# Emission Lines in the Far-UV and Extreme-UV Region

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Abstract. Recent observations with HST and HUT have been made of the spectral region shortward of  $Ly\alpha$  in AGN. Emission lines in this region, such as  $O \vee 1 \lambda 1035$  and  $N e \vee 111 \lambda 774$ , are mainly produced by collisional excitation. Their strengths and the lack of other significant emission lines of higher excitation energy indicate a temperature of  $\sim 6 \times 10^4$  K for the highly ionized gas in line-emitting clouds. Model calculations suggest that the strength of  $O \vee 1$  and  $N e \vee 111$  emission is correlated to the intensity of the high-energy photons above  $100 \, eV$ . Therefore, these lines are useful diagnostics of the continuum shape in the soft X-ray range. The luminosity dependence of emission-line equivalent widths is more prominent for lines of higher ionization level, suggesting that the UV bump in AGN spectra may be weaker at higher luminosities.

## 1. Introduction

Our study of emission lines starts in the optical region, and gradually extends to the UV region and even beyond. A significant number of resonance atomic transitions fall in the UV region. Observations with *IUE* over the past 18 years have provided a wealth of AGN spectra (Kinney, Rivolo, & Koratkar 1990; Lanzetta, Turnshek, & Sandoval 1992) that enable us to study Ly $\alpha$ , C IV, and many other UV lines. In recent years, frequent *IUE* observations have been coordinated with ground-based and X-ray observations to monitor time variations of the continuum and emission lines, Ly $\alpha$  and C IV in particular. Such studies have established more accurate estimates of the size of the broad-line region in Seyfert galaxies and revealed the stratification in the region (Peterson 1993).

The properties of the lines below 1216 Å have not been studied as well as Ly $\alpha$  and C IV. The O VI  $\lambda$ 1035 emission line was measured in the *IUE* spectra of medium-redshift quasars (Kinney et al. 1985). This line is a doublet, blended with Ly $\beta$ . Recently, Laor et al. (1994) used high-resolution *HST* data to study the emission lines in four bright quasars. They were able to deblend Ly $\beta$  from O VI and found Ly $\beta$ /Ly $\alpha = 0.03 - 0.12$ . Weak features such as Ly $\gamma$ , C III  $\lambda$ 977, S VI  $\lambda\lambda$ 933, 945, and N III  $\lambda$ 991 have also been identified. Redshifted O VI emission is a common feature in the optical spectra of many high-z quasars, but its profile suffers from severe Ly $\alpha$  forest absorption. An observational gap existed in the sub-Ly $\alpha$  region of many well-studied low-redshift Seyfert galaxies.

#### 2. HUT Observations

The Hopkins Ultraviolet Telescope (HUT, Davidsen 1993) is a 0.9-m telescope capable of spectroscopy between 850 and 1850 Å, with a resolution of  $\sim 2.5$  Å. It was flown during the Astro-1 mission between 1990 December 2–11. In March 1995 an improved version of the HUT made a second, 16-day flight on the Astro-2 mission aboard the Space Shuttle Endeavour (Kruk et al. 1995). The HUT data base on more than 300 celestial targets includes far-UV spectra of about twenty Seyfert galaxies and quasars.



Figure 1. HUT spectrum of NGC 1068. Significant narrow emission lines are present. Airglow lines are marked with Earth symbols. Adopted from Kriss et al. (1992)

A good example of HUT data is the spectrum of the Seyfert 2 galaxy NGC 1068 (Kriss et al. 1992) shown in Fig. 1. The spectrum reveals strong lines of CIII  $\lambda$ 977, NIII  $\lambda$ 991, Ly $\beta$ , OVI  $\lambda\lambda$ 1032,1038, and HeII  $\lambda$ 1085. The ratios of CIII]  $\lambda$ 1909 to  $\lambda$ 977 and NIII]  $\lambda$ 1750 to  $\lambda$ 991 imply temperatures exceeding 24,000 K, much higher than those typically encountered in photoionized gas with abundant C<sup>+2</sup> and N<sup>+2</sup> populations.

