

Research Paper

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Transfer of AlGaIn/GaN RF-devices onto diamond substrates via van der Waals bonding

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Abstract

We present a novel bonding process for gallium nitride-based electronic devices on diamond heat spreaders. In the proposed technology, GaN devices are transferred from silicon (Si) onto single (SCD) and polycrystalline diamond (PCD) substrates by van der Waals bonding. Load-pull measurements on Si and SCD heat spreaders at 3 GHz and 50 V drain bias show comparable power-added-efficiency and output power (P_{out}) levels. A thermal analysis of the hybrids was performed by comparison of $2 \times 1 \text{ mm}^2$ AlGaIn/GaN Schottky diodes on Si, PCD, and SCD, which exhibit a homogeneous field in the channel in contrast to gated transistors. Significantly different currents are observed due to the temperature dependent mobility in the 2DEG channel. These measurements are supported by a 3D thermal finite element analysis, which suggests a large impact of our transfer technique on the thermal resistance of these devices. In summary, we show a promising new GaN-on-diamond technology for future high-power, microwave GaN device applications.

Introduction

The ongoing development of GaN-based power amplifiers for higher P_{out} and higher frequencies opens a growing market for wide-bandgap semiconductor devices [1,2]. Radio-frequency (RF) communication systems would extremely profit from this trend leading to even higher integration densities, smaller system sizes, and overall better system performance.

Due to the lack of large area GaN substrates, GaN devices are mainly fabricated on three different substrates: Silicon carbide (SiC), Si, and sapphire. GaN-on-SiC achieves currently the highest power densities attributed to its largest thermal conductivity. Despite extremely high radio-frequency (RF) power densities of 40 W mm^{-1} [2], thermal management is still the primary limiting factor. Diamond as substrate would provide four times higher thermal conductivity ($\kappa_{\text{diamond}} > 2000 \text{ W m}^{-1} \text{ K}^{-1}$) than SiC ($\kappa_{\text{SiC}} = 490 \text{ W m}^{-1} \text{ K}^{-1}$), which potentially allows even higher performance (see Fig. 1).

Heterointegration of AlGaIn/GaN layers on diamond as new substrate was realized by growth and by bonding. The growth of AlGaIn/GaN on diamond, on both polycrystalline diamond (PCD) [5,6] and SCD [7–9] usually suffers from a poor electronic quality of the AlGaIn/GaN layers, which require careful interface control and stress engineering during the epitaxial growth and the final cool down. The reverse approach, the growth of diamond on AlGaIn/GaN was demonstrated both on the top [10] or the bottom [11] of the GaN layer and already competes with the GaN-on-SiC technology [12,13]. The most developed GaN-on-diamond process grows PCD on the nitrogen-face GaN buffer layer with a thin interfacial stabilization layer [14]. Extensive reliability testing [15], scaling to 4-inch wafer size [11], and RF-performance of 7.9 W mm^{-1} (P_{out}) at power-added-efficiency (PAE) of 46% [16] (10 GHz, 40 V bias) demonstrated already its efficiency for next-generation high-power semiconductor devices. Conversely, the heterointegration of GaN on diamond by bonding offers more flexibility since both, the AlGaIn/GaN layer and diamond are grown, and optimized in separate processes. Early bonding experiments at high temperatures of about 800°C [17] were limited to small areas ($< 2 \times 2 \text{ cm}^2$). More recently, GaN devices adhesively bonded at low temperatures [18] achieved a current record for GaN-on-diamond with P_{out} of 11 W mm^{-1} and 51% PAE (10 GHz, 40 V bias) [19].

Despite these promising results, further improvement of GaN-on-diamond technology remains challenging. A bottleneck in thermal performance is posed by defect-rich nucleation layers at the buffer/substrate interface. Moreover, the substrate transfer onto diamond introduces thermally poor stabilization [20], or adhesion layers [18], and the nucleation layer of grown diamond contains additional voids and graphitic compounds within several nm thickness [21]. The material specific thermal boundary resistance (TBR) as calculated from diffuse

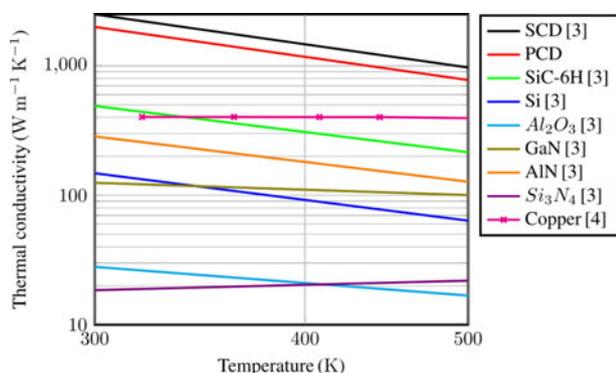


Fig. 1. Temperature dependence of thermal conductivities for different substrates (SCD, PCD, Si, SiC-6H, and Al_2O_3) and AlGaIn/GaN device specific materials (all from [3]) and copper [4].

mismatch models is the theoretical minimum resistance of a GaN/diamond interface, but experimental measurements on GaN-on-diamond devices revealed much higher resistances which was explained by these poor interlayers [22].

