in Section 6), and, in particular, it characterizes when  $\phi_{\xi,s}$  is closed. This is essentially proved in [Har75, Lemma 3.1], but for completeness, we add it here.

**Proposition 3.3.** Let  $\phi_{\xi,s}$  be the degree s extension of  $\phi_{\xi}$ , where  $\phi_{\xi}$  is defined as above except now we allow  $v_{\xi} \in V_{\ell}$  for  $\ell \in \{-2,0,2\}$ , and V is either with the trivial representation or with the adjoint representation. Then

$$d\phi_{\xi,s} = \begin{cases} (s-2n+\ell) \cdot \frac{dt_{\xi}}{t_{\xi}} \wedge \phi_{\xi,s} & \text{if $\rho$ is adjoint representation} \\ (s-2n) \cdot \frac{dt_{\xi}}{t_{\xi}} \wedge \phi_{\xi,s} & \text{if $\rho$ is trivial representation} \end{cases},$$

where  $t_{\xi}: A_{\xi} \to \mathbb{R}$  is the character that corresponds to the positive root on  $\mathfrak{a}_{\xi}$  defined as in (2).

*Proof.* Let  $X_1, \dots, X_{n+1} \in \mathfrak{p}_{\xi} = \mathfrak{t}_{\xi} \oplus \mathfrak{a}_{\xi} \oplus \mathfrak{n}_{\xi}$ , we need to compute  $d\phi_{\xi,s}(X_1, \dots, X_{n+1})$ . Recall that

$$d\phi_{\xi,s}(X_1,\dots,X_{n+1}) = \sum_{i} (-1)^{i+1} (X_i + \rho(X_i)) \phi_{\xi,s}(X_1,\dots,\widehat{X}_i,\dots,X_{n+1})$$
  
+ 
$$\sum_{i < j} (-1)^{i+j} \phi_{\xi,s}([X_i,X_j],X_1,\dots,\widehat{X}_i,\dots,\widehat{X}_j,\dots,X_{n+1}).$$

The form  $\phi_{\xi}$  is a closed *n*-form on the horosphere  $\mathcal{H}_{\xi}(1)$ . It suffices to only consider the term  $d\phi_{\xi,s}(T,X_1,\cdots,X_n)$ , where  $T\in\mathfrak{a}_{\xi}$  and  $X_i\in\mathfrak{n}_{\xi}$  since all the other terms vanish (see [MM63, (4.10)]). We choose  $X_i=u_i\in\mathfrak{n}_{\xi}$  and normalize T, such that  $[T,X_i]=2X_i$ , that is T is dual to  $dt_{\xi}/t_{\xi}$ . Then at any point  $p=na\in\mathcal{N}_{\xi}A_{\xi}$ , we have

$$d\phi_{\xi,s}(T, X_1, \dots, X_n) = (T + \rho(T))\phi_{\xi,s}(X_1, \dots, X_n) + \sum_{i} (-1)^i \phi_{\xi,s}([T, X_i], X_1, \dots, \widehat{X}_i, \dots, X_n).$$

The second term can be simplified as

$$\sum_{i} (-1)^{i} \phi_{\xi,s}([T, X_{i}], X_{1}, \cdots, \widehat{X}_{i}, \cdots, X_{n}) = -2 \sum_{i=1}^{n} \phi_{\xi,s}(X_{1}, \cdots, X_{i}, \cdots, X_{n})$$
$$= -2n \cdot \phi_{\xi,s}(X_{1}, \cdots, X_{n}).$$

Furthermore, we compute

$$\rho(T)\phi_{\xi,s}(X_1,\dots,X_n) = \begin{cases} \ell \cdot \phi_{\xi,s}(X_1,\dots,X_n) & \text{if } \rho \text{ is adjoint representation} \\ 0 & \text{if } \rho \text{ is trivial representation} \end{cases}$$

and

$$T\phi_{\mathcal{E},s}(X_1,\cdots,X_n)=s\cdot\phi_{\mathcal{E},s}(X_1,\cdots,X_n).$$

Therefore.

$$d\phi_{\xi,s} = \begin{cases} (s-2n+\ell) \cdot \frac{dt_{\xi}}{t_{\xi}} \wedge \phi_{\xi,s} & \text{if } \rho \text{ is adjoint representation} \\ (s-2n) \cdot \frac{dt_{\xi}}{t_{\xi}} \wedge \phi_{\xi,s} & \text{if } \rho \text{ is trivial representation} \end{cases}.$$

**Corollary 3.4.** In case of adjoint representation and  $v_{\xi} \in V_{-2}$ ,  $\phi_{\xi,s}$  is closed when s = 2n + 2.

## 3.3. Nontoric cusps

In the case of toric cusps, we know  $\Gamma_{\xi} < N_{\xi}$ , hence, the above constructed differential form  $\phi_{\xi}$  (hence also its extension  $\phi_{\xi,s}$ ) is automatically  $\Gamma_{\xi}$ -invariant, and by the Van Est isomorphism (Theorem 2.4), every cohomology class in  $H^n(\Gamma_{\xi}, \operatorname{Ad})$  can be represented this way. For a general nontoric cusp,  $\Gamma_{\xi}$  is a Bieberbach group. With trivial coefficient, we still have  $H^n(\Gamma_{\xi}, \mathbb{R}) \cong H^n(Z_{\xi}, \mathbb{R}) \cong H^n(\mathfrak{u}, \mathbb{R}) \cong \mathbb{R}$ , and the canonical volume form is  $\Gamma_{\xi}$ -invariant.

However, with adjoint representation it is less clear. Since  $Z_{\xi}$  is a finite index normal subgroup of  $\Gamma_{\xi}$ , the transfer map  $\iota: H^k(Z_{\xi}, \operatorname{Ad}) \to H^k(\Gamma_{\xi}, \operatorname{Ad})$ , given by taking the average over the finite group action of  $Z_{\xi} \setminus \Gamma_{\xi}$ , is a left inverse of the restriction map  $i^*$ . That is, the following composition

$$H^k(\Gamma_{\mathcal{E}}, \mathrm{Ad}) \xrightarrow{i^*} H^k(Z_{\mathcal{E}}, \mathrm{Ad}) \xrightarrow{\iota} H^k(\Gamma_{\mathcal{E}}, \mathrm{Ad})$$

is the identity map. In particular,  $i^*$  is injective and  $\iota$  is surjective. Therefore, we can identify cohomology classes of  $H^k(\Gamma_{\mathcal{E}}, \operatorname{Ad})$  with their images in  $H^k(Z_{\mathcal{E}}, \operatorname{Ad})$ .

**Proposition 3.5.** For any cohomology class  $\alpha \in H^n(\Gamma_{\xi}, \mathrm{Ad})$ , there is a unique  $N_{\xi}$ -invariant,  $\Gamma_{\xi}$ -invariant closed differential form  $\overline{\psi_{\alpha}} \in C^{\infty}(P_{\xi}(1), \mathrm{Hom}(\Lambda^n \mathfrak{n}_{\xi}, V))$  representing  $\alpha$  whose initial value satisfies,

$$\overline{\psi_\alpha}(1) = (u_1^* \wedge \cdots \wedge u_n^*) \otimes v$$

for some  $v \in V_{-2}$ . Moreover,  $\Theta v$  must be fixed by  $\Gamma_{\xi}$  under the adjoint action, where  $\Theta : \mathfrak{g} \to \mathfrak{g}$  is the Cartan involution associated with the basepoint O. Thus,  $H^n(\Gamma_{\xi}, \operatorname{Ad}) \neq 0$  if and only if  $\Gamma_{\xi}$  fixes a nontrivial vector in  $\mathfrak{n}_{\xi}$ .

