



The H and K Family of Mock Theta Functions

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Abstract. In his last letter to Hardy, Ramanujan defined 17 functions $F(q)$, $|q| < 1$, which he called mock θ -functions. He observed that as q radially approaches any root of unity ζ at which $F(q)$ has an exponential singularity, there is a θ -function $T_\zeta(q)$ with $F(q) - T_\zeta(q) = O(1)$. Since then, other functions have been found that possess this property. These functions are related to a function $H(x, q)$, where x is usually q^r or $e^{2\pi ir}$ for some rational number r . For this reason we refer to H as a “universal” mock θ -function. Modular transformations of H give rise to the functions K, K_1, K_2 . The functions K and K_1 appear in Ramanujan’s lost notebook. We prove various linear relations between these functions using Appell–Lerch sums (also called generalized Lambert series). Some relations (mock theta “conjectures”) involving mock θ -functions of even order and H are listed.

1 Introduction

In Ramanujan’s last letter to Hardy ([22, pp. 354–355], [23, pp. 127–131], [26, pp. 56–61]) he observes that the asymptotic expansions of certain q -series with exponential singularities at roots of unity “close” in a striking manner. For example, let

$$G(q) = \sum_{n=0}^{\infty} \frac{q^{n^2}}{(1-q)(1-q^2)\cdots(1-q^n)} = \prod_{m=0}^{\infty} \frac{1}{(1-q^{5m+1})(1-q^{5m+4})}$$

(where the last equality is the first Rogers–Ramanujan identity). If $q = e^{-t}$ and $t \rightarrow 0^+$ (so that q approaches 1 radially from inside the unit circle), then

$$G(q) = \sqrt{\frac{2}{5-\sqrt{5}}} \exp\left(\frac{\pi^2}{15t} - \frac{t}{60}\right) + o(1).$$

In the same letter Ramanujan notes that it is only for some special q -series $f(q)$ that the exponential closes, *i.e.*, its argument terminates with some power t^N . If $f(q)$ is not the sum of a theta function and a function which is $O(1)$ at all roots of unity ζ , and if for each such ζ there is an approximation of the form

$$f(q) = \sum_{\mu=1}^M t^{k_\mu} \exp\left(\sum_{\nu=-1}^N c_{\mu\nu} t^\nu\right) + O(1)$$

Received by the editors November 21, 2010.

Published electronically October 5, 2011.

Support by the Natural Sciences and Engineering Research Council of Canada is gratefully acknowledged.

AMS subject classification: 11B65, 33D15.

Keywords: mock theta function, q -series, Appell–Lerch sum, generalized Lambert series.

as $t \rightarrow 0^+$ with $q = \zeta e^{-t}$, he calls $f(q)$ a *mock θ -function*. It appears from his letter, however, that he was actually concerned with functions having the (possibly) more restrictive property that for every root of unity ζ , there are modular forms $h_j^{(\zeta)}(q)$ and rational numbers α_j , $1 \leq j \leq J(\zeta)$, such that

$$f(q) = \sum_{j=1}^{J(\zeta)} q^{\alpha_j} h_j^{(\zeta)}(q) + O(1)$$

as q radially approaches ζ . For a further description of mock theta functions see [13].

The most well-known infinite family of mock θ -functions is defined by $M(q^r, q)$, where r is a noninteger rational number and

$$M(x, q) = \sum_{n=0}^{\infty} \frac{q^{n(n+1)}}{(x)_{n+1}(q/x)_{n+1}}.$$

In this paper we will use the standard notation for the q -shifted factorial:

$$(a; q^k)_0 = 1, \quad (a; q^k)_n = \prod_{m=0}^{n-1} (1 - aq^{km}), \quad (a; q^k)_\infty = \prod_{m=0}^{\infty} (1 - aq^{km}),$$

where k is a positive integer. When $k = 1$ it is customary to write $(a)_n$ instead of $(a; q)_n$.

The functions $M(q, q^5)$ and $M(q^2, q^5)$ appear in the celebrated Mock Theta Conjectures stated by Ramanujan in the lost notebook [23] and later proved by Hickerson [15]. These conjectures are linear relations involving the fifth order mock θ -functions.

The function

$$N(y, q) = \sum_{n=0}^{\infty} \frac{q^{n^2}}{(yq)_n (y^{-1}q)_n}$$

is related to $M(x, q)$ by a modular transformation law, proved in [12] and restated in Section 4. This function is also known as the rank generating function (see, for example [4]).

In [12] the functions $M(q^r, q)$ and $N(e^{2\pi ir}, q)$ are denoted by $M(r, q)$ and $N(r, q)$, respectively. The product $(e^{2\pi ir})_n (e^{-2\pi ir})_n$ in the definition of $N(r, q)$ in [12] should be $(e^{2\pi ir}q)_n (e^{-2\pi ir}q)_n$. The function $N_1(r, q)$ in [12] is equal to our $M(e^{2\pi ir}q, q^2)$.

In this paper we study another infinite family of mock θ -functions defined by $H(q^r, q)$, where r is a noninteger rational number and

$$H(x, q) = \sum_{n=0}^{\infty} \frac{q^{\frac{1}{2}n(n+1)} (-q)_n}{(x)_{n+1} (q/x)_{n+1}}.$$

The function $H(x, q)$ is defined in a different way by Choi [6, eq. (2.34)]. It is closely related to the functions

$$\begin{aligned}
 K(y, q) &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{n^2} (q; q^2)_n}{(yq^2; q^2)_n (y^{-1}q^2; q^2)_n}, \\
 K_1(y, q) &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{(n+1)^2} (q; q^2)_n}{(yq; q^2)_{n+1} (y^{-1}q; q^2)_{n+1}}, \\
 K_2(y, q) &= \sum_{n=0}^{\infty} \frac{q^{\frac{1}{2}n(n+1)} (-1)_n}{(yq)_n (y^{-1}q)_n}.
 \end{aligned}$$

Of particular interest are the linear and modular relations connecting these functions. One such linear relation involving the functions K and K_1 appears on page 8 of the lost notebook (see also [1, pp. 264–267]). We prove several more in Section 3. The modular relations are studied in Section 4.

For every classical mock θ -function $f(q)$ explicit linear relations involving f , H (or M), and ordinary θ -functions are known [13]. These relations are usually referred to as mock theta “conjectures”, even when their proofs are known. The “conjectures” for the functions of even order involve H and are listed in Section 5.

In Section 2 we show that the function H is a normalized level 2 Appell function (see Section 6 for the definition of an Appell function), whereas the function M is a normalized level 3 Appell function. Appell functions of higher level can often be expressed in terms of those with lower level. A linear relation expressing M in terms of H and a θ -function is given in Section 3. By this relation and the mock theta “conjectures”, every classical mock θ -function is related to H . For this reason we refer to H as a “universal” mock θ -function.

A preprint of this paper was circulated during a conference at the University of Florida in 2004. Subsequently, several results of the preprint were cited by Bringmann, Ono, and Rhoades [5]. Their desire to see a published version is fulfilled here.

2 Appell-Lerch Sums

To prove linear relations and construct transformation laws for these functions it is more convenient to work with the Appell-Lerch sums (also called generalized Lambert series) studied in [18, 19]. Two of these sums defined for positive integers k are

$$(2.1) \quad U_k(x, q) = \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{1 - xq^n},$$

$$(2.2) \quad V_k(y, q) = \frac{1}{1 - y^{-1}} \sum_{n=-\infty}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{1 - yq^n}.$$

It is not difficult to show that

$$(2.3) \quad U_k(x, q) = \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{(1-xq^n)(1-q^{n+1}/x)},$$

$$(2.4) \quad V_k(y, q) = \sum_{n=-\infty}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{(1-yq^n)(1-y^{-1}q^n)}.$$

Observe that

$$\begin{aligned} U_k(x, q) &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{1-xq^n} + \sum_{n=-\infty}^{-1} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{1-xq^n} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{1-xq^n} + \sum_{n=1}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n-1)}}{1-xq^{-n}} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{1-xq^n} + \sum_{n=0}^{\infty} \frac{(-1)^{n+1} q^{\frac{1}{2}kn(n+1)}}{1-xq^{-n-1}} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{1-xq^n} + \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)} q^{n+1}/x}{1-q^{n+1}/x} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)} (1-q^{2n+1})}{(1-xq^n)(1-q^{n+1}/x)} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{(1-xq^n)(1-q^{n+1}/x)} - \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)} q^{2n+1}}{(1-xq^n)(1-q^{n+1}/x)} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{(1-xq^n)(1-q^{n+1}/x)} + \sum_{n=1}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n-1)} q^{2n-1}}{(1-xq^{n-1})(1-q^n/x)} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{(1-xq^n)(1-q^{n+1}/x)} + \sum_{n=-\infty}^{-1} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)} q^{-2n-1}}{(1-xq^{-n-1})(1-q^{-n}/x)} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{(1-xq^n)(1-q^{n+1}/x)} + \sum_{n=-\infty}^{-1} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{(q^{n+1}-x)(q^n-1/x)} \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{(1-xq^n)(1-q^{n+1}/x)}, \end{aligned}$$

which is (2.3). Similarly,

$$\begin{aligned} V_k(y, q) &= \frac{1}{(1-y)(1-y^{-1})} \\ &\quad + \frac{1}{1-y^{-1}} \sum_{n=1}^{\infty} \left(\frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{1-yq^n} + \frac{(-1)^{kn} q^{\frac{1}{2}n(kn-1)}}{1-yq^{-n}} \right) \end{aligned}$$