Figure 2 shows the HUT spectra of Mrk 335 and Fairall 9. The two Seyfert 1 galaxies have different soft X-ray properties. Mrk 335 has a very strong soft X-ray excess (Walter & Fink 1993). Its O VI/C IV ratio of 0.58 is one of the highest among our candidates. Fairall 9 has a weak soft X-ray excess. Its O VI/C IV ratio is only 0.36. There appears to be a correlation between the O VI strength and the soft X-ray intensity. Zheng, Kriss, & Davidsen (1995) found that the objects with stronger X-ray intensity (at 2 keV) exhibit stronger O VI emission. Recent Astro-2 observations, including that of Fairall 9, further suggest that O VI strength is more closely related to the soft X-ray intensity below 1 keV.



Figure 2. *HUT* spectra of Mrk 335 and Fairall 9. Airglow features are marked with Earth symbols.



Figure 3. Quasar extreme-ultraviolet spectrum. The data above 350 Å are from summed HST FOS spectra of high-z quasars. The dotted curve displays the HUT spectrum of the quasar HS1700+64 below 350 Å (Davidsen et al. 1996). A normalization is applied to connect the two curves. The sharp depression in flux below ~ 300 Å is due to intergalactic He II Ly $\alpha$  absorption.

#### 3. HST Observations

Hundreds of HST Faint Object Spectrograph (FOS) spectra reveal the sub-Ly $\alpha$  region of quasars of z > 0.2. Figure 3 displays the spectral region below 1216 Å in the rest frame. The curve above 350 Å is a composite quasar spectrum (Zheng et al. 1996) based on the FOS spectra of quasars of z > 0.33. There are no significant emission lines except for Ne VIII  $\lambda$ 774 (Hamann, Zuo, & Tytler 1995) below 1000 Å. This line has an exceptionally high ionization level and requires very strong soft X-ray emission to power it. The features near 690 Å and 800 Å are not certain due to the limited S/N level. They may be due to the Bowen fluorescence mechanism (see below).

At the very short wavelength end, the HUT spectrum of the quasar HS 1700+64 (Davidsen, Kriss, & Zheng 1996) is normalized and displayed with a dotted curve. There is a prominent HeII Ly $\alpha$  emission line at 304 Å, and its intensity relative to Ly $\alpha$  is consistent with photoionization models. This line is produced by recombination and powered by high-energy photons of E > 54 eV. Note the blue wing of this line is suppressed by the absorption by the intergalactic medium. Abundant HeII Ly $\alpha$  photons in the line-forming region lead to the Bowen fluorescence effect (Eastman & MacAlpine 1985; Netzer, Elitzur, & Ferland 1985). The O<sup>++</sup> and N<sup>++</sup> ions have transitions whose wavelengths coincide very closely with HeII Ly $\alpha$   $\lambda$ 304. Theoretical estimates suggest that the O III/HeII Ly $\alpha$  intensity ratio is of order 0.1 and is only weakly dependent on the ionizing flux. The 690 Å feature is not present in every spectrum that covers the wavelengths between 650 Å and 700 Å, so its identification is only suggestive.

### 4. Discussion

#### 4.1. Cloud Temperature

A typical photoionization model (Kwan & Krolik 1981) suggests that the highionization lines are formed in the ionized region where the temperature is of order  $3 \times 10^4$  K. Several UV metal lines are the major coolant for the region, keeping the temperature in a narrow range of thermal equilibrium, and therefore provide useful temperature estimates.

As shown in Table 1, these collisionally excited lines require a temperature  $T \gtrsim 4 \times 10^4$  K for their presence. The collisional excitation rate is proportional to  $\exp(-E_{exc}/kT)$  where  $E_{exc}$  is the excitation energy (Osterbrock 1989). Only when the cloud temperature is comparable to the excitation temperature, does the collisionally excited line becomes significant. The EUV region contains a large number of resonance metal lines, and many of them have been identified as intervening absorption lines (Vogel & Reimers 1995). The lack of emission lines below 750 Å suggests that the temperature of the highly ionized gas in the broad-line region is not likely to exceed  $6 \times 10^4$  K.