*Proof.* We first choose the harmonic representative of  $i^*(\alpha)$  in  $H^n(Z_\xi, \operatorname{Ad}) \cong H^n(\mathfrak{n}_\xi, \operatorname{ad})$  under the Van Est isomorphism, and we denote it by  $\psi_\alpha$ . Then, by Lemma 2.3, we can choose an  $N_\xi$ -invariant form  $\psi_\alpha \in C^\infty(P_\xi(1), \operatorname{Hom}(\Lambda^n\mathfrak{n}_\xi, V))$ , such that  $\psi_\alpha(1) = (u_1^* \wedge \cdots \wedge u_n^*) \otimes v_\alpha$  for some  $v_\alpha \in V_{-2}$ . Since  $\psi_\alpha$  is a top form, it is closed. We denote  $\overline{\psi_\alpha}$  the image of  $\psi_\alpha$  under the transfer map, hence, by definition, we have

$$\overline{\psi_{\alpha}}(p) = \frac{1}{D} \sum_{[\gamma_i] \in Z_{\mathcal{E}} \setminus \Gamma_{\mathcal{E}}} \psi_{\alpha}(\gamma_i \cdot p),$$

where  $D = [\Gamma_{\xi} : Z_{\xi}]$ . By Lemma 2.2,  $\overline{\psi_{\alpha}}$  is closed. Clearly it is  $\Gamma_{\xi}$ -invariant, and since  $\iota \circ i^* = Id$ ,  $\overline{\psi_{\alpha}}$  represents  $\alpha$ . To see it is also  $N_{\xi}$ -invariant, we compute for any  $u \in N_{\xi}$  and  $p \in P_{\xi}(1)$ 

$$\begin{split} \overline{\psi_{\alpha}}(up) &= \frac{1}{D} \sum_{[\gamma_i] \in Z_{\mathcal{E}} \setminus \Gamma_{\mathcal{E}}} \psi_{\alpha}(\gamma_i \cdot up) \\ &= \frac{1}{D} \sum_{[\gamma_i] \in Z_{\mathcal{E}} \setminus \Gamma_{\mathcal{E}}} \psi_{\alpha}(\gamma_i u \gamma_i^{-1} \cdot \gamma_i p) \\ &= \frac{1}{D} \sum_{[\gamma_i] \in Z_{\mathcal{E}} \setminus \Gamma_{\mathcal{E}}} \psi_{\alpha}(\gamma_i p) \\ &= \overline{\psi_{\alpha}}(p), \end{split}$$

where the third equality uses the  $N_{\xi}$ -invariance of  $\psi_{\alpha}$  together with the fact that  $P_{\xi}(1)$  normalizes  $N_{\xi}$ . To compute  $\overline{\psi_{\alpha}}(1)$ , we can write  $\gamma_i$  uniquely as  $n_i m_i \in P_{\xi}(1)$ , then by the  $N_{\xi}$ -invariance of  $\psi_{\alpha}$ , we obtain

$$\overline{\psi_{\alpha}}(1) = \frac{1}{N} \sum_{[\gamma_i] \in Z_{\mathcal{E}} \setminus \Gamma_{\mathcal{E}}} \psi_{\alpha}(m_i)$$

$$= \frac{1}{N} \sum_{i} \operatorname{Ad}(m_i^{-1}) \otimes \rho(m_i^{-1}) \left( (u_1^* \wedge \dots \wedge u_n^*) \otimes v_{\alpha} \right)$$

$$= (u_1^* \wedge \dots \wedge u_n^*) \otimes \left( \frac{1}{N} \sum_{i} \rho(m_i^{-1}) (v_{\alpha}) \right),$$

where the last equality uses the fact that  $K_{\xi}$  acts isometrically on  $\mathcal{H}_{\xi}(1)$  so, in particular, it preserves its volume form. Since  $A_{\xi}$  commutes with  $K_{\xi}$ , it is clear that

$$v := \left(\frac{1}{N} \sum_{i} \rho(m_i^{-1})(v_\alpha)\right) \in V_{-2}.$$

The uniqueness of v follows from the injectivity of  $i^*$ . Finally, to see why  $\Theta v$  is fixed by  $\Gamma_{\xi}$ , we first note that the collection of  $\{m_i\}$  form a group. Thus,  $\rho(m_i)v = v$  for every  $m_i$ . Since  $\Theta$  fixes the Lie algebra of  $K_{\xi}$ , it commutes with  $\rho(m_i)$  for all  $m_i$ . It follows that  $\rho(m_i)(\Theta v) = \Theta v$  for all  $m_i$ . Since  $\Theta v \in \mathfrak{n}_{\xi}$ ,  $\Gamma_{\xi}$  fixes  $\Theta v$ .

**Remark 3.6.** The proposition is an explicit realization of the isomorphism  $H^n(\Gamma_{\xi}, \operatorname{Ad}) \cong (H^n(Z_{\xi}, \operatorname{Ad}))^{\Gamma_{\xi}}$  obtained, for example, via the spectral sequence. In the example of Kleinian group with infinitely many cusps constructed by Italiano et al. [IMM22], they are toric. However, there are in general nontoric examples (see [FKS21]), where  $\Gamma_{\xi}$  does not fix any nontrivial vector in  $\mathfrak{n}_{\xi}$ , hence,  $H^n(\Gamma_{\xi}, \operatorname{Ad}) = 0$ .

#### 4. Construction of Eisenstein series

In Section 3, we constructed  $\Gamma_{\xi}$ -invariant V-valued  $(V = \mathbb{R} \text{ or } \mathfrak{g})$  n-forms  $\phi_{\xi,s}$  on  $\mathbb{H}^{n+1}$  out from any full rank parabolic fixed point  $\xi \in \partial_{\infty} \mathbb{H}^{n+1}$ . In this section, we will further construct from each  $\phi_{\xi,s}$  a  $\Gamma$ -invariant n-form on  $\mathbb{H}^{n+1}$ , that is an element in  $\Omega^n(\mathbb{H}^{n+1},V)^{\Gamma}$ , by the process of *Eisenstein series*. Since we will need  $\Gamma$ -actions on the differential forms, we want to first extend our definition of  $\phi_{\xi,s}$  to a function that supports on the entire G, but of course this is uniquely determined because  $\phi_{\xi,s}$  is already a form on  $\mathbb{H}^{n+1}$ . By the discussion in Section 2.1 and in view of its value on  $P_{\xi}$ , the unique extension  $\phi_{\xi,s} \in C^{\infty}(G, \operatorname{Hom}(\Lambda^n(\mathfrak{a}_{\xi} \oplus \mathfrak{n}_{\xi}), V))$  satisfies for any  $g \in G$ ,

1. 
$$\phi_{\xi,s}(g) = \phi_{\xi,s}(n \cdot a \cdot k) = \mathrm{Ad}_{\mathfrak{g}}^*(k^{-1}) \otimes \rho(k^{-1})(\phi_{\xi,s}(n)) \cdot t_a^s$$
,  
2.  $\phi_{\xi,s}(n) = \phi_{\xi,s}(1) = (u_1^* \wedge \cdots \wedge u_n^*) \otimes v_{\xi}$ ,

where  $g = nak \in N_{\xi}A_{\xi}K$  is the unique Iwasawa decomposition, and  $v_{\xi}$  is any vector in  $V_{-2}$  if  $V = \mathfrak{g}$ . From Proposition 3.5, we know that any cohomology class in  $H^n(\Gamma_{\xi}, Ad)$  can be represented by the above  $\phi_{\xi,s}$ , which is (left)  $N_{\xi}$ -invariant and  $\Gamma_{\xi}$ -invariant. We now define

$$E(\phi_{\xi}, g, s) = \sum_{\gamma \in \Gamma_{\xi} \setminus \Gamma} \phi_{\xi, s}(\gamma g)$$

to be the Eisenstein series associated with  $\phi_{\xi}$  (with degree s). We will see later there is a unique s, such that  $E(\phi_{\xi}, g, s)$  is close, so we will not emphasize on the dependence of s. By the construction, it is  $\Gamma$ -invariant. But it is unclear whether it converges or not.

**Proposition 4.1.** Let  $\xi$  be a full rank parabolic fixed point. The Eisenstein series  $E(\phi_{\xi}, g, s)$  absolutely converges if the Poincaré series  $P_{s/2}(O) := \sum_{\gamma \in \Gamma} e^{-(s/2)d(O, \gamma O)}$  converges for some  $O \in \mathbb{H}^{n+1}$ .