$$\begin{aligned}
 &= \frac{1}{(1-y)(1-y^{-1})} \\
 &\quad + \frac{1}{1-y^{-1}} \sum_{n=1}^{\infty} \left(\frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{1-yq^n} - \frac{(-1)^{kn} q^{\frac{1}{2}n(kn-1)} y^{-1} q^n}{1-y^{-1}q^n} \right) \\
 &= \frac{1}{(1-y)(1-y^{-1})} \\
 &\quad + \frac{1}{1-y^{-1}} \sum_{n=1}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)} (1-y^{-1})(1+q^n)}{(1-yq^n)(1-y^{-1}q^n)} \\
 &= \frac{1}{(1-y)(1-y^{-1})} + \sum_{n=1}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)} (1+q^n)}{(1-yq^n)(1-y^{-1}q^n)} \\
 &= \sum_{n=0}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{(1-yq^n)(1-y^{-1}q^n)} + \sum_{n=1}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)} q^n}{(1-yq^n)(1-y^{-1}q^n)} \\
 &= \sum_{n=0}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{(1-yq^n)(1-y^{-1}q^n)} + \sum_{n=-\infty}^{-1} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn-1)} q^{-n}}{(1-yq^{-n})(1-y^{-1}q^{-n})} \\
 &= \sum_{n=0}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{(1-yq^n)(1-y^{-1}q^n)} + \sum_{n=-\infty}^{-1} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)} q^{-2n}}{(1-yq^{-n})(1-y^{-1}q^{-n})} \\
 &= \sum_{n=0}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{(1-yq^n)(1-y^{-1}q^n)} + \sum_{n=-\infty}^{-1} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{(q^n - y)(q^n - y^{-1})} \\
 &= \sum_{n=-\infty}^{\infty} \frac{(-1)^{kn} q^{\frac{1}{2}n(kn+1)}}{(1-yq^n)(1-y^{-1}q^n)},
 \end{aligned}$$

which is (2.4).

By (2.3) and (2.4) we see that $U_k(x, q) = U_k(q/x, q)$ and $V_k(y, q) = V_k(y^{-1}, q)$. Also, $V_k(e^{2\pi ir}, q)$ is real when r is a noninteger rational number and q is real with $0 < |q| < 1$. (This function plays an important role in equation (4.2).)

Many of our identities involve the Jacobi θ -function defined by

$$j(x, q) = \sum_{n=-\infty}^{\infty} (-1)^n q^{\frac{1}{2}n(n-1)} x^n = (x)_{\infty} (q/x)_{\infty} (q)_{\infty}$$

(where the last equality is the well-known Jacobi triple-product identity; see, for example [7, p. 12]). Following Hickerson [15], we define a θ -product (or θ -quotient) to be an expression of the form

$$Cq^e x_1^{f_1} \cdots x_r^{f_r} L_1^{g_1} \cdots L_s^{g_s},$$

where C is a complex number, e and f_i are rational numbers, g_j are integers, and each L_j has the form

$$j(Dq^h x_1^{k_1} \cdots x_r^{k_r}, \pm q^m)$$

for some complex number D (usually $D = \pm 1$) and rational numbers h, k_i, m with $m > 0$. A θ -function is a finite sum of θ -products. Thus $(q)_\infty = j(q, q^3)$ is a θ -function, even though it lacks the factor $q^{\frac{1}{24}}$ needed to make it a modular form.

The sums U_1 and V_1 (multiplied by $1 - y$) turn out to be θ -functions, since

$$(2.5) \quad U_1(x, q) = \frac{(q)_\infty^2}{(x)_\infty(q/x)_\infty} = \frac{(q)_\infty^3}{j(x, q)},$$

$$(2.6) \quad V_1(y, q) = \frac{U_1(y, q)}{1 - y^{-1}} = \frac{(q)_\infty^2}{(y)_\infty(y^{-1})_\infty}.$$

Equation (2.5) is the expansion for the reciprocal of a θ -function and is equivalent to the next to last formula on page 1 of the lost notebook (see also [1, p. 264, eq. (12.2.9)]). The function $(1 - z)U_1(z, q)/(q)_\infty$ is the crank statistic of Garvan [8, eq. (1.25)].

At this point we introduce two more θ -functions: Jacobi's $\theta_4(0, q)$ defined by

$$\theta_4(0, q) = \sum_{n=-\infty}^{\infty} (-1)^n q^{n^2} = j(q, q^2) = \frac{(q)_\infty^2}{(q^2; q^2)_\infty} = \frac{(q)_\infty}{(-q)_\infty}$$

and the Gauss function $\psi(q)$ defined by

$$\psi(q) = \sum_{n=0}^{\infty} q^{\frac{1}{2}n(n+1)} = \frac{1}{2}j(-1, q) = j(-q, q^4) = \frac{(q^2; q^2)_\infty}{(q)_\infty} = \frac{(q^2; q^2)_\infty}{(q; q^2)_\infty}.$$

Using these functions we can express the relations between the H and K family and the sums U_2 and V_2 as follows:

$$(2.7) \quad H(x, q) = \frac{U_2(x, q)}{\theta_4(0, q)}$$

$$(2.8) \quad K(y, q) = (1 - y)(1 - y^{-1}) \frac{V_2(y, q^2)}{\psi(q)},$$

$$(2.9) \quad K_1(y, q) = \frac{1}{(1 - y^{-1})\psi(q)} \sum_{n=-\infty}^{\infty} \frac{q^{(n+1)(2n+1)}}{1 - yq^{2n+1}} = \frac{V_2(y, q^2) - V_1(y, q)}{\psi(q)},$$

$$(2.10) \quad K_2(y, q) = \frac{(1 - y)(1 - y^{-1})}{\theta_4(0, q)} \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)}}{(1 - yq^n)(1 - y^{-1}q^n)}$$

$$= \frac{1 - y}{1 + y} \left(1 + 2y \frac{U_2(y, q)}{\theta_4(0, q)} \right).$$

Equation (2.8) is equivalent to the last formula on page 1 of the lost notebook. Other identities of this type are given in [1, Chapter 12].

Equations (2.5)–(2.10) can be proved by the Watson–Whipple transformation [7, p. 242, eq. (III.17)]:

$$(2.11) \quad {}_8\phi_7 \left[\begin{matrix} a, & qa^{\frac{1}{2}}, & -qa^{\frac{1}{2}}, & b, & c, & d, & e, & f \\ & a^{\frac{1}{2}}, & -a^{\frac{1}{2}}, & \frac{aq}{b}, & \frac{aq}{c}, & \frac{aq}{d}, & \frac{aq}{e}, & \frac{aq}{f} \end{matrix} ; q; \frac{a^2q^2}{bcdef} \right] = \frac{(aq)_\infty \left(\frac{aq}{de}\right)_\infty \left(\frac{aq}{df}\right)_\infty \left(\frac{aq}{ef}\right)_\infty}{\left(\frac{aq}{d}\right)_\infty \left(\frac{aq}{e}\right)_\infty \left(\frac{aq}{f}\right)_\infty \left(\frac{aq}{def}\right)_\infty} {}_4\phi_3 \left[\begin{matrix} \frac{aq}{bc}, & d, & e, & f \\ \frac{aq}{b}, & \frac{aq}{c}, & \frac{def}{a} \end{matrix} ; q; q \right].$$

We now prove (2.7). (The proofs of the other identities are similar.) Observe that

$$\frac{(qa^{\frac{1}{2}})_n (-qa^{\frac{1}{2}})_n}{(a^{\frac{1}{2}})_n (-a^{\frac{1}{2}})_n} = \frac{(1 - aq^2)(1 - aq^4) \cdots (1 - aq^{2n})}{(1 - a)(1 - aq^2) \cdots (1 - aq^{2n-2})} = \frac{1 - aq^{2n}}{1 - a}.$$

Let e and f tend to infinity (or equivalently, put $e = 1/e'$, $f = 1/f'$, simplify and then let $e' = f' = 0$). Then $(aq/e)_n$ and $(aq/f)_n$ tend to 1. Also,

$$\begin{aligned} (e)_n &= (1 - e)(1 - eq) \cdots (1 - eq^{n-1}) \\ &= (-e)^n \left(-\frac{1}{e} + 1\right) \left(-\frac{1}{e} + q\right) \cdots \left(-\frac{1}{e} + q^{n-1}\right) \\ &\sim (-e)^n q^{\frac{1}{2}n(n-1)} \end{aligned}$$

as $e \rightarrow \infty$. Similarly, $(f)_n \sim (-f)^n q^{\frac{1}{2}n(n-1)}$ as $f \rightarrow \infty$. Hence in the limit, (2.11) becomes

$$\begin{aligned} \sum_{n=0}^{\infty} \left(\frac{1 - aq^{2n}}{1 - a}\right) \frac{(a)_n (b)_n (c)_n (d)_n}{(q)_n \left(\frac{aq}{b}\right)_n \left(\frac{aq}{c}\right)_n \left(\frac{aq}{d}\right)_n} \left(\frac{a^2}{bcd}\right)^n q^{n(n+1)} &= \\ \frac{(aq)_\infty}{\left(\frac{aq}{d}\right)_\infty} \sum_{n=0}^{\infty} \frac{\left(\frac{aq}{bc}\right)_n (d)_n}{(q)_n \left(\frac{aq}{b}\right)_n \left(\frac{aq}{c}\right)_n} \left(-\frac{a}{d}\right)^n q^{\frac{1}{2}n(n+1)}. \end{aligned}$$

