

ON THE DETERMINATION OF HORIZONTAL FORCES A FLOATING ICE PLATE EXERTS ON A STRUCTURE*

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ABSTRACT. At first, the general approach for calculating the horizontal forces an ice cover exerts on structures is discussed. Ice-force determination consists of two parts: (1) the analysis of the in-plane forces, assuming that the ice cover remains intact, and (2) the use of a failure criterion; because an ice force cannot be larger than the force capable of breaking up the ice cover. For an estimate of the *largest* ice force, an *elastic* plate analysis and a failure criterion are often sufficient. A review of the literature revealed that in the majority of the analyses, it is assumed that the failure load is directly related to a "crushing strength" of the ice cover. Observations in the field and tests in the laboratory show, however, that in some instances the ice cover failed by buckling. Subsequently the ice-force analyses based on the *buckling* failure mechanism are reviewed and their shortcomings are pointed out. A new method of analysis, which is based on the buckling of a floating ice wedge, is then presented.

RÉSUMÉ. Sur la détermination des forces horizontales qu'une plaque de glace flottante exerce sur une structure. Tout d'abord, on discute l'approche générale du calcul des forces horizontales qu'une nappe de glace exerce sur une structure. La détermination des forces dues à la glace comprend deux parties: (1) l'analyse des forces dans un plan dans l'hypothèse où la nappe de glace est sensée rester intacte et (2) l'introduction du seuil de rupture; en effet une force exercée par la glace ne peut pas être supérieure à celle capable de briser la nappe de glace. Pour l'estimation des forces les plus grandes une analyse plane élastique et l'emploi d'un seuil de rupture sont souvent suffisants. Une revue de la littérature révèle que, dans la majorité des analyses, on suppose que la charge de rupture est directement reliée à une "résistance à l'écrasement" de la plaque de glace. Les observations *in situ* et les essais en laboratoire montrent cependant que dans quelques cas la glace se rompt au flambage. Par conséquent les analyses des forces dans la glace basées sur les mécanismes de ruptures au flambage sont revues et leurs insuffisances sont tirées au clair. On présente alors une nouvelle méthode d'analyse basée sur le flambage d'une poutre de glace flottante.

ZUSAMMENFASSUNG. Zur Bestimmung der von einer schwimmenden Eisplatte auf ein Bauwerk ausgeübten Horizontalkräfte. Zunächst wird das generelle Verfahren zur Berechnung der Horizontalkräfte, die eine Eisdecke auf Bauwerke ausübt, dargelegt. Bestimmung von Eiskräften besteht aus zwei Teilen: (1) Die Analyse der inneren Flächenkräfte unter der Annahme, dass die Eisdecke unbeschädigt bleibt; (2) Die Einführung eines Bruchkriteriums, weil keine Eiskraft grösser werden kann als die Kraft, die zum Bruch der Eisdecke führt. Zur Abschätzung der grössten Eiskraft genügt oft eine elastische Plattenanalyse und ein Bruchkriterium. Aus der Durchsicht der Literatur geht hervor, dass die Mehrzahl der Analysen von der Annahme ausgeht, dass die Bruchlast stehe in unmittelbarer Beziehung zur "Quetschfestigkeit" der Eisdecke. Feldbeobachtungen und Laborversuche zeigen dagegen, dass in einigen Fällen die Eisdecke durch Knicken nachgab. Deshalb werden die Eiskraftanalysen auf der Basis des Nachgebens durch einen Knickmechanismus herangezogen und ihre Unzulänglichkeiten hervorgehoben. Schliesslich wird eine neue Methode der Analyse vorgelegt, die auf die Knickung eines schwimmenden Eiskeiles beruht.

I. INTRODUCTION AND STATEMENT OF PROBLEM

When a river, lake, or part of a sea freezes over, the formed ice cover is often in contact with a variety of structures, such as bridge piers, dams, hydroelectric power stations, and off-shore drilling platforms. Subsequently, any attempt to change the position of the ice cover (by changing the water level, or by the movement of the water base beneath it or the air above it, or by the change in the air temperature) creates forces between the structures and the formed ice cover.

For a rational design of a structure which comes in contact with a floating ice cover, the engineer needs to know the forces the ice will exert on this structure; especially the value of the largest possible force.

An early description of problems of this type was given by Barnes (1906, 1928) and by Komarovskiy (1932-33). The effort and progress made in this field during the following few decades was very modest, as evidenced by the results presented at the ASCE symposium on Ice pressure against dams in 1954.