Our approach focuses on the combination of GaN and diamond by van der Waals (vdW) bonding technology. Capillary forces pull both surfaces into close contact, and further annealing creates a strong bond between both surfaces. This approach is similar to the technique first employed for GaAs thin films [23]. Here, GaAs electronics showed improved performance on diamond substrates compared with Si [24]. However, the impact of diamond as new substrate depends on the application. Due to material limitations, AlGaAs/GaAs RF-transistors operate at relatively low power where the impact of diamond might be less pronounced. This paper presents AlGaIn/GaN RF-transistors and large $2\text{ mm} \times 1\text{ mm}$ AlGaIn/GaN Schottky diodes transferred from silicon onto diamond by vdW bonding. In contrast to gated transistors, lateral diodes have a homogeneous electric field in the channel, and thus, uniform heat generation. This allows simplified electric modeling of the temperature dependent electron mobility in order to accurately extract channel temperatures. Based on these results, the TBR was evaluated by thermalfinite element analysis (FEA) simulations, which reveal the excellent thermal performance of SCD compared with Si and PCD.

This paper is an extended version of work published in [25] and includes the detailed thermal analysis of our diodes.

Fabrication of GaN-on-Diamond materials and devices

Diamond substrate fabrication

PCD was grown by microwave chemical vapor deposition in hydrogen/ methane plasma on Si substrates. Diamond nanoparticles of 7 nm mean size initialize the nucleation to grow $350\mu\text{m}$ thick PCD wafers. For the application as a heat spreader, the thermal quality of this starting layer is decisive since this nucleation side is bonded onto the GaN buffer surface. The detailed seeding procedure is described in [26]. After growth, the rough growth side was polished to a final diamond thickness of $300\mu\text{m}$. The nucleation side was shortly polished to remove the first several nm of nanocrystalline diamond. Extended mechanical polishing roughens the nucleation surface due to the anisotropy of the mechanical polishing on polycrystalline material [27]. Mechanical polishing of the growth side of PCD cannot achieve the surface quality required for bonding.

In addition to PCD, SCD (Element Six, $8\text{ mm} \times 8\text{ mm}$) was bonded onto GaN buffer layers, which presents the current reference for our technology. Since direct bonding requires extremely flat surfaces, small surface roughness is necessary to achieve good bonding results. SCD is polished to $\text{RMS} < 0.5\text{ nm}$, which makes it ideal for bonding. Besides, there is no thermally poor nucleation layer but a high-quality diamond directly at the GaN buffer interface. In this work, for the bonding, $8\text{ mm} \times 8\text{ mm}$ SCD and PCD substrates with a diameter of 9 mm were employed.

GaN epitaxy and device fabrication

We demonstrate the transfer of two topologically very different electronic GaN devices. First, RF-transistors with a gate-width of $2 \times 300\mu\text{m}$ quantify the heat-spreading at high local power densities. Second, large $2\text{ mm} \times 1\text{ mm}$ diodes for DC-switching applications demonstrate the homogeneity of our transfer technology on a larger area and are used for the electric and thermal analysis.

Both device types were fabricated on AlGaIn/GaN heterostructures grown by metal-organic chemical vapor deposition on high-resistivity Si substrates. The heterostructure consists of an aluminum nitride (AlN) nucleation layer, a GaN buffer, an AlGaIn barrier and a GaN cap. AlGaIn transition layers are employed to accommodate thermal and lattice mismatch. The device processing is based on standard III-V equipment with lateral device definition by optical stepper lithography. Every process step is optimized for homogeneity and reproducibility to achieve a high yield of GaN-on-Si devices. Ref. [28] presents a detailed description of our GaN epitaxy on Si and subsequent device processing.

GaN transfer and bonding

The combination of GaN and diamond is established by vdW forces. This approach offers the flexibility of a bonding technology combined with a small separation between diamond and the hot transistor region. Bonding temperatures below 300°C minimize thermal stress during bonding and can be performed on a large area without damaging the GaN devices. The bonded layers demonstrated strong vdW forces able to withstand mechanical stress applied by razor blades and extended thermal stress at 250°C .