*Proof.* Fix any  $g \in G$ , by the Iwasawa decomposition, we have  $\gamma g = n_{\gamma} a_{\gamma} k_{\gamma} \in N_{\xi} A_{\xi} K$ , where K is the stabilizer of O. We can write each term

$$\phi_{\xi,s}(\gamma g) = \phi_{\xi,s}(n_\gamma a_\gamma k_\gamma) = \mathrm{Ad}_{\mathfrak{g}}^*(k_\gamma^{-1}) \otimes \rho(k_\gamma^{-1}) \phi_\xi(1) \cdot t_\gamma^s,$$

where we set  $t_{\gamma} = t(a_{\gamma})$ . If we fix any norm on  $\operatorname{Hom}(\Lambda^n(\mathfrak{a}_{\xi} \oplus \mathfrak{n}_{\xi}), V)$ , then we can estimate

$$||E(\phi_{\xi},g,s)|| \leq \sum_{\gamma \in \Gamma_{\xi} \backslash \Gamma} ||\phi_{\xi,s}(\gamma g)|| = \sum_{\gamma \in \Gamma_{\xi} \backslash \Gamma} ||\operatorname{Ad}_{\mathfrak{g}}^{*}(k_{\gamma}^{-1}) \otimes \rho(k_{\gamma}^{-1}) \phi_{\xi}(1)|| \cdot t_{\gamma}^{s}.$$

Since K is compact,  $\operatorname{Ad}_{\mathfrak{g}}^*(k_{\gamma}^{-1}) \otimes \rho(k_{\gamma}^{-1}) \phi_{\xi}(1)$  is uniformly bounded in  $\Lambda^n \mathfrak{g}^* \otimes V$ , hence also uniformly bounded when projected onto  $\Lambda^n(\mathfrak{a}_{\xi}^* \oplus \mathfrak{n}_{\xi}^*) \otimes V$ . So for the absolute convergence of the Eisenstein series, it suffices to consider the series

$$\sum_{\gamma \in \Gamma_{\xi} \setminus \Gamma} t_{\gamma}^{s}$$

and show its convergence. Note that the above series is well-defined, since  $t_{\gamma}$  does not depend on the choice of  $\gamma$  in the coset  $\Gamma_{\xi} \setminus \Gamma$ . Indeed, if  $\gamma_0 \in \Gamma_{\xi} = \Gamma \cap P_{\xi}$  is any element, then  $\gamma_0 \in N_{\xi}K_{\xi}$ , so that we can write  $\gamma_0 = n_0 m_0$ . Since  $m_0$  commutes with  $a_{\gamma}$  and normalizes  $N_{\xi}$ , we have  $\gamma_0 \gamma_0 g = n_0 m_0 n_{\gamma} a_{\gamma} k_{\gamma} = n_0 (m_0 n_{\gamma} m_0^{-1}) a_{\gamma} m_0 k_{\gamma}$ , which gives rises to the  $N_{\xi} A_{\xi} K$ -Iwasawa decomposition. This shows that the  $A_{\xi}$  component does not change when replacing  $\gamma$  by  $\gamma_0 \gamma$ , thus  $t_{\gamma}$  is independent on the choice of  $\gamma$  in the coset.

Next, we want to relate the series to the Poincaré series. Note that since  $\xi$  is a full rank parabolic fixed point,  $\Gamma_{\xi}$  acts cocompactly on  $\mathcal{H}_{\xi}(1)$ , so there exists a constant C such that a fundamental domain of  $\Gamma_{\xi} \setminus \mathcal{H}_{\xi}(1)$  is contained in the metric ball B(O, C/2). Thus, for any coset  $\Gamma_{\xi} \gamma$ , there exists a representative  $\bar{\gamma} \in \Gamma_{\xi} \gamma$ , such that under the Iwasawa decomposition  $\bar{\gamma}g = n_{\gamma}a_{\gamma}k_{\gamma}$ , the unipotent component  $n_{\gamma}$  translates O at most C, that is  $d(n_{\gamma}O, O) \leq C$ . Now we can estimate the Poincaré series

$$\begin{split} P_s(gO) &= \sum_{\gamma \in \Gamma} e^{-sd(O,\gamma gO)} \geq \sum_{[\bar{\gamma}] \in \Gamma_{\mathcal{E}} \backslash \Gamma} e^{-sd(O,n_{\gamma}a_{\gamma}O)} \\ &\geq \sum_{[\bar{\gamma}] \in \Gamma_{\mathcal{E}} \backslash \Gamma} e^{-s(d(O,n_{\gamma}O) + d(n_{\gamma}O,n_{\gamma}a_{\gamma}O))} \\ &\geq e^{-sC} \sum_{[\bar{\gamma}] \in \Gamma_{\mathcal{E}} \backslash \Gamma} e^{-sd(O,a_{\gamma}O)}, \end{split}$$

where the first inequality makes the particular choice of  $\bar{\gamma}$  described above, the second inequality uses the triangle inequality, and the last inequality uses the left invariance of the metric.

Observe that  $a_{\gamma}O$  is on the geodesic connecting O and  $\xi$ . By hyperbolic geometric computation,  $d(O,a_{\gamma}O)=|\ln t_{\gamma}^2|$ , and  $t_{\gamma}>1$  if and only if  $a_{\gamma}O$  lies inside the horoball of  $\mathcal{H}_{\xi}(1)$ . In fact, the computation can be carried out within a totally geodesic copy  $\mathbb{H}^2\subset\mathbb{H}^{n+1}$  which contains  $\xi,O$ , where  $a_{\gamma}$  acts by isometry. Then without loss of generality, we can use the upper half plane model and assume  $O=i, \xi=i\infty$ , and  $a_{\gamma}=\mathrm{diag}(a,a^{-1})\in\mathrm{SL}_2(\mathbb{R})$ . The mobius transformation gives  $a_{\gamma}i=a^2i$ , hence  $d(O,a_{\gamma}O)=d(i,a^2i)=\ln|a^2|$ . On the other hand, from the definition of the character (2), we know that  $t_{\gamma}=t(a_{\gamma})=a$  by the following matrix computation

$$\begin{pmatrix} a \\ a^{-1} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} a \\ a^{-1} \end{pmatrix}^{-1} = a^2 \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}.$$

By the geometry of cusps (essentially the Margulis lemma), we know that there are only finitely many  $\Gamma_{\xi} \backslash \Gamma$ -orbits that lie inside any horoball  $B(\xi)$  at  $\xi$ , that is, the cardinality of  $\Gamma_{\xi} \backslash \Gamma O \cap B(\xi)$  (which makes sense since  $B(\xi)$  is  $\Gamma_{\xi}$ -invariant) is always finite. So are finitely many  $\Gamma_{\xi} \backslash \Gamma g$ -orbits  $(\Gamma_{\xi} \backslash \Gamma gO \cap B(\xi))$  according to triangle inequality. Hence, up to a finite error, we have

$$P_s(gO) \geq e^{-sC} \sum_{[\gamma] \in \Gamma_\xi \backslash \Gamma} e^{2s \ln t_\gamma} = e^{-sC} \sum_{[\gamma] \in \Gamma_\xi \backslash \Gamma} t_\gamma^{2s}.$$

Note that  $d(O, a_{\gamma}O) = -2\ln(t_{\gamma})$  if and only if  $t_{\gamma} \leq 1$ , which holds for all but finitely many  $[\gamma] \in \Gamma_{\xi} \setminus \Gamma$ , since all but finitely many  $[\gamma]gO$  lie outside the horoball of  $\mathcal{H}_{\xi}(1)$ . Hence, up to passing finite many terms, the above inequality holds.

Note that the convergence of the Poincaré series doesn't depend on the basepoint. Thus, if  $P_{s/2}(O)$  converges, then the Eisenstein series  $E(\phi_{\mathcal{E}}, g, s)$  converges.