Now put $a = q$, $b = x$, $c = q/x$, and $d = -q$ to get

$$\sum_{n=0}^{\infty} \frac{(-1)^n q^{n(n+1)} (1 - q^{2n+1})}{(1 - xq^n)(1 - q^{n+1}/x)} = \frac{(q)_\infty}{(-q)_\infty} \sum_{n=0}^{\infty} \frac{q^{\frac{1}{2}n(n+1)} (-q)_n}{(x)_{n+1} (q/x)_{n+1}}.$$

Therefore

$$\begin{aligned} \theta_4(0, q) H(x, q) &= \sum_{n=0}^{\infty} (-1)^n q^{n(n+1)} \left(\frac{xq^n}{1 - xq^n} + \frac{1}{1 - q^{n+1}/x}\right) \\ &= \sum_{n=0}^{\infty} (-1)^n q^{n(n+1)} \left(\frac{1}{1 - q^{n+1}/x} - \frac{1}{1 - q^{-n}/x}\right) \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)}}{1 - q^{n+1}/x} = \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)}}{1 - xq^n} = U_2(x, q), \end{aligned}$$

since $U_2(x, q) = U_2(q/x, q)$. This completes the proof of (2.7).

Subtracting $V_2(y, q^2)$ from $V_1(y, q)$ removes the even terms in $V_1(y, q)$. The second equality in (2.9) easily follows from this observation. The second equality in (2.10) is a consequence of the two identities:

$$(2.12) \quad yU_2(y, q) + y^{-1}U_2(y^{-1}, q) = -\theta_4(0, q),$$

$$(2.13) \quad yU_2(y, q) - y^{-1}U_2(y^{-1}, q) = (y - y^{-1}) \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)}}{(1 - yq^n)(1 - y^{-1}q^n)}.$$

The proof of (2.13) is straightforward. We now prove (2.12). By (2.3) we get $U_2(x, q) = U_2(q/x, q)$. Hence,

$$\begin{aligned} yU_2(y, q) + y^{-1}U_2(y^{-1}, q) &= yU_2(y, q) + y^{-1}U_2(yq, q) \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)} y}{1 - yq^n} + \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)} y^{-1}}{1 - yq^{n+1}} \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)} y}{1 - yq^n} + \sum_{n=-\infty}^{\infty} \frac{(-1)^{n-1} q^{n(n-1)} y^{-1}}{1 - yq^n} \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n^2} (yq^n - y^{-1}q^{-n})}{1 - yq^n} \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n^2} (1 + yq^n)(1 - y^{-1}q^{-n})}{1 - yq^n} \\ &= - \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n^2} (1 + yq^n)(1 - yq^n) y^{-1} q^{-n}}{1 - yq^n} \\ &= - \sum_{n=-\infty}^{\infty} (-1)^n q^{n^2} (1 + y^{-1}q^{-n}) \\ &= - \sum_{n=-\infty}^{\infty} (-1)^n q^{n^2} - y^{-1} \sum_{n=-\infty}^{\infty} (-1)^n q^{n(n-1)} \\ &= -\theta_4(0, q), \end{aligned}$$

since the last sum vanishes.

The M and N analogues of (2.7) and (2.8) are ([12])

$$(2.14) \quad M(x, q) = \frac{U_3(x, q)}{(q)_{\infty}},$$

$$(2.15) \quad N(y, q) = \frac{(1 - y)(1 - y^{-1})}{(q)_{\infty}} V_3(y, q).$$

We see that the H and K family is related to U_2 and V_2 , and the M and N family is related to U_3 and V_3 . In his proof of the Mock Theta Conjectures, Hickerson denotes the function M by g . In view of these observations, in the forthcoming paper *A survey of classical mock theta functions* [13], the functions H, K, M, N are denoted by g_2, h_2, g_3, h_3 , respectively.

3 Linear Relations

In this section we give linear relations for mock θ -functions, where the coefficients are usually θ -functions. Unlike the modular transformation laws in Section 4, convergence of the functions in these relations is not required. Equality is to be interpreted as equality of formal q -series (or Laurent series in q after replacing q by a suitable power of q if necessary).

Since

$$U_2(x, q) + U_2(-x, q) = 2U_1(x^2, q^2) = \frac{2(q^2; q^2)_\infty^3}{j(x^2, q^2)},$$

it follows by (2.7) that

$$(3.1) \quad H(x, q) + H(-x, q) = \frac{2\psi^2(q)}{j(x^2, q^2)}.$$

From (2.10) and (2.7) we obtain

$$(3.2) \quad \frac{1+y}{1-y}K_2(y, q) = 1 + 2yH(y, q).$$

Since $K_2(x, q) = K_2(x^{-1}, q)$, it follows by (3.2) that

$$(3.3) \quad xH(x, q) + x^{-1}H(x^{-1}, q) = -1.$$

Substituting $H(x^{-1}, q) = H(q/x^{-1}, q) = H(xq, q)$ into (3.3), we get

$$xH(x, q) + x^{-1}H(xq, q) = -1,$$

which is equivalent to the functional equation

$$H(xq, q) = -x^2H(x, q) - x.$$

Combining (3.2) and (3.1) gives

$$\frac{1+y}{1-y}K_2(y, q) - \frac{1-y}{1+y}K_2(-y, q) = \frac{4y\psi^2(q)}{j(y^2, q^2)},$$

or equivalently,

$$\frac{K_2(y, q)}{(1-y)(1-y^{-1})} + \frac{K_2(-y, q)}{(1+y)(1+y^{-1})} = \frac{4V_1(y^2, q^2)}{\theta_4(0, q)}.$$

By (2.8) and (2.9),

$$(3.4) \quad (1 - y)^{-1}K(y, q) - (1 - y^{-1})K_1(y, q) = \frac{(1 - y^{-1})V_1(y, q)}{\psi(q)} = \frac{(q)_\infty^3}{\psi(q)j(y, q)}.$$

When $y = -a$ we obtain the fifth formula on page 8 of the lost notebook (see also [1, p. 265, eq. (12.3.2)]):

$$\sum_{n=0}^{\infty} \frac{(-1)^n (q; q^2)_n q^{n^2}}{(-a; q^2)_{n+1} (-q^2/a; q^2)_n} - (1 + 1/a) \sum_{n=0}^{\infty} \frac{(-1)^n (q; q^2)_n q^{(n+1)^2}}{(-aq; q^2)_{n+1} (-q/a; q^2)_{n+1}} = \frac{(q; q^2)_\infty \theta_4(0, q)}{(-a; q)_\infty (-q/a; q)_\infty}.$$

Generalizations of equations (3.1) and (3.4) were later given in [5, Theorem 1.3] and in [17, eq. (1.7), (1.11)].

The functions H and K are related by the identity [14]

$$\frac{qH(x, q)}{x} + \frac{K(-x^2/q, q^2)}{1 + x^2/q} = \frac{(q^2; q^2)_\infty^3}{j(x, q)j(-x^2/q, q^4)}.$$

The analogous relation between M and N is

$$xM(x, q) + 1 = \frac{N(x, q)}{1 - x},$$

which is equivalent to

$$(3.5) \quad x^2U_3(x, q) + (1 - x)V_3(x, q) + x(q)_\infty = 0$$

by (2.14) and (2.15).

We now prove (3.5). By (2.1) and (2.2) we have

$$\begin{aligned} & x^2U_3(x, q) + (1 - x)V_3(x, q) + x(q)_\infty \\ &= x^2 \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{\frac{3}{2}n(n+1)}}{1 - xq^n} + \frac{1 - x}{1 - x^{-1}} \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{\frac{1}{2}n(3n+1)}}{1 - xq^n} + xj(q^2, q^3) \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n x^2 q^{\frac{3}{2}n(n+1)}}{1 - xq^n} - \sum_{n=-\infty}^{\infty} \frac{(-1)^n xq^{\frac{1}{2}n(3n+1)}}{1 - xq^n} + \sum_{n=-\infty}^{\infty} (-1)^n xq^{\frac{1}{2}n(3n+1)} \\ &= \sum_{n=-\infty}^{\infty} \frac{(-1)^n xq^{\frac{1}{2}n(3n+1)}(xq^n - 1 + (1 - xq^n))}{1 - xq^n} = 0. \end{aligned}$$

The functions M and N can be expressed in terms of H . More precisely,

$$M(x^4, q^4) = \frac{qH(x^6q, q^6)}{x^2} + \frac{x^2H(x^6/q, q^6)}{q} - \frac{x^2(q^2; q^2)_\infty^3 (q^{12}; q^{12})_\infty j(x^2q, q^2) j(x^{12}q^6, q^{12})}{q(q^4; q^4)_\infty (q^6; q^6)_\infty^2 j(x^4, q^2) j(x^6/q, q^2)}.$$

This identity is a special case of more general identities expressing each U_k as a combination of $k - 1$ copies of U_2 and a θ -function [13, eq. (6.7), (6.8)]. Identities (6.7) and (6.8) in [13] were discovered by the author and proved by Gordon [14].

All of the functions $K, K_1, K_2, M,$ and N can be expressed in terms of H and θ -functions. In [13] a case is made for considering H as a “universal” mock θ -function.

