

# Fermionic and Bosonic Representations of the Extended Affine Lie Algebra $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$

*Dedicated to Professor Robert Moody on the occasion of his 60th birthday*

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*Abstract.* We construct a class of fermions (or bosons) by using a Clifford (or Weyl) algebra to get two families of irreducible representations for the extended affine Lie algebra  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  of level  $(1, 0)$  (or  $(-1, 0)$ ).

## 0 Introduction

Spinor representations for the affine Kac-Moody Lie algebras were first developed by Frenkel [F1] and Kac-Peterson [KP] independently. As one interesting application, Frenkel [F1] obtained addition formulas for elliptic  $\theta$ -functions. Their ideas were to use an Clifford algebra with infinitely many generators to construct certain quadratic elements. It turns out that these quadratic elements plus the identity element span an orthogonal affine Kac-Moody Lie algebra. Thereafter, Feingold-Frenkel [FF] constructed representations for all classical affine Kac-Moody Lie algebras by using Clifford or Weyl algebras with infinitely many generators. Following [FF] we call these corresponding representations fermionic or bosonic.

The Clifford (or Weyl) algebras have natural representations on the exterior (or symmetric) algebras of polynomials over half of generators. Those representations are also important in quantum and statistical mechanics where the generators are interpreted as operators which create or annihilate particles and satisfy Fermi (or Bose) statistics.

Motivated by [F1], [KP] and [FF], we start from the Clifford (or Weyl) algebras with generators  $a_i(n), a_i^*(n), 1 \leq i \leq N, n \in \mathbb{Z}$ , subject to relations

$$\begin{aligned}\{a_i(n), a_j(m)\}_\rho &= \{a_i^*(n), a_j^*(m)\}_\rho = 0, \\ \{a_i(n), a_j^*(m)\}_\rho &= \rho \delta_{ij} \delta_{n+m, 0},\end{aligned}$$

where  $\{a, b\}_\rho = ab + \rho ba$  is the Jordan (or Lie) bracket depending on  $\rho = \pm 1$ . We then go on to construct a class of fermions (or bosons) and obtain representations for the newly developed extended affine Lie algebra (or double affine Lie algebra)

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$\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  by enlarging the underlying module of the Clifford (or Weyl) algebra, where  $\mathbb{C}_q = \mathbb{C}_q[x^{\pm 1}, y^{\pm 1}]$  is a quantum torus of two variables. More precisely, let  $c_x$  and  $c_y$  denote the two central elements corresponding to the variables  $x$  and  $y$  respectively, the pair  $(\lambda, \mu)$  is called the level for  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  if  $c_x$  and  $c_y$  act as the scalars  $\lambda$  and  $\mu$  respectively. The modules we obtained are completely reducible and we thus got two families of irreducible representations for the extended affine Lie algebra  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  of level  $(1, 0)$  and  $(-1, 0)$  respectively. Extended affine Lie algebras are a high dimensional generalization of affine Kac-Moody Lie algebras which were introduced in [H-KT] and further studied in [BGK] and [AABGP]. A representation for the extended affine Lie algebra  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  of level  $(1, 0)$  has been obtained in [G1] and [BS] by using the principal vertex operator construction and in [G2] by using the homogeneous vertex operator construction. These vertex operator representations for  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  were constructed only when  $q$  is not a root of unity. Our fermionic and bosonic representations are carried out for any non-zero  $q$ . A notion of highest weight modules was also proposed in [G2]. Some unitary representations which are related to our algebra have been given by Jakobsen-Kac [JK] and Wakimoto [W].

Throughout this paper, we denote the field of complex numbers and the ring of integers by  $\mathbb{C}$  and  $\mathbb{Z}$  respectively.

## 1 Extended Affine Lie Algebras

In this section, we will recall some basics on construction of the extended affine Lie algebra  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$ , which is of type  $A_{N-1}$  with nullity 2. For more information on extended affine Lie algebras, see [AABGP], [BGK] and [H-KT].