Fig. 2 shows a summary of the different process steps in GaN-on-diamond technology. After device fabrication on GaN-on-Si wafers and dicing into $12\text{ mm} \times 12\text{ mm}$ reticles (2a), the device layer is transferred onto a sapphire carrier (2b). The Si substrate is wet-chemically removed (2c) by an acetic mixture of hydrofluoric acid (49 wt% HF) and nitric acid (69wt% HNO_3), as described in [29–31]. This etching is very isotropic on Si and highly selective to the AlN nucleation layer. The AlN nucleation layer is bonded onto a polished PCD or SCD substrate (2d) by vdW bonding as described in [23]. After a solid bond is established, the sapphire carrier is removed (2e) at elevated temperatures of about 200°C . Finally, the GaN-on-diamond chip is soldered to a copper heat sink with gold/ indium adhesion layers (2f).

Electrical measurements on RF-transistors and DC-diodes

3 GHz, 50 V load-pull measurements

3 GHz (50 V bias) load-pull measurements of RF-transistors with $2 \times 300\mu\text{m}$ gate-width and 500 nm gate-length have been compared on silicon and on bonded SCD heat spreaders.

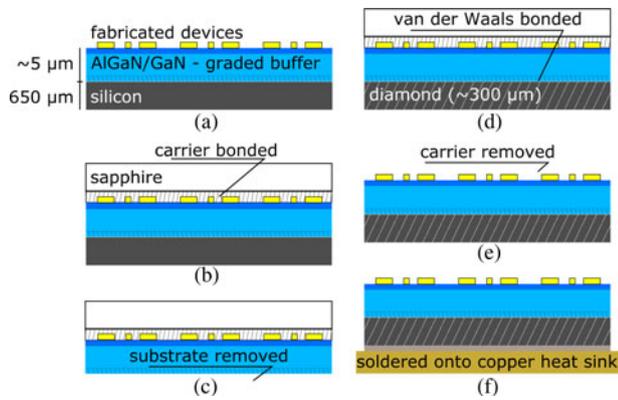


Fig. 2. GaN-on-diamond fabrication steps. After device fabrication on GaN-on-Si wafers (a) the transfer starts with adhesive bonding of the device layer onto a carrier wafer (b). In (c) the Si substrate is etched and the resulting AlN nucleation layer bonded (d) onto a polished diamond substrate (PCD or SCD). The carrier is removed (e) and the GaN-diamond hybrid soldered onto a copper heat sink (f).

First, the output impedance was tuned to achieve maximum P_{out} ($P_{out-max}$) and PAE (PAE_{max}). Table 1 summarizes the best results achieved for GaN-on-Si transistors and its transferred GaN-on-SCD equivalent in continuous wave (cw) measurements. PAE_{max} improved from 50.6% (on Si) to 54.2% (on SCD) at similar P_{out} of 5.65 W mm^{-1} and 5.39 W mm^{-1} , respectively. $P_{out-max}$ was comparable on Si (6.79 W mm^{-1}) and SCD (6.63 W mm^{-1}), whereas PAE increased from 43.6% to 46.5%. However, the measurement of GaN-on-Si transistors on the wafer and GaN-on-SCD on isolated chips requires additional detailed characterizations to guarantee objective comparisons.

As summarized in Table 2 pulsed measurements on the GaN-on-SCD transistor revealed potentially much higher PAE and P_{out} . In pulsed measurements, PAE_{max} increased from 54.2% to 59.1% at P_{out} of 5.39 W mm^{-1} and 5.91 W mm^{-1} , whereas $P_{out-max}$ increased from 6.63 W mm^{-1} to 7.44 W mm^{-1} at PAE of 46.5% and 53.5%, respectively.

With output impedance fixed at maximum PAEs, load-pull measurements between a GaN-on-SCD transistor and the original GaN-on-Si transistors were compared. Fig. 3 shows PAE (blue), gain (green), and P_{out} (red) for different input power levels (P_{in}) on SCD (with symbols) and Si (no symbols). For $P_{in} < 15 \text{ dBm}$ GaN-on-SCD achieved small improvements in PAE, gain and P_{out} compared with all GaN-on-Si equivalent transistors. For $P_{in} > 15 \text{ dBm}$ P_{out} and gain curves approach the values of GaN-on-Si.

The underlying complexity of the large sequence of critical fabrication steps to final GaN-on-SCD devices suggests that further improvements are quite probable. These first experiments successfully demonstrated that RF-transistors can be transferred from Si onto SCD without any observable performance degradation. Small improvements in PAE were recognized at comparable P_{out} levels.