Other linear relations involving H can be constructed using the transformation laws in Section 4 and the hyperbolic function identity

$$\frac{\cosh ax}{\cosh x} = \frac{\sinh(1+a)x}{\sinh 2x} + \frac{\sinh(1-a)x}{\sinh 2x}.$$

Some special values of H, K_2 and K are

(3.6) $H(-1, q) = 1/2,$

(3.7) $H(q, -q^2) = \psi(q^4),$

(3.8) $H(iq, q^2) = \psi(q^4),$

(3.9) $H(i, q) = \theta_4(0, -q)/2 + i/2,$

(3.10) $K_2(i, q) = \theta_4(0, -q),$

(3.11) $K_2(1, q) = 1/\theta_4(0, q),$

(3.12) $K(1, q) = 1/\psi(q).$

Equation (3.6) is obtained from (3.3) with $x = -1$. Observe that (3.7) is (3.8) with q replaced by $-iq$, and (3.10) follows from (3.9) by (3.2). We will now prove (3.8) and (3.9).

By (2.7) and (2.5),

$$\begin{aligned} (3.13) \quad H(iq, q^2) &= \frac{U_2(iq, q^2)}{\theta_4(0, q^2)} = \frac{1}{\theta_4(0, q^2)} \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{2n(n+1)}}{1 - iq^{2n+1}} \\ &= \frac{1}{\theta_4(0, q^2)} \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{2n(n+1)}(1 + iq^{2n+1})}{1 + q^{4n+2}} \\ &= \frac{1}{\theta_4(0, q^2)} \left(U_1(-q^2, q^4) + iq \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{2n(n+2)}}{1 + q^{4n+2}} \right) = \psi(q^4). \end{aligned}$$

The last sum in (3.13) vanishes, since

$$\sum_{n=-\infty}^{-1} \frac{(-1)^n q^{2n(n+2)}}{1 + q^{4n+2}} = - \sum_{n=0}^{\infty} \frac{(-1)^n q^{2n(n+2)}}{1 + q^{4n+2}}.$$

By (3.1) we have

$$(3.14) \quad H(i, q) + H(-i, q) = \frac{2q\psi^2(q)}{j(-1, q^2)} = \theta_4(0, -q),$$

and by (3.3) we have

$$iH(i, q) - iH(-i, q) = -1,$$

which is equivalent to

$$(3.15) \quad H(i, q) - H(-i, q) = i.$$

Adding (3.14) and (3.15) we obtain (3.9).

To prove (3.11), we begin with the first identity of (2.10):

$$\begin{aligned} \theta_4(0, q)K_2(y, q) &= (1 - y)(1 - y^{-1}) \sum_{n=-\infty}^{\infty} \frac{(-1)^n q^{n(n+1)}}{(1 - yq^n)(1 - y^{-1}q^n)} \\ &= 1 + (1 - y)(1 - y^{-1}) \sum_{\substack{n=-\infty \\ n \neq 0}}^{\infty} \frac{(-1)^n q^{n(n+1)}}{(1 - yq^n)(1 - y^{-1}q^n)}. \end{aligned}$$

When $y = 1$, this becomes $\theta_4(0, q)K_2(1, q) = 1$, which is (3.11).

The proof (3.12) is similar to the proof of (3.11). By (2.8) and (2.2) we obtain

$$\begin{aligned} \psi(q)K(y, q) &= (1 - y)(1 - y^{-1})V_2(y, q^2) = (1 - y) \sum_{n=-\infty}^{\infty} \frac{q^{n(2n+1)}}{1 - yq^{2n}} \\ &= 1 + (1 - y) \sum_{\substack{n=-\infty \\ n \neq 0}}^{\infty} \frac{q^{n(2n+1)}}{1 - yq^{2n}}. \end{aligned}$$

When $y = 1$, this becomes $\psi(q)K(1, q) = 1$, which is (3.12).

In the transformation law for $H(q^r, -q)$ (see Section 4) the integral vanishes when $r = 1/2$. This implies that $H(q^{1/2}, -q)$ and $K_2(i, -q)$ are θ -functions. Using computer algebra we found identities (3.7)–(3.10).

4 Transformation Laws

In discussing the approximation of mock θ -functions near roots of unity, we have adhered to the notation $q = e^{-\alpha}$, employed by Ramanujan and his early successors. This maps the right half-plane $\text{Re}(\alpha) > 0$ onto the punctured disc $0 < |q| < 1$. In the classical theory of θ -functions, as expounded for example in [25, 27], it is customary to write instead $q = e^{\pi i \tau}$ with $\text{Im}(\tau) > 0$. Thus $\alpha = -\pi i \tau$. The transformations of mock θ -functions are more complicated than those of θ -functions; they involve Mordell integrals [21]. For example, the θ -functions $(q)_\infty$, $\theta_4(q)$, and $\psi(q)$ satisfy the transformation laws:

$$\begin{aligned}
 (4.1) \quad q^{\frac{1}{24}}(q)_\infty &= \sqrt{\frac{2\pi}{\alpha}} q_1^{\frac{1}{6}}(q_1^4; q_1^4)_\infty, \\
 q^{\frac{1}{24}}(-q; -q)_\infty &= \sqrt{\frac{\pi}{\alpha}} q_1^{\frac{1}{24}}(-q_1; -q_1)_\infty, \\
 \theta_4(0, q) &= \sqrt{\frac{4\pi}{\alpha}} q_1^{\frac{1}{4}} \psi(q_1^2), \\
 \theta_4(0, -q) &= \sqrt{\frac{\pi}{\alpha}} \theta_4(0, -q_1), \\
 q^{\frac{1}{8}} \psi(q) &= \sqrt{\frac{\pi}{2\alpha}} \theta_4(0, q_1^2), \\
 q^{\frac{1}{8}} \psi(-q) &= \sqrt{\frac{\pi}{\alpha}} q_1^{\frac{1}{8}} \psi(-q_1),
 \end{aligned}$$

where $q = e^{-\alpha}$ and $q_1 = e^{-\beta}$ with $\alpha\beta = \pi^2$. Here

$$(a; -q)_n = \prod_{m=0}^{\infty} (1 - a(-q)^m) = (a; q^2)_\infty (-aq; q^2)_\infty.$$

Observe that (4.1) is the functional equation for the Dedekind η -function (see, for example [3, p. 48]); the other five identities above can easily be deduced from it.

The corresponding laws for the mock θ -function $H(q^r, q)$ are

$$\begin{aligned}
 q^{r(1-r)} H(q^r, q) &= \sqrt{\frac{\pi}{4\alpha}} \csc(\pi r) q_1^{-\frac{1}{4}} K(e^{2\pi i r}, q_1^2) \\
 &\quad - \sqrt{\frac{\alpha}{\pi}} \int_0^\infty e^{-\alpha x^2} \frac{\cosh(2r-1)\alpha x}{\cosh \alpha x} dx, \\
 q^{r(1-r)} H(-q^r, q) &= -\sqrt{\frac{4\pi}{\alpha}} \sin(\pi r) q_1^{-\frac{1}{4}} K_1(e^{2\pi i r}, q_1^2) \\
 &\quad + \sqrt{\frac{\alpha}{\pi}} \int_0^\infty e^{-\alpha x^2} \frac{\cosh(2r-1)\alpha x}{\cosh \alpha x} dx,
 \end{aligned}$$

$$\begin{aligned}
 q^{r(1-r)}H(q^r, -q) &= \sqrt{\frac{\pi}{4\alpha}} \cot\left(\frac{\pi r}{2}\right) K_2(e^{\pi ir}, -q_1) \\
 &\quad + \sqrt{\frac{\alpha}{\pi}} \int_0^\infty e^{-\alpha x^2} \frac{\sinh(2r-1)\alpha x}{\sinh \alpha x} dx, \\
 q^{r(1-r)}H(-q^r, -q) &= \sqrt{\frac{\pi}{4\alpha}} \tan\left(\frac{\pi r}{2}\right) K_2(-e^{\pi ir}, -q_1) \\
 &\quad - \sqrt{\frac{\alpha}{\pi}} \int_0^\infty e^{-\alpha x^2} \frac{\sinh(2r-1)\alpha x}{\sinh \alpha x} dx.
 \end{aligned}$$

Observe that the first two transformation laws for H involve the same Mordell integral. Using (3.1), (3.4), and transformation laws for the above θ -functions one can show that these laws are equivalent. Since $H(-q^r, -q) = H(q^{1-r}, -q)$, the last two transformation laws for H are also equivalent. A complete transformation theory of H is found in [5, Theorem 4.3] proved by the same method of contour integration used to prove (4.2) below. This method extends back to the work of Watson [26].