Let  $q$  be a non-zero complex number. A quantum torus associated to  $q$  (see [M]) is the unital associative  $\mathbb{C}$ -algebra  $\mathbb{C}_q[x^{\pm 1}, y^{\pm 1}]$  (or, simply  $\mathbb{C}_q$ ) with generators  $x^{\pm 1}, y^{\pm 1}$  and relations

$$(1.1) \quad xx^{-1} = x^{-1}x = yy^{-1} = y^{-1}y = 1 \quad \text{and} \quad yx = qxy.$$

Then

$$(1.2) \quad x^{m_1} y^{n_1} x^{m_2} y^{n_2} = q^{n_1 m_2} x^{m_1+m_2} y^{n_1+n_2}$$

and

$$(1.3) \quad \mathbb{C}_q = \sum_{m,n \in \mathbb{Z}} \oplus \mathbb{C} x^m y^n.$$

Set  $\Lambda(q) = \{n \in \mathbb{Z} : q^n = 1\}$ .  $q$  is said to be generic if  $\Lambda(q) = \{0\}$ .

Define  $\kappa, \chi: \mathbb{C}_q \rightarrow \mathbb{C}$  to be  $\mathbb{C}$ -linear functions given by

$$(1.4) \quad \kappa(x^m y^n) = \begin{cases} 1, & \text{if } m = 0 \text{ and } n \in \Lambda(q), \\ 0, & \text{otherwise;} \end{cases}$$

$$(1.5) \quad \chi(x^m y^n) = \begin{cases} 1, & \text{if } m = 0 \text{ and } n = 0, \\ 0, & \text{otherwise.} \end{cases}$$

Let  $d_x, d_y$  be the degree operators on  $\mathbb{C}_q$  defined by

$$(1.6) \quad d_x(x^m y^n) = mx^m y^n, \quad d_y(x^m y^n) = nx^m y^n$$

for  $m, n \in \mathbb{Z}$ . Moreover, for any  $n \in \Lambda(q)$ , we define the operator

$$(1.7) \quad d(n) = y^n d_x.$$

We shall identify  $d_x$  with  $d(0)$ . Then the operators  $d_y$  and  $d(n)$ ,  $n \in \Lambda(q)$ , form a derivation subalgebra of  $\mathbb{C}_q$ . They satisfy

$$(1.8) \quad [d(n), d(n')] = 0, \quad [d_y, d(n)] = nd(n),$$

for  $n, n' \in \Lambda(q)$ . Note that  $y^n d_x$  is not a derivation if  $n \in \mathbb{Z} \setminus \Lambda(q)$ .

For the associative algebra  $\mathbb{C}_q$  over  $\mathbb{C}$ , we have the matrix algebra  $M_N(\mathbb{C}_q)$  with entries from  $\mathbb{C}_q$ . We will write  $A(x) \in M_N(\mathbb{C}_q)$  for  $A \in M_N(\mathbb{C})$  and  $x \in \mathbb{C}_q$ . Let  $\mathfrak{gl}_N(\mathbb{C}_q)$  be the Lie algebra  $M_N(\mathbb{C}_q)^-$  as usual. The Lie algebra  $\mathfrak{gl}_N(\mathbb{C}_q)$  has a nondegenerate invariant form given by

$$(1.9) \quad (A(a), B(b)) = \text{tr}(AB)\chi(ab), \quad \text{for } A, B \in M_N(\mathbb{C}), a, b \in \mathbb{C}_q.$$

Note that  $\chi(ab) = \chi(ba)$  for  $a, b \in \mathbb{C}_q$ .

Let  $c(n)$  be symbols indexed by  $n \in \Lambda(q)$ . We sometimes denote  $c(0)$  by  $c_x$ . Now we form a central extension of  $\mathfrak{gl}_N(\mathbb{C}_q)$ ,

$$(1.10) \quad \widehat{\mathfrak{gl}}_N(\mathbb{C}_q) = \mathfrak{gl}_N(\mathbb{C}_q) \oplus \left( \sum_{n \in \Lambda(q)} \oplus \mathbb{C}c(n) \right) \oplus \mathbb{C}c_y$$

with Lie bracket

$$(1.11) \quad \begin{aligned} [A(x^{m_1} y^{n_1}), B(x^{m_2} y^{n_2})] &= A(x^{m_1} y^{n_1})B(x^{m_2} y^{n_2}) - B(x^{m_2} y^{n_2})A(x^{m_1} y^{n_1}) \\ &+ \text{tr}(AB)\kappa((d_x x^{m_1} y^{n_1})x^{m_2} y^{n_2})c(n_1 + n_2) \\ &+ \text{tr}(AB)\chi((d_y x^{m_1} y^{n_1})x^{m_2} y^{n_2})c_y \end{aligned}$$