Table 1. Maximum PAE and P_{out} and corresponding P_{out} and PAE for the $2 \times 300 \mu\text{m}$ gate-width transistor on SCD and Si

	GaN-on-SCD (cw)		GaN-on-Si (cw)	
	PAE (%)	P_{out} (dBm)	PAE (%)	P_{out} (dBm)
PAE_{max}	54.2	35.1	50.6	35.3
$P_{out-max}$	46.5	36.0	43.6	36.1

Table 2. Maximum PAE and P_{out} and corresponding P_{out} and PAE for the $2 \times 300 \mu\text{m}$ gate-width transistor on SCD - cw and pulsed

	cw measurement		pulsed measurement	
	PAE (%)	P_{out} (dBm)	PAE (%)	P_{out} (dBm)
PAE_{max}	54.2	35.1	59.1	35.5
$P_{out-max}$	46.5	36.0	53.5	36.5

DC-Measurements on $2 \times 1 \text{ mm}^2$ diodes

To characterize the GaN-on-diamond technology on larger devices, $2 \text{ mm} \times 1 \text{ mm}$ AlGaIn/GaN-diodes were transferred from Si onto PCD and SCD substrates. In Fig. 4 differential interference contrast (DIC) micrographs show the different appearance of two adjacent AlGaIn/GaN-diodes on Si (Fig. 4a), PCD (Fig. 4b) and SCD (Fig. 4c). Si appears dark, whereas PCD and SCD diffusively scatter the light. PCD is easily distinguished by the visible grains in the background. The diodes ($17 \mu\text{m}$ channel length) have a channel width of 48 mm distributed over 40 fingers. The specific structure of this Schottky diode was reported in [32].

In Fig. 5 we increased the bias voltage (V_{Bias}) from 0 to 5 V on Si (red), PCD (blue) and SCD (black). Each bias point was kept for 10 s to apply a large heat load, whereas the baseplate temperature below the copper block regulated to constant 40°C . In our following electrical and thermal analysis, the 5 V power levels of 32, 43 and 52 W for Si, PCD, and SCD, respectively, are compared. For SCD an increase to 10 V results in a power of 130 W (not shown), which would probably destroy the diode on Si.

Thermal analysis of diodes on SCD, PCD, and si

Temperatures from electrical analysis

In the section 3 electrical measurements, RF-transistors and diodes were shown. Whereas RF-transistors were only transferred onto SCD, our diodes were bonded onto both SCD and PCD. In contrast

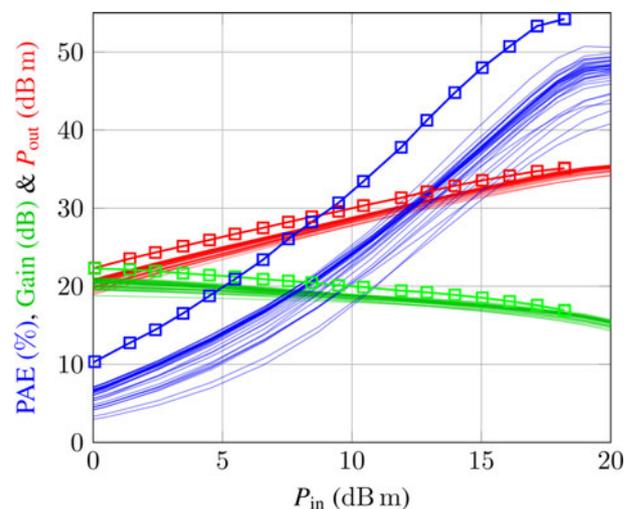


Fig. 3. Load-pull measurement results of PAE (blue), gain (green) and P_{out} (red) on RF-transistors with gate-widths of $2 \times 300 \mu\text{m}$ operated at 3 GHz, 50 V DC bias and constant load (load at maximum PAE). The transistor transferred onto SCD (with symbols) shows higher PAE than any GaN-on-Si transistor (other curves) at any input power level. At low input power, gain and P_{out} is also higher on SCD but approximates the Si devices at higher input power.

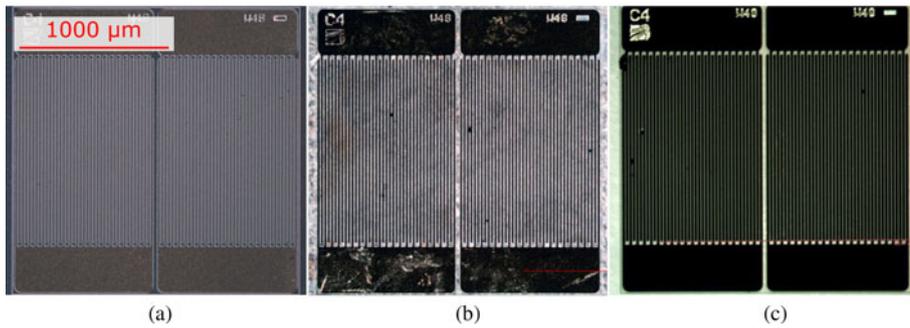


Fig. 4. Optical micrographs (DIC) of 2 mm x 1 mm diodes bonded onto Si (a), PCD, (b) and SCD (c). All diodes are fully operable without any cracks of the GaN device layer.

to our transfer of RF-transistors, the bonding process with diodes is well-engineered. The diodes are optically homogeneous bonded and demonstrate a large increase in diode currents on both SCD and PCD. The simplified structure of diodes compared with transistors facilitates our electrical and thermal analysis and allows the quantification of channel temperatures present in our diodes.