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During the past decade the problem of ice forces gained in importance, partly because of the discovery of oil in the northern regions and the resulting increased activity in this area. The intensified research effort that followed is evidenced by the recent increase of conferences on this subject, such as the conference on Ice pressures against structures (Gold and Williams, 1968), the Ice seminar (1969), the international conferences on Port and ocean engineering under Arctic conditions (Trondheim, Norway, 1971; Reykjavik, Iceland, 1973; Fairbanks, Alaska, U.S.A., 1975), as well as the IAHR symposia on Ice and its action on hydraulic structures (Reykjavik, Iceland, 1970; Leningrad, U.S.S.R., 1972; Hanover, New Hampshire, U.S.A., 1975). In spite of all these efforts, to date there are no reliable methods for predicting the ice forces an ice cover may exert on a structure.

The analyses of these forces are usually grouped into several categories. For example, analyses of forces due to a rise or drop in the water level, analyses of forces because a structure constrains the movement of the ice cover in the horizontal plane, and analyses of forces due to impact of a moving ice flow and a structure. The present paper discusses problems in the second category. Examples of such problems are shown in Figure 1.

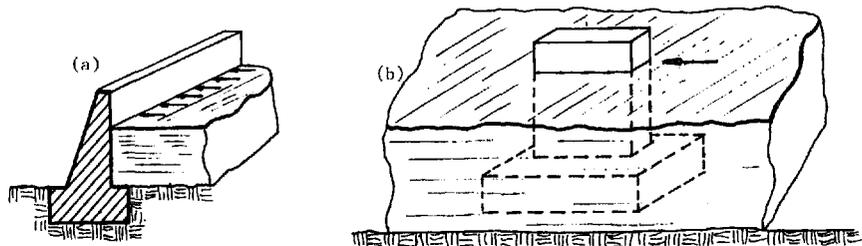


Fig. 1. Typical situations in which forces due to ice constraints are applied to a structure: (a) ice cover pressing against a rigid wall, (b) ice cover pressing against a rectangular pier.

Since ice is a visco-elastic material, the forces an ice cover exerts on a structure are time-dependent. The study of these forces, caused by constrained thermal expansions, was initiated by N. Royen in 1922. The analytical and test results by Royen, as well as the more recent results of a number of other investigators, were reviewed by Korzhavin (1962, chapter 14), by Drouin (1968), and by Michel (1970). Other attempts to solve this problem were presented by Panfilov (1965), Nevel (1968) and Jumppanen ([1974]).

The conclusion of the survey by Drouin was that the determination of ice forces as a function of time is not yet solved satisfactorily. The relevant results of Panfilov (1965), Nevel (1968), and Jumppanen ([1974]) are as yet not conclusive either. A major difficulty in solving this problem is the absence of appropriate constitutive equations for the ice cover. It should also be noted that even if expressions which determine the ice forces were available, the magnitude of these forces may sometimes be limited by a possible failure (crushing or buckling) of the ice cover, as discussed previously.

For the design of a structure which is in contact with an ice cover, it is often sufficient to know the *largest* force the ice cover may exert on the structure. This approach simplifies the necessary analyses.

To demonstrate the essential aspects of this ice-force problem, consider as an example, an ice cover constrained by two parallel "rigid" walls, as shown in Figure 2. Assume that (at a time $t = 0$) the temperature of the ice plate is instantaneously raised by a constant value, ΔT . Then, the corresponding axial compression force (per unit length of wall) is approximately

$$\bar{p}_0 = Eh\alpha\Delta T. \quad (1)$$

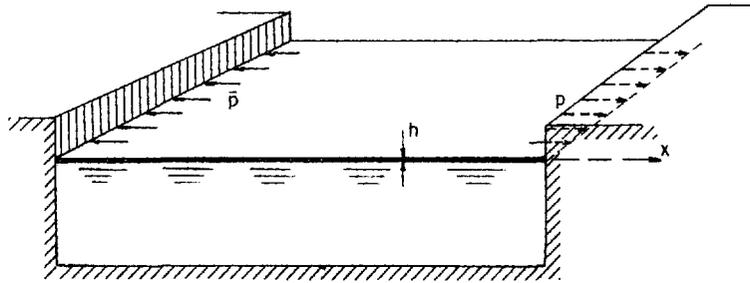


Fig. 2. Floating ice cover constrained by two parallel "rigid" walls.

Equation (1) is based on the assumption that, at the instant of the temperature rise ($t = 0$), the ice obeys Hooke's law, that Young's modulus E is an averaged value across the plate thickness, and that the coefficient of linear expansion α is also an average value which does not vary noticeably in the temperature range of interest.