for  $m_1, m_2, n_1, n_2 \in \mathbb{Z}$ ,  $A, B \in M_N(\mathbb{C})$ , where  $c(n)$ , for  $n \in \Lambda(q)$  and  $c_y$  are central elements of  $\widehat{\mathfrak{gl}}_N(\mathbb{C}_q)$ . Note that

$$(1.12) \quad \kappa(x^{m_1+m_2} y^{n_1+n_2}) = 0 \quad \text{if } n_1 + n_2 \in \mathbb{Z} \setminus \Lambda(q).$$

The derivations  $d_y$  and  $d(n)$  for  $n \in \Lambda(q)$  can be extended to derivations on  $\mathfrak{gl}_N(\mathbb{C}_q)$ . Now we can define the semi-direct product of the Lie algebra  $\widehat{\mathfrak{gl}}_N(\mathbb{C}_q)$  and those derivations:

$$(1.13) \quad \widetilde{\widehat{\mathfrak{gl}}_N(\mathbb{C}_q)} = \widehat{\mathfrak{gl}}_N(\mathbb{C}_q) \oplus \left( \sum_{n \in \Lambda(q)} \oplus \mathbb{C}d(n) \right) \oplus \mathbb{C}d_y,$$

with Lie bracket

$$(1.14) \quad \begin{aligned} [d(n'), A(x^m y^n)] &= mA(x^m y^{n+n'}), & [d_y, A(x^m y^n)] &= nA(x^m y^n), \\ [d(n'), c(n'')] &= [d(n'), c_y] = [d_y, c_y] = 0, & [d_y, c(n')] &= n'c(n'), \end{aligned}$$

for  $m, n \in \mathbb{Z}, A \in M_N(\mathbb{C}), n', n'' \in \Lambda(q)$ .

Next we extend the nondegenerate form on  $\mathfrak{gl}_N(\mathbb{C}_q)$  to a symmetric bilinear form on  $\widetilde{\mathfrak{gl}_N(\mathbb{C}_q)}$  as follows:

$$(1.15) \quad (A(a), B(b)) = \text{tr}(AB)\chi(ab), \quad (c_y, d_y) = (c(n), d(-n)) = 1,$$

all others are zero, for  $A, B \in M_N(\mathbb{C}), a, b \in \mathbb{C}_q, n \in \Lambda(q)$ . Then:

**Lemma 1.16** *The form defined as in (1.15) is nondegenerate and invariant.*

**Proof** It is enough to show the form is invariant in the following case.

$$(1.17) \quad (d(n''), [A(x^{m_1} y^n), B(x^{m_2} y^{n'})]) = ([d(n''), A(x^{m_1} t^n)], B(x^{m_2} y^{n'})),$$

for  $m_1, m_2, n, n' \in \mathbb{Z}, n'' \in \Lambda(q), A, B \in M_N(\mathbb{C})$ . Indeed, the left hand side of (1.17) is

$$\begin{aligned} &(d(n''), m_1 q^{m_2 n} \text{tr}(AB)\kappa(x^{m_1+m_2} y^{n+n'})c(n+n')) \\ &= m_1 q^{m_2 n} \text{tr}(AB)\delta_{m_1+m_2,0}(d(n''), \kappa(y^{n+n'})c(n+n')). \end{aligned}$$

The right hand side of (1.17) is

$$\begin{aligned} (m_1 A(x^{m_1} y^{n+n'}), B(x^{m_2} y^{n'})) &= m_1 q^{m_2(n+n')} \text{tr}(AB)\chi(x^{m_1+m_2} y^{n+n'+n'}) \\ &= m_1 q^{m_2 n} \text{tr}(AB)\delta_{m_1+m_2,0}\chi(y^{n+n'+n'}). \end{aligned}$$

Since  $n'' \in \Lambda(q)$ , we have

$$(d(n''), \kappa(y^{n+n'})c(n+n')) = \chi(y^{n+n'+n'})$$

and the proof is completed. ■

If  $q$  is generic,  $\widetilde{\mathfrak{gl}_1(\mathbb{C}_q)}$  is the so-called  $q$ -analog for two-dimensional Virasoro algebra [KPS]. Its vertex operator representation was given in [G-KL].

