# SOME SPECTRAL PROPERTIES OF AN INTEGRAL OPERATOR IN POTENTIAL THEORY 

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## 1. Introduction

In [7] Plemelj established some fundamental results in two- and three-dimensional potential theory about the eigenvalues of both the double layer potential operator and its adjoint, the normal derivative of the single layer potential operator. In [3] Blumenfeld and Mayer established some additional results concerning the eigenvalues of these integral operators in the case of $\mathbb{R}^{2}$. The spectral properties established by Plemelj [7] and by Blumenfeld and Mayer [3] have had a profound effect in the area of integral equation methods in scattering and potential theory in both $\mathbb{R}^{2}$ and $\mathbb{R}^{3}$.

Some applications that have been made of these results may be found in Colton and Kress [4]. A complete list, however, of all the different uses that have been made of the efforts of Plemelj [7] and of Blumenfeld and Mayer [3] would be a formidable task.

This paper arose from the author's long interest in the spectral properties of the double layer potential integral operator in both $\mathbb{R}^{2}$ and $\mathbb{R}^{3}$. For sufficiently smooth boundaries, it can be shown that for both $\mathbb{R}^{2}$ and $\mathbb{R}^{3}$ the point 0 lies in the spectrum of the integral operator. A fundamental question is what part of the spectrum does 0 lie in?

For the case of $\mathbb{R}^{3}$ some partial results on this topic are known. If the underlying boundary is either a sphere or a prolate spheroid, it can be shown (see [1]) that the eigenvalues of the double layer potential integral operator lie in the interval $[-1,0)$. Furthermore, it can be established that for both geometries 0 lies in the continuous spectrum of the integral operator. As for other geometries, the spectral classification of 0 remains an open question.

For the case of $\mathbb{R}^{2}$, it turns out that the situation is somewhat different for the double layer potential operator, which in this paper we denote by $K$ and define in equation (2.1). After encountering some serious difficulties in an attempt to establish a general theorem about which part of the spectrum of $K$ the point 0 lies in, the author looked at some specific examples, namely, the circle and the ellipse. For the case of the circle, the point 0 lies in the point spectrum of $K$. For the case of an ellipse, 0 lies in the continuous spectrum of $K$. Consequently, these examples demonstrate that 0 does not always lie in the same component of the spectrum of $K$ for all sufficiently smooth boundaries.

In the next section we give some notation, definitions and basic results which we shall use. In Section 3, we consider the case of a circle and compute the spectrum of $K$. It is shown that 0 is an eigenvalue and moreover, that it has infinite geometric multiplicity.

In Section 4 we consider the case of an ellipse. There we compute all the eigenvalues for $K$ and establish that every eigenvalue has a geometric multiplicity of 1 . Furthermore, we prove that the continuous spectrum of $K$ is equal to the set $\{0\}$.

## 2. Notation, definitions and basic results

In this section we give our notation and state some definitions and results which we shall require. Let $D_{i}$ be a bounded, simply connected domain in $\mathbb{R}^{2}$ containing the origin with a $C^{2}$ boundary $\partial D$, and let $D_{e}$ denote the region exterior to $\bar{D}_{i}$. Let $\hat{n}$ denote a unit normal directed out of $D_{i}$. Let $x$ and $y$ denote typical points in $\mathbb{R}^{2}$.

We now define the following integral operators of potential theory:

$$
\begin{gather*}
(K \psi)(x):=\frac{1}{\pi} \int_{\partial D} \psi(y) \frac{\partial}{\partial n(y)} \ln \frac{1}{|x-y|} d s_{y}, \quad x \in \partial D,  \tag{2.1}\\
(D \psi)(x):=\frac{1}{\pi} \int_{\partial D} \psi(y) \frac{\partial}{\partial n(y)} \ln \frac{1}{|x-y|} d s_{y}, \quad x \in \mathbb{R}^{2} \backslash \partial D . \tag{2.2}
\end{gather*}
$$

Here it is understood that the integration is taken with respect to arc length.
A standard result in two-dimensional potential theory (e.g. see [9, pp. 78-80]) states that for closed smooth curves $\partial D$

$$
\begin{equation*}
\lim _{\substack{x \rightarrow y \\ x, y \in \partial D}} \frac{\partial}{\partial n(y)} \ln \frac{1}{|x-y|}=-\frac{1}{2} \kappa(y) \tag{2.3}
\end{equation*}
$$

where $\kappa(y)$ denotes the curvature of $\partial D$ at $y$. Consequently, unlike the weakly singular nature of the double layer kernel in $\mathbb{R}^{3}$, the double layer kernel in $\mathbb{R}^{2}$ is continuous for all points $x$ and $y$ on $\partial D$, including when $x=y$.