The slope of all diode curves (grown on Si and transferred onto SCD/PCD) in Fig. 5 is identical for small bias voltages where almost no heat is present. This slope $\partial U/\partial I$ is related to the on-state resistance

$$R_{on} = \partial U/\partial I = \frac{l}{w} R_{ch} + R_C = \frac{l}{w} \frac{1}{q\mu n_s} + R_C, \quad (1)$$

where l and w are length and width of the channel, respectively. R_C includes contact resistances and R_{ch} is the channel resistance between the contacts. R_{ch} resolves to the inverse product of the elementary charge q , the electron mobility μ and the carrier density n_s . The identical on-state resistance at low bias on all substrates shows that our transfer process does not change the AlGaIn/GaN structure and that different currents at higher bias are predominantly caused by a temperature effect. The current saturation at higher bias is based on a channel pinch-off from the increasing channel potential [33]. Saturation was observed on SCD around 10 V, whereas on silicon the channel saturates already at 5 V with currents of 13 and 6.5 A, respectively. Since electron density is almost constant with temperature [34], improved transport is mainly based on higher mobility.

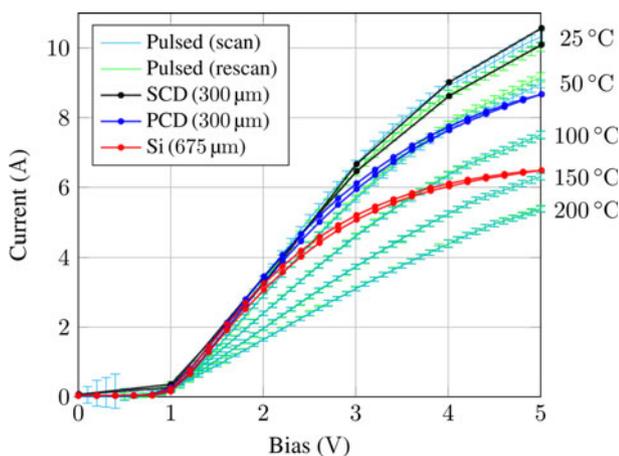


Fig. 5. Currents in GaN-diodes on Si (red), PCD (blue), and SCD (black) and pulsed measurements (1 ms) on identical diodes of a similar wafer for increasing/ decreasing bias (cyan/green) and different baseplate temperatures (25°C-200°C). The standard deviation results from all diodes on the wafer.

The temperature dependence of electron mobility in AlGaIn/GaN devices (> 300 K) is modeled with

$$\mu_0 = \mu_{300} \left(\frac{T}{300} \right)^{-\alpha}, \quad (2)$$

where μ_{300} is the low-field mobility at 300 K and α a positive exponent. Generally, the bulk mobility in GaN is dominated by polar optical phonon scattering giving $\alpha_{bulk} = 1.5$. In the two-dimensional electron gas of AlGaIn/GaN heterostructures, however, the exponent is higher and depends on the carrier density (see compiled data in Fig. 6). Moreover, at large electrical fields > 50 kVcm⁻¹ [35], the mobility-field dependence cannot be neglected. However, the average field in our diodes with $l = 17\mu\text{m}$ at 5 V bias results in electric fields < 3 kVcm⁻¹, thus, mobility can be assumed to be independent of the electrical field. AlGaIn/GaN transistors have highly inhomogeneous electric fields along the channel, which complicates the analysis of the temperature dependent mobility. In contrast, the employed lateral Schottky diodes with only two contacts exhibit an almost constant electric field. Therefore, changes in the current can be assigned directly to the temperature dependence of the mobility, i.e. larger currents result primarily from lower channel temperatures.