Because of the visco-elastic response of the ice cover, the pressure \bar{p} , for $t > 0$, will be smaller than the \bar{p}_0 value given in Equation (1), namely

$$\bar{p}(t) < \bar{p}_0. \quad (2)$$

A similar argument leads to the conclusion that Equation (2) is also valid for the case when the temperature rise is not instantaneous but proceeds gradually, as is the case in the field. Thus Equation (1) yields the largest possible axial compression force due to ΔT . This suggests that, for many engineering problems, it may be sufficient to determine the design forces acting on a structure by assuming that the ice cover responds elastically.

At this point it should be noted that the force estimate given by Equation (1), or the more accurate expression which takes into consideration the visco-elastic response of the ice cover, is valid as long as the ice cover does not fail. Thus, beyond a certain compression force, Equation (1), or its visco-elastic equivalent, yields too high values for the ice force, and it has to be bounded by a failure criterion. This criterion is obtained by noting that *an ice force cannot be larger than the force capable of breaking up the ice cover*. Thus the needed bound for in-plane compression forces is established by determining the intensity of the force at which the ice cover fails. Note, that this criterion is valid for the determination of the largest ice forces that may be caused by a temperature change or due to any other reason such as the movement of the water in the horizontal direction.

Many published analyses are based on this notion. In the majority of these publications it is assumed that this failure load is directly related to a "crushing strength" of the ice. Namely, that the mechanism of failure is due to crushing or splitting of the floating ice plate in the immediate vicinity of the structure. For a review of the relevant literature up to 1962 refer to Korzhavin (1962). For a more recent utilization of this approach refer to Shadrin and Panfilov (1962), Afanas'yev (1968), Michel (1970), Schwarz ([1971]), Afanas'yev (1972) and Assur ([1974]).

However, in laboratory and field tests it was observed by E. Rose (1947), Kheysin (1961), Afanas'yev and others (1971) and Nevel and others (1972), that for relatively thin ice plates (compared to the width of contact of plate and structure) the ice cover failed by buckling in the vicinity of the contact area.

The published analyses of ice forces on isolated structures which are based on the buckling failure mechanism, are not conclusive, however. The purpose of the present paper is first to review these few attempts and then to develop an improved method of analysis for the determination of the largest force that a relatively thin ice cover may exert on an isolated structure. The proposed analysis is based on the buckling mechanism of the floating ice plate.

2. REVIEW OF RELEVANT ANALYSES AND TESTS

An early attempt to determine analytically the buckling force for an ice plate was presented by Albenga (1921). However, because Albenga neglected the effect of the liquid base, the obtained results do not apply to floating ice covers.

On the basis of observations in the field, Rose (1947, p. 883) stated that "arching and buckling of an ice sheet does occur but that there is a limiting thickness of about 1 foot [0.3 m] beyond which buckling will rarely occur". He also claimed (p. 880) that "for ice sheets less than 1 foot [0.3 m] thick the critical buckling load is the limiting factor governing the ice pressure on a dam, whereas for very thick ice sheets (about 2 feet [0.6 m] or more) the crushing strength pressure" is the limiting factor. No analysis was given for the determination of the buckling load.

Shadrin and Panfilov (1962) presented an analysis for the semi-infinite ice cover whose straight edge is pressing against a rigid wall, as shown in Figures 1(a) or 2. The authors assumed that the type of ice cover failure depends on the thickness of the ice cover; namely, buckling for thin covers and local crushing at the contact area for thick covers. This assumption is in agreement with the observations made by Rose (1947), as described above.

Shadrin and Panfilov expressed the failure load (per unit length) as

$$\bar{p}_t = \sigma_{ct}h, \quad (3)$$

where σ_{ct} is a "compression strength" and h is the thickness of the ice cover.

For relatively thin covers the buckling load was determined from the differential equation

$$D \frac{d^4w}{dx^4} + \bar{p} \frac{d^2w}{dx^2} + \gamma w = 0, \quad (4)$$

and the boundary conditions

$$\left. \begin{aligned} w(0) &= 0, \\ \frac{d^2w}{dx^2} \Big|_0 &= 0, \end{aligned} \right\} \quad (5)$$

where $w(x)$ is the vertical deflection at x and at any value $-\infty < y < \infty$, $D = Eh^3/[12(1-\nu^2)]$ is the flexural rigidity of the cover, and γ is the specific weight of the liquid base. The solution of the above formulation yields

$$\bar{p}_b = 2(\gamma D)^{\frac{1}{2}}. \quad (6)$$

It should be noted that the above expression for \bar{p}_b is also valid when the boundary at $x = 0$ is clamped (Hetényi, 1946, ch. 7), as may often be the case in the field.