If  $N \geq 2$ ,  $\widetilde{\mathfrak{gl}_N(\mathbb{C}_q)}$  is an extended affine Lie algebra of type  $A_{N-1}$  with nullity 2. (See [AABGP] and [BGK] for definitions.) Some representations for those Lie algebras with generic  $q$  have been obtained in [BS], [G1] and [G2].

If  $q = 1$ , then

$$\widetilde{\mathfrak{gl}_N(\mathbb{C}_q)} \cong (\widetilde{\mathfrak{gl}_N} \otimes \mathbb{C}[y, y^{-1}]) \oplus \mathbb{C}c_y \oplus \mathbb{C}d_y,$$

where

$$\widetilde{\mathfrak{gl}}_N = \mathfrak{gl}_N(\mathbb{C}[x, x^{-1}]) \oplus \mathbb{C}c_x \oplus \mathbb{C}d_x$$

is the affinization of the general linear Lie algebra  $\mathfrak{gl}_N$ .

Let  $E_{ij}$  be the matrix whose  $(i, j)$ -entry is 1 and 0 elsewhere. Then, in  $\widehat{\mathfrak{gl}}_N(\mathbb{C}_q)$ ,

$$(1.18) \quad \begin{aligned} [E_{ij}(x^{m_1} y^{n_1}), E_{kl}(x^{m_2} y^{n_2})] &= \delta_{jk} q^{n_1 m_2} E_{il}(x^{m_1+m_2} y^{n_1+n_2}) - \delta_{il} q^{n_2 m_1} E_{kj}(x^{m_1+m_2} y^{n_1+n_2}) \\ &\quad + m_1 q^{n_1 m_2} \delta_{jk} \delta_{il} \delta_{m_1+m_2, 0} \kappa(y^{n_1+n_2}) c(n_1 + n_2) \\ &\quad + q^{n_1 m_2} \delta_{jk} \delta_{il} \delta_{m_1+m_2, 0} \delta_{n_1+n_2, 0} n_1 c_y \end{aligned}$$

for  $m_1, m_2, n_1, n_2 \in \mathbb{Z}$ .

## 2 Construction of Fermions and Bosons

In this section, we will give a unified treatment for both fermionic and bosonic constructions as were done in [FF]. However, for simplicity, we use a slightly different normal ordering from [FF] (see also [F2]).

Let  $\mathcal{R}$  be an associative algebra. Let  $\rho = \pm 1$ . We define a  $\rho$ -bracket on  $\mathcal{R}$  as follows:

$$(2.1) \quad \{r_1, r_2\}_\rho = r_1 r_2 + \rho r_2 r_1, \quad r_1, r_2 \in \mathcal{R}.$$

Then, one can easily see that

$$(2.2) \quad \{r_1, r_2\}_\rho = \rho \{r_2, r_1\}_\rho \text{ and } [r_1 r_2, r_3] = r_1 \{r_2, r_3\}_\rho - \rho \{r_1, r_3\}_\rho r_2$$

for  $r_1, r_2, r_3, r_4 \in \mathcal{R}$ , where  $[r_1, r_2] = \{r_1, r_2\}_{-1}$  is the Lie bracket.

Define  $\mathfrak{a}(N, \rho)$  to be the unital associative algebra with infinitely many generators  $a_i(n), a_i^*(n), n \in \mathbb{Z}, 1 \leq i \leq N$ , subject to relations

$$(2.3) \quad \begin{aligned} \{a_i(n), a_j(m)\}_\rho &= \{a_i^*(n), a_j^*(m)\}_\rho = 0, \\ \{a_i(n), a_j^*(m)\}_\rho &= \rho \delta_{ij} \delta_{n+m, 0}. \end{aligned}$$

We now define the normal ordering as follows.

$$(2.4) \quad : a_i(n) a_j^*(m) : = \begin{cases} a_i(n) a_j^*(m) & \text{if } n \leq m, \\ -\rho a_j^*(m) a_i(n) & \text{if } n > m, \end{cases}$$

for  $n, m \in \mathbb{Z}, 1 \leq i, j \leq N$ . Set

$$(2.5) \quad \theta(n) = \begin{cases} 1, & \text{for } n > 0, \\ 0, & \text{for } n \leq 0. \end{cases}$$