The channel temperatures were quantified by comparison of pulsed (1 ms) measurements on identical Schottky diodes on a silicon substrate (on-wafer) with the I - V characteristics of the diodes

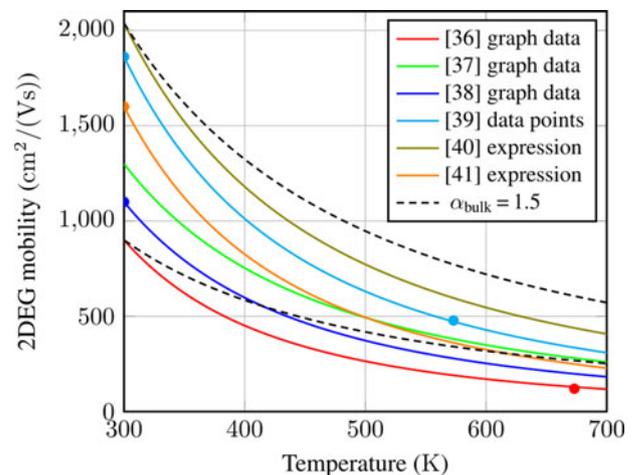


Fig. 6. Experimentally determined AlGaIn/GaN Hall mobilities from different groups [36-41]. If distinct measurement points (marked as ‘data points’) or analytical expressions (‘expression’) were not mentioned, equation (4.2) together with points read from plotted graphs were used (‘graph data’). At 430 K (157°C) mobilities are 61%-83% larger than at 330 K (57°C). As comparison, the bulk mobility (black, dashed) decreases by only 45%.

Table 3. Thermal FEA parameter values of joule heat (P_{sv}), TBR and channel temperature (T_{ch}) for AlGaIn/GaN Schottky diodes on SCD, PCD, and Si. Bold values were calculated through FEA. For SCD and Si, a TBR of $10 \text{ m}^2 \text{ K/GW}$ [22] was assumed. The TBR of PCD is estimated from the temperature difference to SCD in our pulsed electrical measurements

	P_{sv} (W/mm ²)	TBR (m ² K/GW)	T_{ch} (°C)
SCD (300µm)	65	10	52
PCD (300µm)	53	500	74
Si (675µm)	40	10	160
Si (100µm)	65	10	101

on SCD, PCD, and Si (Fig. 5). Within the employed measurement system for these high-voltage diodes, charging effects require comparably long pulse lengths to prevent any electrical measurement errors. The thermal error due to electrical self-heating was corrected through our subsequent thermal simulation.

For pulsed measurements the baseplate and approximately the whole wafer was kept at constant temperatures (25°C–200°C) and it is assumed that the short load generates no significant heat. At low bias, the 25°C curves perfectly overlap with all measured diodes (SCD, PCD, and Si). At larger bias, the currents on PCD and Si reduce due to decreased mobility. In contrast, the diode on SCD follows the 25°C line, thus the generated heat is efficiently removed from the channel. It should be noted, however, that the limited pulse length generates some heat in the channel adding some temperature error. In first approximation this 1 ms pulse error adds some defined temperature offset depending on the bias, but not on the baseplate temperature. The relative error in this analysis therefore shrinks with higher baseplate temperatures. The measurements of our diodes on SCD, PCD, and Si were performed at baseplate temperatures of 40°C, it is therefore reasonable to assume that the channel temperature in our 25°C pulsed curve at 5 V bias equals at least 40°C. The 5 V bias temperatures on PCD and Si are 50°C and 150°C, respectively. To verify these assumptions, the temperature dependence of the mobility in Fig. 6 was taken into account. As shown above, the saturation current on SCD is 100% higher in comparison with the diode on a silicon substrate. The corresponding 100% change of the mobility can easily be explained by a change of the temperature 423 K (channel temperature on silicon as extracted from Fig. 5) to about 315 K, i.e. about 42°C in agreement with the similar discussion above. If the channel temperature on SCD would be significantly higher, Fig. 6 would give unrealistic high channel temperatures on silicon substrates. The extracted channel temperature is the base for the following thermal simulations.

3D thermal simulation

3D FEA (COMSOL) simulates the channel temperatures generated by electrical Joule heat in the diode's channel region. In literature, thermal simulations are widely employed and experimentally verified by Raman thermography [42,43] and thermorefectance [44–46]. As in our electrical analysis, the simpler structure of our large diodes compared with complex RF-transistors allows an accurate analysis based on relatively few assumptions. By changing the TBR between the GaN buffer and the substrate, temperatures extracted from our electrical analysis are reproduced and the required TBRs estimated.

The thermal resistance within the buffer/substrate region is affected by several layers. The accurate thermal modeling of each layer is complex since material quality is gradually improving in growth direction. The close distance between channel and substrate as well as material specific interfacial resistances complicate the analysis and common diffraction limited optical measurement techniques cannot resolve the several nm thin interlayers. It is, therefore, accepted to introduce a transition layer with an effective TBR, which accounts for all nonidealities at the buffer/substrate interface [47]. In GaN-on-Si devices such transition layer adds an effective TBR of $10 \text{ m}^2 \text{ K/GW}$ [22] which we introduced as 100 nm thin TBR layer of $\kappa_{\text{thermal}} = 10 \text{ W m}^{-1} \text{ K}^{-1}$ (COMSOL 'thin-film' boundary condition).

