

# Steady models of optically thin, magnetically supported two-temperature accretion disks around a black hole

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**Abstract.** We obtained steady solutions of optically thin two-temperature magnetized accretion disks around a black hole. We included relativistic bremsstrahlung cooling, synchrotron cooling and inverse Compton effects and assumed that the disk is threaded by toroidal magnetic field. We found that a magnetic pressure dominated new branch, which we call a ‘low- $\beta$  branch’, appears in the thermal equilibrium curves. The luminosity of the optically thin, magnetically supported disk can exceed 10% of the Eddington luminosity ( $0.1L_{\text{Edd}}$ ).

**Keywords.** Accretion: accretion disks – black holes – state transitions – MHD

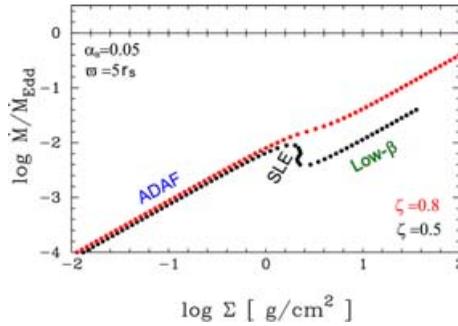
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## 1. Introduction

X-ray hard state (low/hard state) of black hole candidates (BHCs) can be modeled by optically thin advection dominated accretion flows (ADAFs). Abramowicz *et al.* (1995) showed that when the accretion rate  $\dot{M}$  exceeds a critical accretion rate  $\dot{M}_c$ , optically thin solution disappears. Above this critical accretion rate, the hard X-ray emitting optically thin hot accretion disk undergoes a transition to an optically thick cold disk which emit soft X-rays. RXTE observations of galactic BHCs revealed that some BHCs (e.g., XTE J1550-564) stay in X-ray hard state even when  $L > 0.1L_{\text{Edd}}$  (e.g., Gierliński, Newton 2006). This luminosity is higher than the critical luminosity above which thermal instability takes place. Machida *et al.* (2006) carried out global 3D MHD simulations of this cooling instability and showed that magnetically supported disk is created.

## 2. Models

We extended the basic equations for 1D steady, axisymmetric, optically thin black hole accretion flows (e.g., Kato *et al.* 1998) by incorporating the azimuthal magnetic field. General relativistic effects are simulated by using pseudo-Newtonian potential (Paczynski & Wiita 1980). We prescribe the vertically integrated  $\varpi\varphi$ -component of the stress tensor based on the results of local and global 3D MHD simulations (e.g., Machida *et al.* 2006; Hirose *et al.* 2006) as  $T_{\varpi\varphi} = -\alpha_B W$  and we set  $\alpha_B = 0.05$ . Machida *et al.* (2006) showed that  $\int_{-\infty}^{\infty} v_{\varpi} \langle B_{\varphi} \rangle dz \propto \varpi^{-1}$  where  $\langle B_{\varphi} \rangle$  is the azimuthally averaged



**Figure 1.** Local thermal equilibrium curves of optically thin disks at  $\varpi = 5r_s$  when  $M = 10M_\odot$ ,  $\alpha_B = 0.05$  and  $\zeta = 0.5$  (black) and  $\zeta = 0.8$  (red). We obtained a low- $\beta$  branch in addition to the ADAF and SLE branches.  $\dot{M}_{\text{Edd}} = 4\pi GM/(\eta_e \kappa_{\text{es}} c)$  is the Eddington accretion rate where  $\eta_e$  is the energy conversion efficiency taken to be  $\eta_e = 0.1$ .

toroidal field. Here we assume the radial advection rate of the toroidal magnetic field  $\dot{\Phi}$  as

$$\dot{\Phi} = - \int_{-\infty}^{\infty} v_\varpi \langle B_\varphi \rangle dz = -v_\varpi B_0(\varpi) \sqrt{4\pi} H \equiv \dot{\Phi}_{\text{out}} \left( \frac{\varpi}{\varpi_{\text{out}}} \right)^{-\zeta}$$

where  $B_0(\varpi) = 2^{5/4} \pi^{1/4} (\mathcal{R}T/\mu)^{1/2} \Sigma^{1/2} H^{-1/2} \beta^{-1/2}$  is the strength of equatorial toroidal magnetic field and  $\dot{\Phi}_{\text{out}}$  is the radial advection rate of the toroidal magnetic field at the outer boundary.

### 3. Results

We obtained local thermal equilibrium curves of optically thin, two-temperature black hole accretion flows at  $\varpi = 5r_s$  (figure 1). We found a low- $\beta$  branch in which the magnetic heating balances with the radiative cooling. The low- $\beta$  branch exists even when  $\dot{M} \sim 0.5\dot{M}_{\text{Edd}}$  and explains the luminous, X-ray hard state in BHCs observed during the transition from a low/hard state to a high/soft state. Since  $T_e < 10^9\text{K}$  when  $\dot{M} > \dot{M}_{\text{Edd}}$  in low- $\beta$  branch, the power-law X-ray spectrum in such disks becomes steeper and can explain the steep power law spectra observed in very high/steep power law state.

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