Then

$$(2.6) \quad a_i(n) a_j^*(m) = : a_i(n) a_j^*(m) : + \rho \delta_{ij} \delta_{n+m, 0} \theta(n - m).$$

It follows from (2.2) that

$$(2.7) \quad \begin{aligned} [a_i(m)a_j^*(n), a_k(p)] &= \delta_{jk}\delta_{n+p,0}a_i(m), \\ [a_i(m)a_j^*(n), a_k^*(p)] &= -\delta_{ik}\delta_{m+p,0}a_j^*(n), \end{aligned}$$

for  $m, n, p \in \mathbb{Z}, 1 \leq i, j, k \leq N$ .

Let  $\mathfrak{a}(N, \rho)^+$  be the subalgebra generated by  $a_i(n), a_j^*(m), a_k^*(0)$ , for  $n, m > 0$ , and  $1 \leq i, j, k \leq N$ . Let  $\mathfrak{a}(N, \rho)^-$  be the subalgebra generated by  $a_i(n), a_j^*(m), a_k(0)$ , for  $n, m < 0$ , and  $1 \leq i, j, k \leq N$ . Those generators in  $\mathfrak{a}(N, \rho)^+$  are called annihilation operators while those in  $\mathfrak{a}(N, \rho)^-$  are called creation operators. Let  $V(N, \rho)$  be a simple  $\mathfrak{a}(N, \rho)$ -module containing an element  $v_0$ , called a “vacuum vector”, and satisfying

$$(2.8) \quad \mathfrak{a}(N, \rho)^+ v_0 = 0.$$

So all annihilation operators kill  $v_0$  and

$$(2.9) \quad V(N, \rho) = \mathfrak{a}(N, \rho)^- v_0.$$

Now we may construct a class of fermions (if  $\rho = 1$ ) or bosons (if  $\rho = -1$ ) on  $V(N, \rho)$ . For any  $m, n \in \mathbb{Z}, 1 \leq i, j \leq N$ , set

$$(2.10) \quad f_{ij}(m, n) = \sum_{s \in \mathbb{Z}} q^{-ns} a_i(m-s) a_j^*(s) \text{.}$$

We also set

$$(2.11) \quad D = \sum_{i=1}^N \sum_{s \in \mathbb{Z}} s a_i(s) a_i^*(-s) \text{.}$$

Although  $f_{ij}(m, n)$  and  $D$  are infinite sums, they are well-defined as operators on  $V(N, \rho)$ . Indeed, for any vector  $v \in V(N, \rho) = \mathfrak{a}(N, \rho)^- v_0$ , only finitely many terms in (2.10) and (2.11) can make a non-zero contribution to  $f_{ij}(m, n)v$  and  $Dv$ .

**Lemma 2.12** *We have*

$$(2.13) \quad [f_{ij}(m, n), a_k(p)] = \delta_{jk} q^{np} a_i(m+p),$$

$$(2.14) \quad [f_{ij}(m, n), a_k^*(p)] = -\delta_{ik} q^{-n(m+p)} a_j^*(m+p),$$

$$(2.15) \quad [f_{ij}(m, n), a_k(p) a_l^*(s)] = \delta_{jk} q^{np} a_i(m+p) a_l^*(s) - \delta_{il} q^{-n(m+s)} a_k(p) a_j^*(m+s).$$

for  $m, n, p, s \in \mathbb{Z}$  and  $1 \leq i, j, k, l \leq N$ .

**Proof**

$$\begin{aligned} [f_{ij}(m, n), a_k(p)] &= \sum_{s \in \mathbb{Z}} q^{-ns} [ : a_i(m-s)a_j^*(s) :, a_k(p) ] \\ &= \sum_{s \in \mathbb{Z}} q^{-ns} [ a_i(m-s)a_j^*(s), a_k(p) ] \\ &= \sum_{s \in \mathbb{Z}} q^{-ns} \delta_{jk} \delta_{s+p,0} a_i(m-s) = q^{np} \delta_{jk} a_i(m+p). \end{aligned}$$

So (2.13) holds true. The proof of (2.14) is similar, and (2.15) follows from (2.13) and (2.14). ■

In what follows we shall mean  $\frac{q^{mn}-1}{q^n-1} = m$  if  $q^n = 1$ . This will make our formula more concise.