Symmetries were used to reduce the device geometry to one quarter of our multi-finger diode, which consisted of a 22 nm AlGaIn cap layer, 4.5µm GaN buffer and the substrate (675µm silicon or 300µm PCD/SCD) soldered onto a large copper block. The copper block base temperature was set to constant 40°C in accordance with our electrical measurements. The measured electrical power at 5 V bias (cp. Fig. 5) was converted to heat flux applied to the AlGaIn/GaN channel region. A heat load of 65, 53, 40 W/mm² was allocated to the 17µm long and 1.2 mm wide channel region for SCD, PCD, and Si, respectively. The thermal conductivity of the 100 nm thin transition layer was varied between $10 \text{ W m}^{-1} \text{ K}^{-1}$ and $0.2 \text{ W m}^{-1} \text{ K}^{-1}$ to reproduce the channel temperatures from our electrical analysis on devices with different substrates.

Fig. 7 shows the temperature distribution of 20 diode fingers on SCD (7a), PCD (7b), and Si (7c, d) with respective heat load at 5 V bias. Fig. 7c shows a hypothetical diode on 100µm thin Si to compare a thermally optimized Si chip against SCD (heat-flux therefore adapted from SCD). On 675µm Si our assumptions reveal channel temperatures of 160°C, which fairly matches the 150°C as deduced from our pulsed current measurements. As discussed, the minimum channel temperature of our diodes on SCD and PCD must be larger than the 40°C baseplate temperature. Also, we expect TBRs larger than on Si since all defect rich boundary layers are transferred with the devices onto our diamond substrates and the vdW bond is expected to add some additional thermal resistance compared with the grown interface on Si. For SCD, even the optimum TBR conductivity of $10 \text{ W m}^{-1} \text{ K}^{-1}$ (equals TBR_{Si}) still results in temperatures of 52°C, which is slightly higher than the ~40°C expected from our electrical analysis. We therefore assume that the transfer onto SCD does not add any significant thermal resistance compared with the already present GaN transition layer resistance.

The difference between the SCD temperature extracted from electrical measurements and the thermal simulation suggests a similar offset for the PCD temperature. For SCD the simulated temperature was ~25°C higher, which is explained by the neglected limited current pulse. We, therefore, assume that the actual temperature on PCD should be also 25°C higher, i.e. 75°C. In our thermal simulation a temperature of 74°C is simulated for a low thermal conductivity of $0.2 \text{ W m}^{-1} \text{ K}^{-1}$ for the transition layer, which is 50x larger than for SCD and Si (both $10 \text{ W m}^{-1} \text{ K}^{-1}$). The large TBR on PCD is explained from the thermally poor diamond nucleation layer. A hypothetical 100µm thick Si substrate (Fig. 7c) as compared with SCD shows a temperature difference of ~50°C, which clearly demonstrates the outstanding performance of our GaN devices on SCD (cp. Tab. 3).

A 100 nm TBR layer of $\kappa_{\text{thermal}} = 10 \text{ W m}^{-1} \text{ K}^{-1}$ equals an interfacial TBR of $10 \text{ m}^2 \text{ K/GW}$ (TBR = $100 \text{ nm} / \kappa_{\text{thermal}}$ [47]),