In view of Proposition 3.3 and Corollary 3.4, we obtain

**Corollary 4.2.** Let  $\xi \in \mathbb{H}^{n+1}$  be any full rank parabolic fixed point.

- 1. If  $V = \mathfrak{g}$  is the adjoint representation, then the Eisenstein series  $E(\phi_{\xi}, g, 2n + 2)$  is always an absolutely convergent closed form.
- 2. If  $V = \mathbb{R}$  is the trivial representation, and in addition if either  $\Gamma$  is of convergence type or  $\delta(\Gamma) < n$ , then the Eisenstein series  $E(\phi_{\mathcal{E}}, g, 2n)$  is an absolutely convergent closed form.

*Proof.* The absolute convergence follows from Proposition 4.1. To see it is closed, we have

$$dE(\phi_{\xi},g,s) = d\left(\sum_{[\gamma] \in \Gamma_{\xi} \backslash \Gamma} \phi_{\xi,s}(\gamma g)\right) = \sum_{[\gamma] \in \Gamma_{\xi} \backslash \Gamma} d\phi_{\xi,s}(\gamma g) = 0.$$

Here,  $\phi_{\xi,s}$  is closed by Corollary 3.4 if  $V = \mathfrak{g}$ . If  $V = \mathbb{R}$ ,  $\phi_{\xi,s}$  is automatically closed since it is a top form. The absolute convergence of both series ensures the interchanging of the differential operator d with the infinite sum in the second equality. The last equality follows from Lemma 2.2.

### 5. Intertwining operators

In this section, we assume that all the cusps are full rank, and  $V = \mathbb{R}$  or  $\mathfrak{g}$ . Let  $E(\phi_{\xi'}, g, s)$  be the Eisenstein series corresponding to the parabolic fixed point  $\xi'$ . In order to see which cohomology class the Eisenstein series  $E(\phi_{\xi'}, g, s)$  restricts to in  $H^n(\Gamma_{\xi}, V) \cong H^n(\mathfrak{u}_{\xi}, V)$ , we need to look at the image of  $E(\phi_{\xi'}, g, s)$  under the map  $I_2 \circ r_1$  (see the following commutative diagram), which is equivalent to trace along  $r_2 \circ I_1$ .

$$\Omega^*(\mathbb{H}^{n+1}, V)^{\Gamma} \xrightarrow{I_1} \Omega^*(\mathbb{H}^{n+1}, V)^{N_{\xi}}$$

$$\downarrow^{r_1} \qquad \qquad \downarrow^{r_2}$$

$$\Omega^*(\mathcal{H}_{\xi}(1), V)^{\Gamma_{\xi}} \xrightarrow{I_2} \Omega^*(\mathcal{H}_{\xi}(1), V)^{N_{\xi}}.$$

To obtain the image under  $I_1$ , we define the intertwining operator from  $\xi'$  to  $\xi$  as

$$E^{\xi}(\phi_{\xi'},g,s) = \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} E(\phi_{\xi'},ug,s) du,$$

where du is the Haar measure normalized, such that  $\operatorname{vol}(\Gamma_{\xi} \setminus N_{\xi}) = 1$ . Then, by restricting the integral to the horosphere  $\mathcal{H}_{\xi}(1)$ , we obtain an element in  $\Omega^*(\mathcal{H}_{\xi}(1), V)^{N_{\xi}}$  which is exactly the image of  $E(\phi_{\xi'}, g, s)$  under  $r_2 \circ I_1$ . The main goal of this section is to compute explicitly  $E^{\xi}(\phi_{\xi'}, g, s)$ . We will follow the general approach in [HC68, Chapter II]. For the convenience, we introduce the following notation: for any  $g, h \in G$ , we write  ${}^gh$  to denote the conjugate of h by g, that is,  $[{}^g]h = ghg^{-1}$ . We will need the following two lemmas.

**Lemma 5.1.** For any  $g \in G$  and  $a \in A_{\xi}$ , one has

$$\phi_{\xi,s}(ag) = t_a^s \cdot \phi_{\xi,s}(g).$$

*Proof.* Let  $g = n_g a_g k_g$ . By definition,

$$\phi_{\xi,s}(ag) = \phi_{\xi,s}(an_g a_g k_g) = \phi_{\xi,s}(^a n_g a a_g k_g) = t_a^s t_g^s \phi_{\xi}(^a n_g k_g) = t_a^s t_g^s \phi_{\xi}(k_g).$$

where the last equality follows from the  $N_{\xi}$ -invariance of  $\phi_{\xi}$ . We also have

$$\phi_{\xi,s}(g) = \phi_{\xi,s}(n_g a_g k_g) = t_g^s \phi_{\xi}(n_g k_g) = t_g^s \phi_{\xi}(k_g).$$

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Therefore,  $\phi_{\xi,s}(ag) = t_a^s \phi_{\xi,s}(g)$ .

**Lemma 5.2.** (Geometric Bruhat decomposition) Let  $G = SO^+(n+1,1)$ , and fix a basepoint  $O \in \mathbb{H}^{n+1}$ . For any  $\xi, \xi' \in \partial_{\infty} \mathbb{H}^{n+1}$ , G decomposes as a disjoint union

$$G = \begin{cases} P_{\xi} \cup N_{\xi} w P_{\xi} & \text{if } \xi = \xi' \\ k P_{\xi} \cup N_{\xi'} k w P_{\xi} & \text{if } \xi \neq \xi' \end{cases},$$

where  $w \in K$  is any isometry that reflects the geodesic  $O\xi$ , and  $k \in K$  is any isometry that sends  $\xi$  to  $\xi'$ .

*Proof.* The case  $\xi = \xi'$  is just the classical Bruhat decomposition. For  $\xi \neq \xi'$ , and for any  $k \in K$  that sends  $\xi$  to  $\xi'$ , using the classical Bruhat decomposition we have,

$$G=kG=kP_{\xi}\cup kN_{\xi}wP_{\xi}=kP_{\xi}\cup N_{\xi'}kwP_{\xi},$$

where the last equality uses  $kN_{\mathcal{E}}k^{-1} = N_{\mathcal{E}'}$ .

Although we started with  $\phi_{\xi,s}$  whose value is in  $C^{\infty}(P_{\xi}, \operatorname{Hom}(\Lambda^n \mathfrak{n}_{\xi}, V))$ , after taking the Eisenstein series and intertwining operator, even though it is still  $N_{\xi}$ -invariant, the value could lie in a bigger space  $C^{\infty}(P_{\xi}, \operatorname{Hom}(\Lambda^n(\mathfrak{n}_{\xi} \oplus \mathfrak{a}_{\xi}), V))$ . For this reason, we introduce the extended space  $\overline{\Omega}^n(\mathcal{H}_{\xi}(1), V) \supset \Omega^n(\mathcal{H}_{\xi}(1), V)$  to be the restriction of V-valued forms on  $\mathbb{H}^{n+1}$  to the horosphere  $\mathcal{H}_{\xi}(1)$ , without projecting to its (co)tangent space. In other words, forms in  $\overline{\Omega}^n(\mathcal{H}_{\xi}(1), V)$  may look like  $(dy \land dx_1 \land \cdots \land dx_{n-1}) \otimes v$ . Now, we can state our theorem.

**Theorem 5.3.** Let  $E(\phi_{\xi}, g, s)$ ,  $E(\phi_{\xi'}, g, s)$  denote the Eisenstein series corresponding to the rank n parabolic fixed points  $\xi, \xi'$ , respectively. Then for any  $g \in P_{\xi}$ , we have

$$E^{\xi}(\phi_{\xi'}, g, s) = \epsilon \phi_{\xi', s}(g) + c_s(\phi_{\xi'})_{-s+2n}(g),$$

where

$$\epsilon = \begin{cases} 0 & \text{if } \Gamma_{\xi} \text{ and } \Gamma_{\xi'} \text{ are not } \Gamma\text{-conjugate} \\ 1 & \text{if } \xi = \xi' \end{cases},$$

and  $c_s: \Omega^n(\mathcal{H}_{\xi'}(1), V)^{N_{\xi'}} \to \overline{\Omega}^n(\mathcal{H}_{\xi}(1), V)^{N_{\xi}}$  is a linear operator.