The analogous transformation laws for $M(q^r, q)$ are [12]

$$\begin{aligned}
 q^{\frac{3}{2}r(1-r)-\frac{1}{24}}M(q^r, q) &= \sqrt{\frac{\pi}{2\alpha}} \csc(\pi r) q_1^{-\frac{1}{6}} N(e^{2\pi ir}, q_1^4) - \sqrt{\frac{3\alpha}{2\pi}} J(r, \alpha), \\
 q^{\frac{3}{2}r(1-r)-\frac{1}{24}}M(-q^r, q) &= -\sqrt{\frac{2\pi}{\alpha}} q_1^{\frac{4}{3}} M(e^{2\pi ir} q_1^2, q_1^4) - \sqrt{\frac{3\alpha}{2\pi}} J_1(r, \alpha), \\
 q^{\frac{3}{2}r(1-r)-\frac{1}{24}}M(q^r, -q) &= \sqrt{\frac{\pi}{4\alpha}} \csc\left(\frac{\pi r}{2}\right) q_1^{-\frac{1}{24}} N(e^{\pi ir}, -q_1) - \sqrt{\frac{3\alpha}{2\pi}} J_2(r, \alpha), \\
 q^{\frac{3}{2}r(1-r)-\frac{1}{24}}M(-q^r, -q) &= \sqrt{\frac{\pi}{4\alpha}} \sec\left(\frac{\pi r}{2}\right) q_1^{-\frac{1}{24}} N(-e^{\pi ir}, -q_1) - \sqrt{\frac{3\alpha}{2\pi}} J_2(1-r, \alpha),
 \end{aligned}$$

where the Mordell integrals J, J_1, J_2 are defined by

$$\begin{aligned}
 J(r, \alpha) &= \int_0^\infty e^{-\frac{3}{2}\alpha x^2} \frac{\cosh(3r-2)\alpha x + \cosh(3r-1)\alpha x}{\cosh \frac{3}{2}\alpha x} dx, \\
 J_1(r, \alpha) &= \int_0^\infty e^{-\frac{3}{2}\alpha x^2} \frac{\sinh(3r-2)\alpha x - \sinh(3r-1)\alpha x}{\sinh \frac{3}{2}\alpha x} dx, \\
 J_2(r, \alpha) &= \int_0^\infty e^{-\frac{3}{2}\alpha x^2} \left(\cosh\left(3r - \frac{7}{2}\right)\alpha x + \cosh\left(3r - \frac{5}{2}\right)\alpha x \right. \\
 &\quad \left. + \cosh\left(3r - \frac{1}{2}\right)\alpha x - \cosh\left(3r + \frac{1}{2}\right)\alpha x \right) / \cosh 3\alpha x dx.
 \end{aligned}$$

A complete transformation theory of M is found in [4, Theorems 2.1, 2.2].

In [13] we deduce the first transformation laws for H and M from the following

transformation law for U_k :

$$\begin{aligned}
 (4.2) \quad q^{\frac{1}{2}kr(1-r)}U_k(q^r, q) &= \frac{4\pi}{\alpha} \sin(\pi r) V_k(e^{2\pi ir}, q_1^4) \\
 &\quad - \sum_{m=1}^{k-1} \theta_1\left(\frac{m\pi}{k}, q_1^{\frac{2}{k}}\right) \int_0^\infty e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr-m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx \\
 &= \frac{4\pi}{\alpha} \sin(\pi r) V_k(e^{2\pi ir}, q_1^4) \\
 &\quad - \sqrt{\frac{k\alpha}{2\pi}} \sum_{m=1}^{k-1} q^{\frac{(k-2m)^2}{8k}} j(q^m, q^k) \int_0^\infty e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr-m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx,
 \end{aligned}$$

where the Jacobi θ -function θ_1 is defined by

$$\theta_1(z; \tau) = \theta_1(z, q) = 2 \sum_{n=0}^\infty (-1)^n q^{\frac{(2n+1)^2}{4}} \sin(2n+1)z.$$

As usual for Jacobi θ -functions $q = e^{\pi i \tau}$.

We will now prove (4.2) by contour integration and the saddle-point method. By analytic continuation, it suffices to prove the identity for real $\alpha > 0$. Put $q = e^{-\alpha}$ and consider the contour integral

$$\begin{aligned}
 I = I_1 + I_2 &= \frac{1}{2\pi i} \int_{-\infty-\epsilon i}^{+\infty-\epsilon i} \frac{\pi}{\sin \pi z} \frac{e^{-\frac{1}{2}k\alpha z(z+1)}}{1 - e^{-\alpha(z+r)}} dz \\
 &\quad + \frac{1}{2\pi i} \int_{+\infty+\epsilon i}^{-\infty+\epsilon i} \frac{\pi}{\sin \pi z} \frac{e^{-\frac{1}{2}k\alpha z(z+1)}}{1 - e^{-\alpha(z+r)}} dz,
 \end{aligned}$$

where $\epsilon > 0$ is sufficiently small. By Cauchy’s residue theorem, I is equal to the sum of the residues of the poles of the integrand inside the contour. Now $\pi / \sin \pi z$ has a simple pole of residue $(-1)^n$ at each integer n and $1/(1 - e^{-\alpha(z+r)})$ has a simple pole of residue $1/\alpha$ at $z = -r$. If ϵ is sufficiently small, there are no other poles inside the contour. Hence

$$(4.3) \quad I = \sum_{n=-\infty}^\infty \frac{(-1)^n q^{\frac{1}{2}kn(n+1)}}{1 - q^{n+r}} + \frac{\pi}{\sin(-\pi r)} \frac{q^{-\frac{1}{2}kr(1-r)}}{\alpha} = U_k(q^r, q) + \frac{\pi q^{-\frac{1}{2}kr(1-r)}}{\alpha \sin(-\pi r)}.$$

We now consider I_2 . In the upper half plane we have

$$\frac{1}{\sin \pi z} = -2i \sum_{n=0}^\infty e^{(2n+1)\pi iz},$$

so

$$I_2 = \sum_{n=0}^\infty \int_{-\infty+\epsilon i}^{+\infty+\epsilon i} \frac{e^{(2n+1)\pi iz - \frac{1}{2}k\alpha z(z+1)}}{1 - e^{-\alpha(z+r)}} dz = \sum_{n=0}^\infty J_n,$$

say. The integrand of J_n has poles in the upper half plane at the points z where $1 - e^{-\alpha(z+r)} = 0$, that is, at the points

$$z_m = -r + \frac{2\pi im}{\alpha}$$

for $m = 1, 2, \dots$. The residue at z_m (multiplied by $2\pi i$) is

$$\begin{aligned} \mu_{n,m} &= 2\pi i \frac{e^{(2n+1)\pi iz_m - \frac{1}{2}k\alpha z_m(z_m+1)}}{\alpha} \\ &= \frac{2\pi i}{\alpha} e^{-(2n+1)\pi ir} q_1^{(2n+1)2m} q^{-\frac{1}{2}kr(1-r)} e^{-k(1-2r)\pi im} q_1^{-2km^2}, \end{aligned}$$

where $q_1 = e^{-\pi^2/\alpha}$. Next, we symmetrize the denominator of the integrand of J_n by using the identity

$$\frac{1}{1-t} = \frac{t^{-\frac{1}{2}k} + t^{-\frac{1}{2}k+1} + t^{-\frac{1}{2}k+2} + \dots + t^{\frac{1}{2}k-1}}{t^{-\frac{1}{2}k} - t^{\frac{1}{2}k}}.$$

Applying this with $t = e^{-\alpha(z+r)}$, we find that the integrand of J_n is

$$\frac{e^{\frac{1}{2}k\alpha(z+r)} + e^{(\frac{1}{2}k-1)\alpha(z+r)} + \dots + e^{(-\frac{1}{2}k+1)\alpha(z+r)}}{e^{\frac{1}{2}k\alpha(z+r)} - e^{-\frac{1}{2}k\alpha(z+r)}} e^{-\frac{1}{2}k\alpha z} e^{(2n+1)\pi iz - \frac{1}{2}k\alpha z^2}.$$

To find the saddle point, we set the derivative of the last factor equal to 0, getting $(2n + 1)\pi i - k\alpha z = 0$ or

$$z = \frac{(2n + 1)\pi i}{k\alpha} = w_n,$$

say. We move the upper contour of J_n up to the horizontal line through w_n , getting J'_n . By the residue theorem,

$$J_n = J'_n + \text{sum of residues of poles of integrand between the two contours.}$$

These poles are the points $z_m = -r + \frac{2\pi im}{\alpha}$ for which $0 < 2m < \frac{2n+1}{k}$, or equivalently, $0 < m \leq \frac{n}{k}$. Hence,

$$J_n = J'_n + \sum_{0 < m \leq \frac{n}{k}} \mu_{n,m}.$$

Summing over n , we obtain

$$I_2 = \sum_{n=0}^{\infty} J'_n + \sum_{m=1}^{\infty} \sum_{n=km}^{\infty} \mu_{n,m}.$$

Now

$$\mu_{n+1,m} = e^{2\pi iz_m} \mu_{n,m} = e^{-2\pi ir} q_1^{4m} \mu_{n,m}.$$

Hence,

$$\begin{aligned} \sum_{m=1}^{\infty} \sum_{n=km}^{\infty} \mu_{n,m} &= \sum_{m=1}^{\infty} \frac{\mu_{km,m}}{1 - e^{-2\pi ir} q_1^{4m}} \\ &= \sum_{m=1}^{\infty} \frac{2\pi i e^{-(2km+1)\pi ir} q_1^{(2km+1)2m} q^{-\frac{1}{2}kr(1-r)} e^{-k(1-2r)\pi im} q_1^{-2km^2}}{\alpha (1 - e^{-2\pi ir} q_1^{4m})} \\ &= \frac{2\pi i}{\alpha} q^{-\frac{1}{2}kr(1-r)} e^{-\pi ir} \sum_{m=1}^{\infty} \frac{(-1)^{km} q_1^{2km^2+2m}}{1 - e^{-2\pi ir} q_1^{4m}}, \end{aligned}$$

so

$$(4.4) \quad I_2 = \frac{2\pi i}{\alpha} q^{-\frac{1}{2}kr(1-r)} e^{-\pi ir} \sum_{m=1}^{\infty} \frac{(-1)^{km} q_1^{2km^2+2m}}{1 - e^{-2\pi ir} q_1^{4m}} + \sum_{n=0}^{\infty} J'_n.$$