Let $C(\partial D)$ denote the Banach space of complex-valued, continuous functions defined on $\partial D$ equipped with the maximum norm. Since the integral operator $K$ has a continuous kernel, it follows that $K$ is a compact linear operator on $C(\partial D)$ (see [4, Theorem 1.10]).

Let + and - denote the limits obtained for the double layer potential $(D \psi)(x)$ by approaching the boundary $\partial D$ from $D_{e}$ and $D_{i}$, respectively, that is

$$
\begin{align*}
& \left(D_{+} u\right)(x)=\lim _{\substack{x \\
x_{e} \in D_{e}}}(D u)\left(x_{e}\right), \quad x \in \partial D,  \tag{2.4}\\
& \left(D_{-} u\right)(x)=\lim _{\substack{x_{i} \rightarrow x \\
x_{i} \in D_{i}}}(D u)\left(x_{i}\right), \quad x \in \partial D . \tag{2.5}
\end{align*}
$$

It can be shown (e.g. see [5, p. 49] or [8, p. 392]) that

$$
\begin{equation*}
\left(D_{ \pm} u\right)(x)=(K u)(x) \pm u(x), \quad x \in \partial D . \tag{2.6}
\end{equation*}
$$

Let $A$ denote any bounded linear operator mapping a Banach space $X$ into itself. By an eigenvalue of $A$ we mean a complex number $\lambda$ such that the nullspace $N(\lambda I-$ $A) \neq\{0\}$ where $I$ denotes the identity operator. Let $\rho(A)$ denote the resolvent set of $A$. Let $\sigma(A)$ denote the spectrum of $A$. Let $\sigma_{C}(A), \sigma_{P}(A)$, and $\sigma_{R}(A)$ denote the continuous spectrum, point spectrum, and residual spectrum of $A$, respectively. It is known (e.g. see [2, Chapter 18] or [4, Theorem 1.34]) that if $X$ is an infinite dimensional Banach space and if $A$ is a compact linear operator then $\lambda=0$ lies in $\sigma(A)$ and $\sigma(A) \backslash\{0\}$ consists of at most a countable number of eigenvalues, with $\lambda=0$ the only possible limit point.

It can be shown (see [3]) that the eigenvalues of the integral operator $K$, defined in equation (2.1), lie in the interval $[-1,1)$ and are symmetric with respect to the origin. The only exception is the eigenvalue -1 corresponding to constant eigenfunctions.

Finally, we shall denote the set of positive integers by $\mathbb{N}$.

## 3. The Circle

In this section $\partial D$ is taken to be a circle of radius $a$. Under this assumption, we compute the spectrum of the integral operator $K$ and also determine the spectral properties of the point $\lambda=0$ for $K$.

With respect to polar coordinates, let the points $x$ and $y$ be given by ( $r_{x}, \phi_{x}$ ) and $(r, \phi)$, respectively. Then

$$
\begin{equation*}
\ln \frac{1}{|x-y|}=-\frac{1}{2} \ln \left[r_{x}^{2}+r^{2}-2 r_{x} r \cos \left(\phi-\phi_{x}\right)\right] \tag{3.1}
\end{equation*}
$$

From equation (3.1), for $x, y \in \partial D$, we have the following known result (e.g. see [5, p. 52])

$$
\begin{align*}
\frac{\partial}{\partial n(y)} \ln \frac{1}{|x-y|} & =\frac{\partial}{\partial r}\left\{-\frac{1}{2} \ln \left[a^{2}+r^{2}-2 a r \cos \left(\phi-\phi_{x}\right)\right]\right\}_{r=a} \\
& =-\frac{1}{2 a} \tag{3.2}
\end{align*}
$$

Before proceeding to the stated purposes of this section, it is worthy of note to examine the result in equation (3.2) in the context of equation (2.3). It is a well known result in differential geometry that the curvature of a circle of radius $a$ is $1 / a$. Consequently, the results in equations (2.3) and (3.2) are seen to be compatible.

From equations (2.1) and (3.2) it follows that

$$
\begin{equation*}
K \psi(x)=-\frac{1}{2 \pi} \int_{0}^{2 \pi} \psi(a, \phi) d \phi \tag{3.3}
\end{equation*}
$$

Letting $\psi$ equal $1, \cos m \phi$, and $\sin m \phi, m \in \mathbb{N}$, respectively, in equation (3.3) we have

$$
\begin{equation*}
K(1)=-1, \tag{3.4}
\end{equation*}
$$

$$
\begin{equation*}
K(\cos m \phi)=K(\sin m \phi)=0 . \tag{3.5}
\end{equation*}
$$