Proof. By the definition,

$$E^{\xi}(\phi_{\xi'},g,s) = \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} E(\phi_{\xi'},ug,s) du = \sum_{\gamma \in \Gamma_{\xi'} \backslash \Gamma} \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} \phi_{\xi',s}(\gamma ug) du.$$

By assumption, the Eisenstein series is absolutely convergent, and this ensures the interchanging between the integral and the infinite summation. This shows up several times in the rest computations, but we emphasize that all the summations that appear are absolutely dominated by the Eisenstein series (which is absolutely convergent), thus interchanging the summation with the integral simply causes no problem by Fubini's theorem. We first prove the case when  $\xi = \xi'$ . By the Bruhat decomposition, we have

$$G = SO^+(n+1,1) = P_{\mathcal{E}} \cup N_{\mathcal{E}} w P_{\mathcal{E}}.$$

Hence, elements in  $\Gamma$  are either represented by elements in  $P_{\xi}$  or by the ones in  $N_{\xi}wP_{\xi}$ . We denote the set of former elements by  $\Gamma_1$  which is exactly  $\Gamma_{\xi}$  and the set of the latter one by  $\Gamma_w$ . Now we split the summation of  $E^{\xi}(\phi_{\xi}, g, s)$  into two kinds,

$$\begin{split} E^{\xi}(\phi_{\xi},g,s) &= \sum_{\gamma \in \Gamma_{\xi} \backslash \Gamma_{1}} \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} \phi_{\xi,s}(\gamma ug) du + \sum_{\gamma \in \Gamma_{\xi} \backslash \Gamma_{w}} \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} \phi_{\xi,s}(\gamma ug) du \\ &= \mathfrak{E}_{1} + \mathfrak{E}_{2}, \end{split}$$

where  $\Gamma_{\xi} \setminus \Gamma_{w}$  makes sense since  $\Gamma = \Gamma_{\xi} \cup \Gamma_{w}$ , and it denotes the set of all nontrivial right cosets of  $\Gamma_{\xi}$  in  $\Gamma$ . It is clear that

$$\mathfrak{E}_1 = \int_{u \in \Gamma_{\xi} \setminus N_{\xi}} \phi_{\xi,s}(ug) du = \phi_{\xi,s}(g),$$

where the last equality follows from the  $N_{\xi}$ -invariance of  $\phi_{\xi,s}$ . To simplify  $\mathfrak{E}_2$ , we will need the following lemma.

**Lemma 5.4.**  $\Gamma_{\xi}$  acts on  $\Gamma_{\xi} \backslash \Gamma_{w}$  from the right with trivial stabilizer, hence the quotient is the double coset  $\Gamma_{\xi} \backslash \Gamma_{w} / \Gamma_{\xi}$ .

*Proof.* Given any  $\gamma_w \in \Gamma_w$  and  $a \in \Gamma_{\xi}$ , such that  $\Gamma_{\xi} \gamma_w \cdot a = \Gamma_{\xi} \gamma_w$ , it suffices to show a = 1. From the identity, we obtain  $\gamma_w a \gamma_w^{-1} \in \Gamma_{\xi}$ , or that  $a \in \Gamma_{\gamma_w \xi}$ . Hence,

$$a\in \Gamma_{\xi}\cap \Gamma_{\gamma_w\,\xi}.$$

Since  $\gamma_w \notin \Gamma_{\xi}$ , we know  $\gamma_w \xi \neq \xi$ , hence  $\Gamma_{\xi} \cap \Gamma_{\gamma_w \xi} = 1$ . Therefore, a = 1 and the action has trivial stabilizer.

Now we can simplify  $\mathfrak{E}_2$  as

$$\begin{split} \sum_{\gamma \in \Gamma_{\xi} \backslash \Gamma_{w}} \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} \phi_{\xi,s}(\gamma u g) du &= \sum_{\gamma \in \Gamma_{\xi} \backslash \Gamma_{w} / \Gamma_{\xi}} \sum_{\gamma_{\xi} \in \Gamma_{\xi}} \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} \phi_{\xi,s}(\gamma \gamma_{\xi} u g) du \\ &= \sum_{\gamma \in \Gamma_{\xi} \backslash \Gamma_{w} / \Gamma_{\xi}} \int_{u \in N_{\xi}} \phi_{\xi,s}(\gamma u g) du, \end{split}$$

where the last equality uses the fact that  $N_{\xi}$  is abelian and du is the biinvariant Haar measure. We compute further  $\mathfrak{C}_2$ .

$$\begin{split} \mathfrak{E}_2 &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(\gamma u g) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(u_{\gamma} w p_{\gamma} u g) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(u_{\gamma} w p_{\gamma} u g) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w p_{\gamma} u n_{g} a_{g} m_{g}) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} a_{\gamma} n_{\gamma} u n_{g} a_{g} m_{g}) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} a_{\gamma} u' a_{g} m_{g}) du' \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w a_{\gamma} a_{g} m_{\gamma} u'' a_{g} m_{g}) du' \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w a_{\gamma} a_{g} m_{\gamma} u'' m_{g}) du' \cdot t_{g}^{2n} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w a_{\gamma} a_{g} m_{\gamma} u'' m_{g}) du'' \cdot t_{g}^{2n} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} u'' m_{g}) du'' \cdot t_{g}^{2n} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} u'' m_{g}) du'' \cdot t_{g}^{2n} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} u'' m_{g}) du'' \cdot t_{g}^{2n} \\ &= Ad_{\mathfrak{B}}^{*}(m_{\mathfrak{B}}^{-1}) \otimes \rho(m_{\mathfrak{B}}^{-1}) \binom{\sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} u'' m_{g}) du'' t_{\gamma}^{-s} t_{g}^{2n-s} \\ &= Ad_{\mathfrak{B}}^{*}(m_{\mathfrak{B}}^{-1}) \otimes \rho(m_{\mathfrak{B}}^{-1}) \binom{\sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} u'' m_{g}) du'' t_{\gamma}^{-s} t_{g}^{2n-s} \\ &= Ad_{\mathfrak{B}}^{*}(m_{\mathfrak{B}}^{-1}) \otimes \rho(m_{\mathfrak{B}}^{-1}) \binom{\sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} u'' m_{g}) du'' t_{\gamma}^{-s} t_{g}^{2n-s} \\ &= Ad_{\mathfrak{B}}^{*}(m_{\mathfrak{B}}^{-1}) \otimes \rho(m_{\mathfrak{B}}^{-1}) \binom{\sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E},s}(w m_{\gamma} u'' m_{g}) du'' t_{\gamma}^{-s} t_{g}^{2n-s} \\ &= Ad_{\mathfrak{B}}^{*}(m_{\mathfrak{B}}^{-1}) \otimes \rho(m_{\mathfrak{B}}^{-1}) \binom{\sum_{\gamma \in \Gamma_{\mathcal{E}} \backslash \Gamma_{w} / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}$$

where in the last line  $C_0(\phi_{\xi,s})$  is just a constant (independent of g) in  $\operatorname{Hom}(\Lambda^n(\mathfrak{a}_\xi\oplus\mathfrak{p}_\xi),V)$ . Hence, if we write the variable  $g=n_ga_gm_g$ , then we can regard  $\mathfrak{E}_2$  as the differential form constructed via the extension (with degree 2n-s) of the  $N_\xi$ -invariant form whose initial value is exactly  $C_0(\phi_{\xi,s})$ . We denote  $c_s:\Omega^n(\mathcal{H}_\xi(1),V)^{N_\xi}\to\overline{\Omega}^n(\mathcal{H}_\xi(1),V)^{N_\xi}$  the unique operator given by the initial value  $c_s(\phi_\xi)(1)=C_0(\phi_{\xi,s})$ . From the expression, it is clear that  $c_s$  is linear. Therefore, we obtain,

$$E^{\xi}(\phi_{\xi}, g, s) = \int_{u \in \Gamma_{\xi} \setminus N_{\xi}} E(\phi_{\xi}, ug, s) du = \phi_{\xi}(n_{g}m_{g}) t_{g}^{s} + c_{s}(\phi_{\xi})(n_{g}m_{g}) t_{g}^{2n-s}$$
$$= \phi_{\xi, s}(g) + c_{s}(\phi_{\xi})(2n-s)(g).$$

Next, we assume  $\xi$  and  $\xi'$  are not  $\Gamma$ -conjugate. Then, similar to Lemma 5.4, we have,

**Lemma 5.5.** Suppose  $\xi$  and  $\xi'$  are not  $\Gamma$ -conjugate, then  $\Gamma_{\xi}$  acts on  $\Gamma_{\xi'} \setminus \Gamma$  from the right with trivial stabilizer, hence the quotient is the double set  $\Gamma_{\xi'} \setminus \Gamma/\Gamma_{\xi}$ .