Before going on to evaluate the integral J'_n , we remark that the integral I_1 over the lower contour can be handled similarly. This time the expansion

$$\frac{1}{\sin \pi z} = 2i \sum_{n=0}^{\infty} e^{-(2n+1)\pi iz}$$

is employed. Note that this is just the complex conjugate of the expansion used in the upper half plane. Thus $I_1 = \sum_{n=0}^{\infty} K_n$, where

$$K_n = \int_{-\infty - \epsilon i}^{+\infty - \epsilon i} \frac{e^{-(2n+1)\pi iz - \frac{1}{2}k\alpha z(z+1)}}{1 - e^{-\alpha(z+r)}} dz.$$

The lower contour is moved down to the horizontal line through \bar{w}_n , giving

$$K_n = \bar{J}'_n + \sum_{0 < m \leq \frac{n}{2}} \bar{\mu}_{n,m}.$$

The sum here is just the complex conjugate of the one evaluated above, so from (4.4) it follows that

$$(4.5) \quad I_1 = -\frac{2\pi i}{\alpha} q^{-\frac{1}{2}kr(1-r)} e^{\pi ir} \sum_{m=1}^{\infty} \frac{(-1)^{km} q_1^{2km^2+2m}}{1 - e^{2\pi ir} q_1^{4m}} + \sum_{n=0}^{\infty} \bar{J}'_n.$$

Adding (4.4) and (4.5), we obtain

$$\begin{aligned} I &= I_1 + I_2 \\ &= \frac{2\pi i}{\alpha} q^{-\frac{1}{2}kr(1-r)} \sum_{m=1}^{\infty} (-1)^{km} q_1^{2km^2+2m} \left[\frac{e^{-\pi ir}}{1 - e^{-2\pi ir} q_1^{4m}} - \frac{e^{\pi ir}}{1 - e^{2\pi ir} q_1^{4m}} \right] \\ &\quad + \sum_{n=0}^{\infty} (J'_n + \bar{J}'_n) \\ &= \frac{4\pi}{\alpha} q^{-\frac{1}{2}kr(1-r)} \sin(\pi r) \sum_{m=1}^{\infty} \frac{(-1)^{km} q_1^{2km^2+2m} (1 + q_1^{4m})}{(1 - e^{2\pi ir} q_1^{4m})(1 - e^{-2\pi ir} q_1^{4m})} + \sum_{n=0}^{\infty} (J'_n + \bar{J}'_n). \end{aligned}$$

It now follows from equation (4.3) that

$$\begin{aligned}
 (4.6) \quad U_k(q^r, q) &= I - \frac{\pi}{\sin(-\pi r)} \frac{q^{-\frac{1}{2}kr(1-r)}}{\alpha} \\
 &= \frac{4\pi \sin(\pi r) q^{-\frac{1}{2}kr(1-r)}}{\alpha} \left[\frac{1}{4 \sin^2(\pi r)} + \sum_{m=1}^{\infty} \frac{(-1)^{km} q_1^{2km^2+2m} (1 + q_1^{4m})}{(1 - e^{2\pi i r} q_1^{4m})(1 - e^{-2\pi i r} q_1^{4m})} \right] \\
 &\quad + \sum_{n=0}^{\infty} (J'_n + \bar{J}'_n) \\
 &= \frac{4\pi \sin(\pi r) q^{-\frac{1}{2}kr(1-r)}}{\alpha} V_k(e^{2\pi i r}, q_1^4) + \sum_{n=0}^{\infty} (J'_n + \bar{J}'_n).
 \end{aligned}$$

We now evaluate $\sum_{n=0}^{\infty} (J'_n + \bar{J}'_n)$. In the integral J'_n put $z = -r + p + x$, where $p = \frac{(2n+1)\pi i}{k\alpha}$ and x is a real variable running from $-\infty$ to ∞ . This gives

$$J'_n = q^{-\frac{1}{2}kr} \int_{-\infty}^{\infty} ABC \, dx,$$

where

$$\begin{aligned}
 A &= e^{(2n+1)\pi i(-r+p+x)}, \\
 B &= \frac{1 + e^{-\alpha(p+x)} + e^{-2\alpha(p+x)} + \dots + e^{(-k+1)\alpha(p+x)}}{e^{\frac{1}{2}k\alpha(p+x)} - e^{-\frac{1}{2}k\alpha(p+x)}}, \\
 C &= e^{-\frac{1}{2}k\alpha(-r+p+x)^2}.
 \end{aligned}$$

Simplifying, we obtain

$$\begin{aligned}
 J'_n &= q^{-\frac{1}{2}kr(1-r)} q_1^{\frac{(2n+1)^2}{2k}} \int_{-\infty}^{\infty} \frac{e^{k\alpha r x - \frac{1}{2}k\alpha x^2}}{2i(-1)^n \cosh \frac{1}{2}k\alpha x} \sum_{m=1}^{k-1} e^{-\frac{m(2n+1)\pi i}{k}} e^{-m\alpha x} \, dx \\
 &\quad + q^{-\frac{1}{2}kr(1-r)} q_1^{\frac{(2n+1)^2}{2k}} \int_{-\infty}^{\infty} \frac{e^{k\alpha r x - \frac{1}{2}k\alpha x^2}}{2i(-1)^n \cosh \frac{1}{2}k\alpha x} \, dx \\
 &= P_n + Q_n,
 \end{aligned}$$

say. Since Q_n is purely imaginary, we have $J'_n + \bar{J}'_n = P_n + \bar{P}_n$. Hence,

$$\begin{aligned}
 J'_n + \bar{J}'_n &= -q^{-\frac{1}{2}kr(1-r)} \sum_{m=1}^{k-1} (-1)^n q_1^{\frac{(2n+1)^2}{2k}} \sin\left(\frac{m(2n+1)\pi}{k}\right) \int_{-\infty}^{\infty} \frac{e^{(kr-m)\alpha x - \frac{1}{2}k\alpha x^2}}{\cosh \frac{1}{2}k\alpha x} \, dx \\
 &= -2q^{-\frac{1}{2}kr(1-r)} \sum_{m=1}^{k-1} (-1)^n q_1^{\frac{(2n+1)^2}{2k}} \sin\left(\frac{m(2n+1)\pi}{k}\right) \\
 &\quad \cdot \int_0^{\infty} e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr-m)\alpha x}{\cosh \frac{1}{2}k\alpha x} \, dx
 \end{aligned}$$

and so

(4.7)

$$\begin{aligned} \sum_{n=0}^{\infty} (J'_n + \bar{J}'_n) &= -q^{-\frac{1}{2}kr(1-r)} \sum_{m=1}^{k-1} \sum_{n=0}^{\infty} 2(-1)^n q_1^{\frac{(2n+1)^2}{2k}} \sin\left(\frac{m(2n+1)\pi}{k}\right) \\ &\quad \cdot \int_0^{\infty} e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr-m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx \\ &= -q^{-\frac{1}{2}kr(1-r)} \sum_{m=1}^{k-1} \theta_1\left(\frac{m\pi}{k}, q_1^{\frac{2}{k}}\right) \int_0^{\infty} e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr-m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx, \end{aligned}$$

where the Jacobi θ -function θ_1 is defined by

$$\theta_1(z; \tau) = \theta_1(z, q) = 2 \sum_{n=0}^{\infty} (-1)^n q^{\frac{(2n+1)^2}{4}} \sin(2n+1)z = ie^{-iz} q^{\frac{1}{4}} j(e^{2iz}, q^2)$$

and satisfies the transformation law

$$-i\sqrt{-i\tau} \exp\left(\frac{iz^2}{\pi\tau}\right) \theta_1(z; \tau) = \theta_1\left(\frac{z}{\tau}; -\frac{1}{\tau}\right).$$

Hence,

$$-i\sqrt{-i\tau} \exp\left(\frac{iz^2}{\pi\tau}\right) \theta_1(z, q) = \theta_1\left(\frac{z}{\tau}, q_1\right).$$

Replacing q by $q^{\frac{k}{2}}$ (so $\tau \rightarrow \frac{k}{2}\tau, q_1 \rightarrow q_1^{\frac{k}{2}}$), we obtain

$$-i\sqrt{\frac{-ik\tau}{2}} \exp\left(\frac{2iz^2}{k\pi\tau}\right) \theta_1(z, q^{\frac{k}{2}}) = \theta_1\left(\frac{2z}{k\tau}, q_1^{\frac{k}{2}}\right).$$

When $z = m\pi\tau/2$, this becomes

$$\begin{aligned} (4.8) \quad \theta_1\left(\frac{m\pi}{k}, q_1^{\frac{k}{2}}\right) &= -i\sqrt{\frac{-ik\tau}{2}} \exp\left(\frac{m^2\pi i\tau}{2k}\right) \theta_1\left(\frac{m\pi\tau}{2}, q^{\frac{k}{2}}\right) \\ &= -i\sqrt{\frac{-ik\tau}{2}} q^{\frac{m^2}{2k}} \theta_1\left(\frac{m\pi\tau}{2}, q^{\frac{k}{2}}\right) \\ &= -i\sqrt{\frac{k\alpha}{2\pi}} q^{\frac{m^2}{2k}} \theta_1\left(\frac{m\pi\tau}{2}, q^{\frac{k}{2}}\right) = \sqrt{\frac{k\alpha}{2\pi}} q^{\frac{(k-2m)^2}{8k}} j(q^m, q^k). \end{aligned}$$