**Proposition 2.16**

$$\begin{aligned} [f_{ij}(m_1, n_1), f_{kl}(m_2, n_2)] &= \delta_{jk} q^{n_1 m_2} f_{il}(m_1 + m_2, n_1 + n_2) \\ &\quad - \delta_{il} q^{n_2 m_1} f_{kj}(m_1 + m_2, n_1 + n_2) \\ &\quad + \rho \delta_{jk} \delta_{il} q^{n_1 m_2} \delta_{m_1+m_2,0} \frac{q^{m_1(n_1+n_2)} - 1}{q^{n_1+n_2} - 1} q^{n_1+n_2} \end{aligned}$$

for all  $m_1, m_2, n_1, n_2 \in \mathbb{Z}$  and  $1 \leq i, j, k, l \leq N$ .

**Proof** It follows from (2.15) and (2.6) that

$$\begin{aligned} [f_{ij}(m_1, n_1), q^{-n_2 t} : a_k(m_2 - t)a_l^*(t) :] &= \delta_{jk} q^{n_1 m_2 - (n_1+n_2)t} a_i(m_1 + m_2 - t) a_l^*(t) \\ &\quad - \delta_{il} q^{n_2 m_1 - (n_1+n_2)(m_1+t)} a_k(m_2 - t) a_j^*(m_1 + t) \\ &= \delta_{jk} q^{n_1 m_2 - (n_1+n_2)t} ( : a_i(m_1 + m_2 - t) a_l^*(t) : + \rho \delta_{il} \delta_{m_1+m_2,0} \theta(m_1 + m_2 - 2t) ) \\ &\quad - \delta_{il} q^{n_2 m_1 - (n_1+n_2)(m_1+t)} ( : a_k(m_2 - t) a_j^*(m_1 + t) : \\ &\quad \quad \quad + \rho \delta_{jk} \delta_{m_1+m_2,0} \theta(m_2 - t - m_1 - t) ) \\ &= \delta_{jk} q^{n_1 m_2} q^{-(n_1+n_2)t} : a_i(m_1 + m_2 - t) a_l^*(t) : \\ &\quad - \delta_{il} q^{n_2 m_1} q^{-(n_1+n_2)(m_1+t)} : a_k(m_2 - t) a_j^*(m_1 + t) : \\ &\quad + \rho \delta_{jk} \delta_{il} \delta_{m_1+m_2,0} q^{n_1 m_2} q^{-(n_1+n_2)t} ( \theta(-2t) - \theta(-2m_1 - 2t) ). \end{aligned}$$

Since

$$\begin{aligned} & \sum_{t \in \mathbb{Z}} q^{-(n_1+n_2)t} (\theta(-2t) - \theta(-2m_1 - 2t)) \\ &= \begin{cases} 0, & \text{if } m_1 = 0, \\ \sum_{t=-m_1}^{-1} q^{-(n_1+n_2)t}, & \text{if } m_1 > 0, \\ \sum_{t=0}^{-m_1-1} q^{-(n_1+n_2)t}, & \text{if } m_1 < 0, \end{cases} \\ &= \frac{q^{m_1(n_1+n_2)} - 1}{q^{n_1+n_2} - 1} q^{n_1+n_2}, \end{aligned}$$

we see that (2.16) holds true and the proof is completed. ■

**Lemma 2.17**

(2.18)  $[D, a_k(p)] = pa_k(p),$

(2.19)  $[D, a_k^*(p)] = pa_k^*(p),$

(2.20)  $[D, f_{ij}(m, n)] = mf_{ij}(m, n),$

for  $p, m, n \in \mathbb{Z}, 1 \leq k, i, j \leq N.$

**Proof** Let  $D_i = \sum_{s \in \mathbb{Z}} s: a_i(s)a_i^*(-s) :$ , for  $1 \leq i \leq N.$  As in (2.13), we may obtain that  $[D_i, a_k(p)] = \delta_{ik}pa_k(p)$  and hence (2.18) follows. (2.19) can be proved similarly.