*Proof.* For any  $\gamma \in \Gamma$  and  $a \in \Gamma_{\xi}$ , such that  $\Gamma_{\xi'}\gamma \cdot a = \Gamma_{\xi'}\gamma$ , we have  $\gamma a \gamma^{-1} \in \Gamma_{\xi'}$ , or that  $a \in \Gamma_{\gamma \xi'}$ . Hence,

$$a \in \Gamma_{\xi} \cap \Gamma_{\gamma \xi'}$$
.

Since  $\xi$  and  $\xi'$  are not  $\Gamma$ -conjugate, we know  $\gamma \xi' \neq \xi$ , hence  $\Gamma_{\xi} \cap \Gamma_{\gamma \xi'} = 1$ . Therefore, a = 1 and the action has trivial stabilizer.

Using the lemma, we can simplify

$$\begin{split} E^{\xi}(\phi_{\xi'},g,s) &= \sum_{\gamma \in \Gamma_{\xi'} \backslash \Gamma} \int_{u \in \Gamma_{\xi} \backslash N_{\xi}} \phi_{\xi',s}(\gamma ug) du \\ &= \sum_{\gamma \in \Gamma_{\xi'} \backslash \Gamma / \Gamma_{\xi}} \int_{u \in N_{\xi}} \phi_{\xi',s}(\gamma ug) du \\ &= \mathfrak{G}. \end{split}$$

Also by the Bruhat decomposition (Lemma 5.2), for any  $\gamma \in \Gamma$ , either  $\gamma \in kP_{\xi}$  or  $\gamma \in N_{\xi'}kwP_{\xi}$ . If  $\gamma \in kP_{\xi}$ , then  $k^{-1}\gamma \in P_{\xi}$ , so  $k^{-1}\gamma\xi = \xi$ , or that  $\gamma\xi = k\xi = \xi'$ , contradiction to that  $\xi$  and  $\xi'$  are not  $\Gamma$ -conjugate. Thus,  $\gamma \in N_{\xi'}kwP_{\xi}$ . We now further compute  $\mathfrak{E}$ .

$$\begin{split} \mathfrak{E} &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (\gamma u g) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (n'_{\gamma} k w p_{\gamma} u g) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w p_{\gamma} u g) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w p_{\gamma} u g) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w m_{\gamma} a_{\gamma} n_{\gamma} u n_{g} a_{g} m_{g}) du \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w a_{\gamma} a_{g} m_{\gamma} u'' a_{g} m_{g}) du' \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w a_{\gamma} a_{g} m_{\gamma} u'' m_{g}) du' \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u'' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w a_{\gamma} a_{g} m_{\gamma} u'' m_{g}) du'' \cdot t_{g}^{2n} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u''' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w a_{\gamma} a_{g} m_{\gamma} u'' m_{g}) du'' \cdot t_{g}^{2n} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u''' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w m_{\gamma} u'' m_{g}) du'' \cdot t_{g}^{2n} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u''' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w m_{\gamma} u'' m_{g}) du'' t_{\gamma}^{-s} t_{g}^{2n-s} \\ &= \sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u''' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w m_{\gamma} u'' m_{g}) du'' t_{\gamma}^{-s} t_{g}^{2n-s} \end{aligned} \qquad (2)$$

$$= \operatorname{Ad}^{*}(m_{g}^{-1}) \otimes \rho(m_{g}^{-1}) \left(\sum_{\gamma \in \Gamma_{\mathcal{E}'} \backslash \Gamma / \Gamma_{\mathcal{E}}} \int_{u''' \in N_{\mathcal{E}}} \phi_{\mathcal{E}',s} (k w m_{\gamma} u'' m_{g}) du'' t_{\gamma}^{-s} t_{g}^{2n-s} \right)$$

$$= \operatorname{Ad}^{*}(m_{g}^{-1}) \otimes \rho(m_{g}^{-1}) (C_{0}(\phi_{\mathcal{E}',s})) \cdot t_{g}^{2n-s},$$

where in (1), we apply the Langlands decompositions  $p_{\gamma} = m_{\gamma} a_{\gamma} n_{\gamma} \in M_{\xi} A_{\xi} N_{\xi}, g = n_{g} a_{g} m_{g} \in$  $N_{\xi}A_{\xi}M_{\xi}$ . In (2), we use  $ka_{\gamma}^{-1}a_{g}^{-1}k^{-1} \in A_{\xi'}$  and the fact that k is an isometry which send character  $t_{\xi}$ to  $t_{\xi'}$ . Then we apply Lemma 5.1. In (3), we use the property of  $\phi_{\xi,s}$  and the fact that  $m_g \in K$ . Hence, if we denote  $c_s: \Omega^n(\mathcal{H}_{\mathcal{E}'}(1), V)^{N_{\mathcal{E}'}} \to \overline{\Omega}^n(\mathcal{H}_{\mathcal{E}}(1), V)^{N_{\mathcal{E}}}$  the unique operator given by the initial value  $c_s(\phi_{\mathcal{E}'})(1) = C_0(\phi_{\mathcal{E}',s})$ , then  $E^{\xi}(\phi_{\mathcal{E}'},g,s) = c_s(\phi_{\mathcal{E}'})(n_g m_g) t_g^{-s+2n}$ .

### 6. Cohomology classes associated to cusps

In this section, we study the cohomology class of the restriction of the Eisenstein series  $E(\phi_{\mathcal{E}'}, g, s)$ to the cusp associated to the parabolic fixed point  $\xi$ . Throughout the section, we assume s = 2n + 2 if  $V = \mathfrak{g}$  and s = 2n if  $V = \mathbb{R}$ . We also assume the Eisenstein series converges at s. By Corollary 4.2, we have  $dE(\phi_{\xi'}, g, s) = 0$ . Then

$$dE^{\xi}(\phi_{\xi'},g,s) = d\int_{u \in \Gamma_{\xi} \setminus N_{\xi}} E(\phi_{\xi'},ug,s) du = \int_{u \in \Gamma_{\xi} \setminus N_{\xi}} dE(\phi_{\xi'},ug,s) du = 0.$$

As before, the absolute convergence of the Eisenstein series ensures the interchanging of differential and integral, and the last equality follows from Lemma 2.2. We will use the computation of  $E^{\xi}(\phi_{\xi}, g, s)$  in Section 5 to see which cohomology class  $[E^{\xi}(\phi_{\xi'},g,s)]|_{\mathcal{H}_{\xi}(1)}$  represents in  $H^{n}(\Gamma_{\xi},V)$ .