Substituting (4.8) into (4.7) gives

(4.9)

$$\begin{aligned} \sum_{n=0}^{\infty} (J'_n + \bar{J}'_n) &= -q^{-\frac{1}{2}kr(1-r)} \sum_{m=1}^{k-1} \theta_1\left(\frac{m\pi}{k}, q_1^{\frac{2}{k}}\right) \int_0^{\infty} e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr-m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx \\ &= -q^{-\frac{1}{2}kr(1-r)} \sqrt{\frac{k\alpha}{2\pi}} \sum_{m=1}^{k-1} q^{\frac{(k-2m)^2}{8k}} j(q^m, q^k) \\ &\quad \cdot \int_0^{\infty} e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr-m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx. \end{aligned}$$

Finally, by (4.6) and (4.9) we get the transformation law

$$\begin{aligned}
 q^{\frac{1}{2}kr(1-r)}U_k(q^r, q) &= \frac{4\pi}{\alpha} \sin(\pi r) V_k(e^{2\pi ir}, q_1^4) \\
 &\quad - \sum_{m=1}^{k-1} \theta_1\left(\frac{m\pi}{k}, q_1^{\frac{2}{k}}\right) \int_0^\infty e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr - m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx \\
 &= \frac{4\pi}{\alpha} \sin(\pi r) V_k(e^{2\pi ir}, q_1^4) \\
 &\quad - \sqrt{\frac{k\alpha}{2\pi}} \sum_{m=1}^{k-1} q^{\frac{(k-2m)^2}{8k}} j(q^m, q^k) \int_0^\infty e^{-\frac{1}{2}k\alpha x^2} \frac{\cosh(kr - m)\alpha x}{\cosh \frac{1}{2}k\alpha x} dx,
 \end{aligned}$$

which completes the proof of (4.2). ■

5 Mock Theta Conjectures for Functions of Even Order

Hickerson [15,16] proved that Ramanujan’s fifth and seventh order mock θ -functions are related to the function M . The third order mock θ -function $\omega(q)$ is $M(q, q^2)$, and the third order mock θ -function $\psi(q)$ is equal to $qM(q, q^4)$. Ramanujan gave relations, later proved by Watson [26], between $\omega(q)$ or $\psi(q)$ and some of the other third order mock θ -functions. A complete list of relations between all of the third order mock θ -functions and the function M is given in [13].

It turns out that the mock θ -functions of even order are related to the function H . Lists of all of these relations (referred to as mock theta “conjectures” even after their proofs are known) are found in [13]. We will discuss some of these relations.

The second order mock θ -function $B(q)$ is $H(q, q^2)$ (see [20]). The function $V_1(q)$ in [11, 20] is equal to $qH(q, q^4)$.

The sixth order mock θ -functions $\phi(q)$ and $\psi(q)$ defined by

$$\phi(q) = \sum_{n=0}^\infty \frac{(-1)^n q^{n^2} (q; q^2)_n}{(-q; q)_{2n}}, \quad \psi(q) = \sum_{n=0}^\infty \frac{(-1)^n q^{(n+1)^2} (q; q^2)_n}{(-q; q)_{2n+1}}$$

are related to H by

$$\begin{aligned}
 \phi(q^4) &= \frac{(q^2; q^2)_\infty^3 (q^3; q^3)_\infty^2 (q^{12}; q^{12})_\infty^3}{(q)_\infty^2 (q^6; q^6)_\infty^3 (q^8; q^8)_\infty (q^{24}; q^{24})_\infty} - 2qH(q, q^6), \\
 \psi(q^4) &= \frac{q^3 (q^2; q^2)_\infty^2 (q^4; q^4)_\infty (q^{24}; q^{24})_\infty^2}{(q)_\infty (q^3; q^3)_\infty (q^8; q^8)_\infty^2} - q^3H(q^3, q^6).
 \end{aligned}$$

These identities were discovered using transformation laws and computer algebra. Some proofs for these and similar identities for the eighth and tenth order functions below will appear in a forthcoming paper. Relations between other sixth order mock θ -functions and $\phi(q)$ or $\psi(q)$ are found in [2].

Gordon and the author [11] discovered the eighth order mock θ -functions by applying the half-shift method to the θ -functions appearing in the Göllnitz–Gordon identities [9],[10], [24, eq. (36), (34)]. Two of the eighth order mock θ -functions are

$$S_0(q) = \sum_{n=0}^{\infty} \frac{q^{n^2}(-q; q^2)_n}{(-q^2; q^2)_n}, \quad S_1(q) = \sum_{n=0}^{\infty} \frac{q^{n(n+2)}(-q; q^2)_n}{(-q^2; q^2)_n}.$$

They are related to H by

$$S_0(-q^2) = \frac{j(-q, q^2)j(q^6, q^{16})}{j(q^2; q^8)} - 2qH(q, q^8),$$

$$S_1(-q^2) = \frac{j(-q, q^2)j(q^2, q^{16})}{j(q^2, q^8)} - 2qH(q^3, q^8).$$

On page 9 of the lost notebook Ramanujan defined four functions that came to be known as the tenth order mock θ -functions. These functions are

$$\phi(q) = \sum_{n=0}^{\infty} \frac{q^{\frac{1}{2}n(n+1)}}{(q; q^2)_{n+1}}, \quad \psi(q) = \sum_{n=0}^{\infty} \frac{q^{\frac{1}{2}(n+1)(n+2)}}{(q; q^2)_{n+1}},$$

$$X(q) = \sum_{n=0}^{\infty} \frac{(-1)^n q^{n^2}}{(-q; q)_{2n}}, \quad \chi(q) = \sum_{n=0}^{\infty} \frac{(-1)^n q^{(n+1)^2}}{(-q; q)_{2n+1}}.$$

The mock theta “conjectures” of order 10 are

$$\phi(q) = \frac{(q^{10}; q^{10})_{\infty}^2 j(-q^2, q^5)}{(q^5; q^5)_{\infty} j(q^2, q^{10})} + 2qH(q^2, q^5),$$

$$\psi(q) = -\frac{q(q^{10}; q^{10})_{\infty}^2 j(-q, q^5)}{(q^5; q^5)_{\infty} j(q^4, q^{10})} + 2qH(q, q^5),$$

$$X(-q^2) = \frac{(q^4; q^4)_{\infty}^2 (j(-q^2, q^{20})^2 j(q^{12}, q^{40}) + 2q(q^{40}; q^{40})_{\infty}^3)}{(q^2; q^2)_{\infty} (q^{20}; q^{20})_{\infty} (q^{40}; q^{40})_{\infty} j(q^8, q^{40})} - 2qH(q, q^{20}) + 2q^5H(q^9, q^{20}),$$

$$\chi(-q^2) = \frac{q^2(q^4; q^4)_{\infty}^2 (2q(q^{40}; q^{40})_{\infty}^3 - j(-q^6, q^{20})^2 j(q^4, q^{40}))}{(q^2; q^2)_{\infty} (q^{20}; q^{20})_{\infty} (q^{40}; q^{40})_{\infty} j(q^{16}, q^{40})} - 2q^3H(q^3, q^{20}) - 2q^5H(q^7, q^{20}).$$

The first two were stated and proved by Choi [6, pp. 533–534], and the last two were discovered by the author [13] by matching the Mordell integrals in their transformation laws (obtained using computer algebra) with the Mordell integrals in the transformation laws for $H(q^r, q)$. A rigorous proof has yet to be worked out.

6 Concluding Remarks

Unlike that for θ -functions, transformation laws for mock θ -functions are not unique. For example,

$$(6.1) \quad q^{\frac{3}{2}r(1-r)}U_3(q^r, q) = \frac{4\pi}{\alpha} \sin(\pi r) V_3(e^{2\pi ir}, q_1^4) + \text{an integral}$$

$$(6.2) \quad = \frac{-2\pi i}{\alpha} e^{3\pi ir} U_3(e^{2\pi ir}, q_1^4) + \text{some integrals.}$$

Equation (6.1) is (4.2) with $k = 3$. To prove (6.2) we first introduce the Appell function.

The Appell function of level l (not to be confused with the level of a modular form) is defined by (see, for example [30])

$$A_l(u, v; \tau) = e^{l\pi iu} \sum_{n=-\infty}^{\infty} \frac{(-1)^n e^{l\pi i(n^2+n)\tau+2\pi inv}}{-e^{2\pi in\tau+2\pi iu}}, \quad \tau \in \mathcal{H}, \quad v \in \mathbf{C}, \quad u \in \mathbf{C} \setminus (\mathbf{Z}\tau + \mathbf{Z}),$$

and satisfies the transformation law [30]

$$A_l(u, v; \tau) = \frac{e^{\pi i(lu-2v)u/\tau}}{\tau} A_l(u/\tau, v/\tau; -1/\tau) + \text{some integrals.}$$

When $l = 3$ and $v = 0$, this law becomes

$$(6.3) \quad e^{3\pi iu} U_3(e^{2\pi iu}, e^{2\pi i\tau}) = \frac{e^{3\pi iu^2/\tau}}{\tau} e^{3\pi iu/\tau} U_3(e^{2\pi iu/\tau}, e^{-2\pi i/\tau}) + \text{some integrals.}$$

Recall that $q = e^{-\alpha} = e^{\pi i\tau}$ and $q_1 = e^{-\pi^2/\alpha}$. So $\alpha = -\pi i\tau$ and $q_1 = e^{-\pi i/\tau}$. If we put $u = r\tau$, then (6.3) simplifies to

$$q^{3r} U_3(q^{2r}, q^2) = \frac{-\pi i}{\alpha} q^{3r^2} e^{3\pi ir} U_3(e^{2\pi ir}, q_1^2) + \text{some integrals.}$$

Replacing q by $q^{\frac{1}{2}}$ (hence $\alpha \rightarrow \frac{1}{2}\alpha, q_1 \rightarrow q_1^2$), this becomes (6.2).