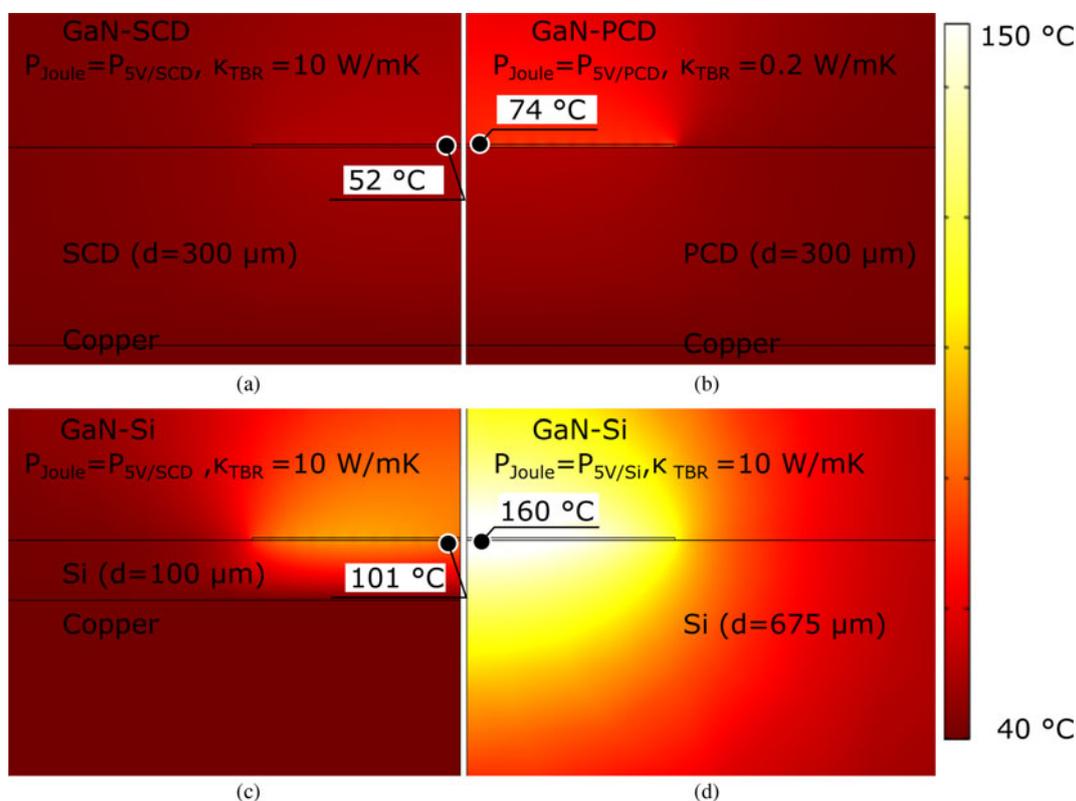


Fig. 7. 3D thermal simulations (COMSOL) visualize temperature distributions in our $2\text{ mm} \times 1\text{ mm}$ diodes - shown is only the channel cross section for clarity. (a), (b), and (d) show our electrically measured structures with adapted κ_{TBR} to account for measured current differences on SCD (a), PCD (b), and Si (d). (c) is a hypothetical $100\mu\text{m}$ thin Si substrate with heat load from SCD (a). The direct comparison shows that even thinned Si is easily outperformed by SCD.

which is in range of commonly measured TBRs on Si and SiC [48]. The scatter in compiled values, and the uncertainty in individual measurements is considerably high. For GaN-on-diamond fabricated by Element Six, TBRs of $18\text{ m}^2\text{K/GW}$ – $27\text{ m}^2\text{K/GW}$ at 50 nm to 324 nm thickness are reported [22], which match the TBR of Si. In their process, the negative impact of the diamond nucleation layer might be compensated by the removal of low-quality GaN nucleation layers before diamond growth, whereas our process still includes these thermally poor GaN buffer layers. In contrast, our transfer process on SCD substrates does not impose an additional diamond nucleation layer but high-quality diamond directly at the buffer/substrate interface.

Summary

Our low-temperature GaN-on-diamond transfer process with its seamless integration into GaN-on-Si device fabrication offers significant advantages in flexibility over diamond growth techniques on GaN. Pulsed load-pull measurements achieved $7.43\text{ W mm}^{-1} P_{\text{out}}$ with 53.5% PAE (3 GHz, 50 V bias) in these first results.

We quantified the cooling performance of hybrids fabricated by our transfer technique with simpler, large diode structures. The electrical analysis based on temperature dependent mobility data (cp. Fig. 6) and pulsed (1 ms) measurements (cp. Fig. 5) consistently predict channel temperature differences of $>100\text{ K}$ and an absolute temperature of $\sim 150^\circ\text{C}$ on Si at 5 V bias. 3D thermal simulations with an effective TBR_{Si} of $10\text{ m}^2\text{ K/GW}$ [48] result in a similar channel temperature of 160°C . The larger temperature in our thermal simulation is explained from the relative difference between the pulsed heat load in the channel and the measured

substrate temperature. Based on this thermal model, SCD and PCD were introduced as alternative substrates and the TBR adjusted to reproduce the electrically predicted temperatures. For SCD, even a perfect TBR (= TBR_{Si} of $10\text{ m}^2\text{ K/GW}$) reveals temperatures larger than the predicted $\sim 40^\circ\text{C}$ (minimum temperature set by baseplate temperature), which suggests that the thermal performance of SCD is excellently exploited and the vdW bonding adds no significant thermal resistance to our devices. On PCD, in contrast, the present diamond nucleation layer increases the TBR consequently reducing the currents, which demonstrates the benefits in SCD compared to PCD as high-power GaN substrates.

Our temperature measurements are based on several consistent assumptions, however, more detailed quantification is needed to remove uncertainties regarding the SCD and PCD thermal performance. Future work will focus on the transfer of RF-devices and quantify channel and interface temperatures by Raman thermography or thermoreflectance to allow a direct measurement of lattice and surface temperatures and compare our devices with state-of-the-art technology.

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