Next, Equations (3) and (6) are compared by noting that according to tests, if the corresponding $\bar{p}_b < \bar{p}_t$ ice cover fails by buckling and if $\bar{p}_t < \bar{p}_b$ ice cover fails by crushing. Substituting Equations (3) and (6) in the first inequality, it follows that the ice cover will fail by buckling when

$$h < \frac{3(1-\nu^2) \sigma_{ct}^2}{\gamma E}. \quad (7)$$

As an example, if $\nu = 0.34$, $\gamma = 0.001 \text{ kg/cm}^3$, $E = 40\,000 \text{ kg/cm}^2$, and $\sigma_{ct} = 20 \text{ kg/cm}^2$ then according to Equation (7) the ice cover will buckle when $h < 27 \text{ cm} \approx 10.5 \text{ inches}$. This result agrees reasonably well with the field observations reported by Rose (1947).

According to the above analysis when Inequality (7) is satisfied, the largest force that the ice cover may exert on the long rigid structure is the buckling load $\bar{p}_b = 2(\gamma D)^{\frac{1}{2}}$; otherwise the largest force is $\bar{p}_t = \sigma_{ct}h$.

Next we review the calculation of the largest force an ice cover exerts on a ship hull, presented by Kheysin (1961); a problem closely related to the determination of ice forces on

isolated structures. Kheysin reduced this problem to the determination of a uniform pressure p at which a floating elastic semi-infinite plate, shown in Figure 3, will fail. He based his analysis on observations in the field that ice cover failures often occur in two stages: the formation of cracks which emanate from the points 1 and 2, followed by the buckling of the inner part of the ice cover which is formed by these cracks.

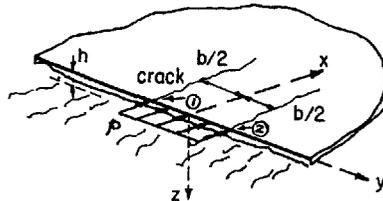


Fig. 3. Ice pressing against a wall of limited length.

For his analysis Kheysin assumed that the cracks are initiated by the largest shear stress along the edge

$$(\sigma_{xy})_{\max} = \sigma_{xy}(0, \pm b/2) = \frac{p}{\pi h}. \quad (8)$$

Choosing for the "shear strength" of the (sea-)ice cover $\sigma_{st} = 30 \text{ Mg/m}^2$, and using the equation

$$(\sigma_{xy})_{\max} = \sigma_{st}, \quad (9)$$

as a crack criterion, he obtained from Equation (8) the pressure intensity which causes the formation of the cracks as

$$p_{st} \approx (94 \text{ Mg/m}^2) h. \quad (10)$$

To simplify the buckling analysis, Kheysin assumed that the two cracks which emanate from the points $(0, \pm b/2)$ are parallel to the x -axis, and that the formed ice strip buckles when

$$p_b|_{\text{cyl bend}} = (\gamma D)^{\frac{1}{2}}. \quad (11)$$

Assuming $\nu = 0.36$, $\gamma = 1 \text{ Mg/m}^3$, and $E = 40\,000 \text{ Mg/m}^2$, the above equation reduces to

$$p_b|_{\text{cyl bend}} = (62 \text{ Mg/m}^{\frac{3}{2}}) h^{\frac{1}{2}}. \quad (12)$$

Comparing Equations (10) and (12), Kheysin found that for $0 < h < 2.2 \text{ m}$,* which is the range of practical interest, the inequality

$$p_{st} > p_b|_{\text{cyl bend}}, \quad (13)$$

holds (although, according to Kheysin $p_{st} < p_b$, where p_b is the buckling load of the uncracked, semi-infinite plate). Based on observations that an ice cover cracks prior to its destruction, Kheysin (1961) concluded that the largest ice force is given by p_{st} , namely Equation (10). According to the above analysis, after the plate cracks, buckling takes place immediately, since $p_b|_{\text{cyl bend}} < p_{st}$.

The above method is based on a number of questionable assumptions. Namely, that Equation (9) is a valid criterion, that the cracks are parallel and normal to the free edge, etc. It should also be noted that p_b given in Equation (11) is the critical buckling pressure when the edge is free to rotate and displace vertically, which is usually not the case when the ice cover is in contact with a fixed vertical structure or the hull of a ship. Other shortcomings of

* Note that in the previous example (Shadrin and Panfilov, 1962) $E = 400\,000 \text{ Mg/m}^2$ was used. For this value of E the corresponding interval is $0 < h < 0.23 \text{ m}$.