Thus $\lambda=-1$ is an eigenvalue of $K$ with corresponding eigenfunction 1 , and $\lambda=0$ is also an eigenvalue of $K$ with corresponding eigenfunctions $\{\cos m \phi, \sin m \phi: m \in \mathbb{N}\}$. From the completeness of the orthogonal set of eigenfunctions $\{1, \cos m \phi, \sin m \phi: m \in \mathbb{N}\}$ in $L^{2}[0,2 \pi]$, it follows, by a similar argument as used in [1], that with respect to the underlying Banach space $C(\partial D)$

$$
\begin{equation*}
\sigma_{P}(K)=\{-1,0\} . \tag{3.6}
\end{equation*}
$$

That is, $\lambda=-1$ and $\lambda=0$ are the only eigenvalues of $K$ for the case when $\partial D$ is a circle.
In view of the fact that $K$ is a compact linear operator on $C(\partial D)$, it follows that if $\lambda \neq 0$ then either $\lambda \in \rho(K)$ or $\lambda \in \sigma_{P}(K)$. Consequently, $\lambda=-1$ and $\lambda=0$ are the only elements of $\sigma(K)$. Furthermore, since $\cos m \phi$ and $\sin m \phi, m \in \mathbb{N}$, are all eigenfunctions of $K$ corresponding to $\lambda=0$, it follows that

$$
\begin{equation*}
\operatorname{dim} N(K)=\infty . \tag{3.7}
\end{equation*}
$$

## 4. The ellipse

The elliptical coordinates ( $\mu, \phi$ ) are related to the rectangular Cartesian coordinates ( $y_{1}, y_{2}$ ) by the transformation

$$
\begin{align*}
& y_{1}=\frac{1}{2} c \cosh \mu \cos \phi, \\
& y_{2}=\frac{1}{2} c \sinh \mu \sin \phi \tag{4.1}
\end{align*}
$$

where $0 \leqq \mu<\infty, 0 \leqq \phi \leqq 2 \pi$. The closed curves corresponding to $\mu=$ constant, $0 \leqq \phi \leqq 2 \pi$ are confocal ellipses of interfocal distance $c$, eccentricity $e=(\cosh \mu)^{-1}$, major axis $c \cosh \mu$ and minor axis $c \sinh \mu$. The limiting case $\mu=0$ represents the line segment between the foci.

In this section $\partial D$ will denote the ellipse corresponding to $\mu=b, 0 \leqq \phi \leqq 2 \pi$, where $b$ is some constant. To avoid the degenerate case, we will assume that $b>0$.

In terms of elliptical coordinates it can be shown that the gradient of a scalar function $\Phi(\mu, \phi)$ is given by

$$
\begin{equation*}
\nabla \Phi(\mu, \phi)=\frac{2}{c \tau}\left(\frac{\partial \Phi}{\partial \mu} \hat{e}_{\mu}+\frac{\partial \Phi}{\partial \phi} \hat{e}_{\phi}\right) \tag{4.2}
\end{equation*}
$$

where

$$
\begin{equation*}
\tau:=\left[\cosh ^{2} \mu \sin ^{2} \phi+\sinh ^{2} \mu \cos ^{2} \phi\right]^{1 / 2} \tag{4.3}
\end{equation*}
$$

and where $\hat{e}_{\mu}$ and $\hat{e}_{\phi}$ denote the orthonormal vectors

$$
\begin{align*}
& \hat{e}_{\mu}:=(\sinh \mu \cos \phi \hat{i}+\cosh \mu \sin \phi \hat{j}) / \tau \\
& \hat{e}_{\phi}:=(-\cosh \mu \sin \phi \hat{i}+\sinh \mu \cos \phi \hat{j}) / \tau \tag{4.4}
\end{align*}
$$

Furthermore, it can be shown that the element of arc length $d s$ is given by

$$
\begin{equation*}
d s=\frac{c}{2} \tau d \phi \tag{4.5}
\end{equation*}
$$

From [6, p. 1202] we have

$$
\begin{align*}
\ln \frac{1}{|x-y|}= & -\left(\mu_{>}+\ln \frac{c}{4}\right)+\sum_{n=1}^{\infty} \frac{2}{n} e^{-n \mu_{>}}\left[\cosh n \mu_{<} \cos n \phi \cos n \phi_{x}\right. \\
& \left.+\sinh n \mu_{<} \sin n \phi \sin n \phi_{x}\right] \tag{4.6}
\end{align*}
$$

where $\mu_{>}=\max \left\{\mu_{x}, \mu_{y}\right\}, \mu_{<}=\min \left\{\mu_{x}, \mu_{y}\right\}$, and $\left(\mu_{x}, \phi_{x}\right)$ and ( $\mu_{y}, \phi$ ) denote the elliptical coordinates of the points $x$ and $y$, respectively.