**Proposition 6.1.** Let  $\xi$  and  $\xi'$  be full rank parabolic fixed points. Then

$$[E^{\xi}(\phi_{\xi'},g,s)]|_{\mathcal{H}_{\mathcal{E}}(1)} = [\phi_{\xi}] \in H^{n}(\Gamma_{\xi},V)$$

if  $\xi = \xi'$ . If  $\xi$  and  $\xi'$  are not  $\Gamma$ -conjugate, then

$$[E^\xi(\phi_{\xi'},g,s)]|_{\mathcal{H}_{\mathcal{E}}(1)}=0\in H^n(\Gamma_\xi,V).$$

*Proof.* By Theorem 5.3,

$$E^{\xi}(\phi_{\mathcal{E}'}, g, s) = \epsilon \phi_{\mathcal{E}', s}(g) + c_s(\phi_{\mathcal{E}'})_{-s+2n}(g).$$

Since  $dE^{\xi}(\phi_{\xi'}, g, s) = 0$ , and  $d\phi_{\xi', s}(g) = 0$ , we obtain

$$d(c_s(\phi_{\mathcal{E}'})_{-s+2n})(g) = 0.$$

Hence, it suffices to prove that

$$[c_s(\phi_{\xi'})_{-s+2n}(g)]|_{\mathcal{H}_{\mathcal{E}}(1)}=0\in H^n(\Gamma_{\xi},V).$$

We first consider the case that  $V = \mathfrak{g}$ . Recall that we can regard the extended form  $c_s(\phi_{\mathcal{E}'}) \in$  $\overline{\Omega}(\mathcal{H}_{\mathcal{E}}(1),V)^{N_{\mathcal{E}}}$  as a function in  $C^{\infty}(P_{\mathcal{E}}(1),\operatorname{Hom}(\Lambda^{n}(\mathfrak{a}_{\mathcal{E}}\oplus\mathfrak{n}_{\mathcal{E}}),V))$ . Since  $V=V_{-2}\oplus V_{0}\oplus V_{2}$  and  $c_s(\phi_{\mathcal{E}'})$  is  $N_{\mathcal{E}}$ -invariant, we can decompose  $c_s(\phi_{\mathcal{E}'})(1)$  as

$$c_s(\phi_{\xi'})(1) = A_2 + A_0 + A_{-2} + \frac{dt}{t} \wedge A_2' + \frac{dt}{t} \wedge A_0' + \frac{dt}{t} \wedge A_{-2}',$$

where  $A_i \in \Lambda^n \mathfrak{n}_{\xi}^* \otimes V_i$ , and  $A_i' \in \Lambda^{n-1} \mathfrak{n}_{\xi}^* \otimes V_i$  for  $i \in \{-2,0,2\}$ . For each i, we construct an  $N_{\xi}$ -invariant V-valued form  $c_i(\phi_{\xi'})$  (or  $c_i'(\phi_{\xi'})$ ) on the horosphere  $\mathcal{H}_{\xi}(1)$  whose initial value is  $A_i$  (or  $A_i'$ ), that is,  $c_i(\phi_{\mathcal{E}'}) \in C^{\infty}(P_{\mathcal{E}}(1), \text{Hom}(\Lambda^n \mathfrak{n}_{\mathcal{E}}, V))$  satisfies:

1.  $c_i(\phi_{\xi'})(n) = c_i(\phi_{\xi'})(1) = A_i$ , 2.  $c_i(\phi_{\xi'})(pm) = \operatorname{Ad}^*_{\mathfrak{a}_{\xi} \oplus \mathfrak{n}_{\xi}}(m^{-1}) \otimes \rho(m^{-1})(c_i(\phi_{\xi'})(p))$  for all  $m \in K_{\xi}$  and  $p \in P_{\xi}(1)$ .

And  $c_i'(\phi_{\xi'})$  is defined similarly. Following the construction in Section 3, we can define degree 2n-s extensions of  $c_i(\phi_{\xi'})$  and  $c_i'(\phi_{\xi'})$  in  $C^{\infty}(P_{\xi}, \operatorname{Hom}(\Lambda^n \mathfrak{p}_{\xi}, V))$ , denoted by  $(c_i(\phi_{\xi'}))_{2n-s}$  and  $(c_i'(\phi_{\xi'}))_{2n-s}$ , respectively. It follows that

$$c_s(\phi_{\xi'})_{2n-s} = \sum_i (c_i(\phi_{\xi'}))_{2n-s} + \sum_i \frac{dt}{t} \wedge (c_i'(\phi_{\xi'}))_{2n-s}.$$

By Proposition 3.3,

$$0 = d(c_s(\phi_{\xi'})_{2n-s}) = (2n - s + 2 - 2n)\frac{dt}{t} \wedge c_2(\phi_{\xi'})_{2n-s} + (2n - s - 2n)\frac{dt}{t} \wedge c_0(\phi_{\xi'})_{2n-s} + (2n - s - 2n)\frac{dt}{t} \wedge c_0(\phi_{\xi'})_{2n-s} + \sum_i \frac{dt}{t} \wedge (dc_i'(\phi_{\xi'})_{2n-s}),$$

where  $i \in \{-2, 0, 2\}$ . Observe that

$$\frac{dt}{t} \wedge c_i(\phi_{\xi'})_{2n-s}(1) \in \frac{dt}{t} \wedge \operatorname{Hom}(\Lambda^n \mathfrak{n}_{\xi}, V_i),$$

and by the proof of Lemma 2.3,

$$\frac{dt}{t} \wedge d(c'_i(\phi_{\xi'})_{2n-s})(1) \in \frac{dt}{t} \wedge \operatorname{Hom}(\Lambda^n \mathfrak{n}_{\xi}, V_{i+2})$$

for  $i \in \{-2, 0, 2\}$ . Therefore, by comparing their components in  $V_2, V_0$ , and  $V_{-2}$ , respectively, we obtain

$$\begin{split} \frac{dt}{t} \wedge (2n - s + 2 - 2n)c_2(\phi_{\xi'})_{2n - s}(1) + \frac{dt}{t} \wedge d(c_0'(\phi_{\xi'}))_{2n - s}(1) &= 0, \\ \frac{dt}{t} \wedge (2n - s - 2n)c_0(\phi_{\xi'})_{2n - s}(1) + \frac{dt}{t} \wedge d(c_{-2}'(\phi_{\xi'}))_{2n - s}(1) &= 0, \\ \frac{dt}{t} \wedge (2n - s - 2 - 2n)c_{-2}(\phi_{\xi'})_{2n - s}(1) &= 0. \end{split}$$

Note that s = 2 + 2n, the coefficients in the above equalities are nonzero. Thus, we have

$$c_2(\phi_{\xi'})(1) = \frac{d(c'_0(\phi_{\xi'}))(1)}{2 - s},$$

$$c_0(\phi_{\xi'})(1) = \frac{d(c'_{-2}(\phi_{\xi'}))(1)}{-s},$$

$$c_{-2}(\phi_{\xi'})(1) = 0.$$

By Lemma 2.2,

$$c_2(\phi_{\xi'}) = \frac{d(c_0'(\phi_{\xi'}))}{2 - s},$$

$$c_0(\phi_{\xi'}) = \frac{d(c_{-2}'(\phi_{\xi'}))}{-s},$$

$$c_{-2}(\phi_{\xi'}) = 0.$$

We see that  $c_0(\phi_{\xi'}) + c_2(\phi_{\xi'})$  is a coboundary, which equals the restriction (in the strong sense, i.e., also project the (co)tangent space) of  $c_s(\phi_{\xi'})_{2n-s}$  to the horosphere  $\mathcal{H}_{\xi}(1)$ . Hence

$$[c_s(\phi_{\xi'})_{-s+2n}]|_{\mathcal{H}_{\mathcal{E}}(1)}=0\in H^n(\Gamma_{\xi},\mathrm{Ad}).$$

The argument is similar for the case that  $V = \mathbb{R}$ . In this case, we write

$$c_s(\phi_{\xi'}) = c_1(\phi_{\xi'}) + \frac{dt}{t} \wedge c_2(\phi_{\xi'}),$$

where  $c_1(\phi_{\xi'}) \in C^{\infty}(P_{\xi}(1), \operatorname{Hom}(\Lambda^n \mathfrak{n}_{\xi}, \mathbb{R}))$  and  $c_2(\phi_{\xi'}) \in C^{\infty}(P_{\xi}(1), \operatorname{Hom}(\Lambda^{n-1} \mathfrak{n}_{\xi}, \mathbb{R}))$ . By Proposition 3.3,

$$0 = d(c_s(\phi_{\xi'})_{2n-s}) = (2n - s - 2n)\frac{dt}{t} \wedge c_1(\phi_{\xi'})_{2n-s},$$

which implies that  $c_1(\phi_{\mathcal{E}'}) = 0$ . Hence

$$[c_s(\phi_{\mathcal{E}'})_{-s+2n}]|_{\mathcal{H}_{\mathcal{E}}(1)} = 0 \in H^n(\Gamma_{\mathcal{E}}, \mathbb{R}).$$

Now we are ready to prove Theorems 1.1, 1.2, and Corollary 1.3.