Equation (11) are that the ice strip formed in the field is not of constant width and is constrained along the entire length of the formed edges by the surrounding ice cover. Also, according to Nevel and others (1972), who conducted many closely related tests, the formation of radial cracks is *not* followed immediately by ice cover failure. Therefore, the results obtained by Kheysin, which are also reproduced in the book by Ya. N. Popov and others (1967, p. 60–62) should be used with caution.

Afanas'yev and others (1971) published results of an extensive laboratory test program in which a number of thin floating ice plates ($2.6 \text{ cm} \leq h \leq 3.4 \text{ cm}$) were subjected to a horizontal force, by pressing a circular or a rectangular model pier against the free edge of the plate. According to these results, when $b/h > 4$ and the ice cover was not made more rigid by the proximity of the basin walls or the presence of simulated ridges, the ice plate failed by buckling. During these tests it was observed (p. 64) that at a certain intensity of the load, radial cracks formed in the plate. At a further increase of the load, the ice plate lifted along a nearly circular curve at a distance of about $14h$ from the pier. This was followed by a breakdown of the ice field.

Afanas'yev and others (1971) then presented an analysis for the semi-infinite ice plate pressing uniformly against a rigid wall, which is similar to the one given by Shadrin and Panfilov (1962) and discussed previously. The obtained results are also similar, namely, that for fresh-water ice the ice cover buckles when $h < 30 \text{ cm}$ and for salt-water ice when $h < 20 \text{ cm}$. Thus, for cases for which these inequalities are satisfied, the buckling load is the largest force the ice cover can exert against the wall.

Nevel and others (1972) presented results of extensive ice force tests conducted at CRREL. The test set-up was similar to the one described by Afanas'yev and others (1971), except that the used basin ($6.3 \text{ m} \times 6.3 \text{ m}$) and the thickness of the tested plates ($6.5 \text{ cm} \leq h \leq 22.4 \text{ cm}$) were larger in size. The description of the test results is also more complete. The test results with *rectangular* piers showed that for an average ice temperature of -4.3°C and $h \approx 12 \text{ cm}$, the ice plate failed in buckling for $b/h > 2$. For *circular* piers, instead of the b -value, the pier diameter D was used. From the presented test results it follows that for cold ice (-4.0°C) the ice plate buckled for $D/h > 4.85$ whereas for warmer ice (-2.7°C) buckling occurred for larger D/h ratios, namely $D/h > 8.2$. In some of these tests, it was observed that, as the load increased, cracks occurred as indicated in Figure 4. For rectangular piles (Fig. 4(a)) the cracks formed a funnel-shaped plate which buckled in the wider part.

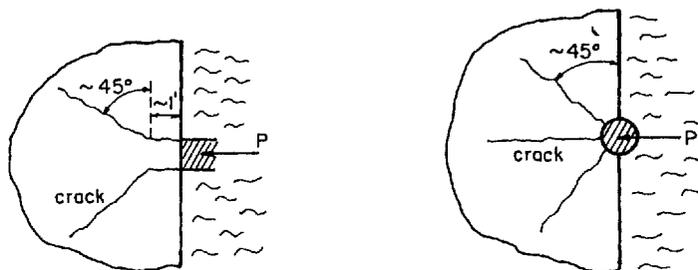


Fig. 4. Form of cracking observed when an ice cover presses against a pier: (a) rectangular pier (left), (b) circular pier (right).

According to the above test results by Afanas'yev and others (1971) and Nevel and others (1972), for rectangular piers a thin ice plate may fail in buckling for $b/h > 2$, and for circular piers for $D/h > 5$, ranges of practical interest. Thus, for a variety of isolated structures and relatively thin ice covers, the largest ice forces appear to be related to the buckling load of the ice cover and not to the local "crushing strength".

The results of the reviewed analyses for the determination of the largest force an ice cover may exert on a *long vertical wall* (for example a long dam), presented by Shadrin and Panfilov (1962) and Afanas'yev and others (1971), are in agreement with observations in the field, regarding occurrence of buckling failures. Therefore, it is reasonable to expect that if the material parameters are chosen properly, the largest ice force for problems of this type can be estimated using these analyses, namely, Equations (4) to (7).

