## A GENERALIZATION OF SONINE'S FIRST FINITE INTEGRAL

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In this note I show that

$$
\begin{align*}
J_{\mu+v+2 n+1}(z)= & \frac{z^{v+1} \Gamma(\mu+n+1)}{2^{\nu} \Gamma(\mu+1) \Gamma(v+n+1)} \\
& \times \int_{0}^{\frac{1 \pi}{2}} J_{\mu}(z \sin \theta)_{2} F_{1}\left(-n, \mu+v+n+1 ; \mu+1 ; \sin ^{2} \theta\right) \sin ^{\mu+1} \theta \cos ^{2 v+1} \theta d \theta \tag{1}
\end{align*}
$$

where $J$ denotes the Bessel function of the first kind of the orders and arguments indicated, $n=0,1,2,3, \ldots$ and the real parts of both $\mu$ and $v$ exceed -1 . This is a generalization of Sonine's first finite integral [1, p. 373] to which it reduces in the special case $n=0$.

I start with the Weber-Schafheitlin integral

$$
\begin{equation*}
I(\mu, v, n, r)=\int_{0}^{\infty} z^{-v} J_{\mu+v+2 n+1}(z) J_{\mu}(r z) d z \tag{2}
\end{equation*}
$$

with the conditions on $n, \mu$ and $v$ as given above. The integral is convergent and [1, p. 401] its value is given by

$$
I(\mu, v, n, r)=\left\{\begin{array}{l}
\frac{r^{\mu} \Gamma(\mu+n+1)}{2^{\nu} \Gamma(\mu+1) \Gamma(v+n+1)}{ }_{2} F_{1}\left(\mu+n+1, n-v ; \mu+1 ; r^{2}\right) \quad(0<r<1)  \tag{3}\\
0(1<r<\infty)
\end{array}\right.
$$

the integral vanishing when $r>1$ because of a factor $\Gamma(-n)$ in the denominator of the term multiplying the hypergeometric function. Applying Hankel's inversion formula to (2), we obtain

$$
z^{-v-1} J_{\mu+v+2 n+1}(z)=\int_{0}^{\infty} r I(\mu, v, n, r) J_{\mu}(z r) d r
$$

and substitution from (3) gives
$z^{-v-1} J_{\mu+v+2 n+1}(z)=\frac{\Gamma(\mu+n+1)}{2^{v} \Gamma(\mu+1) \Gamma(v+n+1)} \int_{0}^{1} r^{\mu+1}{ }_{2} F_{1}\left(\mu+n+1,-n-v ; \mu+1 ; r^{2}\right) J_{\mu}(z r) d r$.

Using the well-known transformation formula [2, p. 8],

$$
{ }_{2} F_{1}\left(\mu+n+1,-n-v ; \mu+1 ; r^{2}\right)=\left(1-r^{2}\right)_{2}^{v} F_{1}\left(-n, \mu+v+n+1 ; \mu+1 ; r^{2}\right),
$$

and writing $r=\sin \theta$, we obtain the required result (1) directly from (4).

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As well as Sonine's first finite integral, there are some further interesting special cases of the general formula (1). Thus the two modifications of Bessel's integral [1, pp. 20, 21],

$$
\begin{aligned}
J_{2 n}(z) & =\frac{2}{\pi} \int_{0}^{\frac{1}{2} \pi} \cos 2 n \theta \cos (z \sin \theta) d \theta, \\
J_{2 n+1}(z) & =\frac{2}{\pi} \int_{0}^{\frac{1}{2 \pi}} \sin (2 n+1) \theta \sin (z \sin \theta) d \theta,
\end{aligned}
$$

are obtained by writing $v=-\frac{1}{2}$ and $\mu=\mp \frac{1}{2}$ respectively in (1). Again, taking $v=0, \mu=-\frac{1}{2}$ in (1), expressing the hypergeometric function in terms of a Legendre polynomial [2, p. 50], making a few reductions and writing $x=\sin \theta$, we have

$$
J_{2 n+\frac{t}{}}(z)=(-1)^{n} /\left(\frac{2 z}{\pi}\right) \int_{0}^{1} P_{2 n}(x) \cos z x d x
$$

and this formula gives, in effect, the so-called even Legendre transform of $\cos z x$ [3, p. 97]. In a similar way, substitution of $v=0, \mu=\frac{1}{2}$ in (1) leads to

$$
J_{2 n+\frac{z}{2}}(z)=(-1)^{n} \sqrt{\left(\frac{2 z}{\pi}\right) \int_{0}^{1} P_{2 n+1}(x) \sin z x d x}
$$

and hence to the odd Legendre transform of $\sin z x$.

## REFERENCES

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