Now it follows from (2.18) and (2.19) that

$$\begin{aligned} [D, q^{-ns}: a_k(m-s)a_l^*(s) :] &= mq^{-ns}a_k(m-s)a_l^*(s) \\ &= mq^{-ns}(: a_k(m-s)a_l^*(s) : + \rho\delta_{kl}\delta_{m,0}\theta(m-2s)) \\ &= mq^{-ns}: a_k(m-s)a_l^*(s) : \end{aligned}$$

which clearly yields (2.20). ■

Now we see that the operators  $f_{ij}(m, n)$  for  $m, n \in \mathbb{Z}, 1 \leq i, j \leq N,$  together with 1 and  $D,$  form a Lie algebra  $\mathcal{L}.$  Next we will consider two subalgebras  $\mathcal{L}_v$  and  $\mathcal{L}_h$  of  $\mathcal{L}.$  To this end, we first let

(2.21) 
$$\mathcal{J} = \sum_{i=1}^N f_{ii}(0, 0),$$

then one can easily show that

(2.22)  $[J, a_j(n)] = a_j(n) \quad \text{and} \quad [J, a_j^*(n)] = -a_j^*(n),$

for  $n \in \mathbb{Z}, 1 \leq j \leq N.$  For any

(2.23)  $v = a_{i_1}(n_1) \cdots a_{i_s}(n_s)a_{j_1}^*(m_1) \cdots a_{j_r}^*(m_r)v_0$

from  $V(N, \rho)$ , one has

$$(2.24) \quad \mathfrak{J}v = (s - t)v.$$

Let  $\mathcal{L}_v$  be the subalgebra spanned by operators  $f_{ij}(m, 0)$  for  $m \in \mathbb{Z}$ ,  $1 \leq i, j \leq N$ . Then  $\mathcal{L}_v$  is isomorphic to the affinization of the general linear Lie algebra  $\mathfrak{gl}_N$ . Moreover, as (5.56) in [FF], we have:

**Proposition 2.25** *As  $\mathcal{L}_v$ -module,  $V(N, \rho)$  is completely reducible. Moreover,*

$$V(N, \rho) = \sum_{k \in \mathbb{Z}} \oplus V_k,$$

where  $V_k$  is the  $k$ -eigenspace of the operator  $\mathfrak{J}$ , and each  $V_k$  is an irreducible  $\mathcal{L}_v$ -module.

**Remark 2.26** Let  $\mathcal{L}_h$  be the subalgebra spanned by operators  $f_{ij}(0, n)$  for  $n \in \mathbb{Z}$ ,  $1 \leq i, j \leq N$ . Then  $\mathcal{L}_h$  is isomorphic to the Lie algebra  $\mathfrak{gl}_N(\mathbb{C}[Z/\Lambda(0)])$ , where  $\mathbb{C}[Z/\Lambda(0)]$  is the group algebra. In particular, if  $q$  is generic, then  $\mathcal{L}_h$  is isomorphic to the loop algebra  $\mathfrak{gl}_N(\mathbb{C}[Z])$ . Namely, each  $\mathcal{L}_v$ -module  $V_k$  allows  $\mathcal{L}_h$ -action. This shows that the level- $\rho$  module for the affine Lie algebra  $\hat{\mathfrak{gl}}_N$  allows a level-0 action for  $\hat{\mathfrak{gl}}_N$ .

### 3 Lifting to $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$

In this section, we shall obtain the  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$ -module action by enlarging the Fock space  $V(N, \rho)$ .

From (2.16), we see that, if  $n_1 + n_2 \in \Lambda(q)$ ,

$$(3.1) \quad \begin{aligned} [f_{ij}(m_1, n_1), f_{kl}(m_2, n_2)] &= \delta_{jk}q^{n_1m_2} f_{il}(m_1 + m_2, n_1 + n_2) \\ &\quad - \delta_{il}q^{n_2m_1} f_{kj}(m_1 + m_2, n_1 + n_2) \\ &\quad + \rho\delta_{jk}\delta_{il}q^{n_1m_2}\delta_{m_1+m_2,0}m_1, \end{aligned}$$

for  $m_1, m_2, n_1, n_2 \in \mathbb{Z}$ ,  $1 \leq i, j, k, l \leq N$ .