For the determination of the largest ice forces on *isolated structures* based on the buckling failure mode, only the analysis of Kheysin (1961) seems to be available. As shown above, this analysis is based on questionable assumptions. According to Kheysin, the failure load is the load at which the plate cracks, since buckling takes place *immediately* after the cracks form, which is not substantiated by the observations described by Afanas'yev and others (1971) and by Nevel and others (1972).

The above findings suggest the need for an improved analysis for the determination of the largest forces an ice cover exerts on isolated structures, based on the buckling failure mechanism. Such an analysis is presented in the following section.

3. DETERMINATION OF THE LARGEST ICE FORCE ON AN ISOLATED STRUCTURE

Preliminaries

According to the laboratory tests by Nevel and others (1972), when a pier with a flat surface presses against the free edge of a floating ice plate (or vice versa) then at a certain intensity of load radial cracks emanate from the loaded region, as indicated in Figure 4(a), and a further increase of the load leads to more radial cracks and finally to the buckling of the cracked region. A similar observation was also made by Afanas'yev and others (1971). These observations suggest that the largest force a relatively thin ice cover may exert on an isolated structure may be related to the *buckling load of floating wedge-shaped plates*.

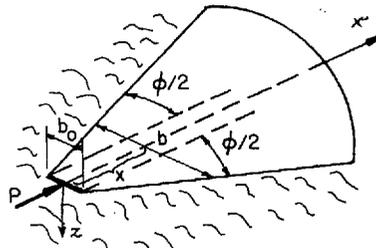


Fig. 5. Geometry of a floating wedge with a force P tending to buckle it.

The buckling analysis of a floating wedge

Consider a floating truncated wedge, as shown in Figure 5. It is assumed that the buckling load P_b may be determined approximately from the eigenvalue problem consisting of the differential equation

$$[EI(x) w'']'' + Pw'' + k(x) w = 0, \quad 0 < x < \infty, \quad (14)$$

the boundary conditions at the contact area

$$\left. \begin{aligned} w(0) &= 0, \\ w''(0) &= 0, \end{aligned} \right\} \quad (15)$$

or

$$\left. \begin{aligned} w(0) &= 0, \\ w'(0) &= 0, \end{aligned} \right\} \quad (16)$$

and regularity conditions at $x = \infty$. In the above formulation, $w(x)$ is a vertical perturbation of the wedge axis at x . Furthermore

$$\left. \begin{aligned} k(x) &= \gamma(b_0 + x\phi^*), \\ EI(x) &= D(b_0 + x\phi^*), \end{aligned} \right\} \quad (17)$$

where

$$D = \frac{Eh^3}{12(1-\nu^2)}, \quad \phi^* = 2 \tan \frac{\phi}{2}, \quad (18)$$

and γ is the specific weight of the liquid base. The term $(1-\nu^2)$ was inserted in the denominator of $EI(x)$ to account for plate action, as done by Papkovich (1962, p. 424-26) in connection with a related problem and discussed by Kerr (1976).

Substituting Equations (17) into Equation (14) we obtain

$$[D(b_0 + x\phi^*) w'''] + Pw'' + \gamma(b_0 + x\phi^*) w = 0, \quad 0 < x < \infty, \quad (19)$$

a fourth-order equation with variable coefficients. For $\phi^* = 0$ (i.e. for $\phi = 0$) and with $b_0 = 1$ the above equation reduces to the equation with constant coefficients used by Kheysin (1961).

Because of the nature of the variable coefficients, no exact closed-form solution is available. A numerical solution can be easily obtained using, for example, a finite-difference scheme, since the buckling mode is restricted to the vicinity of the contact region at $x = 0$. However, the solution obtained will be valid only for the set of parameters assumed.

For the method to determine P_{\max} proposed in the following, an analytical expression which shows the dependence of the buckling load P_b on the various geometrical and material parameters of the floating ice wedge and the width of the structure, is more suitable. Noting that the representation of the buckling of a cracked ice cover by the buckling of a floating wedge is, at best, an approximation, it appears sufficient to solve the wedge formulation approximately. In the following such an analytical solution is obtained first using the Stodola-Vianello method (Hildebrand, 1962, p. 200).