Comparing (6.1) and (6.2), we cannot conclude that

$$\frac{4\pi}{\alpha} \sin(\pi r) V_3(e^{2\pi ir}, q_1^4) = \frac{-2\pi i}{\alpha} e^{3\pi ir} U_3(e^{2\pi ir}, q_1^4),$$

which is equivalent to

$$2i \sin(\pi r) V_3(e^{2\pi ir}, q_1^4) = e^{3\pi ir} U_3(e^{2\pi ir}, q_1^4).$$

By (3.5) (with $x \rightarrow e^{2\pi ir}, q \rightarrow q_1^4$, then divide by $e^{\pi ir}$) we obtain

$$2i \sin(\pi r) V_3(e^{2\pi ir}, q_1^4) - e^{3\pi ir} U_3(e^{2\pi ir}, q_1^4) = e^{\pi ir} (q_1^4; q_1^4)_{\infty}.$$

Hence $2i \sin(\pi\tau) V_3(e^{2\pi i\tau}, q_1^4)$ and $e^{3\pi i\tau} U_3(e^{2\pi i\tau}, q_1^4)$ differ by a θ -function.

In general, the mock θ -functions on the right-hand sides of transformation laws similar to (6.1) and (6.2) (with the same left side) differ by a θ -function, because they have the same shadow [28].

The presence of a nonzero Mordell integral in a transformation formula for a function $f(q)$ does not always indicate that $f(q)$ is a mock θ -function. We provide an example using Zwegers' μ -function [29] (this function is a normalized level 1 Appell function):

$$\mu(a, b, q) = \mu(u, v; \tau) = ib^{\frac{1}{2}} q^{-\frac{1}{8}} A_1(u, v; \tau) / j(b, q),$$

where $a = e^{2\pi iu}$, $b = e^{2\pi iv}$, and $q = e^{2\pi i\tau}$.

If we put $u = 1/2 + \tau/2$ and $v = 1/2$, then

$$\mu(u, v; \tau) = \mu(-q^{\frac{1}{2}}, -1, q) = q^{\frac{1/8}{2i}}, \quad \mu\left(\frac{u}{\tau}, \frac{v}{\tau}; -\frac{1}{\tau}\right) = \mu(-q_1^{-\frac{1}{2}}, q_1^{-\frac{1}{2}}, q_1) = 0,$$

where $q_1 = e^{-2\pi i/\tau}$. The transformation law for $f(q) = \mu(-q^{\frac{1}{2}}, -1, q)$ becomes

$$q^{1/8} = e^{-\alpha/4} = 2 \int_0^\infty e^{-\alpha x^2} \frac{\cos \alpha x}{\cosh \pi x} dx,$$

where $\alpha = -\pi i\tau$.

Equation (2.7) expresses the function $H(x, q)$ as a normalized level 2 Appell function. In particular,

$$H(x, q) = \frac{U_2(x, q)}{j(q, q^2)} = \frac{\tilde{A}_2(x, -1, q)}{j(q, q^2)},$$

where

$$\tilde{A}_l(a, b, q) = \sum_{n=-\infty}^\infty \frac{(-1)^n q^{\frac{1}{2} \ln(n+1)} b^n}{1 - aq^n}.$$

By the identity

$$\frac{1}{1-x} = \frac{1+x+x^2+\dots+x^{l-1}}{1-x^l},$$

where $x = aq^n$, it is not difficult to show that

$$(6.4) \quad \tilde{A}_l(a, b, q) = \sum_{m=0}^{l-1} a^m \tilde{A}_1(a^l, (-1)^{l-1} b q^m, q^l).$$

Kang [17] used this to prove that

$$(6.5) \quad iaH(a, q) = \frac{\eta^4(2\tau)}{\eta^2(\tau)\vartheta(2u; 2\tau)} + aq^{-\frac{1}{4}} \mu(2u, \tau; 2\tau)$$

and

$$ia^{\frac{3}{2}} q^{-\frac{1}{24}} M(a, q) = \frac{\eta^3(3\tau)}{\eta(\tau)\vartheta(3u; 3\tau)} + aq^{-\frac{1}{6}} \mu(3u, \tau; 3\tau) + a^2 q^{-\frac{2}{3}} \mu(3u, 2\tau; 3\tau),$$

where $a = e^{2\pi iu}$, $q = e^{2\pi i\tau}$, and $\eta(\tau) = q^{\frac{1}{24}}(q)_{\infty}$ is the Dedekind η -function.

The transformation laws for H and μ can be combined to eliminate the Mordell integrals. This resulting transformation law is

$$q^{-\frac{1}{2}r^2} \left(q^{\frac{1}{8}} H(a^{\frac{1}{2}} b^{-\frac{1}{2}} q^{\frac{1}{4}}, q^{\frac{1}{2}}) + i\mu(u, v; \tau) \right) = \frac{1}{\sqrt{-i\tau}} \left(\frac{1}{2} \sec(\pi r) q_1^{-\frac{1}{8}} K(-q_1^{v-u}, q_1) - i\mu\left(\frac{u}{\tau}, \frac{v}{\tau}; -\frac{1}{\tau}\right) \right),$$

where $q = e^{2\pi i\tau}$, $q_1 = e^{-2\pi i/\tau}$, $a = e^{2\pi iu}$, $b = e^{2\pi iv}$, and $r = (u - v)/\tau$. Hence,

$$(6.6) \quad i\mu(u, v; \tau) + q^{\frac{1}{8}} H(a^{\frac{1}{2}} b^{-\frac{1}{2}} q^{\frac{1}{4}}, q^{\frac{1}{2}})$$

and

$$(6.7) \quad i\mu(u, v; \tau) - \frac{a^{\frac{1}{2}} b^{-\frac{1}{2}} q^{-\frac{1}{8}} K(-a/b, q)}{1 + a/b}$$

are Jacobi forms; they behave like θ -functions. A proof that (6.6) vanishes when $u + v = \tau/2$ and (6.7) vanishes when $u + v = 1/2$ is given in [14]. Therefore

$$(6.8) \quad H(a, q) = -iq^{-\frac{1}{4}} \mu(u, \tau - u; 2\tau)$$

and

$$K(a, q) = (a^{\frac{1}{2}} - a^{-\frac{1}{2}}) q^{\frac{1}{8}} \mu\left(\frac{u}{2}, \frac{1-u}{2}; \tau\right) = 2i \sin(\pi u) q^{\frac{1}{8}} \mu\left(\frac{u}{2}, \frac{1-u}{2}; \tau\right),$$

or equivalently,

$$H(x, q) = -iq^{-\frac{1}{4}} \mu(x, q/x, q^2) = \frac{\tilde{A}_1(x, q/x, q^2)}{j(q/x, q^2)}$$

and

$$\frac{K(y, q)}{1 - y} = -y^{-\frac{1}{2}} q^{\frac{1}{8}} \mu(y^{\frac{1}{2}}, -y^{-\frac{1}{2}}, q) = \frac{\tilde{A}_1(y^{\frac{1}{2}}, -y^{-\frac{1}{2}}, q)}{y^{\frac{1}{2}} j(-y^{-\frac{1}{2}}, q)}.$$

Observe that (6.8) removes the θ -quotient from (6.5). This has a nice extension to higher level Appell functions. It follows from (6.4) and (2.5) that

$$(6.9) \quad \begin{aligned} \tilde{A}_l(a, (-1)^{l-1}, q) &= \frac{(q^l; q^l)_{\infty}^3}{j(a^l, q^l)} + \sum_{m=1}^{l-1} a^m \tilde{A}_1(a^l, q^m, q^l) \\ &= \frac{(q^l; q^l)_{\infty}^3}{j(a^l, q^l)} - i \sum_{m=1}^{l-1} a^{m-\frac{1}{2}} q^{\frac{l-m}{8}} j(q^m, q^l) \mu(lu, m\tau; l\tau). \end{aligned}$$

The θ -quotient in (6.9) is removed by the conjectured identity

$$\begin{aligned}\tilde{A}_l(a, (-1)^{l-1}, q) &= \sum_{m=1}^{l-1} \frac{j(q^m, q^l)}{j(q^m/a, q^l)} a^{m-1} \tilde{A}_l(a^{l-1}, q^m/a, q^l) \\ &= -i \sum_{m=1}^{l-1} a^{m-\frac{1}{2}} q^{\frac{l}{8}-\frac{m}{2}} j(q^m, q^l) \mu(lu - u, m\tau - \tau; l\tau)\end{aligned}$$

for $l \geq 2$.

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