If  $n_1 + n_2 \in \mathbb{Z} \setminus \Lambda(q)$ , then

$$(3.2) \quad \begin{aligned} [f_{ij}(m_1, n_1), f_{kl}(m_2, n_2)] &= \delta_{jk}q^{n_1m_2} f_{il}(m_1 + m_2, n_1 + n_2) - \delta_{il}q^{n_2m_1} f_{kj}(m_1 + m_2, n_1 + n_2) \\ &\quad + \rho\delta_{jk}\delta_{il}\delta_{m_1+m_2,0} \frac{q^{(n_1+n_2)}}{q^{n_1+n_2} - 1} (q^{n_2m_1} - q^{n_1m_2}) \\ &= \delta_{jk}q^{n_1m_2} \left( f_{il}(m_1 + m_2, n_1 + n_2) - \rho\delta_{il}\delta_{m_1+m_2,0} \frac{q^{n_1+n_2}}{q^{n_1+n_2} - 1} \right) \\ &\quad - \delta_{il}q^{n_2m_1} \left( f_{kj}(m_1 + m_2, n_1 + n_2) - \rho\delta_{kj}\delta_{m_1+m_2,0} \frac{q^{n_1+n_2}}{q^{n_1+n_2} - 1} \right), \end{aligned}$$

for  $m_1, m_2, n_1, n_2 \in \mathbb{Z}, 1 \leq i, j, k, l \leq N$ . Therefore, if we define

$$(3.3) \quad F_{ij}(m, n) = \begin{cases} f_{ij}(m, n), & \text{for } n \in \Lambda(q) \\ f_{ij}(m, n) - \rho \delta_{ij} \delta_{m,0} \frac{q^n}{q^n - 1}, & \text{for } n \in \mathbb{Z} \setminus \Lambda(q). \end{cases}$$

Then we have

$$(3.4) \quad [F_{ij}(x^{m_1} y^{n_1}), F_{kl}(x^{m_2} y^{n_2})] = \delta_{jk} q^{n_1 m_2} F_{il}(x^{m_1+m_2} y^{n_1+n_2}) - \delta_{il} q^{n_2 m_1} F_{kj}(x^{m_1+m_2} y^{n_1+n_2}) + m_1 q^{n_1 m_2} \delta_{jk} \delta_{il} \delta_{m_1+m_2,0} \kappa(y^{n_1+n_2}) \rho$$

for  $m_1, m_2, n_1, n_2 \in \mathbb{Z}$ .

Set

$$(3.5) \quad W(N, \rho) = V(N, \rho) \otimes \mathbb{C}[y, y^{-1}],$$

and define operators

$$(3.6) \quad e_{ij}(m, n) = F_{ij}(m, n) \otimes y^n$$

which act on  $W(N, \rho)$  as follows

$$(3.7) \quad e_{ij}(m, n)(v \otimes y^k) = F_{ij}(m, n)v \otimes y^{n+k},$$

for  $m, n, k \in \mathbb{Z}, 1 \leq i, j \leq N$  and  $v \in V(N, \rho)$ . Then following from (1.18), (3.4) and Proposition 2.25 one can easily prove that:

**Theorem 3.8**  $W(N, \rho)$  is a module for the Lie algebra  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  under the action given by

$$\begin{aligned} \pi(E_{ij}(x^m y^n)) &= e_{ij}(m, n), \quad \text{for } 1 \leq i, j \leq N, m, n \in \mathbb{Z}; \\ \pi(c(n)) &= \rho 1 \otimes y^n, \quad \pi(d(n)) = D \otimes y^n, \quad \text{for } n \in \Lambda(q), \\ \pi(c_y) &= 0, \quad \pi(d_y) = 1 \otimes d_y. \end{aligned}$$

Moreover,  $W(N, \rho)$  is completely reducible and each component  $V_k \otimes \mathbb{C}[y, y^{-1}]$  is irreducible.

**Remark 3.9** The above theorem gives us two families of irreducible representations for the extended affine Lie algebra  $\widetilde{\mathfrak{gl}}_N(\mathbb{C}_q)$  of level  $(\rho, 0)$ , for  $N \geq 2$ . In particular, if  $q$  is not a root of unity (equivalently  $\Lambda(q) = \{0\}$ ), then the center of the central extension in (1.10) is two dimensional. In this case, an irreducible vertex operator representation has been constructed in [BS], [G1] and [G2].

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