#### 4.2. Effect of the Ionizing Continuum

The intensity of each emission line depends on the number of photons above the line's ionization threshold. While the intensities of high-ionization lines are related to the continuum shape (Krolik & Kallman 1988), they are also sensitive

Line (Å)	Ionization Energy (eV)	Characteristic Temperatureª(K)
C iv λ1549	8.0	$4.0 \times 10^{4}$
$O vi \lambda 1035$	12.0	$6.0 imes10^4$
N 111 $\lambda 991$	12.5	$6.3 imes10^4$
$ m C{}^{III}\lambda 977$	12.7	$6.4 imes10^4$
${\rm Neviii}\lambda774$	16.0	$8.1 imes10^4$

Table 1. Major FUV and EUV Emission Lines.

<sup>a</sup>The value at which  $\exp(-E_{exc}/kT) = 0.1$  and the line strength starts to become significant.



Figure 4. Theoretical line ratios produced with a power law plus different EUV bumps. Solid curve: With a bump centered at 100 eV; dotted curve: 50 eV, and dashed curve: 25 eV.



Figure 5. EW of Ly $\alpha$ , C IV, and O VI vs. specific luminosity of UV continuum L $_{\nu}$ , and their best fit. Squares: *HUT* data; circles: *IUE* data; triangles: *HST* data; pentagons: optical data. Solid symbols are for radio-quiet objects.

to the level of the ionizing flux. Namely, a stronger flux may mimic the effect of a harder ionizing continuum.

Take the example of O v1. It is a collisionally excited line and highly sensitive to the temperature. Would an ionizing source with abundant photons below 113 eV (the ionization threshold of O v1) produce excessive heat in the line-forming region, thus boosting the O v1 emission? A series of photoionization calculations are carried out using the code CLOUDY (Ferland 1993). The basic ionizing continuum is a power law of  $F_{\nu} \propto \nu^{-2}$  between 1 eV and 200 eV and  $F_{\nu} \propto \nu^{-0.7}$  between 200 eV and 20 keV. Three models assume an additional EUV bump centered at 25, 50 and 100 eV, respectively. The bump is assumed to be a Gaussian, with its width as 20% of the peak value. It is normalized so that its total energy is equal to that of the underlying power-law continuum. For each continuum model, several calculations are carried out with different ionization parameters for a fixed nucleon density of  $10^{10}$  cm<sup>-3</sup>. As shown in Fig. 4, only the model with 100 eV bump can produce both the O v1/Ly $\alpha$  and O v1/C IV ratios that are close to the observed values. The other two models produces over predict the C IV emission because of the bump at lower energies.

This example demonstrates that the number of photons with energy above  $113 \,\mathrm{eV}$  dictates the temperature and volume of the region in which O VI emission

is formed, and therefore it is possible to use multi-line ratios to diagnose the continuum shape in the EUV and soft X-ray band.

### 4.3. Luminosity Dependence

The intensities of high-ionization lines are known to exhibit an inverse luminosity dependence (Baldwin 1977; Kinney et al. 1990). Zheng, Fang, & Binette (1992) found that such a relation is more significant for lines of higher ionization level. Zheng et al. (1995) found that, over five decades of luminosity, the O VI equivalent width shows a very strong anti-correlation with luminosity (Fig. 5). Expressing such a correlation with  $\log(EW) = A \log L_{UV} + constant$ , we find that, for Ly $\alpha$ , C IV, and O VI, the coefficient A is -0.076, -0.14, and -0.25 respectively. Between Seyfert 1 galaxies and high-z quasars the equivalent widths of Ly $\alpha$ , C IV, and O VI decrease approximately by factors of 3, 6, and 25, respectively. The significant change in high-ionization line intensity with luminosity is consistent with the assumption that the soft X-ray continuum between 0.1 and 1 keV in high-luminosity quasars is weaker than that in low-luminosity AGNs.

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