For this purpose Equation (19) is rewritten as

$$w'''' + 4\kappa^4 w = -\frac{\phi^*}{b_0} \left[(xw_1'')'' + \frac{\gamma}{D} xw_1 \right] - \lambda^2 w_1'', \quad (20)$$

where

$$\kappa^4 = \frac{\gamma}{4D}, \quad \lambda^2 = \frac{P}{Db_0}. \quad (21)$$

Assuming for w_1 the expression

$$w_1(x) = A_1 \exp(-\kappa x) \sin \kappa x, \quad (22)$$

and substituting it into Equation (20), we obtain a non-homogeneous ordinary differential equation with constant coefficients for $w(x)$. Its solution, which satisfies the boundary conditions of a simply supported edge, Equations (15), is

$$w(x) = A_1 \exp(-\kappa x) [A \sin(\kappa x) + (\kappa A + B) x \sin(\kappa x) + Bx \cos(\kappa x)], \quad (23)$$

where

$$A = \frac{-\phi^*}{2\kappa b_0}, \quad B = \frac{P}{8D\kappa b_0}. \quad (24)$$

The first eigenvalue is obtained by substituting Equations (24) into the condition (Hildebrand, 1963, p. 216, equation 129)

$$\int_0^\infty w_1^2(x) dx = \int_0^\infty w_1(x) w(x) dx. \tag{25}$$

Performing the necessary integrations and noting that $\lambda^2 = P/(Db_0)$, we obtain the approximate buckling load

$$(P_b)_{ss} = 5.3D\kappa[\kappa b_0 + \phi^*]. \tag{26}$$

If instead of the boundary conditions for a simply supported end the ones for a clamped end at $x = 0$ are used, the above procedure yields

$$(P_b)_c = 8D\kappa[2\kappa b_0 + \phi^*]. \tag{27}$$

Note, that in both cases the respective P_b expression consists of a sum of two terms:

$$(D\kappa^2 b_0) \quad \text{and} \quad (D\kappa\phi^*),$$

each multiplied by a constant coefficient. Thus, for the P_b expressions obtained, the different boundary conditions at $x = 0$ affect only these coefficients.

To check if the above two terms are representative of the buckling solution for the compressed wedge, the wedge problem is next solved again using the energy method.

With the notation used above, the second variation of the total potential energy for the wedge problem shown in Figure 5, is

$$V_2 = \int_0^\infty \left(\frac{EI}{2} w''^2 + \frac{k}{2} w^2 - \frac{P}{2} w'^2 \right) dx, \tag{28}$$

where EI and k are given in Equations (17). Assuming

$$w(x) = A \exp(-\beta x) \sin \beta x, \tag{29}$$

as an approximate shape function, substituting it into V_2 , and performing the necessary integrations, we obtain

$$V_2 = \frac{A^2}{8\beta^2} \left[\left(6Db_0\beta^4 - P\beta^2 + \frac{\gamma b_0}{2} \right) \beta + \phi^* \left(2D\beta^4 + \frac{\gamma}{2} \right) \right]. \tag{30}$$

If β is chosen *a priori* as

$$\beta = \kappa' = \left(\frac{\gamma}{D} \right)^{\frac{1}{4}} = 4^{\frac{1}{4}}\kappa, \tag{31}$$

the value which corresponds to the compressed semi-infinite plate, it follows that $V_2 = V_2(A)$. The condition $\delta V_2 = 0$ then reduces to

$$\frac{\partial V_2}{\partial A} = 0. \tag{32}$$

The above equation yields

$$(6Db_0\kappa'^4 + \gamma b_0/2 - P_b\kappa'^2) \kappa' + \phi^*(2D\kappa'^4 + \gamma/2) = 0, \tag{33}$$

and thus

$$P_b = 6D\kappa' \left[b_0\kappa' \left(1 + \frac{\gamma}{12D\kappa'^4} \right) + \phi^* \left(\frac{1}{3} + \frac{\gamma}{12D\kappa'^4} \right) \right]. \tag{34}$$

Noting that $\gamma/(12D\kappa'^4) = 1/12$ and $\kappa' = 4^{\frac{1}{4}}\kappa$, the expression for P_b reduces to

$$P_b = 3.5D\kappa[3.7b_0\kappa + \phi^*]. \tag{35}$$

If β is not fixed *a priori*, then $V_2 = V_2(A, \beta)$. The condition $\delta V_2 = 0$ then reduces to

$$\frac{\partial V_2}{\partial A} = 0, \quad \frac{\partial V_2}{\partial \beta} = 0. \quad (36)$$

The first equation is identical to Equation (32) and yields Equation (34) with $\kappa' = \beta$. Namely,

$$P_b = 6D\beta \left[b_0\beta \left(1 + \frac{\gamma}{12D\beta^4} \right) + \phi^* \left(\frac{1}{3} + \frac{\gamma}{12D\beta^4} \right) \right]. \quad (37)$$

The second equation in (36) yields, after eliminating P_b by using Equation (37),

$$\beta(4Db_0\beta^4 - b_0\gamma/3) + \phi^*(2D\beta^4/3 - \gamma/2) = 0. \quad (38)$$

This is the equation for the determination of β which appears in Equation (37).