# Proof of Theorems 1.1, 1.2

Since every full rank parabolic subgroup corresponds to a parabolic fixed point  $\xi \in \mathbb{H}^{n+1}$ , it follows immediately from Proposition 6.1. Also, the harmonicity of the Eisenstein series  $E(\phi_{\xi})$  in the case of trivial coefficient is clear, since it is the (absolutely convergent) sum of harmonic forms (see Remark 3.2).

#### **Proof of Corollary 1.3**

Let  $C_1, ..., C_N$  be N toric cusps. For each cusp  $C_i$ , choose a corresponding parabolic subgroup  $\Gamma_i \cong \mathbb{Z}^n$  (unique up to conjugate). By Lemma 2.3 and Theorem 2.4, the dimension of the cohomology group  $H^n(\Gamma_i, \operatorname{Ad})$  is n. Thus, in view of Theorem 1.2, they in total correspond to nN linearly independent Eisenstein cohomology classes in  $H^n(\Gamma, \operatorname{Ad})$ . So the corollary follows immediately.

### 7. Further discussions

Our work seems to suggest that the classical method of Eisenstein series should also fit in a broader context for certain problems, and this paper only focuses on a specific aspect of that, namely, the cusp counting problem for hyperbolic manifolds. Before discussing possible directions of extensions of our results, we first give some examples where our theorems can be applied. Nevertheless, we would like to point out that the dimension of group cohomology is often very hard to compute, so it is unclear how sharp the inequality in Corollary 1.3 is in general.

### Nonuniform lattices

As is mentioned in the Introduction, nonuniform lattices do not satisfy the assumption and conclusion of Theorem 1.1. In this case,  $\delta(\Gamma) = n$ ,  $\Gamma$  is of divergent type, and cohomology classes (with trivial coefficient) arising from cusps *cannot* be linearly independent. However, our Theorem 1.2 and Corollary 1.3 both apply, thus giving an explicit upper bound on the number of cusps. One type of explicit construction of lattices comes from arithmetic (e.g., take  $\Gamma = SO(n+1,1;\mathbb{Z}) < SO(n+1,1)$  up to finite index to kill the torsion). The number of cusps thus is closely related to the arithmetic feature of the construction (e.g., the ideal class group of the corresponding number field). There are also

nonarithmetic constructions of lattices due to the work of Gromov–Piatetski-Shapiro [GPS88], and our results are potentially more useful in these examples.

### Geometrically finite Kleinian groups

The simplest such example arises from the Schottky-type construction. Take a rank n parabolic subgroup  $\Gamma_1 \cong \mathbb{Z}^n < \operatorname{Isom}^+(\mathbb{H}^{n+1})$  and an elementary subgroup generated by a single hyperbolic translation  $\Gamma_2 \cong \mathbb{Z} < \operatorname{Isom}^+(\mathbb{H}^{n+1})$ , then up to a choice of conjugates of  $\Gamma_1, \Gamma_2$ , the group  $\Gamma := < \Gamma_1, \Gamma_2 >$  is isomorphic to the free amalgamation  $\Gamma_1 * \Gamma_2$  by Maskit's Klein combination theorem [Mas64]. The resulting Kleinian group is then geometrically finite and have exactly one cusp. On the other hand, the critical exponent satisfies  $\delta(\Gamma) < n$ , so in this case, both our theorems hold. It is clear that the n-th Betti number is 1, but  $H^n(\Gamma, \operatorname{Ad})$  will depend on the representation of  $\Gamma$  in  $\operatorname{Isom}^+(\mathbb{H}^{n+1})$ . Thus, our corollary gives a uniform lower bound on  $\dim H^n(\Gamma, \operatorname{Ad})$ , that is,  $\dim H^n(\Gamma, \operatorname{Ad}) \ge n$ , and in fact, it holds for all geometrically finite Kleinian groups  $\Gamma \cong \mathbb{Z}^n * \mathbb{Z}$ .

# Geometrically infinite Kleinian groups

As we mentioned in Remark 3.6, the Kleinian groups  $\Gamma < \mathrm{Isom}^+(\mathbb{H}^{n+1})$  ( $4 \le n \le 7$ ) constructed in [IMM22] has infinitely many full rank toric cusps. It is a finitely generated (also finitely presented in the case n=6,7) normal subgroup of a lattice, which is constructed from an algebraic fibration of the lattice over  $\mathbb{Z}$ . It follows that  $\delta(\Gamma) = n$  [Rob05]. As a consequence of Corollary 1.3, the cohomology group  $H^n(\Gamma, \mathrm{Ad})$  has infinite dimension. On the other hand, Kapovich's example [Kap95] of Kleinian group  $\Gamma < \mathrm{Isom}^+(\mathbb{H}^4)$  has infinitely many rank 1 cusps, and the critical exponent satisfies  $\delta(\Gamma) \in [2,3]$ . Our results do not cover this example since the cusp is not full rank, but it would be very interesting to see a generalization of the results to lower degrees.

By carefully examining all these examples, we believe that there might be interesting general relations between the value of critical exponent and the number of cusps for a Kleinian group.

**Conjecture 7.1.** Given a finitely generated Kleinian group  $\Gamma < \mathrm{Isom}^+(\mathbb{H}^{n+1})$ , if  $\delta(\Gamma) < k$  for a positive integer  $k \le n$ , then  $\Gamma$  has only finitely many rank k cusps.

We now proceed to discuss possible extensions of our results.

# Other coefficient modules

To possibly extend our main results, one can try to consider other  $\Gamma$ -coefficient modules V. We believe this should be straightforward. First, to construct a cohomology classes in  $H^n(\Gamma_{\xi}, V)$ , one takes the  $V_{-\lambda}$  sections where  $\lambda$  denotes the highest weight of the representation. Second, the absolute convergence of Eisenstein series follows a similar argument provided  $\lambda > 0$ , which is automatic if V is not the trivial representation. Finally, the computation of the intertwining operators in Section 5 works verbatim, and the argument in Section 6 also works very similarly. Our Theorem 1.2 is thus expected to hold for any nontrivial coefficient module. However, one cautionary point is that our Lemma 2.3 may not hold anymore.

# Other degrees

One can also consider lower degree Eisenstein cohomology classes in  $H^k(\Gamma, V)$ . But the construction of cohomology classes in  $H^k(\Gamma_\xi, V)$  is more delicate. For example, the Lie algebra cohomology  $H^k(\mathfrak{u}, V)$  might be zero. When it is nonzero, then one can construct the Eisenstein series, and it is absolutely convergent if the weights of the constructed cohomology classes are small enough (negative large) in terms of k and the critical exponent of  $\Gamma$ . The rest computation of the intertwining operators follows similarly. It would be very interesting to see if one can construct an absolutely convergent Eisenstein

cohomology class in  $H^1(\Gamma, V)$ . This would lead to a cusp finiteness theorem, since finitely generated groups always satisfy dim  $H^1(\Gamma, V) < \infty$ .

# Rank k cusps

When  $\Gamma_{\xi}$  is a rank k parabolic group, there are natural classes in  $H^k(\Gamma_{\xi}, V)$ . Under the same construction, the Eisenstein series is comparable to the Poincaré series if the parabolic fixed point is *bounded*. Thus, the absolute convergence again depends on the weights of the class constructed in  $H^k(\Gamma_{\xi}, V)$ , the value k, and the critical exponent of  $\Gamma$ . However, the most difficult part is the computation of the intertwining operators. Our method fails when  $\phi_{\xi,s}$  is not  $N_{\xi}$ -invariant. One needs to further analyze the behavior of the function  $\phi_{\xi,s}$  in the orthogonal directions to the subspace where it is invariant.

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