At the point $(b, \phi) \in \partial D$ the unit tangent vector $\hat{T}$ and the outer unit normal vector $\hat{n}$ are given, respectively, by

$$
\begin{equation*}
\hat{T}=\hat{e}_{\phi}, \quad \hat{n}=\hat{e}_{\mu} . \tag{4.7}
\end{equation*}
$$

For $y=(b, \phi) \in \partial D$ and $x=\left(\mu_{x} \phi_{x}\right) \in D_{i}$ it follows from equations (4.2) and (4.6) that

$$
\begin{gather*}
\frac{\partial}{\partial n(y)} \ln \frac{1}{|x-y|}=\frac{-2}{c \tau}\left\{1+2 \sum_{n=1}^{\infty} e^{-n b}\left[\cosh n \mu_{x} \cos n \phi \cos n \phi_{x}\right.\right. \\
+  \tag{4.8}\\
\left.\left.+\sinh n \mu_{x} \sin n \phi \sin n \phi_{x}\right]\right\} .
\end{gather*}
$$

Consequently, from equations (2.4), (4.5) and (4.8) we have

$$
\begin{gather*}
D \psi(x)=-\frac{1}{\pi} \int_{0}^{2 \pi} \psi(\mu, \phi)\left\{1+2 \sum_{n=1}^{\infty} e^{-n b}\left[\cosh n \mu_{x} \cos n \phi \cos n \phi_{x}\right.\right. \\
 \tag{4.9}\\
\left.\left.+\sinh n \mu_{x} \sin n \phi \sin n \phi_{x}\right]\right\} d \phi, \quad x \in D_{i} .
\end{gather*}
$$

Letting $\psi$ equal $1, \cos m \phi$, and $\sin m \phi$, where $m \in \mathbb{N}$, respectively, in equation (4.9), then using the orthogonality of the trigonometric functions, and finally letting
$\mu_{x} \rightarrow b$, we have from equation (2.6)

$$
\begin{gather*}
K(1)=-1,  \tag{4.10}\\
K(\cos m \phi)=-e^{-2 m b} \cos m \phi, \quad m \in \mathbb{N},  \tag{4.11}\\
K(\sin m \phi)=e^{-2 m b} \sin m \phi, \quad m \in \mathbb{N} . \tag{4.12}
\end{gather*}
$$

Thus it is seen that $-1,-\mathrm{e}^{-2 m b}$, and $e^{-2 m b}, m \in \mathbb{N}$, are eigenvalues of $K$ with corresponding eigenfunctions $1, \cos m \phi$, and $\sin m \phi$, respectively. From the completeness of the orthogonal set of eigenfunctions $\{1, \cos m \phi, \sin m \phi: m \in \mathbb{N}\}$ in $L^{2}[0,2 \pi]$, it follows, by an argument similar to one used in [1], that with respect to the underlying Banach space $C(\partial D)$

$$
\begin{equation*}
\sigma_{P}(K)=\left\{-1, \pm e^{-2 m b}: m \in \mathbb{N}\right\} . \tag{4.13}
\end{equation*}
$$

That is, the above eigenvalues are the only eigenvalues of $K$. Consequently, unlike the situation when $\partial D$ is a circle,

$$
\begin{equation*}
0 \notin \sigma_{P}(K) \tag{4.14}
\end{equation*}
$$

when $\partial D$ is an ellipse. Furthermore, it is seen that

$$
\begin{equation*}
\operatorname{dim} N(-I-K)=\operatorname{dim} N\left( \pm e^{-2 m b} I-K\right)=1 \tag{4.15}
\end{equation*}
$$

for each $m \in \mathbb{N}$. Therefore each eigenvalue has a geometric multiplicity of 1 .
To complete the analysis of this section we establish the following result which determines the spectral nature of the point $\lambda=0$.

Theorem 4.1 Let $\partial D$ denote the ellipse corresponding to $\mu=b, 0 \leqq \phi \leqq 2 \pi$, where $b$ is some positive constant. Then

$$
\{0\}=\sigma_{c}(K) .
$$

Proof. Since $K$ is a compact linear operator on $C(\partial D)$,

$$
\begin{equation*}
0 \in \sigma(K) . \tag{4.16}
\end{equation*}
$$

Furthermore, since each eigenfunction must necessarily lie in the range of $K, R(K)$, we have

$$
\begin{equation*}
\{1, \cos m \phi, \sin m \phi: m \in \mathbb{N}\} \subset R(K) . \tag{4.17}
\end{equation*}
$$

It follows that $R(K)$ is dense in $C(\partial D)$. Consequently, from equations (4.14) and (4.16) it
follows that

$$
\begin{equation*}
0 \in \sigma_{c}(K) \tag{4.18}
\end{equation*}
$$

Finally, by using the fact that $K$ is a compact linear operator on $C(\partial D)$, we have from equation (4.18)

$$
\{0\}=\sigma_{C}(K) .
$$

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