Comparing Equation (35) with Equations (26) and (27) it follows that it is of the same functional form as these two equations except for the constant coefficients.

In order to compare the approximate expressions obtained for the buckling load P_b with each other, they were evaluated for $E = 30\,000 \text{ kg cm}^2$, $\nu = 0.34$, $\gamma = 0.001 \text{ kg/cm}^2$, $b_0 = 20 \text{ cm}$, and $h = 10 \text{ cm}$. The results are shown in Figure 6.

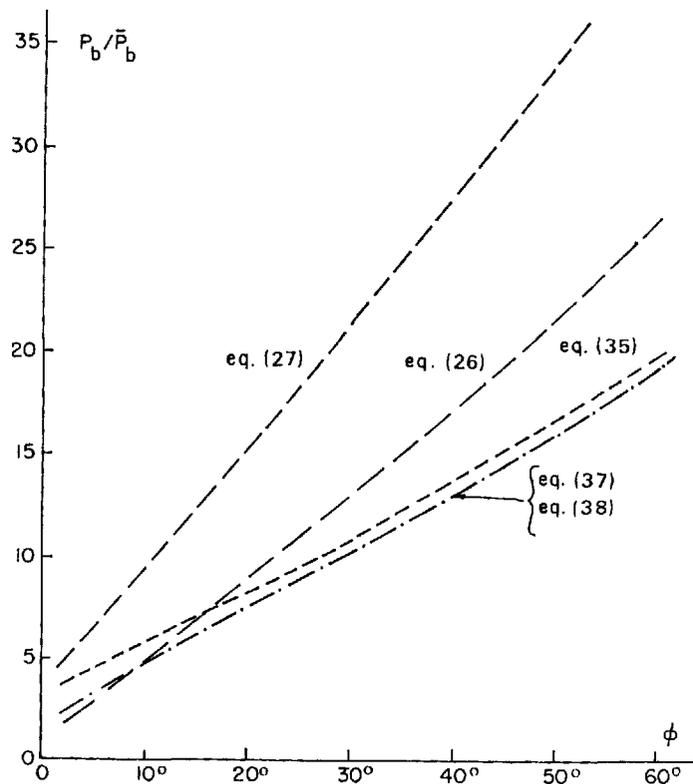


Fig. 6. Critical buckling loads for a floating wedge. $\bar{P}_b = 2b_0(\gamma D)^{1/2}$. The numbers by each curve indicate equation numbers in the paper.

* Note that this equation may be obtained directly from the condition $\partial P_b / \partial \beta = 0$.

Proposed method to determine P_{\max}

When an actual ice cover presses against a bridge pier or the leg of a drilling platform, the boundary conditions at the contact area between ice and structure will fluctuate, depending on a number of factors. Also, the value ϕ is usually not a measurable entity since according to test observations, at the onset of buckling, the ice cover exhibits many radial (and some non-radial) cracks.

It is therefore proposed to express the largest force an ice cover exerts on an isolated structure (when buckling dominates) by the expression

$$P_{\max} = n'D\kappa(\kappa b_0 + m'), \quad (39)$$

where according to Equations (18) and (21),

$$D = \frac{Eh^3}{12(1-\nu^2)}, \quad \kappa = \left(\frac{\gamma}{4D}\right)^{\frac{1}{2}},$$

and n' and m' are parameters to be determined from field tests.

Since the effective values of E and ν will depend upon the nature of the pre-cracked plate, the term $E/[12(1-\nu^2)] = E^*$ should be determined from the same test data. Thus, the proposed expression for P_{\max} is

$$P_{\max} = n(\gamma E^* h^3)^{\frac{1}{2}} \left[\left(\frac{\gamma}{E^* h^3}\right)^{\frac{1}{2}} b_0 + m \right], \quad (40)$$

where γ is the specific weight of the liquid base, h is the ice-cover thickness, and b_0 is the width of pier. The parameters m , n , and E^* are to be determined from appropriate test results.

It is anticipated that m , n , and E^* will vary for different categories of problems and different ice temperatures (in particular in the vicinity of the melting temperature). This point, as well as the validity of Equation (40) to represent P_{\max} (when the buckling failure mode predominates), can only be established by comparing it with the relevant results of laboratory and field tests.

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